


Measurement of the Sensitivity of Two-Particle Correlations in pp Collisions to the Presence of Hard Scatterings

G. Aad *et al.**
(ATLAS Collaboration)

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A key open question in the study of multiparticle production in high-energy pp collisions is the relationship between the “ridge”—i.e., the observed azimuthal correlations between particles in the underlying event that extend over all rapidities—and hard or semihard scattering processes. In particular, it is not known whether jets or their soft fragments are correlated with particles in the underlying event. To address this question, two-particle correlations are measured in pp collisions at $\sqrt{s} = 13$ TeV using data collected by the ATLAS experiment at the LHC, with an integrated luminosity of 15.8 pb^{-1} , in two different configurations. In the first case, charged particles associated with jets are excluded from the correlation analysis, while in the second case, correlations are measured between particles within jets and charged particles from the underlying event. Second-order flow coefficients, v_2 , are presented as a function of event multiplicity and transverse momentum. These measurements show that excluding particles associated with jets does not affect the measured correlations. Moreover, particles associated with jets do not exhibit any significant azimuthal correlations with the underlying event, ruling out hard processes contributing to the ridge.

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In heavy-ion collisions, two-particle correlations (2PC) in relative azimuthal angle with large pseudorapidity [1] separation show distinct long-range correlations [2–12]. These long-range correlations are a simple manifestation of the single-particle anisotropies, v_n , that originate from the hydrodynamic expansion of the quark-gluon plasma produced in these collisions. The v_n are defined by parametrizing the azimuthal distribution of produced particles as

$$\frac{dN}{d\phi} \propto \left(1 + 2 \sum_{n=1}^{\infty} v_n \cos[n(\phi - \Psi_n)] \right), \quad (1)$$

where ϕ is the azimuthal angle of the particle momentum and v_n and Ψ_n are the magnitude and phase of the n th-order anisotropy; see Refs. [4,10] and references therein.

Because of their hydrodynamic origin in nucleus-nucleus ($A + A$) collisions, such long-range correlations were not expected in smaller colliding systems such as proton-nucleus ($p + A$) or proton-proton (pp) collisions, where collective phenomena were not commonly expected to develop. However, measurements by CMS showed the

presence of such long-range correlations, known as the “ridge,” in high-multiplicity pp collisions [13]. Further investigations by ATLAS [9,14,15] have demonstrated that these long-range correlations in pp collisions are produced from single-particle anisotropies similar to those in heavy-ion collisions. These long-range correlations have been interpreted as evidence of collective effects similar to those seen in heavy-ion collisions. However, some authors have proposed that the ridge primarily results from correlated production of partons in the presence of dense gluonic initial states (i.e., the “glasma”) [16–20], implying that much of the correlation structure associated with the ridge should be associated with hard- or semihard scattering processes. Previous measurements [21] have shown that the ridge is unmodified in pp collisions producing a Z boson, but no direct measurement in pp collisions of the correlation between jets or their fragments and the underlying event has yet been performed, while such a correlation has been observed in $p + \text{Pb}$ collisions [22,23].

This Letter presents 2PC measurements in pp collisions at a center-of-mass energy (\sqrt{s}) of 13 TeV, using the ATLAS detector at the LHC. The measurements are performed with two different particle-pair selections. The first case explores correlations between tracks that are not jet constituents, while the second case measures correlations between tracks that are constituents of jets and tracks that are well-separated from jets. Similar measurements in $p + \text{Pb}$ collisions have shown significant nonzero v_2 for low [23] and high [22] transverse momentum (p_T) particles generated in hard processes. Correlations are also measured

*Full author list given at the end of the Letter.

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in events that are explicitly selected by requiring the presence or absence of low- p_T jets. These measurements can address whether or not the presence of jets affects the ridge, and if the particles from jets exhibit azimuthal correlations with particles from the underlying event and therefore contribute to the ridge.

The measurements presented here are performed using the ATLAS [24] inner detector (ID), minimum-bias trigger scintillators, calorimeters, and the trigger and data acquisition systems [25]. The ID records charged-particle trajectories within the pseudorapidity range $|\eta| < 2.5$ using a combination of silicon pixel detectors including the “insertable B-layer” [26,27], silicon microstrip detectors, and a straw-tube transition radiation tracker, all immersed in a 2 T axial magnetic field [1,28]. The ATLAS calorimeter system consists of a liquid argon (LAr) electromagnetic calorimeter covering $|\eta| < 3.2$, a steel-scintillator sampling hadronic calorimeter covering $|\eta| < 1.7$, a LAr hadronic calorimeter covering $1.5 < |\eta| < 3.2$, and two LAr electromagnetic and hadronic forward calorimeters covering $3.2 < |\eta| < 4.9$. The ATLAS trigger system [29] consists of a Level-1 trigger implemented using a combination of dedicated electronics and programmable logic, and a software-based high-level trigger. An extensive software suite [30] is used in data simulation, in the reconstruction and analysis of real and simulated data, in detector operations, and in the trigger and data acquisition systems of the experiment.

The data were collected during Run 2 of the LHC (2015–2018), with an average collision rate per bunch crossing (μ) of less than 3, and an integrated luminosity of 15.8 pb^{-1} . The data used here were recorded using multiple minimum-bias, high-multiplicity, and jet triggers, which are described in Ref. [31]. Additional offline requirements are imposed on the events selected by the triggers. The events are required to have a reconstructed vertex with $|z| < 100 \text{ mm}$. To suppress events with more than one pp collision in the same bunch crossing, events are required to have only one reconstructed vertex. Pileup events where the vertices from multiple collisions are sufficiently close such that they are reconstructed as a single vertex are not removed by the one vertex requirement. However, such merged events typically have a broader distribution for the longitudinal impact parameter of tracks relative to the vertex ($|z_0 \sin(\theta)|$). Such events are reduced by requiring that the standard deviation of $|z_0 \sin(\theta)|$ for all tracks in an event is less than 0.25 mm.

The reconstruction and performance of tracks and primary vertices in the ID are described in Refs. [32–34]. The specific track selection criteria can be found in Ref. [31]. The track reconstruction efficiencies $\epsilon(p_T, \eta)$ are obtained using Monte Carlo (MC) generated events that are passed through a Geant4 [35] simulation [36] of the ATLAS detector and reconstructed using the procedures applied to the data. The efficiency varies between 69% and 87% as a function of η and p_T .

Jets used in this analysis are reconstructed using the anti- k_t algorithm [37] with a radius parameter of 0.4. The inputs to jet reconstruction are “particle flow objects” as detailed in Ref. [38]. Jets are calibrated to the hadronic scale using scale factors obtained from MC simulations specifically derived for low- μ data. Additional *in situ* corrections [39] are applied, which account for differences in the jet response between the MC samples and data. One issue in this analysis is that the modulation in the soft particles in the event [Eq. (1)] biases the jet p_T in a manner that depends on its orientation relative to the Ψ_n . This affects the measurements of the correlations between jet fragments and the underlying event (UE) particles (discussed in detail below). To mitigate this effect, instead of selecting jets based on their p_T , selections are made on the following groomed quantity:

$$p_T^G = \left| \sum_{\text{constituents}} p_T^{>4 \text{ GeV}} \right|, \quad (2)$$

where the sum runs over all the jet constituents with $p_T > 4 \text{ GeV}$, which considerably reduces the number of UE particles within the jet, and makes this bias negligible, as shown in Ref. [31].

In previous ATLAS measurements of 2PCs in $p + \text{Pb}$ [40,41] and pp [14,15,21] collisions, events were quantified by $N_{\text{ch}}^{\text{rec}}$: the total number of reconstructed tracks with $p_T > 0.4 \text{ GeV}$, passing the track selections discussed above. In this analysis, a slight modification is made to ensure that the event activity is not biased by the presence of jets and only reflects the soft multiplicity in the event. The number of constituent tracks in jets with $p_T^G > 15 \text{ GeV}$ is subtracted from the measured multiplicity, and the corrected quantity, $N_{\text{ch}}^{\text{rec,corr}}$, is used to represent the event activity. While counting the constituent tracks of jets, the $p_T > 4 \text{ GeV}$ requirement is not imposed on the tracks. Additionally, this correction is offset by the average number of UE tracks within the jet cone. This offset is estimated by measuring the average number of tracks, as a function of η and ϕ , that are in a $R = 0.4$ cone in events with similar multiplicity and trigger conditions.

In 2PC measurements, the distribution of particle pairs in relative azimuthal angle $\Delta\phi = \phi^a - \phi^b$ are measured. The labels a and b denote the two particles in the pair. In evaluating the correlation functions, the tracks are weighted by the inverse of their reconstruction efficiency, $1/\epsilon(p_T, \eta)$. To suppress short-range correlations, the particles are required to have a pseudorapidity separation of $|\Delta\eta| > 2$. In pp collisions, back-to-back dijets also make a significant contribution to the 2PCs. To remove this contribution, a template-fit method [14,15,21] is employed in which the measured 2PC is described by a fit having two components. The first component accounts for the dijet contribution, $C^{\text{periph}}(\Delta\phi)$, which is measured using low-multiplicity events (called the “peripheral reference”). This analysis

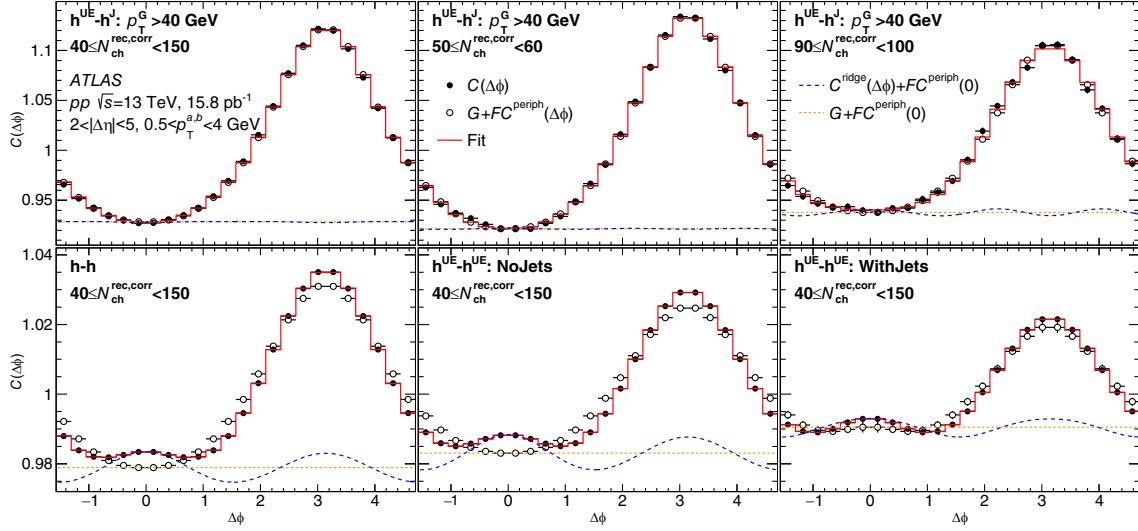


FIG. 1. Template fits to the two-particle correlations in $\Delta\phi$. Events with $10 \leq N_{\text{ch}}^{\text{rec,corr}} < 30$ are used as the peripheral reference. The solid points indicate the measured 2PC, the open circles show the scaled and shifted peripheral reference, and the continuous line shows the fit. The dashed line shows the second-order harmonic component, and the dotted line shows the pedestal of the fit shifted up by $FC^{\text{periph}}(0)$. The top row corresponds to different multiplicity intervals for the $h^{\text{UE}} - h^J$ class. The left, center, and right panels in the bottom row correspond to the $h-h$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ classes, respectively, for the 40–150 multiplicity interval.

uses the $N_{\text{ch}}^{\text{rec,corr}}$ interval of 10–30 to build C^{periph} . The second component accounts for the bulk contribution with a relative harmonic modulation, $C^{\text{ridge}}(\Delta\phi)$. The 2PC can then be described as

$$C(\Delta\phi) = FC^{\text{periph}}(\Delta\phi) + G \left(1 + 2 \sum_{n=2} v_{n,n} \cos(n\Delta\phi) \right) \equiv FC^{\text{periph}}(\Delta\phi) + C^{\text{ridge}}(\Delta\phi), \quad (3)$$

where F and $v_{n,n}$ are fit parameters and G is fixed by the requirement that the integrals of the fit and $C(\Delta\phi)$ are equal. The Fourier moments, $v_{n,n}$, obtained from the template fit quantify the strength of the long-range correlation. It is demonstrated in Refs. [14,15] that the $v_{n,n}$ in pp collisions obtained from Eq. (3) factorize as $v_{n,n}(\phi_T^a, \phi_T^b) = v_n(\phi_T^a) v_n(\phi_T^b)$, where v_n is the single particle anisotropy [Eq. (1)]. Thus, $v_n(\phi_T^b)$ is obtained as $v_n(\phi_T^b) = v_{n,n}(\phi_T^a, \phi_T^b) / \sqrt{v_{n,n}(\phi_T^a, \phi_T^a)}$.

The tracks used in this analysis are categorized as follows: those that are separated from all $\phi_T^G > 15$ GeV jets by at least one unit in η [22] and having $0.5 < p_T < 4$ GeV are considered to be UE tracks (h^{UE}); tracks that are included as particle-flow constituents of jets having $p_T^G > 40$ GeV (called “trigger jets” henceforth) are considered to be jet constituents (h^J). Five classes of correlations are studied in this Letter: (1) standard 2PC [14,15] without applying any rejection of tracks around jets; (2) 2PC where both tracks are h^{UE} —about 14% of $h-h$ 2PC pairs are removed by the above-mentioned rejection; (3) 2PC using events with no jets with $p_T^G > 15$ GeV; (4) 2PC using events with

at least one jet with $p_T^G > 15$ GeV; (5) 2PC performed between h^{UE} and h^J . These five classes are referred to as $h-h$, $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$, and $h^{\text{UE}} - h^J$, respectively, in the text below.

These classes are not mutually exclusive. Specifically, the $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$ and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ classes add up to the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$ class. The $h^{\text{UE}} - h^J$ class has no overlapping particle-pairs with the ones in the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ classes. The $h-h$ class is identical to the measurements performed in the previous ATLAS publications [14,15], and is used as a reference with which other classes are compared.

For the $h^{\text{UE}} - h^J$ case, additional requirements are imposed on the trigger jets to avoid distortions of the 2PC. They must have no other jet with $p_T^G > 15$ GeV within $\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2} = 1$ and they must have a balancing jet with $p_T^G > 15$ GeV and with $|\Delta\phi| > 5\pi/6$. The first requirement removes distortions of the 2PC at smaller $\Delta\phi$ while the second requirement ensures that fragments of the balancing jet are excluded from h^{UE} .

It may happen that some constituents of jets originate in the UE, leading to a contribution of combinatorial pairs in the 2PC. These combinatorial pairs, by construction, have the same correlation as those where both the tracks are from the UE. The contribution of such pairs is removed by the following technique. For each event that contributes to the $h^{\text{UE}} - h^J$ correlation, a separate 2PC is made using another event with similar vertex position and multiplicity. In this event, one track is picked from an η - ϕ region that is

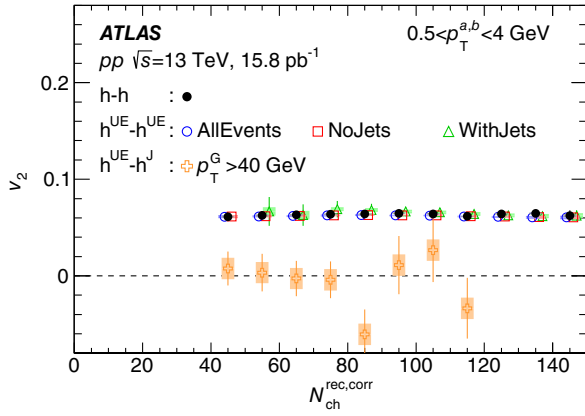


FIG. 2. The multiplicity dependence of v_2 for $2 < |\Delta\eta| < 5$. Events with $10 \leq N_{\text{ch}}^{\text{rec,corr}} < 30$ are used as the peripheral reference. Jets with $p_T^G > 15$ GeV are used to classify the $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$ and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ samples. The data point for the $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ case has a particularly large statistical uncertainty in the 40–50 multiplicity interval and is not shown. The data points for the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ samples are slightly shifted along the x axis for clarity. The error bars and bands correspond to statistical and systematic uncertainties, respectively.

within $R = 0.4$ cone of the jet axis and the other track is picked from the same η range as in the $h^{\text{UE}} - h^J$ event. This combinatorial 2PC is then subtracted from the $h^{\text{UE}} - h^J$ 2PC.

Statistical uncertainties in the measured 2PCs are evaluated using a bootstrapping procedure previously used in Ref. [42]. Systematic uncertainties in the v_2 measurements are estimated by varying different aspects of the analysis. For the template-fit procedure, the $N_{\text{ch}}^{\text{rec,corr}}$ multiplicity range for the peripheral reference selection was varied from the nominal 10–30 to 10–40 and 20–40 [31] and the change in the v_2 values is included as a systematic uncertainty. For the multiplicity dependence, this uncertainty for the v_2 is 0.01 (absolute) for the $h^{\text{UE}} - h^J$ class and is typically within 2% for the other classes. This uncertainty is fully correlated across all multiplicity intervals and is the dominant uncertainty for the $h^{\text{UE}} - h^J$ class. Uncertainties in the tracking efficiency are propagated into the measured v_2 . This uncertainty on the v_2 is less than 0.5%, and is estimated by varying the efficiency up and down within its uncertainties ($\sim \pm 3\%$) [43], and re-evaluating the v_2 . The systematic uncertainty due to nonprimary tracks is estimated by varying the selection criteria for transverse and longitudinal impact parameters, resulting in a 0.5% change in v_2 . The 2PC analyses often use event mixing [4,10] to estimate and correct the 2PCs for the detector's pair acceptance. This correction is quite small, and the full effect of the correction is included as a systematic uncertainty. As discussed previously, the events used in this analysis are required to have the standard deviation of $|z_0 \sin(\theta)|$ for the tracks in an event to be smaller than

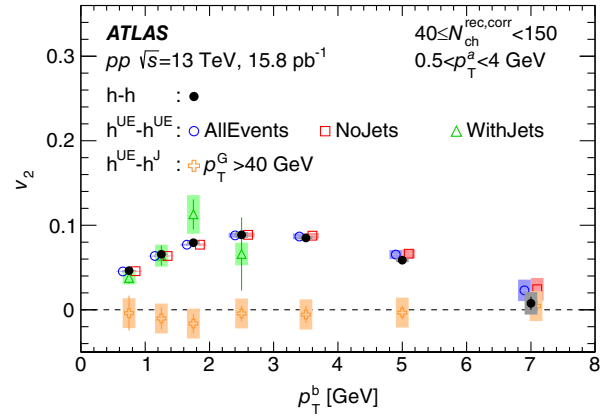


FIG. 3. The p_T^b dependence of the v_2 obtained for the 40–150 multiplicity interval for $2 < |\Delta\eta| < 5$. Events with $10 \leq N_{\text{ch}}^{\text{rec,corr}} < 30$ are used as the peripheral reference. Jets with $p_T^G > 15$ GeV are used to classify the $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$ and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ samples. The data points for the $h-h$ sample are drawn at the nominal values while the data points for the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and the highest p_T^b point of $h^{\text{UE}} - h^J$ samples are shifted slightly for clarity. The error bars and bands correspond to statistical and systematic uncertainties, respectively.

0.25 mm, to reduce pileup. Conservatively, the entire effect of this selection, which varies with multiplicity but is typically within 1%, is taken to be a systematic uncertainty associated with pileup effects.

Figure 1 compares the 2PCs for all classes, except the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$ class. The figure also shows the template fits including the components of the fits. In general, the template fits describe the 2PCs quite well. A near-side ridge is visible for the $h-h$, $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$ cases, while the $C^{\text{periph}}(\Delta\phi)$ appears to describe the full distribution in the $h^{\text{UE}} - h^J$ case.

Figure 2 shows the multiplicity dependence of the v_2 for all five 2PC classes. The v_2 values for the $h-h$ case vary weakly with multiplicity, as previously reported in Refs. [14,15]. The v_2 values in the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ cases, are all consistent with the $h-h$ result. This demonstrates that removing tracks associated with jets does not impact the long-range UE correlations, and nor does the presence (or absence) of jets in an event. Within uncertainties, the v_2 values in the $h^{\text{UE}} - h^J$ case are consistent with zero. The mean v_2 for the $h^{\text{UE}} - h^J$ correlations over the 40–150 multiplicity range is $-0.009 \pm 0.010(\text{statistical}) \pm 0.014(\text{systematic})$. This indicates that particles produced in hard scattering processes (with $p_T^G > 40$ GeV) do not contribute significantly to the long-range correlation observed in pp collisions. Figure 3 shows the p_T dependence of the v_2 . The differential $v_2(p_T)$ values in the $h^{\text{UE}} - h^{\text{UE}}(\text{AllEvents})$, $h^{\text{UE}} - h^{\text{UE}}(\text{NoJets})$, and $h^{\text{UE}} - h^{\text{UE}}(\text{WithJets})$ cases are found to be consistent with the

h - h case. Again, within uncertainties, the $h^{\text{UE}} - h^J v_2$ values are consistent with zero, across the entire measured p_T range. The findings drawn from the p_T dependence are consistent with those from the multiplicity dependence, and similarly demonstrate that the presence or absence of jets has no influence on the flow of the UE and that there are no correlations between jet fragments and the UE. The features of the v_2 values discussed above do not show any systematic variation with the jet selections—for example, the p_T^G thresholds used in the analysis, as discussed in Ref. [31].

In conclusion, this Letter studies long-range 2PCs in pp collisions when rejecting tracks in the vicinity of jets, and the correlations between jet constituent tracks and tracks from the UE. The 2PCs are analyzed using a template-fit procedure, previously developed by ATLAS [15], which extracts second-order Fourier coefficients (v_2) of the anisotropy. These results demonstrate that the magnitude of the v_2 is not affected when removing tracks associated with jets, or by the presence or absence of jets in the event. The v_2 measured with correlations between jet constituents with $p_T < 8$ GeV and UE tracks are consistent with zero within uncertainties. These features are observed both in the v_2 multiplicity and p_T dependence.

The observation that fragments of high- p_T jets in pp collisions do not have measurable long-range azimuthal correlations with the UE and that the production of Z bosons [21] or jets does not significantly influence the long-range correlations between UE particles, suggest a complete “factorization” between hard-scattering processes and the physics responsible for the ridge. Further studies are needed to extend this measurement to higher p_T to compare with previous measurements in $p + \text{Pb}$ collisions [22] where such factorization is broken. This Letter provides important insights into the origin of the long-range correlations observed in pp collisions and offers new fundamental input to theoretical models.

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 H. Abreu¹⁵⁰, Y. Abulaiti¹¹⁷, A. C. Abusleme Hoffman^{137a}, B. S. Acharya^{69a,69b,b}, C. Adam Bourdarios⁴,
 L. Adamczyk^{85a}, L. Adamek¹⁵⁵, S. V. Addepalli²⁶, M. J. Addison¹⁰¹, J. Adelman¹¹⁵, A. Adiguzel^{21c}, T. Aye¹³⁴,
 A. A. Affolder¹³⁶, Y. Afik³⁶, M. N. Agaras¹³, J. Agarwala^{73a,73b}, A. Aggarwal¹⁰⁰, C. Agheorghiesei^{27c},
 A. Ahmad³⁶, F. Ahmadov^{38,c}, W. S. Ahmed¹⁰⁴, S. Ahuja⁹⁵, X. Ai^{62a}, G. Aielli^{76a,76b}, M. Ait Tamlihat^{35e},
 B. Aitbenchikh^{35a}, I. Aizenberg¹⁶⁹, M. Akbiyik¹⁰⁰, T. P. A. Åkesson⁹⁸, A. V. Akimov³⁷, D. Akiyama¹⁶⁸,
 N. N. Akolkar²⁴, K. Al Houry⁴¹, G. L. Alberghi^{23b}, J. Albert¹⁶⁵, P. Albicocco⁵³, G. L. Albouy⁶⁰,
 S. Alderweireldt⁵², M. Aleksa³⁶, I. N. Aleksandrov³⁸, C. Alexa^{27b}, T. Alexopoulos¹⁰, A. Alfonsi¹¹⁴,
 F. Alfonsi^{23b}, M. Algren⁵⁶, M. Alhroob¹²⁰, B. Ali¹³², H. M. J. Ali⁹¹, S. Ali¹⁴⁸, S. W. Alibocus⁹², M. Aliev³⁷,
 G. Alimonti^{71a}, W. Alkahi⁵⁵, C. Allaire⁶⁶, B. M. M. Allbrooke¹⁴⁶, J. F. Allen⁵², C. A. Allendes Flores^{137f},
 P. P. Allport²⁰, A. Aloisio^{72a,72b}, F. Alonso⁹⁰, C. Alpigiani¹³⁸, M. Alvarez Estevez⁹⁹, A. Alvarez Fernandez¹⁰⁰,
 M. G. Alviggi^{72a,72b}, M. Aly¹⁰¹, Y. Amaral Coutinho^{82b}, A. Ambler¹⁰⁴, C. Amelung³⁶, M. Amerl¹⁰¹,
 C. G. Ames¹⁰⁹, D. Amidei¹⁰⁶, S. P. Amor Dos Santos^{130a}, K. R. Amos¹⁶³, V. Ananiev¹²⁵, C. Anastopoulos¹³⁹,
 T. Andeen¹¹, J. K. Anders³⁶, S. Y. Andreadou^{47a,47b}, A. Andreazza^{71a,71b}, S. Angelidakis⁹, A. Angerami^{41,d},
 A. V. Anisenkov³⁷, A. Annovi^{74a}, C. Antel⁵⁶, M. T. Anthony¹³⁹, E. Antipov¹⁴⁵, M. Antonelli⁵³,
 D. J. A. Antrim^{17a}, F. Anulli^{75a}, M. Aoki⁸³, T. Aoki¹⁵³, J. A. Aparisi Pozo¹⁶³, M. A. Aparo¹⁴⁶, L. Aperio Bella⁴⁸,
 C. Appelt¹⁸, A. Apyan²⁶, N. Aranzabal³⁶, C. Arcangeletti⁵³, A. T. H. Arce⁵¹, E. Arena⁹², J-F. Arguin¹⁰⁸,
 S. Argyropoulos⁵⁴, J.-H. Arling⁴⁸, A. J. Armbruster³⁶, O. Arnaez⁴, H. Arnold¹¹⁴, Z. P. Arrubarrena Tame¹⁰⁹,
 G. Artoni^{75a,75b}, H. Asada¹¹¹, K. Asai¹¹⁸, S. Asai¹⁵³, N. A. Asbah⁶¹, J. Assahsah^{35d}, K. Assamagan²⁹,
 R. Astalos^{28a}, S. Atashi¹⁶⁰, R. J. Atkin^{33a}, M. Atkinson¹⁶², N. B. Atlay¹⁸, H. Atmani^{62b}, P. A. Atmasiddha¹⁰⁶,
 K. Augsten¹³², S. Auricchio^{72a,72b}, A. D. Auriol²⁰, V. A. Austrup¹⁰¹, G. Avolio³⁶, K. Axiotis⁵⁶, G. Azuelos^{108,e},
 D. Babal^{28b}, H. Bachacou¹³⁵, K. Bachas^{152,f}, A. Bachiou³⁴, F. Backman^{47a,47b}, A. Badea⁶¹, P. Bagnaia^{75a,75b},
 M. Bahmani¹⁸, A. J. Bailey¹⁶³, V. R. Bailey¹⁶², J. T. Baines¹³⁴, L. Baines⁹⁴, C. Bakalis¹⁰, O. K. Baker¹⁷²,
 E. Bakos¹⁵, D. Bakshi Gupta⁸, R. Balasubramanian¹¹⁴, E. M. Baldwin³⁷, P. Balek^{85a}, E. Ballabene^{23b,23a},
 F. Balli¹³⁵, L. M. Baltes^{63a}, W. K. Balunas³², J. Balz¹⁰⁰, E. Banas⁸⁶, M. Bandieramonte¹²⁹,
 A. Bandyopadhyay²⁴, S. Bansal²⁴, L. Barak¹⁵¹, M. Barakat⁴⁸, E. L. Barberio¹⁰⁵, D. Barberis^{57b,57a},
 M. Barbero¹⁰², G. Barbour⁹⁶, K. N. Barends^{33a}, T. Barillari¹¹⁰, M-S. Barisits³⁶, T. Barklow¹⁴³, P. Baron¹²²,
 D. A. Baron Moreno¹⁰¹, A. Baroncelli^{62a}, G. Barone²⁹, A. J. Barr¹²⁶, J. D. Barr⁹⁶, L. Barranco Navarro^{47a,47b},
 F. Barreiro⁹⁹, J. Barreiro Guimarães da Costa^{14a}, U. Barron¹⁵¹, M. G. Barros Teixeira^{130a}, S. Barsov³⁷,
 F. Bartels^{63a}, R. Bartoldus¹⁴³, A. E. Barton⁹¹, P. Bartos^{28a}, A. Basan¹⁰⁰, M. Baselga⁴⁹, A. Bassalat^{66,g},
 M. J. Basso^{156a}, C. R. Basson¹⁰¹, R. L. Bates⁵⁹, S. Batlamous^{35e}, J. R. Batley³², B. Batool¹⁴¹, M. Battaglia¹³⁶,
 D. Battulga¹⁸, M. Bauce^{75a,75b}, M. Bauer³⁶, P. Bauer²⁴, L. T. Bazzano Hurrell³⁰, J. B. Beacham⁵¹, T. Beau¹²⁷,
 P. H. Beauchemin¹⁵⁸, F. Becherer⁵⁴, P. Bechtel²⁴, H. P. Beck^{19,h}, K. Becker¹⁶⁷, A. J. Beddall^{21d},
 V. A. Bednyakov³⁸, C. P. Bee¹⁴⁵, L. J. Beemster¹⁵, T. A. Beermann³⁶, M. Begalli^{82d}, M. Begel²⁹, A. Behera¹⁴⁵,
 J. K. Behr⁴⁸, J. F. Beirer⁵⁵, F. Beisiegel²⁴, M. Belfkir¹⁵⁹, G. Bella¹⁵¹, L. Bellagamba^{23b}, A. Bellerive³⁴,
 P. Bellos²⁰, K. Beloborodov³⁷, N. L. Belyaev³⁷, D. Benckroun^{35a}, F. Bendebba^{35a}, Y. Benhammou¹⁵¹,
 M. Benoit²⁹, J. R. Bensinger²⁶, S. Bentvelsen¹¹⁴, L. Beresford⁴⁸, M. Beretta⁵³, E. Bergeas Kuutmann¹⁶¹,
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 T. Berry⁹⁵, P. Berta¹³³, A. Berthold⁵⁰, I. A. Bertram⁹¹, S. Bethke¹¹⁰, A. Betti^{75a,75b}, A. J. Bevan⁹⁴,
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 G. Bianco^{23b,23a}, O. Biebel¹⁰⁹, R. Bielski¹²³, M. Biglietti^{77a}, T. R. V. Billoud¹³², M. Bindi⁵⁵, A. Bingul^{21b},
 C. Bini^{75a,75b}, A. Biondini⁹², C. J. Birch-sykes¹⁰¹, G. A. Bird^{20,134}, M. Birman¹⁶⁹, M. Biros¹³³, T. Bisanz⁴⁹

E. Bisceglie^{43b,43a} D. Biswas¹⁴¹ A. Bitadze¹⁰¹ K. Bjørke¹²⁵ I. Bloch⁴⁸ C. Blocker²⁶ A. Blue⁵⁹
 U. Blumenschein⁹⁴ J. Blumenthal¹⁰⁰ G. J. Bobbink¹¹⁴ V. S. Bobrovnikov³⁷ M. Boehler⁵⁴ B. Boehm¹⁶⁶
 D. Bogavac³⁶ A. G. Bogdanchikov³⁷ C. Bohm^{47a} V. Boisvert⁹⁵ P. Bokan⁴⁸ T. Bold^{85a} M. Bomben⁵
 M. Bona⁹⁴ M. Boonekamp¹³⁵ C. D. Booth⁹⁵ A. G. Borbély⁵⁹ I. S. Bordulev³⁷ H. M. Borecka-Bielska¹⁰⁸
 L. S. Borgna⁹⁶ G. Borissov⁹¹ D. Bortoletto¹²⁶ D. Boscherini^{23b} M. Bosman¹³ J. D. Bossio Sola³⁶
 K. Bouaouda^{35a} N. Bouchhar¹⁶³ J. Boudreau¹²⁹ E. V. Bouhova-Thacker⁹¹ D. Boumediene⁴⁰ R. Bouquet⁵
 A. Boveia¹¹⁹ J. Boyd³⁶ D. Boye²⁹ I. R. Boyko³⁸ J. Bracnik²⁰ N. Brahimi^{62d} G. Brandt¹⁷¹ O. Brandt³²
 F. Braren⁴⁸ B. Brau¹⁰³ J. E. Brau¹²³ R. Brenner¹⁶⁹ L. Brenner¹¹⁴ R. Brenner¹⁶¹ S. Bressler¹⁶⁹ D. Britton⁵⁹
 D. Britzger¹¹⁰ I. Brock²⁴ G. Brooijmans⁴¹ W. K. Brooks^{137f} E. Brost²⁹ L. M. Brown^{165j} L. E. Bruce⁶¹
 T. L. Bruckler¹²⁶ P. A. Bruckman de Renstrom⁸⁶ B. Brüers⁴⁸ D. Bruncko^{28b,a} A. Bruni^{23b} G. Bruni^{23b}
 M. Bruschi^{23b} N. Bruscinò^{75a,75b} T. Buanes¹⁶ Q. Buat¹³⁸ D. Buchin¹¹⁰ A. G. Buckley⁵⁹ M. K. Bugge¹²⁵
 O. Bulekov³⁷ B. A. Bullard¹⁴³ S. Burdin⁹² C. D. Burgard⁴⁹ A. M. Burger⁴⁰ B. Burghgrave⁸
 O. Burlayenko⁵⁴ J. T. P. Burr³² C. D. Burton¹¹ J. C. Burzynski¹⁴² E. L. Busch⁴¹ V. Büscher¹⁰⁰ P. J. Bussey⁵⁹
 J. M. Butler²⁵ C. M. Buttar⁵⁹ J. M. Butterworth⁹⁶ W. Buttinger¹³⁴ C. J. Buxo Vazquez¹⁰⁷ A. R. Buzykaev³⁷
 G. Cabras^{23b} S. Cabrera Urbán¹⁶³ L. Cadamuro⁶⁶ D. Caforio⁵⁸ H. Cai¹²⁹ Y. Cai^{14a,14e} V. M. M. Cairo³⁶
 O. Cakir^{3a} N. Calace³⁶ P. Calafiura^{17a} G. Calderini¹²⁷ P. Calfayan⁶⁸ G. Callea⁵⁹ L. P. Caloba^{82b} D. Calvet⁴⁰
 S. Calvet⁴⁰ T. P. Calvet¹⁰² M. Calvetti^{74a,74b} R. Camacho Toro¹²⁷ S. Camarda³⁶ D. Camarero Munoz²⁶
 P. Camarri^{76a,76b} M. T. Camerlingo^{72a,72b} D. Cameron¹²⁵ C. Camincher¹⁶⁵ M. Campanelli⁹⁶ A. Camplani⁴²
 V. Canale^{72a,72b} A. Canesse¹⁰⁴ M. Cano Bret⁸⁰ J. Cantero¹⁶³ Y. Cao¹⁶² F. Capocasa²⁶ M. Capua^{43b,43a}
 A. Carbone^{71a,71b} R. Cardarelli^{76a} J. C. J. Cardenas⁸ F. Cardillo¹⁶³ T. Carli³⁶ G. Carlino^{72a} J. I. Carlotto¹³
 B. T. Carlson^{129,k} E. M. Carlson^{165,156a} L. Carminati^{71a,71b} A. Carnelli¹³⁵ M. Carnesale^{75a,75b} S. Caron¹¹³
 E. Carquin^{137f} S. Carrá^{71a,71b} G. Carratta^{23b,23a} F. Carrio Argos^{33g} J. W. S. Carter¹⁵⁵ T. M. Carter⁵²
 M. P. Casado^{13,l} M. Caspar⁴⁸ E. G. Castiglia¹⁷² F. L. Castillo⁴ L. Castillo Garcia¹³ V. Castillo Gimenez¹⁶³
 N. F. Castro^{130a,130e} A. Catinaccio³⁶ J. R. Catmore¹²⁵ V. Cavaliere²⁹ N. Cavalli^{23b,23a} V. Cavalinni^{74a,74b}
 Y. C. Cekmecelioglu⁴⁸ E. Celebi^{21a} F. Celli¹²⁶ M. S. Centonze^{70a,70b} K. Cerny¹²² A. S. Cerqueira^{82a}
 A. Cerri¹⁴⁶ L. Cerrito^{76a,76b} F. Cerutti^{17a} B. Cervato¹⁴¹ A. Cervelli^{23b} G. Cesarini⁵³ S. A. Cetin^{21d}
 Z. Chadi^{35a} D. Chakraborty¹¹⁵ M. Chala^{130f} J. Chan¹⁷⁰ W. Y. Chan¹⁵³ J. D. Chapman³² E. Chapon¹³⁵
 B. Chargeishvili^{149b} D. G. Charlton²⁰ T. P. Charman⁹⁴ M. Chatterjee¹⁹ C. Chauhan¹³³ S. Chekanov⁶
 S. V. Chekulaev^{156a} G. A. Chelkov^{38,m} A. Chen¹⁰⁶ B. Chen¹⁵¹ B. Chen¹⁶⁵ H. Chen^{14c} H. Chen²⁹
 J. Chen^{62c} J. Chen¹⁴² M. Chen¹²⁶ S. Chen¹⁵³ S. J. Chen^{14c} X. Chen^{62c} X. Chen^{14b,n} Y. Chen^{62a}
 C. L. Cheng¹⁷⁰ H. C. Cheng^{64a} S. Cheong¹⁴³ A. Cheplakov³⁸ E. Cheremushkina⁴⁸ E. Cherepanova¹¹⁴
 R. Cherkaoui El Moursli^{35e} E. Cheu⁷ K. Cheung⁶⁵ L. Chevalier¹³⁵ V. Chiarella⁵³ G. Chiarelli^{74a}
 N. Chiedde¹⁰² G. Chiodini^{70a} A. S. Chisholm²⁰ A. Chitan^{27b} M. Chitishvili¹⁶³ M. V. Chizhov³⁸ K. Choi¹¹
 A. R. Chomont^{75a,75b} Y. Chou¹⁰³ E. Y. S. Chow¹¹⁴ T. Chowdhury^{33g} K. L. Chu¹⁶⁹ M. C. Chu^{64a} X. Chu^{14a,14e}
 J. Chudoba¹³¹ J. J. Chwastowski⁸⁶ D. Cieri¹¹⁰ K. M. Ciesla^{85a} V. Cindro⁹³ A. Ciocio^{17a} F. Ciotto^{72a,72b}
 Z. H. Citron^{169,o} M. Citterio^{71a} D. A. Ciubotaru^{27b} B. M. Ciungu¹⁵⁵ A. Clark⁵⁶ P. J. Clark⁵²
 J. M. Clavijo Columbie⁴⁸ S. E. Clawson⁴⁸ C. Clement^{47a,47b} J. Clercx⁴⁸ L. Clissa^{23b,23a} Y. Coadou¹⁰²
 M. Cobal^{69a,69c} A. Cocco^{57b} R. F. Coelho Barrue^{130a} R. Coelho Lopes De Sa¹⁰³ S. Coelli^{71a} H. Cohen¹⁵¹
 A. E. C. Coimbra^{71a,71b} B. Cole⁴¹ J. Collot⁶⁰ P. Conde Muiño^{130a,130g} M. P. Connell^{33c} S. H. Connell^{33c}
 I. A. Connelly⁵⁹ E. I. Conroy¹²⁶ F. Conventi^{72a,p} H. G. Cooke²⁰ A. M. Cooper-Sarkar¹²⁶
 A. Cordeiro Oudot Choi¹²⁷ F. Cormier¹⁶⁴ L. D. Corpe⁴⁰ M. Corradi^{75a,75b} F. Corriveau^{104,q}
 A. Cortes-Gonzalez¹⁸ M. J. Costa¹⁶³ F. Costanza⁴ D. Costanzo¹³⁹ B. M. Cote¹¹⁹ G. Cowan⁹⁵ K. Cranmer¹⁷⁰
 D. Cremonini^{23b,23a} S. Crépe-Renaudin⁶⁰ F. Crescioli¹²⁷ M. Cristinziani¹⁴¹ M. Cristoforetti^{78a,78b} V. Croft¹¹⁴
 J. E. Crosby¹²¹ G. Crosetti^{43b,43a} A. Cueto⁹⁹ T. Cuhadar Donszelmann¹⁶⁰ H. Cui^{14a,14e} Z. Cui⁷
 W. R. Cunningham⁵⁹ F. Curcio^{43b,43a} P. Czodrowski³⁶ M. M. Czurylo^{63b} M. J. Da Cunha Sargedas De Sousa^{62a}
 J. V. Da Fonseca Pinto^{82b} C. Da Via¹⁰¹ W. Dabrowski^{85a} T. Dado⁴⁹ S. Dahbi^{33g} T. Dai¹⁰⁶ C. Dallapiccola¹⁰³
 M. Dam⁴² G. D'amen²⁹ V. D'Amico¹⁰⁹ J. Damp¹⁰⁰ J. R. Dandoy¹²⁸ M. F. Daneri³⁰ M. Danninger¹⁴²
 V. Dao³⁶ G. Darbo^{57b} S. Darmora⁶ S. J. Das^{29,i} S. D'Auria^{71a,71b} C. David^{156b} T. Davidek¹³³
 B. Davis-Purcell³⁴ I. Dawson⁹⁴ H. A. Day-hall¹³² K. De⁸ R. De Asmundis^{72a} N. De Biase⁴⁸

S. De Castro^{23b,23a} N. De Groot¹¹³ P. de Jong¹¹⁴ H. De la Torre¹⁰⁷ A. De Maria^{14c} A. De Salvo^{75a}
 U. De Sanctis^{76a,76b} A. De Santo¹⁴⁶ J. B. De Vivie De Regie⁶⁰ D. V. Dedovich³⁸ J. Degens¹¹⁴ A. M. Deiana⁴⁴
 F. Del Corso^{23b,23a} J. Del Peso⁹⁹ F. Del Rio^{63a} F. Deliot¹³⁵ C. M. Delitzsch⁴⁹ M. Della Pietra^{72a,72b}
 D. Della Volpe⁵⁶ A. Dell'Acqua³⁶ L. Dell'Asta^{71a,71b} M. Delmastro⁴ P. A. Delsart⁶⁰ S. Demers¹⁷²
 M. Demichev³⁸ S. P. Denisov³⁷ L. D'Eramo⁴⁰ D. Derendarz⁸⁶ F. Derue¹²⁷ P. Dervan⁹² K. Desch²⁴
 C. Deutsch²⁴ F. A. Di Bello^{57b,57a} A. Di Ciaccio^{76a,76b} L. Di Ciaccio⁴ A. Di Domenico^{75a,75b}
 C. Di Donato^{72a,72b} A. Di Girolamo³⁶ G. Di Gregorio⁵ A. Di Luca^{78a,78b} B. Di Micco^{77a,77b} R. Di Nardo^{77a,77b}
 C. Diaconu¹⁰² F. A. Dias¹¹⁴ T. Dias Do Vale¹⁴² M. A. Diaz^{137a,137b} F. G. Diaz Capriles²⁴ M. Didenko¹⁶³
 E. B. Diehl¹⁰⁶ L. Diehl⁵⁴ S. Díez Cornell⁴⁸ C. Díez Pardos¹⁴¹ C. Dimitriadi^{24,161} A. Dimitrievska^{17a}
 J. Dingfelder²⁴ I-M. Dinu^{27b} S. J. Dittmeier^{63b} F. Dittus³⁶ F. Djama¹⁰² T. Djobava^{149b} J. I. Djuvsland¹⁶
 C. Doglioni^{101,98} J. Dolejsi¹³³ Z. Dolezal¹³³ M. Donadelli^{82c} B. Dong¹⁰⁷ J. Donini⁴⁰ A. D'Onofrio^{77a,77b}
 M. D'Onofrio⁹² J. Dopke¹³⁴ A. Doria^{72a} N. Dos Santos Fernandes^{130a} M. T. Dova⁹⁰ A. T. Doyle⁵⁹
 M. A. Draguet¹²⁶ E. Dreyer¹⁶⁹ I. Drivas-koulouris¹⁰ A. S. Drobac¹⁵⁸ M. Drozdova⁵⁶ D. Du^{62a}
 T. A. du Pree¹¹⁴ F. Dubinin³⁷ M. Dubovsky^{28a} E. Duchovni¹⁶⁹ G. Duckeck¹⁰⁹ O. A. Ducu^{27b} D. Duda⁵²
 A. Dudarev³⁶ E. R. Duden²⁶ M. D'uffizi¹⁰¹ L. Dufлот⁶⁶ M. Dührssen³⁶ C. Dülßen¹⁷¹ A. E. Dumitriu^{27b}
 M. Dunford^{63a} S. Dungs⁴⁹ K. Dunne^{47a,47b} A. Duperrin¹⁰² H. Duran Yildiz^{3a} M. Düren⁵⁸ A. Durglishvili^{149b}
 B. L. Dwyer¹¹⁵ G. I. Dyckes^{17a} M. Dyndal^{85a} S. Dysch¹⁰¹ B. S. Dziedzic⁸⁶ Z. O. Earnshaw¹⁴⁶
 G. H. Eberwein¹²⁶ B. Eckerova^{28a} S. Eggebrecht⁵⁵ M. G. Eggleston⁵¹ E. Egidio Purcino De Souza¹²⁷
 L. F. Ehrke⁵⁶ G. Eigen¹⁶ K. Einsweiler^{17a} T. Ekelof¹⁶¹ P. A. Ekman⁹⁸ S. El Farkh^{35b} Y. El Ghazali^{35b}
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 N. Ellis³⁶ J. Elmsheuser²⁹ M. Elsing³⁶ D. Emelianov¹³⁴ Y. Enari¹⁵³ I. Ene^{17a} S. Epari¹³ J. Erdmann⁴⁹
 P. A. Erland⁸⁶ M. Errenst¹⁷¹ M. Escalier⁶⁶ C. Escobar¹⁶³ E. Etzion¹⁵¹ G. Evans^{130a} H. Evans⁶⁸
 L. S. Evans⁹⁵ M. O. Evans¹⁴⁶ A. Ezhilov³⁷ S. Ezzarqtouni^{35a} F. Fabbri⁵⁹ L. Fabbri^{23b,23a} G. Facini⁹⁶
 V. Fadeyev¹³⁶ R. M. Fakhrutdinov³⁷ S. Falciano^{75a} L. F. Falda Ulhoa Coelho³⁶ P. J. Falke²⁴ J. Faltova¹³³
 C. Fan¹⁶² Y. Fan^{14a} Y. Fang^{14a,14e} M. Fanti^{71a,71b} M. Faraj^{69a,69b} Z. Farazpay⁹⁷ A. Farbin⁸ A. Farilla^{77a}
 T. Farooque¹⁰⁷ S. M. Farrington⁵² F. Fassi^{35e} D. Fassouliotis⁹ M. Fauci Giannelli^{76a,76b} W. J. Fawcett³²
 L. Fayard⁶⁶ P. Federic¹³³ P. Federicova¹³¹ O. L. Fedin^{37,m} G. Fedotov³⁷ M. Feickert¹⁷⁰ L. Feligioni¹⁰²
 D. E. Fellers¹²³ C. Feng^{62b} M. Feng^{14b} Z. Feng¹¹⁴ M. J. Fenton¹⁶⁰ A. B. Fenyuk³⁷ L. Ferencz⁴⁸
 R. A. M. Ferguson⁹¹ S. I. Fernandez Luengo^{137f} M. J. V. Fernoux¹⁰² J. Ferrando⁴⁸ A. Ferrari¹⁶¹ P. Ferrari^{114,113}
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 M. C. N. Fiolhais^{130a,130c,r} L. Fiorini¹⁶³ W. C. Fisher¹⁰⁷ T. Fitschen¹⁰¹ P. M. Fitzhugh¹³⁵ I. Fleck¹⁴¹
 P. Fleischmann¹⁰⁶ T. Flick¹⁷¹ L. Flores¹²⁸ M. Flores^{33d,s} L. R. Flores Castillo^{64a} L. Flores Sanz De Acedo³⁶
 F. M. Follega^{78a,78b} N. Fomin¹⁶ J. H. Foo¹⁵⁵ B. C. Forland⁶⁸ A. Formica¹³⁵ A. C. Forti¹⁰¹ E. Fortin³⁶
 A. W. Fortman⁶¹ M. G. Foti^{17a} L. Fountas^{9,t} D. Fournier⁶⁶ H. Fox⁹¹ P. Francavilla^{74a,74b} S. Francescato⁶¹
 S. Franchellucci⁵⁶ M. Franchini^{23b,23a} S. Franchino^{63a} D. Francis³⁶ L. Franco¹¹³ L. Franconi⁴⁸ M. Franklin⁶¹
 G. Frattari²⁶ A. C. Freegard⁹⁴ W. S. Freund^{82b} Y. Y. Frid¹⁵¹ N. Fritzsche⁵⁰ A. Froch⁵⁴ D. Froidevaux³⁶
 J. A. Frost¹²⁶ Y. Fu^{62a} M. Fujimoto¹¹⁸ E. Fullana Torregrosa^{163,a} K. Y. Fung^{64a} E. Furtado De Simas Filho^{82b}
 M. Furukawa¹⁵³ J. Fuster¹⁶³ A. Gabrielli^{23b,23a} A. Gabrielli¹⁵⁵ P. Gadow⁴⁸ G. Gagliardi^{57b,57a}
 L. G. Gagnon^{17a} E. J. Gallas¹²⁶ B. J. Gallop¹³⁴ K. K. Gan¹¹⁹ S. Ganguly¹⁵³ J. Gao^{62a} Y. Gao⁵²
 F. M. Garay Walls^{137a,137b} B. Garcia^{29,i} C. García¹⁶³ A. Garcia Alonso¹¹⁴ A. G. Garcia Caffaro¹⁷²
 J. E. García Navarro¹⁶³ M. Garcia-Sciveres^{17a} G. L. Gardner¹²⁸ R. W. Gardner³⁹ N. Garelli¹⁵⁸ D. Garg⁸⁰
 R. B. Garg^{143,u} J. M. Gargan⁵² C. A. Garner¹⁵⁵ S. J. Gasiorowski¹³⁸ P. Gaspar^{82b} G. Gaudio^{73a} V. Gautam¹³
 P. Gauzzi^{75a,75b} I. L. Gavrilenko³⁷ A. Gavriluk³⁷ C. Gay¹⁶⁴ G. Gaycken⁴⁸ E. N. Gazis¹⁰ A. A. Geanta^{27b}
 C. M. Gee¹³⁶ C. Gemme^{57b} M. H. Genest⁶⁰ S. Gentile^{75a,75b} S. George⁹⁵ W. F. George²⁰ T. Geralis⁴⁶
 P. Gessinger-Befurt³⁶ M. E. Geyik¹⁷¹ M. Ghneimat¹⁴¹ K. Ghorbanian⁹⁴ A. Ghosal¹⁴¹ A. Ghosh¹⁶⁰
 A. Ghosh⁷ B. Giacobbe^{23b} S. Giagu^{75a,75b} P. Giannetti^{74a} A. Giannini^{62a} S. M. Gibson⁹⁵ M. Gignac¹³⁶
 D. T. Gil^{85b} A. K. Gilbert^{85a} B. J. Gilbert⁴¹ D. Gillberg³⁴ G. Gilles¹¹⁴ N. E. K. Gillwald⁴⁸ L. Ginabat¹²⁷
 D. M. Gingrich^{2,e} M. P. Giordani^{69a,69c} P. F. Giraud¹³⁵ G. Giugliarelli^{69a,69c} D. Giugni^{71a} F. Giuli³⁶
 I. Gkialas^{9,t} L. K. Gladilin³⁷ C. Glasman⁹⁹ G. R. Gledhill¹²³ M. Glisic¹²³ I. Gnesi^{43b,v} Y. Go^{29,i}

M. Goblirsch-Kolb³⁶ B. Gocke⁴⁹ D. Godin,¹⁰⁸ B. Gokturk^{21a} S. Goldfarb¹⁰⁵ T. Golling⁵⁶ M. G. D. Gololo,^{33g}
D. Golubkov³⁷ J. P. Gombas¹⁰⁷ A. Gomes^{130a,130b} G. Gomes Da Silva¹⁴¹ A. J. Gomez Delegido¹⁶³
R. Gonçalo^{130a,130c} G. Gonella¹²³ L. Gonella²⁰ A. Gongadze³⁸ F. Gonnella²⁰ J. L. Gonski⁴¹
R. Y. González Andana⁵² S. González de la Hoz¹⁶³ S. Gonzalez Fernandez¹³ R. Gonzalez Lopez⁹²
C. Gonzalez Renteria^{17a} R. Gonzalez Suarez¹⁶¹ S. Gonzalez-Sevilla⁵⁶ G. R. Gonzalvo Rodriguez¹⁶³
L. Goossens³⁶ P. A. Gorbounov³⁷ B. Gorini³⁶ E. Gorini^{70a,70b} A. Gorišek⁹³ T. C. Gosart¹²⁸ A. T. Goshaw⁵¹
M. I. Gostkin³⁸ S. Goswami¹²¹ C. A. Gottardo³⁶ M. Goughri^{35b} V. Goumarre⁴⁸ A. G. Goussiou¹³⁸
N. Govender^{33c} I. Grabowska-Bold^{85a} K. Graham³⁴ E. Gramstad¹²⁵ S. Grancagnolo^{70a,70b} M. Grandi¹⁴⁶
V. Gratchev,^{37,a} P. M. Gravila^{27f} F. G. Gravili^{70a,70b} H. M. Gray^{17a} M. Greco^{70a,70b} C. Grefe²⁴ I. M. Gregor⁴⁸
P. Grenier¹⁴³ C. Grieco¹³ A. A. Grillo¹³⁶ K. Grimm³¹ S. Grinstein^{13,w} J.-F. Grivaz⁶⁶ E. Gross¹⁶⁹
J. Grosse-Knetter⁵⁵ C. Grud,¹⁰⁶ J. C. Grundy¹²⁶ L. Guan¹⁰⁶ W. Guan¹⁷⁰ C. Gubbels¹⁶⁴
J. G. R. Guerrero Rojas¹⁶³ G. Guerrieri^{69a,69b} F. Guescini¹¹⁰ R. Gugel¹⁰⁰ J. A. M. Guhit¹⁰⁶ A. Guida¹⁸
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H. K. Hadavand⁸ A. Hadeef¹⁰⁰ S. Hadzic¹¹⁰ J. J. Hahn¹⁴¹ E. H. Haines⁹⁶ M. Haleem¹⁶⁶ J. Haley¹²¹
J. J. Hall¹³⁹ G. D. Hallewell¹⁰² L. Halser¹⁹ K. Hamano¹⁶⁵ H. Hamdaoui^{35e} M. Hamer²⁴ G. N. Hamity⁵²
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P. H. Hansen⁴² K. Hara¹⁵⁷ D. Harada⁵⁶ T. Harenberg¹⁷¹ S. Harkusha³⁷ M. L. Harris¹⁰³ Y. T. Harris¹²⁶
J. Harrison¹³ N. M. Harrison¹¹⁹ P. F. Harrison¹⁶⁷ N. M. Hartman¹⁴³ N. M. Hartmann¹⁰⁹ Y. Hasegawa¹⁴⁰
A. Hasib⁵² S. Haug¹⁹ R. Hauser¹⁰⁷ C. M. Hawkes²⁰ R. J. Hawkins³⁶ Y. Hayashi¹⁵³ S. Hayashida¹¹¹
D. Hayden¹⁰⁷ C. Hayes¹⁰⁶ R. L. Hayes¹¹⁴ C. P. Hays¹²⁶ J. M. Hays⁹⁴ H. S. Hayward⁹² F. He^{62a}
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C. Heidegger⁵⁴ K. K. Heidegger⁵⁴ W. D. Heidorn⁸¹ J. Heilman³⁴ S. Heim⁴⁸ T. Heim^{17a} J. G. Heinlein¹²⁸
J. J. Heinrich¹²³ L. Heinrich^{110,x} J. Hejbal¹³¹ L. Helary⁴⁸ A. Held¹⁷⁰ S. Hellesund¹⁶ C. M. Helling¹⁶⁴
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M. E. Hesping¹⁰⁰ N. P. Hesse^{156a} H. Hibi⁸⁴ S. J. Hillier²⁰ J. R. Hinds¹⁰⁷ F. Hinterkeuser²⁴ M. Hirose¹²⁴
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M. C. Hodgkinson¹³⁹ B. H. Hodgkinson³² A. Hoecker³⁶ J. Hofer⁴⁸ T. Holm²⁴ M. Holzbock¹¹⁰
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Y. Hori¹¹¹ S. Hou¹⁴⁸ A. S. Howard⁹³ J. Howarth⁵⁹ J. Hoya⁶ M. Hrabovsky¹²² A. Hrynevich⁴⁸
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M. Huhtinen³⁶ S. K. Huiberts¹⁶ R. Hulsken¹⁰⁴ N. Huseynov^{12,m} J. Huston¹⁰⁷ J. Huth⁶¹ R. Hyneman¹⁴³
G. Iacobucci⁵⁶ G. Iakovidis²⁹ I. Ibragimov¹⁴¹ L. Iconomidou-Fayard⁶⁶ P. Iengo^{72a,72b} R. Iguchi¹⁵³
T. Iizawa⁸³ Y. Ikegami⁸³ N. Ilic¹⁵⁵ H. Imam^{35a} M. Ince Lezki⁵⁶ T. Ingebretsen Carlson^{47a,47b}
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K. A. Johns⁷ J. W. Johnson¹³⁶ D. M. Jones³² E. Jones⁴⁸ P. Jones³² R. W. L. Jones⁹¹ T. J. Jones⁹²
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S. Kabana^{137e} A. Kaczmarek⁸⁶ M. Kado¹¹⁰ H. Kagan¹¹⁹ M. Kagan¹⁴³ A. Kahn⁴¹ A. Kahn¹²⁸ C. Kahra¹⁰⁰
T. Kaji¹⁶⁸ E. Kajomovitz¹⁵⁰ N. Kakati¹⁶⁹ I. Kalaitzidou⁵⁴ C. W. Kalderon²⁹ A. Kamenshchikov¹⁵⁵
S. Kanayama¹⁵⁴ N. J. Kang¹³⁶ D. Kar^{33g} K. Karava¹²⁶ M. J. Kareem^{156b} E. Karentzos⁵⁴ I. Karkanas¹⁵²

O. Karkout¹¹⁴ S. N. Karpov³⁸ Z. M. Karpova³⁸ V. Kartvelishvili⁹¹ A. N. Karyukhin³⁷ E. Kasimi¹⁵²
 J. Katzy⁴⁸ S. Kaur³⁴ K. Kawade¹⁴⁰ T. Kawamoto¹³⁵ E. F. Kay³⁶ F. I. Kaya¹⁵⁸ S. Kazakos¹⁰⁷
 V. F. Kazanin³⁷ Y. Ke¹⁴⁵ J. M. Keaveney^{33a} R. Keeler¹⁶⁵ G. V. Kehris⁶¹ J. S. Keller³⁴ A. S. Kelly⁹⁶
 J. J. Kempster¹⁴⁶ K. E. Kennedy⁴¹ P. D. Kennedy¹⁰⁰ O. Kepka¹³¹ B. P. Kerridge¹⁶⁷ S. Kersten¹⁷¹
 B. P. Kerševan⁹³ S. Keshri⁶⁶ L. Keszeghova^{28a} S. Ketabchi Haghight¹⁵⁵ M. Khandoga¹²⁷ A. Khanov¹²¹
 A. G. Kharlamov³⁷ T. Kharlamova³⁷ E. E. Khoda¹³⁸ T. J. Khoo¹⁸ G. Khorauli¹⁶⁶ J. Khubua^{149b}
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 C. Kirfel²⁴ F. Kirfel²⁴ J. Kirk¹³⁴ A. E. Kiryunin¹¹⁰ C. Kitsaki¹⁰ O. Kivernyk²⁴ M. Klassen^{63a} C. Klein³⁴
 L. Klein¹⁶⁶ M. H. Klein¹⁰⁶ M. Klein⁹² S. B. Klein⁵⁶ U. Klein⁹² P. Klimek³⁶ A. Klimentov²⁹
 T. Klioutchnikova³⁶ P. Kluit¹¹⁴ S. Kluth¹¹⁰ E. Kneringer⁷⁹ T. M. Knight¹⁵⁵ A. Knue⁵⁴ R. Kobayashi⁸⁷
 S. F. Koch¹²⁶ M. Kocian¹⁴³ P. Kodyš¹³³ D. M. Koeck¹²³ P. T. Koenig²⁴ T. Koffas³⁴ M. Kolb¹³⁵
 I. Koletsou⁴ T. Komarek¹²² K. Köneke⁵⁴ A. X. Y. Kong¹ T. Kono¹¹⁸ N. Konstantinidis⁹⁶ B. Konya⁹⁸
 R. Kopeliansky⁶⁸ S. Koperny^{85a} K. Korcyl⁸⁶ K. Kordas^{152,aa} G. Koren¹⁵¹ A. Korn⁹⁶ S. Korn⁵⁵
 I. Korolkov¹³ N. Korotkova³⁷ B. Kortman¹¹⁴ O. Kortner¹¹⁰ S. Kortner¹¹⁰ W. H. Kostecka¹¹⁵
 V. V. Kostyukhin¹⁴¹ A. Kotsokechagia¹³⁵ A. Kotwal⁵¹ A. Koulouris³⁶ A. Kourkoumeli-Charalampidi^{73a,73b}
 C. Kourkoumelis⁹ E. Kourlitis⁶ O. Kovanda¹⁴⁶ R. Kowalewski¹⁶⁵ W. Kozanecki¹³⁵ A. S. Kozhin³⁷
 V. A. Kramarenko³⁷ G. Kramberger⁹³ P. Kramer¹⁰⁰ M. W. Krasny¹²⁷ A. Krasznahorkay³⁶ J. W. Kraus¹⁷¹
 J. A. Kremer¹⁰⁰ T. Kresse⁵⁰ J. Kretschmar⁹² K. Kreul¹⁸ P. Krieger¹⁵⁵ S. Krishnamurthy¹⁰³ M. Krivos¹³³
 K. Krizka²⁰ K. Kroeninger⁴⁹ H. Kroha¹¹⁰ J. Kroll¹³¹ J. Kroll¹²⁸ K. S. Krowpman¹⁰⁷ U. Kruchonak³⁸
 H. Krüger²⁴ N. Krumnack⁸¹ M. C. Kruse⁵¹ J. A. Krzysiak⁸⁶ O. Kuchinskaia³⁷ S. Kuday^{3a} S. Kuehn³⁶
 R. Kuesters⁵⁴ T. Kuhl⁴⁸ V. Kukhtin³⁸ Y. Kulchitsky^{37,m} S. Kuleshov^{137d,137b} M. Kumar^{33g} N. Kumari¹⁰²
 A. Kupco¹³¹ T. Kupfer⁴⁹ A. Kupich³⁷ O. Kuprash⁵⁴ H. Kurashige⁸⁴ L. L. Kurchaninov^{156a} O. Kurdysh⁶⁶
 Y. A. Kurochkin³⁷ A. Kurova³⁷ M. Kuze¹⁵⁴ A. K. Kvam¹⁰³ J. Kvita¹²² T. Kwan¹⁰⁴ N. G. Kyriacou¹⁰⁶
 L. A. O. Laatu¹⁰² C. Lacasta¹⁶³ F. Lacava^{75a,75b} H. Lacker¹⁸ D. Lacour¹²⁷ N. N. Lad⁹⁶ E. Ladygin³⁸
 B. Laforge¹²⁷ T. Lagouri^{137e} S. Lai⁵⁵ I. K. Lakomic^{85a} N. Lalloue⁶⁰ J. E. Lambert^{165,j} S. Lammers⁶⁸
 W. Lampl⁷ C. Lampoudis^{152,aa} A. N. Lancaster¹¹⁵ E. Lançon²⁹ U. Landgraf⁵⁴ M. P. J. Landon⁹⁴
 V. S. Lang⁵⁴ R. J. Langenberg¹⁰³ O. K. B. Langrekken¹²⁵ A. J. Lankford¹⁶⁰ F. Lanni³⁶ K. Lantzsche²⁴
 A. Lanza^{73a} A. Lapertosa^{57b,57a} J. F. Laporte¹³⁵ T. Lari^{71a} F. Lasagni Manghi^{23b} M. Lassnig³⁶ V. Latonova¹³¹
 A. Laudrain¹⁰⁰ A. Laurier¹⁵⁰ S. D. Lawlor⁹⁵ Z. Lawrence¹⁰¹ M. Lazzaroni^{71a,71b} B. Le¹⁰¹
 E. M. Le Boulicaut⁵¹ B. Leban⁹³ A. Lebedev⁸¹ M. LeBlanc³⁶ F. Ledroit-Guillon⁶⁰ A. C. A. Lee⁹⁶ S. C. Lee¹⁴⁸
 S. Lee^{47a,47b} T. F. Lee⁹² L. L. Leeuw^{33c} H. P. Lefebvre⁹⁵ M. Lefebvre¹⁶⁵ C. Leggett^{17a} G. Lehmann Miotto³⁶
 M. Leigh⁵⁶ W. A. Leight¹⁰³ W. Leinonen¹¹³ A. Leisos^{152,bb} M. A. L. Leite^{82c} C. E. Leitgeb⁴⁸ R. Leitner¹³³
 K. J. C. Leney⁴⁴ T. Lenz²⁴ S. Leone^{74a} C. Leonidopoulos⁵² A. Leopold¹⁴⁴ C. Leroy¹⁰⁸ R. Les¹⁰⁷
 C. G. Lester³² M. Levchenko³⁷ J. Levêque⁴ D. Levin¹⁰⁶ L. J. Levinson¹⁶⁹ M. P. Lewicki⁸⁶ D. J. Lewis⁴
 A. Li⁵ B. Li^{62b} C. Li^{62a} C-Q. Li^{62c} H. Li^{62a} H. Li^{62b} H. Li^{14c} H. Li^{62b} K. Li¹³⁸ L. Li^{62c} M. Li^{14a,14e}
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 M. Liberatore⁴⁸ B. Liberti^{76a} K. Lie^{64c} J. Lieber Marin^{82b} H. Lien⁶⁸ K. Lin¹⁰⁷ R. E. Lindley⁷
 J. H. Lindon² A. Linss⁴⁸ E. Lipeles¹²⁸ A. Lipniacka¹⁶ A. Lister¹⁶⁴ J. D. Little⁴ B. Liu^{14a} B. X. Liu¹⁴²
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 X. Liu^{62a} Y. Liu^{14d,14e} Y. L. Liu¹⁰⁶ Y. W. Liu^{62a} J. Llorente Merino¹⁴² S. L. Lloyd⁹⁴ E. M. Lobodzinska⁴⁸
 P. Loch⁷ S. Loffredo^{76a,76b} T. Lohse¹⁸ K. Lohwasser¹³⁹ E. Loiacono⁴⁸ M. Lokajicek^{131,a} J. D. Lomas²⁰
 J. D. Long¹⁶² I. Longarini¹⁶⁰ L. Longo^{70a,70b} R. Longo¹⁶² I. Lopez Paz⁶⁷ A. Lopez Solis⁴⁸ J. Lorenz¹⁰⁹
 N. Lorenzo Martinez⁴ A. M. Lory¹⁰⁹ O. Loseva³⁷ X. Lou^{47a,47b} X. Lou^{14a,14e} A. Lounis⁶⁶ J. Love⁶
 P. A. Love⁹¹ G. Lu^{14a,14e} M. Lu⁸⁰ S. Lu¹²⁸ Y. J. Lu⁶⁵ H. J. Lubatti¹³⁸ C. Luci^{75a,75b} F. L. Lucio Alves^{14c}
 A. Lucotte⁶⁰ F. Luehring⁶⁸ I. Luise¹⁴⁵ O. Lukianchuk⁶⁶ O. Lundberg¹⁴⁴ B. Lund-Jensen¹⁴⁴ N. A. Luongo¹²³
 M. S. Lutz¹⁵¹ D. Lynn²⁹ H. Lyons⁹² R. Lysak¹³¹ E. Lytken⁹⁸ V. Lyubushkin³⁸ T. Lyubushkina³⁸
 M. M. Lyukova¹⁴⁵ H. Ma²⁹ K. Ma^{62a} L. L. Ma^{62b} Y. Ma¹²¹ D. M. Mac Donell¹⁶⁵ G. Maccarrone⁵³
 J. C. MacDonald¹⁰⁰ R. Madar⁴⁰ W. F. Mader⁵⁰ J. Maeda⁸⁴ T. Maeno²⁹ M. Maerker⁵⁰ H. Maguire¹³⁹
 V. Maiboroda¹³⁵ A. Maio^{130a,130b,130d} K. Maj^{85a} O. Majersky⁴⁸ S. Majewski¹²³ N. Makovec⁶⁶

V. Maksimovic¹⁵ B. Malaescu¹²⁷ Pa. Malecki⁸⁶ V. P. Maleev³⁷ F. Malek⁶⁰ M. Mali⁹³ D. Malito^{95,ee}
 U. Mallik⁸⁰ S. Maltezos¹⁰ S. Malyukov³⁸ J. Mamuzic¹³ G. Mancini⁵³ G. Manco^{73a,73b} J. P. Mandalia⁹⁴
 I. Mandić⁹³ L. Manhaes de Andrade Filho^{82a} I. M. Maniatis¹⁶⁹ J. Manjarres Ramos^{102,ff} D. C. Mankad¹⁶⁹
 A. Mann¹⁰⁹ B. Mansoulie¹³⁵ S. Manzoni³⁶ A. Marantis^{152,bb} G. Marchiori⁵ M. Marcisovsky¹³¹
 C. Marcon^{71a,71b} M. Marinescu²⁰ M. Marjanovic¹²⁰ E. J. Marshall⁹¹ Z. Marshall^{17a} S. Marti-Garcia¹⁶³
 T. A. Martin¹⁶⁷ V. J. Martin⁵² B. Martin dit Latour¹⁶ L. Martinelli^{75a,75b} M. Martinez^{13,w} P. Martinez Agullo¹⁶³
 V. I. Martinez Outschoorn¹⁰³ P. Martinez Suarez¹³ S. Martin-Haugh¹³⁴ V. S. Martoiu^{27b} A. C. Martyniuk⁹⁶
 A. Marzin³⁶ D. Mascione^{78a,78b} L. Masetti¹⁰⁰ T. Mashimo¹⁵³ J. Masik¹⁰¹ A. L. Maslennikov³⁷ L. Massa^{23b}
 P. Massarotti^{72a,72b} P. Mastrandrea^{74a,74b} A. Mastroberardino^{43b,43a} T. Masubuchi¹⁵³ T. Mathisen¹⁶¹
 J. Matousek¹³³ N. Matsuzawa¹⁵³ J. Maurer^{27b} B. Maček⁹³ D. A. Maximov³⁷ R. Mazini¹⁴⁸ I. Maznas¹⁵²
 M. Mazza¹⁰⁷ S. M. Mazza¹³⁶ E. Mazzeo^{71a,71b} C. Mc Ginn^{29,i} J. P. Mc Gowan¹⁰⁴ S. P. Mc Kee¹⁰⁶
 E. F. McDonald¹⁰⁵ A. E. McDougall¹¹⁴ J. A. Mcfayden¹⁴⁶ R. P. McGovern¹²⁸ G. Mchedlidze^{149b}
 R. P. Mckenzie^{33g} T. C. Mclachlan⁴⁸ D. J. Mclaughlin⁹⁶ K. D. McLean¹⁶⁵ S. J. McMahon¹³⁴
 P. C. McNamara¹⁰⁵ C. M. Mepartland⁹² R. A. McPherson^{165,q} S. Mehlhase¹⁰⁹ A. Mehta⁹² D. Melini¹⁵⁰
 B. R. Mellado Garcia^{33g} A. H. Melo⁵⁵ F. Meloni⁴⁸ A. M. Mendes Jacques Da Costa¹⁰¹ H. Y. Meng¹⁵⁵
 L. Meng⁹¹ S. Menke¹¹⁰ M. Mentink³⁶ E. Meoni^{43b,43a} C. Merlassino¹²⁶ L. Merola^{72a,72b} C. Meroni^{71a}
 G. Merz¹⁰⁶ O. Meshkov³⁷ J. Metcalfe⁶ A. S. Mete⁶ C. Meyer⁶⁸ J-P. Meyer¹³⁵ R. P. Middleton¹³⁴
 L. Mijović⁵² G. Mikenberg¹⁶⁹ M. Mikesstikova¹³¹ M. Mikuž⁹³ H. Mildner¹⁰⁰ A. Milic³⁶ C. D. Milke⁴⁴
 D. W. Miller³⁹ L. S. Miller³⁴ A. Milov¹⁶⁹ D. A. Milstead^{47a,47b} T. Min^{14c} A. A. Minaenko³⁷ I. A. Minashvili^{149b}
 L. Mince⁵⁹ A. I. Mincer¹¹⁷ B. Mindur^{85a} M. Mineev³⁸ Y. Mino⁸⁷ L. M. Mir¹³ M. Miralles Lopez¹⁶³
 M. Mironova^{17a} A. Mishima¹⁵³ M. C. Missio¹¹³ T. Mitani¹⁶⁸ A. Mitra¹⁶⁷ V. A. Mitsou¹⁶³ O. Miu¹⁵⁵
 P. S. Miyagawa⁹⁴ Y. Miyazaki⁸⁹ A. Mizukami⁸³ T. Mkrtychyan^{63a} M. Mlinarevic⁹⁶ T. Mlinarevic⁹⁶
 M. Mlynarikova³⁶ S. Mobius¹⁹ K. Mochizuki¹⁰⁸ P. Moder⁴⁸ P. Mogg¹⁰⁹ A. F. Mohammed^{14a,14e}
 S. Mohapatra⁴¹ G. Mokgatitwane^{33g} L. Moleri¹⁶⁹ B. Mondal¹⁴¹ S. Mondal¹³² G. Monig¹⁴⁶ K. Mönig⁴⁸
 E. Monnier¹⁰² L. Monsonis Romero¹⁶³ J. Montejo Berlingen^{13,83} M. Montella¹¹⁹ F. Montekali^{77a,77b}
 F. Monticelli⁹⁰ S. Monzani^{69a,69c} N. Morange⁶⁶ A. L. Moreira De Carvalho^{130a} M. Moreno Llacer¹⁶³
 C. Moreno Martinez⁵⁶ P. Morettini^{57b} S. Morgenstern³⁶ M. Morii⁶¹ M. Morinaga¹⁵³ A. K. Morley³⁶
 F. Morodei^{75a,75b} L. Morvaj³⁶ P. Moschovakos³⁶ B. Moser³⁶ M. Mosidze^{149b} T. Moskalets⁵⁴ P. Moskvitina¹¹³
 J. Moss^{31,gg} E. J. W. Moyse¹⁰³ O. Mtintsilana^{33g} S. Muanza¹⁰² J. Mueller¹²⁹ D. Muenstermann⁹¹
 R. Müller¹⁹ G. A. Mullier¹⁶¹ A. J. Mullin³² J. J. Mullin¹²⁸ D. P. Mungo¹⁵⁵ D. Munoz Perez¹⁶³
 F. J. Munoz Sanchez¹⁰¹ M. Murin¹⁰¹ W. J. Murray^{167,134} A. Murrone^{71a,71b} J. M. Muse¹²⁰ M. Muškinja^{17a}
 C. Mwewa²⁹ A. G. Myagkov^{37,m} A. J. Myers⁸ A. A. Myers¹²⁹ G. Myers⁶⁸ M. Myska¹³² B. P. Nachman^{17a}
 O. Nackenhorst⁴⁹ A. Nag⁵⁰ K. Nagai¹²⁶ K. Nagano⁸³ J. L. Nagle^{29,i} E. Nagy¹⁰² A. M. Nairz³⁶
 Y. Nakahama⁸³ K. Nakamura⁸³ K. Nakkalil⁵ H. Nanjo¹²⁴ R. Narayan⁴⁴ E. A. Narayanan¹¹² I. Naryshkin³⁷
 M. Naseri³⁴ S. Nasri¹⁵⁹ C. Nass²⁴ G. Navarro^{22a} J. Navarro-Gonzalez¹⁶³ R. Nayak¹⁵¹ A. Nayaz¹⁸
 P. Y. Nechaeva³⁷ F. Nechansky⁴⁸ L. Nedic¹²⁶ T. J. Neep²⁰ A. Negri^{73a,73b} M. Negrini^{23b} C. Nellist¹¹⁴
 C. Nelson¹⁰⁴ K. Nelson¹⁰⁶ S. Nemecek¹³¹ M. Nessi^{36,hh} M. S. Neubauer¹⁶² F. Neuhaus¹⁰⁰ J. Neundorf⁴⁸
 R. Newhouse¹⁶⁴ P. R. Newman²⁰ C. W. Ng¹²⁹ Y. W. Y. Ng⁴⁸ B. Ngair^{35e} H. D. N. Nguyen¹⁰⁸
 R. B. Nickerson¹²⁶ R. Nicolaidou¹³⁵ J. Nielsen¹³⁶ M. Niemeyer⁵⁵ J. Niermann^{55,36} N. Nikiforou³⁶
 V. Nikolaenko^{37,m} I. Nikolic-Audit¹²⁷ K. Nikolopoulos²⁰ P. Nilsson²⁹ I. Ninca⁴⁸ H. R. Nindhito⁵⁶
 G. Ninio¹⁵¹ A. Nisati^{75a} N. Nishu² R. Nisius¹¹⁰ J-E. Nitschke⁵⁰ E. K. Nkadimeng^{33g} S. J. Noacco Rosende⁹⁰
 T. Nobe¹⁵³ D. L. Noel³² T. Nommensen¹⁴⁷ M. B. Norfolk¹³⁹ R. R. B. Norisam⁹⁶ B. J. Norman³⁴ J. Novak⁹³
 T. Novak⁴⁸ L. Novotny¹³² R. Novotny¹¹² L. Nozka¹²² K. Ntekas¹⁶⁰ N. M. J. Nunes De Moura Junior^{82b}
 E. Nurse⁹⁶ J. Ocariz¹²⁷ A. Ochi⁸⁴ I. Ochoa^{130a} S. Oerdek¹⁶¹ J. T. Offermann³⁹ A. Ogrodnik¹³³ A. Oh¹⁰¹
 C. C. Ohm¹⁴⁴ H. Oide⁸³ R. Oishi¹⁵³ M. L. Ojeda⁴⁸ Y. Okazaki⁸⁷ M. W. O'Keefe⁹² Y. Okumura¹⁵³
 L. F. Oleiro Seabra^{130a} S. A. Olivares Pino^{137d} D. Oliveira Damazio²⁹ D. Oliveira Goncalves^{82a} J. L. Oliver¹⁶⁰
 M. J. R. Olsson¹⁶⁰ A. Olszewski⁸⁶ Ö. O. Öncel⁵⁴ D. C. O'Neil¹⁴² A. P. O'Neill¹⁹ A. Onofre^{130a,130e}
 P. U. E. Onyisi¹¹ M. J. Oreglia³⁹ G. E. Orellana⁹⁰ D. Orestano^{77a,77b} N. Orlando¹³ R. S. Orr¹⁵⁵ V. O'Shea⁵⁹
 L. M. Osojnak¹²⁸ R. Ospanov^{62a} G. Otero y Garzon³⁰ H. Otono⁸⁹ P. S. Ott^{63a} G. J. Ottino^{17a} M. Ouchrif^{35d}

J. Ouellette²⁹ F. Ould-Saada¹²⁵ M. Owen⁵⁹ R. E. Owen¹³⁴ K. Y. Oyulmaz^{21a} V. E. Ozcan^{21a} N. Ozturk⁸
S. Ozturk^{21d} H. A. Pacey³² A. Pacheco Pages¹³ C. Padilla Aranda¹³ G. Padovano^{75a,75b} S. Pagan Griso^{17a}
G. Palacino⁶⁸ A. Palazzo^{70a,70b} S. Palestini³⁶ J. Pan¹⁷² T. Pan^{64a} D. K. Panchal¹¹ C. E. Pandini¹¹⁴
J. G. Panduro Vazquez⁹⁵ H. Pang^{14b} P. Pani⁴⁸ G. Panizzo^{69a,69c} L. Paolozzi⁵⁶ C. Papadatos¹⁰⁸ S. Parajuli⁴⁴
A. Paramonov⁶ C. Paraskevopoulos¹⁰ D. Paredes Hernandez^{64b} T. H. Park¹⁵⁵ M. A. Parker³² F. Parodi^{57b,57a}
E. W. Parrish¹¹⁵ V. A. Parrish⁵² J. A. Parsons⁴¹ U. Parzefall⁵⁴ B. Pascual Dias¹⁰⁸ L. Pascual Dominguez¹⁵¹
F. Pasquali¹¹⁴ E. Pasqualucci^{75a} S. Passaggio^{57b} F. Pastore⁹⁵ P. Pasuwan^{47a,47b} P. Patel⁸⁶ U. M. Patel⁵¹
J. R. Pater¹⁰¹ T. Pauly³⁶ J. Pearkes¹⁴³ M. Pedersen¹²⁵ R. Pedro^{130a} S. V. Peleganchuk³⁷ O. Penc³⁶
E. A. Pender⁵² H. Peng^{62a} K. E. Pensi¹⁰⁹ M. Penzin³⁷ B. S. Peralva^{82d} A. P. Pereira Peixoto⁶⁰
L. Pereira Sanchez^{47a,47b} D. V. Perepelitsa^{29,i} E. Perez Codina^{156a} M. Perganti¹⁰ L. Perini^{71a,71b,a}
H. Pernegger³⁶ A. Perrevoort¹¹³ O. Perrin⁴⁰ K. Peters⁴⁸ R. F. Y. Peters¹⁰¹ B. A. Petersen³⁶ T. C. Petersen⁴²
E. Petit¹⁰² V. Petousis¹³² C. Petridou^{152,aa} A. Petrukhin¹⁴¹ M. Pettee^{17a} N. E. Pettersson³⁶ A. Petukhov³⁷
K. Petukhova¹³³ A. Peyaud¹³⁵ R. Pezoa^{137f} L. Pezzotti³⁶ G. Pezzullo¹⁷² T. M. Pham¹⁷⁰ T. Pham¹⁰⁵
P. W. Phillips¹³⁴ G. Piacquadio¹⁴⁵ E. Pianori^{17a} F. Piazza^{71a,71b} R. Piegai³⁰ D. Pietreanu^{27b}
A. D. Pilkington¹⁰¹ M. Pinamonti^{69a,69c} J. L. Pinfold² B. C. Pinheiro Pereira^{130a} A. E. Pinto Pinoargote¹³⁵
K. M. Piper¹⁴⁶ A. Pirttikoski⁵⁶ C. Pitman Donaldson⁹⁶ D. A. Pizzi³⁴ L. Pizzimento^{76a,76b} A. Pizzini¹¹⁴
M.-A. Pleier²⁹ V. Plesanovs⁵⁴ V. Pleskot¹³³ E. Plotnikova³⁸ G. Poddar⁴ R. Poettgen⁹⁸ L. Poggioli¹²⁷
I. Pokharel⁵⁵ S. Polacek¹³³ G. Polesello^{73a} A. Poley^{142,156a} R. Polifka¹³² A. Polini^{23b} C. S. Pollard¹⁶⁷
Z. B. Pollock¹¹⁹ V. Polychronakos²⁹ E. Pompa Pacchi^{75a,75b} D. Ponomarenko¹¹³ L. Pontecorvo³⁶ S. Popa^{27a}
G. A. Popeneciu^{27d} A. Poreba³⁶ D. M. Portillo Quintero^{156a} S. Pospisil¹³² M. A. Postill¹³⁹ P. Postolache^{27c}
K. Potamianos¹⁶⁷ P. P. Potepa^{85a} I. N. Potrap³⁸ C. J. Potter³² H. Potti¹ T. Poulsen⁴⁸ J. Poveda¹⁶³
M. E. Pozo Astigarraga³⁶ A. Prades Ibanez¹⁶³ J. Pretel⁵⁴ D. Price¹⁰¹ M. Primavera^{70a} M. A. Principe Martin⁹⁹
R. Privara¹²² T. Procter⁵⁹ M. L. Proffitt¹³⁸ N. Proklova¹²⁸ K. Prokofiev^{64c} G. Proto¹¹⁰ S. Protopopescu²⁹
J. Proudfoot⁶ M. Przybycien^{85a} W. W. Przygoda^{85b} J. E. Puddefoot¹³⁹ D. Pudzha³⁷ D. Pyatiizbyantseva³⁷
J. Qian¹⁰⁶ D. Qichen¹⁰¹ Y. Qin¹⁰¹ T. Qiu⁵² A. Quadt⁵⁵ M. Queitsch-Maitland¹⁰¹ G. Quetant⁵⁶
G. Rabanal Bolanos⁶¹ D. Rafanoharana⁵⁴ F. Ragusa^{71a,71b} J. L. Rainbolt³⁹ J. A. Raine⁵⁶ S. Rajagopalan²⁹
E. Ramakoti³⁷ K. Ran^{48,14e} N. P. Rapheeha^{33g} H. Rasheed^{27b} V. Raskina¹²⁷ D. F. Rassloff^{63a} S. Rave¹⁰⁰
B. Ravina⁵⁵ I. Ravinovich¹⁶⁹ M. Raymond³⁶ A. L. Read¹²⁵ N. P. Readioff¹³⁹ D. M. Rebuzzi^{73a,73b}
G. Redlinger²⁹ A. S. Reed¹¹⁰ K. Reeves²⁶ J. A. Reidelsturz^{171,ii} D. Reikher¹⁵¹ A. Rej¹⁴¹ C. Rembser³⁶
A. Renardi⁴⁸ M. Renda^{27b} M. B. Rendel¹¹⁰ F. Renner⁴⁸ A. G. Rennie⁵⁹ S. Resconi^{71a} M. Ressegotti^{57b,57a}
S. Rettie³⁶ J. G. Reyes Rivera¹⁰⁷ B. Reynolds¹¹⁹ E. Reynolds^{17a} O. L. Rezanova³⁷ P. Reznicek¹³³ N. Ribaric⁹¹
E. Ricci^{78a,78b} R. Richter¹¹⁰ S. Richter^{47a,47b} E. Richter-Was^{85b} M. Ridel¹²⁷ S. Ridouani^{35d} P. Rieck¹¹⁷
P. Riedler³⁶ M. Rijssenbeek¹⁴⁵ A. Rimoldi^{73a,73b} M. Rimoldi⁴⁸ L. Rinaldi^{23b,23a} T. T. Rinn²⁹
M. P. Rinnagel¹⁰⁹ G. Ripellino¹⁶¹ I. Riu¹³ P. Rivadeneira⁴⁸ J. C. Rivera Vergara¹⁶⁵ F. Rizatdinova¹²¹
E. Rizvi⁹⁴ B. A. Roberts¹⁶⁷ B. R. Roberts^{17a} S. H. Robertson^{104,q} M. Robin⁴⁸ D. Robinson³²
C. M. Robles Gajardo^{137f} M. Robles Manzano¹⁰⁰ A. Robson⁵⁹ A. Rocchi^{76a,76b} C. Roda^{74a,74b}
S. Rodriguez Bosca^{63a} Y. Rodriguez Garcia^{22a} A. Rodriguez Rodriguez⁵⁴ A. M. Rodriguez Vera^{156b} S. Roe³⁶
J. T. Roemer¹⁶⁰ A. R. Roepe-Gier¹³⁶ J. Roggel¹⁷¹ O. Røhne¹²⁵ R. A. Rojas¹⁰³ C. P. A. Roland⁶⁸ J. Roloff²⁹
A. Romaniouk³⁷ E. Romano^{73a,73b} M. Romano^{23b} A. C. Romero Hernandez¹⁶² N. Rompotis⁹² L. Roos¹²⁷
S. Rosati^{75a} B. J. Rosser³⁹ E. Rossi¹²⁶ E. Rossi^{72a,72b} L. P. Rossi^{57b} L. Rossini⁴⁸ R. Rosten¹¹⁹
M. Rotaru^{27b} B. Rottler⁵⁴ C. Rougier^{102,ff} D. Rousseau⁶⁶ D. Rousso³² A. Roy¹⁶² S. Roy-Garand¹⁵⁵
A. Rozanov¹⁰² Y. Rozen¹⁵⁰ X. Ruan^{33g} A. Rubio Jimenez¹⁶³ A. J. Ruby⁹² V. H. Ruelas Rivera¹⁸
T. A. Ruggeri¹ A. Ruggiero¹²⁶ A. Ruiz-Martinez¹⁶³ A. Rummeler³⁶ Z. Rurikova⁵⁴ N. A. Rusakovich³⁸
H. L. Russell¹⁶⁵ G. Russo^{75a,75b} J. P. Rutherford⁷ S. Rutherford Colmenares³² K. Rybacki⁹¹ M. Rybar¹³³
E. B. Rye¹²⁵ A. Ryzhov⁴⁴ J. A. Sabater Iglesias⁵⁶ P. Sabatini¹⁶³ L. Sabetta^{75a,75b} H. F.-W. Sadrozinski¹³⁶
F. Safai Tehrani^{75a} B. Safarzadeh Samani¹⁴⁶ M. Safdari¹⁴³ S. Saha¹⁶⁵ M. Sahinsoy¹¹⁰ M. Saimpert¹³⁵
M. Saito¹⁵³ T. Saito¹⁵³ D. Salamani³⁶ A. Salmikov¹⁴³ J. Salt¹⁶³ A. Salvador Salas¹³ D. Salvatore^{43b,43a}
F. Salvatore¹⁴⁶ A. Salzburger³⁶ D. Sammel⁵⁴ D. Sampsonidis^{152,aa} D. Sampsonidou¹²³ J. Sánchez¹⁶³
A. Sanchez Pineda⁴ V. Sanchez Sebastian¹⁶³ H. Sandaker¹²⁵ C. O. Sander⁴⁸ J. A. Sandesara¹⁰³ M. Sandhoff¹⁷¹

C. Sandoval^{22b}, D. P. C. Sankey¹³⁴, T. Sano⁸⁷, A. Sansoni⁵³, L. Santi^{75a,75b}, C. Santoni⁴⁰, H. Santos^{130a,130b}, S. N. Santpur^{17a}, A. Santra¹⁶⁹, K. A. Saoucha¹³⁹, J. G. Saraiva^{130a,130d}, J. Sardain⁷, O. Sasaki⁸³, K. Sato¹⁵⁷, C. Sauer^{63b}, F. Sauerburger⁵⁴, E. Sauvan⁴, P. Savard^{155,e}, R. Sawada¹⁵³, C. Sawyer¹³⁴, L. Sawyer⁹⁷, I. Sayago Galvan¹⁶³, C. Sbarra^{23b}, A. Sbrizzi^{23b,23a}, T. Scanlon⁹⁶, J. Schaarschmidt¹³⁸, P. Schacht¹¹⁰, D. Schaefer³⁹, U. Schäfer¹⁰⁰, A. C. Schaffer^{66,44}, D. Schaile¹⁰⁹, R. D. Schamberger¹⁴⁵, C. Scharf¹⁸, M. M. Schefer¹⁹, V. A. Schegelsky³⁷, D. Scheirich¹³³, F. Schenck¹⁸, M. Schernau¹⁶⁰, C. Scheulen⁵⁵, C. Schiavi^{57b,57a}, E. J. Schioppa^{70a,70b}, M. Schioppa^{43b,43a}, B. Schlag^{143,u}, K. E. Schleicher⁵⁴, S. Schlenker³⁶, J. Schmeing¹⁷¹, M. A. Schmidt¹⁷¹, K. Schmieden¹⁰⁰, C. Schmitt¹⁰⁰, S. Schmitt⁴⁸, L. Schoeffel¹³⁵, A. Schoening^{63b}, P. G. Scholer⁵⁴, E. Schopf¹²⁶, M. Schott¹⁰⁰, J. Schovancova³⁶, S. Schramm⁵⁶, F. Schroeder¹⁷¹, T. Schroer⁵⁶, H.-C. Schultz-Coulon^{63a}, M. Schumacher⁵⁴, B. A. Schumm¹³⁶, Ph. Schune¹³⁵, A. J. Schuy¹³⁸, H. R. Schwartz¹³⁶, A. Schwartzman¹⁴³, T. A. Schwarz¹⁰⁶, Ph. Schwemling¹³⁵, R. Schwienhorst¹⁰⁷, A. Sciandra¹³⁶, G. Sciolla²⁶, F. Scuri^{74a}, C. D. Sebastiani⁹², K. Sedlaczek¹¹⁵, P. Seema¹⁸, S. C. Seidel¹¹², A. Seiden¹³⁶, B. D. Seidlitz⁴¹, C. Seitz⁴⁸, J. M. Seixas^{82b}, G. Sekhniaidze^{72a}, S. J. Sekula⁴⁴, L. Selem⁶⁰, N. Semprini-Cesari^{23b,23a}, D. Sengupta⁵⁶, V. Senthikumar¹⁶³, L. Serin⁶⁶, L. Serkin^{69a,69b}, M. Sessa^{76a,76b}, H. Severini¹²⁰, F. Sforza^{57b,57a}, A. Sfyrla⁵⁶, E. Shabalina⁵⁵, R. Shaheen¹⁴⁴, J. D. Shahinian¹²⁸, D. Shaked Renous¹⁶⁹, L. Y. Shan^{14a}, M. Shapiro^{17a}, A. Sharma³⁶, A. S. Sharma¹⁶⁴, P. Sharma⁸⁰, S. Sharma⁴⁸, P. B. Shatalov³⁷, K. Shaw¹⁴⁶, S. M. Shaw¹⁰¹, A. Shcherbakova³⁷, Q. Shen^{62c,5}, P. Sherwood⁹⁶, L. Shi⁹⁶, X. Shi^{14a}, C. O. Shimmin¹⁷², Y. Shimogama¹⁶⁸, J. D. Shinner⁹⁵, I. P. J. Shipsey¹²⁶, S. Shirabe^{56,hh}, M. Shiyakova³⁸, J. Shlomi¹⁶⁹, M. J. Shochet³⁹, J. Shojaii¹⁰⁵, D. R. Shope¹²⁵, S. Shrestha^{119,jj}, E. M. Shrif^{33g}, M. J. Shroff¹⁶⁵, P. Sicho¹³¹, A. M. Sickles¹⁶², E. Sideras Haddad^{33g}, A. Sidoti^{23b}, F. Siegert⁵⁰, Dj. Sijacki¹⁵, R. Sikora^{85a}, F. Sili⁹⁰, J. M. Silva²⁰, M. V. Silva Oliveira²⁹, S. B. Silverstein^{47a}, S. Simion⁶⁶, R. Simoniello³⁶, E. L. Simpson⁵⁹, H. Simpson¹⁴⁶, L. R. Simpson¹⁰⁶, N. D. Simpson⁹⁸, S. Simsek^{21d}, S. Sindhu⁵⁵, P. Sinervo¹⁵⁵, S. Singh¹⁵⁵, S. Sinha⁴⁸, S. Sinha¹⁰¹, M. Sioli^{23b,23a}, I. Siral³⁶, E. Sitnikova⁴⁸, S. Yu. Sivoklov^{37,a}, J. Sjölin^{47a,47b}, A. Skaf⁵⁵, E. Skorda⁹⁸, P. Skubic¹²⁰, M. Slawinska⁸⁶, V. Smakhtin¹⁶⁹, B. H. Smart¹³⁴, J. Smiesko³⁶, S. Yu. Smirnov³⁷, Y. Smirnov³⁷, L. N. Smirnova^{37,m}, O. Smirnova⁹⁸, A. C. Smith⁴¹, E. A. Smith³⁹, H. A. Smith¹²⁶, J. L. Smith⁹², R. Smith¹⁴³, M. Smizanska⁹¹, K. Smolek¹³², A. A. Snesarev³⁷, S. R. Snider¹⁵⁵, H. L. Snoek¹¹⁴, S. Snyder²⁹, R. Sobie^{165,q}, A. Soffer¹⁵¹, C. A. Solans Sanchez³⁶, E. Yu. Soldatov³⁷, U. Soldevila¹⁶³, A. A. Solodkov³⁷, S. Solomon²⁶, A. Soloshenko³⁸, K. Solovieva⁵⁴, O. V. Solovyanov⁴⁰, V. Solovyevev³⁷, P. Sommer³⁶, A. Sonay¹³, W. Y. Song^{156b}, J. M. Sonneveld¹¹⁴, A. Sopczak¹³², A. L. Soppio⁹⁶, F. Sopkova^{28b}, V. Sothilingam^{63a}, S. Sottocornola⁶⁸, R. Soualah^{116b}, Z. Soumami^{35e}, D. South⁴⁸, S. Spagnolo^{70a,70b}, M. Spalla¹¹⁰, D. Sperlich⁵⁴, G. Spigo³⁶, M. Spina¹⁴⁶, S. Spinali⁹¹, D. P. Spiteri⁵⁹, M. Spousta¹³³, E. J. Staats³⁴, A. Stabile^{71a,71b}, R. Stamen^{63a}, M. Stamenkovic¹¹⁴, A. Stampekis²⁰, M. Standke²⁴, E. Stanecka⁸⁶, M. V. Stange⁵⁰, B. Stanislaus^{17a}, M. M. Stanitzki⁴⁸, B. Stapf⁴⁸, E. A. Starchenko³⁷, G. H. Stark¹³⁶, J. Stark^{102,ff}, D. M. Starko^{156b}, P. Staroba¹³¹, P. Starovoitov^{63a}, S. Stärz¹⁰⁴, R. Staszewski⁸⁶, G. Stavropoulos⁴⁶, J. Steentoft¹⁶¹, P. Steinberg²⁹, B. Stelzer^{142,156a}, H. J. Stelzer¹²⁹, O. Stelzer-Chilton^{156a}, H. Stenzel⁵⁸, T. J. Stevenson¹⁴⁶, G. A. Stewart³⁶, J. R. Stewart¹²¹, M. C. Stockton³⁶, G. Stoicea^{27b}, M. Stolarski^{130a}, S. Stonjek¹¹⁰, A. Straessner⁵⁰, J. Strandberg¹⁴⁴, S. Strandberg^{47a,47b}, M. Strauss¹²⁰, T. Streblner¹⁰², P. Strizenec^{28b}, R. Ströhmer¹⁶⁶, D. M. Strom¹²³, L. R. Strom⁴⁸, R. Stroynowski⁴⁴, A. Strubig^{47a,47b}, S. A. Stucci²⁹, B. Stugu¹⁶, J. Stupak¹²⁰, N. A. Styles⁴⁸, D. Su¹⁴³, S. Su^{62a}, W. Su^{62d}, X. Su^{62a,66}, K. Sugizaki¹⁵³, V. V. Sulin³⁷, M. J. Sullivan⁹², D. M. S. Sultan^{78a,78b}, L. Sultanaliyeva³⁷, S. Sultansoy^{3b}, T. Sumida⁸⁷, S. Sun¹⁰⁶, S. Sun¹⁷⁰, O. Sunneborn Gudnadottir¹⁶¹, N. Sur¹⁰², M. R. Sutton¹⁴⁶, H. Suzuki¹⁵⁷, M. Svatos¹³¹, M. Swiatlowski^{156a}, T. Swirski¹⁶⁶, I. Sykora^{28a}, M. Sykora¹³³, T. Sykora¹³³, D. Ta¹⁰⁰, K. Tackmann^{48,kk}, A. Taffard¹⁶⁰, R. Tafirout^{156a}, J. S. Tafoya Vargas⁶⁶, R. Takashima⁸⁸, E. P. Takeva⁵², Y. Takubo⁸³, M. Talby¹⁰², A. A. Talyshev³⁷, K. C. Tam^{64b}, N. M. Tamir¹⁵¹, A. Tanaka¹⁵³, J. Tanaka¹⁵³, R. Tanaka⁶⁶, M. Tanasini^{57b,57a}, Z. Tao¹⁶⁴, S. Tapia Araya^{137f}, S. Tapprogge¹⁰⁰, A. Tarek Abouelfadl Mohamed¹⁰⁷, S. Tarem¹⁵⁰, K. Tariq^{14a}, G. Tarna^{102,27b}, G. F. Tartarelli^{71a}, P. Tas¹³³, M. Tasevsky¹³¹, E. Tassi^{43b,43a}, A. C. Tate¹⁶², G. Tateno¹⁵³, Y. Tayalati^{35e,ll}, G. N. Taylor¹⁰⁵, W. Taylor^{156b}, H. Teagle⁹², A. S. Tee¹⁷⁰, R. Teixeira De Lima¹⁴³, P. Teixeira-Dias⁹⁵, J. J. Teoh¹⁵⁵, K. Terashi¹⁵³, J. Terron⁹⁹, S. Terzo¹³, M. Testa⁵³, R. J. Teuscher^{155,q}, A. Thaler⁷⁹, O. Theiner⁵⁶, N. Themistokleous⁵², T. Thevenaux-Pelzer¹⁰², O. Thielmann¹⁷¹

D. W. Thomas,⁹⁵ J. P. Thomas,²⁰ E. A. Thompson,^{17a} P. D. Thompson,²⁰ E. Thomson,¹²⁸ Y. Tian,⁵⁵
V. Tikhomirov,^{37,m} Yu. A. Tikhonov,³⁷ S. Timoshenko,³⁷ D. Timoshyn,¹³³ E. X. L. Ting,¹ P. Tipton,¹⁷²
S. H. Tlou,^{33g} A. Tnourji,⁴⁰ K. Todome,^{23b,23a} S. Todorova-Nova,¹³³ S. Todt,⁵⁰ M. Togawa,⁸³ J. Tojo,⁸⁹
S. Tokár,^{28a} K. Tokushuku,⁸³ O. Toldaiev,⁶⁸ R. Tombs,³² M. Tomoto,^{83,111} L. Tompkins,^{143,u}
K. W. Topolnicki,^{85b} E. Torrence,¹²³ H. Torres,^{102,ff} E. Torró Pastor,¹⁶³ M. Toscani,³⁰ C. Toscirì,³⁹ M. Tost,¹¹
D. R. Tovey,¹³⁹ A. Traeet,¹⁶ I. S. Trandafir,^{27b} T. Trefzger,¹⁶⁶ A. Tricoli,²⁹ I. M. Trigger,^{156a} S. Trincaz-Duvoid,¹²⁷
D. A. Trischuk,²⁶ B. Trocmé,⁶⁰ C. Troncon,^{71a} L. Truong,^{33c} M. Trzebinski,⁸⁶ A. Trzupek,⁸⁶ F. Tsai,¹⁴⁵
M. Tsai,¹⁰⁶ A. Tsiamis,^{152,aa} P. V. Tsiareshka,³⁷ S. Tsigaridas,^{156a} A. Tsirigotis,^{152,bb} V. Tsiskaridze,¹⁵⁵
E. G. Tskhadadze,^{149a} M. Tsopoulou,^{152,aa} Y. Tsujikawa,⁸⁷ I. I. Tsukerman,³⁷ V. Tsulaia,^{17a} S. Tsuno,⁸³ O. Tsur,¹⁵⁰
K. Tsurì,¹¹⁸ D. Tsybychev,¹⁴⁵ Y. Tu,^{64b} A. Tudorache,^{27b} V. Tudorache,^{27b} A. N. Tuna,³⁶ S. Turchikhin,³⁸
I. Turk Cakir,^{3a} R. Turra,^{71a} T. Turtuvshin,^{38,mm} P. M. Tuts,⁴¹ S. Tzamarías,^{152,aa} P. Tzaniş,¹⁰ E. Tzovara,¹⁰⁰
K. Uchida,¹⁵³ F. Ukegawa,¹⁵⁷ P. A. Ulloa Poblete,^{137c,137b} E. N. Umaka,²⁹ G. Unal,³⁶ M. Unal,¹¹ A. Undrus,²⁹
G. Unel,¹⁶⁰ J. Urban,^{28b} P. Urquijo,¹⁰⁵ G. Usai,⁸ R. Ushioda,¹⁵⁴ M. Usman,¹⁰⁸ Z. Uysal,^{21b} L. Vacavant,¹⁰²
V. Vacek,¹³² B. Vachon,¹⁰⁴ K. O. H. Vadla,¹²⁵ T. Vafeiadis,³⁶ A. Vaitkus,⁹⁶ C. Valderanis,¹⁰⁹
E. Valdes Santurio,^{47a,47b} M. Valente,^{156a} S. Valentinetti,^{23b,23a} A. Valero,¹⁶³ E. Valiente Moreno,¹⁶³ A. Vallier,^{102,ff}
J. A. Valls Ferrer,¹⁶³ D. R. Van Arneinan,¹¹⁴ T. R. Van Daalen,¹³⁸ A. Van Der Graaf,⁴⁹ P. Van Gemmeren,⁶
M. Van Rijnbach,^{125,36} S. Van Stroud,⁹⁶ I. Van Vulpen,¹¹⁴ M. Vanadia,^{76a,76b} W. Vandelli,³⁶ M. Vandenbroucke,¹³⁵
E. R. Vandewall,¹²¹ D. Vannicola,¹⁵¹ L. Vannoli,^{57b,57a} R. Vari,^{75a} E. W. Varnes,⁷ C. Varni,^{17a} T. Varol,¹⁴⁸
D. Varouchas,⁶⁶ L. Varriale,¹⁶³ K. E. Varvell,¹⁴⁷ M. E. Vasile,^{27b} L. Vaslin,⁴⁰ G. A. Vasquez,¹⁶⁵ F. Vazeille,⁴⁰
T. Vazquez Schroeder,³⁶ J. Veatch,³¹ V. Vecchio,¹⁰¹ M. J. Veen,¹⁰³ I. Velisek,¹²⁶ L. M. Veloce,¹⁵⁵
F. Veloso,^{130a,130c} S. Veneziano,^{75a} A. Ventura,^{70a,70b} A. Verbytskyi,¹¹⁰ M. Verducci,^{74a,74b} C. Vergis,²⁴
M. Verissimo De Araujo,^{82b} W. Verkerke,¹¹⁴ J. C. Vermeulen,¹¹⁴ C. Vernieri,¹⁴³ P. J. Verschuuren,⁹⁵
M. Vessella,¹⁰³ M. C. Vetterli,^{142,e} A. Vgenopoulos,^{152,aa} N. Viaux Maira,^{137f} T. Vickey,¹³⁹
O. E. Vickey Boeriu,¹³⁹ G. H. A. Viehhauser,¹²⁶ L. Viganì,^{63b} M. Villa,^{23b,23a} M. Villaplana Perez,¹⁶³
E. M. Villhauer,⁵² E. Vilucchi,⁵³ M. G. Vincker,³⁴ G. S. Virdee,²⁰ A. Vishwakarma,⁵² A. Visibile,¹¹⁴ C. Vittori,³⁶
I. Vivarelli,¹⁴⁶ V. Vladimirov,¹⁶⁷ E. Voevodina,¹¹⁰ F. Vogel,¹⁰⁹ P. Vokac,¹³² J. Von Ahnen,⁴⁸ E. Von Toerne,²⁴
B. Vormwald,³⁶ V. Vorobel,¹³³ K. Vorobev,³⁷ M. Vos,¹⁶³ K. Voss,¹⁴¹ J. H. Vosseveld,⁹² M. Vozak,¹¹⁴
L. Vozdecky,⁹⁴ N. Vranjes,¹⁵ M. Vranjes Milosavljevic,¹⁵ M. Vreeswijk,¹¹⁴ N. K. Vu,^{62d,62c} R. Vuillermet,³⁶
O. Vujanovic,¹⁰⁰ I. Vukotic,³⁹ S. Wada,¹⁵⁷ C. Wagner,¹⁰³ J. M. Wagner,^{17a} W. Wagner,¹⁷¹ S. Wahdan,¹⁷¹
H. Wahlberg,⁹⁰ R. Wakasa,¹⁵⁷ M. Wakida,¹¹¹ J. Walder,¹³⁴ R. Walker,¹⁰⁹ W. Walkowiak,¹⁴¹ A. Wall,¹²⁸
T. Wamorkar,⁶ A. Z. Wang,¹⁷⁰ C. Wang,¹⁰⁰ C. Wang,^{62c} H. Wang,^{17a} J. Wang,^{64a} R.-J. Wang,¹⁰⁰ R. Wang,⁶¹
R. Wang,⁶ S. M. Wang,¹⁴⁸ S. Wang,^{62b} T. Wang,^{62a} W. T. Wang,⁸⁰ W. Wang,^{14a} X. Wang,^{14c} X. Wang,¹⁶²
X. Wang,^{62c} Y. Wang,^{62d} Y. Wang,^{14c} Z. Wang,¹⁰⁶ Z. Wang,^{62d,51,62c} Z. Wang,¹⁰⁶ A. Warburton,¹⁰⁴
R. J. Ward,²⁰ N. Warrack,⁵⁹ A. T. Watson,²⁰ H. Watson,⁵⁹ M. F. Watson,²⁰ E. Watton,^{59,134} G. Watts,¹³⁸
B. M. Waugh,⁹⁶ C. Weber,²⁹ H. A. Weber,¹⁸ M. S. Weber,¹⁹ S. M. Weber,^{63a} C. Wei,^{62a} Y. Wei,¹²⁶
A. R. Weidberg,¹²⁶ E. J. Weik,¹¹⁷ J. Weingarten,⁴⁹ M. Weirich,¹⁰⁰ C. Weiser,⁵⁴ C. J. Wells,⁴⁸ T. Wenaus,²⁹
B. Wendland,⁴⁹ T. Wengler,³⁶ N. S. Wenke,¹¹⁰ N. Wermes,²⁴ M. Wessels,^{63a} K. Whalen,¹²³ A. M. Wharton,⁹¹
A. S. White,⁶¹ A. White,⁸ M. J. White,¹ D. Whiteson,¹⁶⁰ L. Wickremasinghe,¹²⁴ W. Wiedenmann,¹⁷⁰ C. Wiel,⁵⁰
M. Wielers,¹³⁴ C. Wiglesworth,⁴² D. J. Wilbern,¹²⁰ H. G. Wilkens,³⁶ D. M. Williams,⁴¹ H. H. Williams,¹²⁸
S. Williams,³² S. Willocq,¹⁰³ B. J. Wilson,¹⁰¹ P. J. Windischhofer,³⁹ F. I. Winkel,³⁰ F. Winklmeier,¹²³
B. T. Winter,⁵⁴ J. K. Winter,¹⁰¹ M. Wittgen,¹⁴³ M. Wobisch,⁹⁷ Z. Wolfs,¹¹⁴ R. Wölker,¹²⁶ J. Wollrath,¹⁶⁰
M. W. Wolter,⁸⁶ H. Wolters,^{130a,130c} A. F. Wongel,⁴⁸ S. D. Worm,⁴⁸ B. K. Wosiek,⁸⁶ K. W. Woźniak,⁸⁶
S. Wozniewski,⁵⁵ K. Wraight,⁵⁹ C. Wu,²⁰ J. Wu,^{14a,14e} M. Wu,^{64a} M. Wu,¹¹³ S. L. Wu,¹⁷⁰ X. Wu,⁵⁶ Y. Wu,^{62a}
Z. Wu,¹³⁵ J. Wuerzinger,¹¹⁰ T. R. Wyatt,¹⁰¹ B. M. Wynne,⁵² S. Xella,⁴² L. Xia,^{14c} M. Xia,^{14b} J. Xiang,^{64c}
X. Xiao,¹⁰⁶ M. Xie,^{62a} X. Xie,^{62a} S. Xin,^{14a,14e} J. Xiong,^{17a} D. Xu,^{14a} H. Xu,^{62a} L. Xu,^{62a} R. Xu,¹²⁸ T. Xu,¹⁰⁶
Y. Xu,^{14b} Z. Xu,⁵² Z. Xu,^{14a} B. Yabsley,¹⁴⁷ S. Yacoob,^{33a} N. Yamaguchi,⁸⁹ Y. Yamaguchi,¹⁵⁴ E. Yamashita,¹⁵³
H. Yamauchi,¹⁵⁷ T. Yamazaki,^{17a} Y. Yamazaki,⁸⁴ J. Yan,^{62c} S. Yan,¹²⁶ Z. Yan,²⁵ H. J. Yang,^{62c,62d} H. T. Yang,^{62a}
S. Yang,^{62a} T. Yang,^{64c} X. Yang,^{62a} X. Yang,^{14a} Y. Yang,⁴⁴ Y. Yang,^{62a} Z. Yang,^{62a} W.-M. Yao,^{17a} Y. C. Yap,⁴⁸
H. Ye,^{14c} H. Ye,⁵⁵ J. Ye,⁴⁴ S. Ye,²⁹ X. Ye,^{62a} Y. Yeh,⁹⁶ I. Yeletsikh,³⁸ B. K. Yeo,^{17a} M. R. Yexley,⁹⁶

P. Yin⁴¹, K. Yorita¹⁶⁸, S. Younas^{27b}, C. J. S. Young⁵⁴, C. Young¹⁴³, Y. Yu^{62a}, M. Yuan¹⁰⁶, R. Yuan^{62b,nn},
 L. Yue⁹⁶, M. Zaazoua^{62a}, B. Zabinski⁸⁶, E. Zaid⁵², T. Zakareishvili^{149b}, N. Zakharchuk³⁴, S. Zambito⁵⁶,
 J. A. Zamora Saa^{137d,137b}, J. Zang¹⁵³, D. Zanzi⁵⁴, O. Zaplatilek¹³², C. Zeitnitz¹⁷¹, H. Zeng^{14a}, J. C. Zeng¹⁶²,
 D. T. Zenger Jr.²⁶, O. Zenin³⁷, T. Ženiš^{28a}, S. Zenz⁹⁴, S. Zerradi^{35a}, D. Zerwas⁶⁶, M. Zhai^{14a,14e}, B. Zhang^{14c},
 D. F. Zhang¹³⁹, J. Zhang^{62b}, J. Zhang⁶, K. Zhang^{14a,14e}, L. Zhang^{14c}, P. Zhang^{14a,14e}, R. Zhang¹⁷⁰, S. Zhang¹⁰⁶,
 T. Zhang¹⁵³, X. Zhang^{62c}, X. Zhang^{62b}, Y. Zhang^{62c,5}, Y. Zhang⁹⁶, Z. Zhang^{17a}, Z. Zhang⁶⁶, H. Zhao¹³⁸,
 P. Zhao⁵¹, T. Zhao^{62b}, Y. Zhao¹³⁶, Z. Zhao^{62a}, A. Zhemchugov³⁸, K. Zheng¹⁶², X. Zheng^{62a}, Z. Zheng¹⁴³,
 D. Zhong¹⁶², B. Zhou¹⁰⁶, H. Zhou⁷, N. Zhou^{62c}, Y. Zhou⁷, C. G. Zhu^{62b}, J. Zhu¹⁰⁶, Y. Zhu^{62c}, Y. Zhu^{62a},
 X. Zhuang^{14a}, K. Zhukov³⁷, V. Zhulanov³⁷, N. I. Zimine³⁸, J. Zinsser^{63b}, M. Ziolkowski¹⁴¹, L. Živković¹⁵,
 A. Zoccoli^{23b,23a}, K. Zoch⁵⁶, T. G. Zorbas¹³⁹, O. Zormpa⁴⁶, W. Zou⁴¹ and L. Zwalinski³⁶

(ATLAS Collaboration)

¹Department of Physics, University of Adelaide, Adelaide, Australia

²Department of Physics, University of Alberta, Edmonton, Alberta, Canada

^{3a}Department of Physics, Ankara University, Ankara, Türkiye

^{3b}Division of Physics, TOBB University of Economics and Technology, Ankara, Türkiye

⁴LAPP, Université Savoie Mont Blanc, CNRS/IN2P3, Annecy, France

⁵APC, Université Paris Cité, CNRS/IN2P3, Paris, France

⁶High Energy Physics Division, Argonne National Laboratory, Argonne, Illinois, USA

⁷Department of Physics, University of Arizona, Tucson, Arizona, USA

⁸Department of Physics, University of Texas at Arlington, Arlington, Texas, USA

⁹Physics Department, National and Kapodistrian University of Athens, Athens, Greece

¹⁰Physics Department, National Technical University of Athens, Zografou, Greece

¹¹Department of Physics, University of Texas at Austin, Austin, Texas, USA

¹²Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹³Institut de Física d'Altes Energies (IFAE), Barcelona Institute of Science and Technology, Barcelona, Spain

^{14a}Institute of High Energy Physics, Chinese Academy of Sciences, Beijing, China

^{14b}Physics Department, Tsinghua University, Beijing, China

^{14c}Department of Physics, Nanjing University, Nanjing, China

^{14d}School of Science, Shenzhen Campus of Sun Yat-sen University, Guangzhou, China

^{14e}University of Chinese Academy of Science (UCAS), Beijing, China

¹⁵Institute of Physics, University of Belgrade, Belgrade, Serbia

¹⁶Department for Physics and Technology, University of Bergen, Bergen, Norway

^{17a}Physics Division, Lawrence Berkeley National Laboratory, Berkeley, California, USA

^{17b}University of California, Berkeley, California, USA

¹⁸Institut für Physik, Humboldt Universität zu Berlin, Berlin, Germany

¹⁹Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

²⁰School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom

^{21a}Department of Physics, Bogazici University, Istanbul, Türkiye

^{21b}Department of Physics Engineering, Gaziantep University, Gaziantep, Türkiye

^{21c}Department of Physics, Istanbul University, Istanbul, Türkiye

^{21d}Istinye University, Sariyer, Istanbul, Türkiye

^{22a}Facultad de Ciencias y Centro de Investigaciones, Universidad Antonio Nariño, Bogotá, Colombia

^{22b}Departamento de Física, Universidad Nacional de Colombia, Bogotá, Colombia

^{22c}Pontificia Universidad Javeriana, Bogota, Colombia

^{23a}Dipartimento di Fisica e Astronomia A. Righi, Università di Bologna, Bologna, Italy

^{23b}INFN Sezione di Bologna, Bologna, Italy

²⁴Physikalisches Institut, Universität Bonn, Bonn, Germany

²⁵Department of Physics, Boston University, Boston, Massachusetts, USA

²⁶Department of Physics, Brandeis University, Waltham, Massachusetts, USA

^{27a}Transilvania University of Brasov, Brasov, Romania

^{27b}Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania

^{27c}Department of Physics, Alexandru Ioan Cuza University of Iasi, Iasi, Romania

^{27d}National Institute for Research and Development of Isotopic and Molecular Technologies,
 Physics Department, Cluj-Napoca, Romania

- ^{27e}University Politehnica Bucharest, Bucharest, Romania
^{27f}West University in Timisoara, Timisoara, Romania
^{27g}Faculty of Physics, University of Bucharest, Bucharest, Romania
^{28a}Faculty of Mathematics, Physics and Informatics, Comenius University, Bratislava, Slovak Republic
^{28b}Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
²⁹Physics Department, Brookhaven National Laboratory, Upton, New York, USA
³⁰Universidad de Buenos Aires, Facultad de Ciencias Exactas y Naturales, Departamento de Física, y CONICET, Instituto de Física de Buenos Aires (IFIBA), Buenos Aires, Argentina
³¹California State University, Sacramento, California, USA
³²Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
^{33a}Department of Physics, University of Cape Town, Cape Town, South Africa
^{33b}iThemba Labs, Western Cape, South Africa
^{33c}Department of Mechanical Engineering Science, University of Johannesburg, Johannesburg, South Africa
^{33d}National Institute of Physics, University of the Philippines Diliman (Philippines), Quezon City, Philippines
^{33e}University of South Africa, Department of Physics, Pretoria, South Africa
^{33f}University of Zululand, KwaDlangezwa, South Africa
^{33g}School of Physics, University of the Witwatersrand, Johannesburg, South Africa
³⁴Department of Physics, Carleton University, Ottawa, Ontario, Canada
^{35a}Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca, Morocco
^{35b}Faculté des Sciences, Université Ibn-Tofail, Kénitra, Morocco
^{35c}Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech, Morocco
^{35d}LPMR, Faculté des Sciences, Université Mohamed Premier, Oujda, Morocco
^{35e}Faculté des sciences, Université Mohammed V, Rabat, Morocco
^{35f}Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir, Morocco
³⁶CERN, Geneva, Switzerland
³⁷Affiliated with an institute covered by a cooperation agreement with CERN
³⁸Affiliated with an international laboratory covered by a cooperation agreement with CERN
³⁹Enrico Fermi Institute, University of Chicago, Chicago, Illinois, USA
⁴⁰LPC, Université Clermont Auvergne, CNRS/IN2P3, Clermont-Ferrand, France
⁴¹Nevis Laboratory, Columbia University, Irvington, New York, USA
⁴²Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
^{43a}Dipartimento di Fisica, Università della Calabria, Rende, Italy
^{43b}INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Rende, Italy
⁴⁴Physics Department, Southern Methodist University, Dallas, Texas, USA
⁴⁵Physics Department, University of Texas at Dallas, Richardson, Texas, USA
⁴⁶National Centre for Scientific Research “Demokritos”, Agia Paraskevi, Greece
^{47a}Department of Physics, Stockholm University, Stockholm, Sweden
^{47b}Oskar Klein Centre, Stockholm, Sweden
⁴⁸Deutsches Elektronen-Synchrotron DESY, Hamburg and Zeuthen, Germany
⁴⁹Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany
⁵⁰Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
⁵¹Department of Physics, Duke University, Durham, North Carolina, USA
⁵²SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
⁵³INFN e Laboratori Nazionali di Frascati, Frascati, Italy
⁵⁴Physikalisches Institut, Albert-Ludwigs-Universität Freiburg, Freiburg, Germany
⁵⁵II. Physikalisches Institut, Georg-August-Universität Göttingen, Göttingen, Germany
⁵⁶Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland
^{57a}Dipartimento di Fisica, Università di Genova, Genova, Italy
^{57b}INFN Sezione di Genova, Genova, Italy
⁵⁸II. Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
⁵⁹SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
⁶⁰LPSC, Université Grenoble Alpes, CNRS/IN2P3, Grenoble INP, Grenoble, France
⁶¹Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA
^{62a}Department of Modern Physics and State Key Laboratory of Particle Detection and Electronics, University of Science and Technology of China, Hefei, China
^{62b}Institute of Frontier and Interdisciplinary Science and Key Laboratory of Particle Physics and Particle Irradiation (MOE), Shandong University, Qingdao, China
^{62c}School of Physics and Astronomy, Shanghai Jiao Tong University, Key Laboratory for Particle Astrophysics and Cosmology (MOE), SKLPPC, Shanghai, China

- ^{62d}*Tsung-Dao Lee Institute, Shanghai, China*
- ^{63a}*Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ^{63b}*Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany*
- ^{64a}*Department of Physics, Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China*
- ^{64b}*Department of Physics, University of Hong Kong, Hong Kong, China*
- ^{64c}*Department of Physics and Institute for Advanced Study, Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China*
- ⁶⁵*Department of Physics, National Tsing Hua University, Hsinchu, Taiwan*
- ⁶⁶*IJCLab, Université Paris-Saclay, CNRS/IN2P3, 91405, Orsay, France*
- ⁶⁷*Centro Nacional de Microelectrónica (IMB-CNM-CSIC), Barcelona, Spain*
- ⁶⁸*Department of Physics, Indiana University, Bloomington, Indiana, USA*
- ^{69a}*INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy*
- ^{69b}*ICTP, Trieste, Italy*
- ^{69c}*Dipartimento Politecnico di Ingegneria e Architettura, Università di Udine, Udine, Italy*
- ^{70a}*INFN Sezione di Lecce, Lecce, Italy*
- ^{70b}*Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy*
- ^{71a}*INFN Sezione di Milano, Milano, Italy*
- ^{71b}*Dipartimento di Fisica, Università di Milano, Milano, Italy*
- ^{72a}*INFN Sezione di Napoli, Napoli, Italy*
- ^{72b}*Dipartimento di Fisica, Università di Napoli, Napoli, Italy*
- ^{73a}*INFN Sezione di Pavia, Pavia, Italy*
- ^{73b}*Dipartimento di Fisica, Università di Pavia, Pavia, Italy*
- ^{74a}*INFN Sezione di Pisa, Pisa, Italy*
- ^{74b}*Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy*
- ^{75a}*INFN Sezione di Roma, Roma, Italy*
- ^{75b}*Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy*
- ^{76a}*INFN Sezione di Roma Tor Vergata, Roma, Italy*
- ^{76b}*Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy*
- ^{77a}*INFN Sezione di Roma Tre, Roma, Italy*
- ^{77b}*Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy*
- ^{78a}*INFN-TIFPA, Trento, Italy*
- ^{78b}*Università degli Studi di Trento, Trento, Italy*
- ⁷⁹*Universität Innsbruck, Department of Astro and Particle Physics, Innsbruck, Austria*
- ⁸⁰*University of Iowa, Iowa City, Iowa, USA*
- ⁸¹*Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA*
- ^{82a}*Departamento de Engenharia Elétrica, Universidade Federal de Juiz de Fora (UFJF), Juiz de Fora, Brazil*
- ^{82b}*Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro, Brazil*
- ^{82c}*Instituto de Física, Universidade de São Paulo, São Paulo, Brazil*
- ^{82d}*Rio de Janeiro State University, Rio de Janeiro, Brazil*
- ⁸³*KEK, High Energy Accelerator Research Organization, Tsukuba, Japan*
- ⁸⁴*Graduate School of Science, Kobe University, Kobe, Japan*
- ^{85a}*AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland*
- ^{85b}*Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- ⁸⁶*Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- ⁸⁷*Faculty of Science, Kyoto University, Kyoto, Japan*
- ⁸⁸*Kyoto University of Education, Kyoto, Japan*
- ⁸⁹*Research Center for Advanced Particle Physics and Department of Physics, Kyushu University, Fukuoka, Japan*
- ⁹⁰*Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina*
- ⁹¹*Physics Department, Lancaster University, Lancaster, United Kingdom*
- ⁹²*Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom*
- ⁹³*Department of Experimental Particle Physics, Jožef Stefan Institute and Department of Physics, University of Ljubljana, Ljubljana, Slovenia*
- ⁹⁴*School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom*
- ⁹⁵*Department of Physics, Royal Holloway University of London, Egham, United Kingdom*
- ⁹⁶*Department of Physics and Astronomy, University College London, London, United Kingdom*
- ⁹⁷*Louisiana Tech University, Ruston, Louisiana, USA*
- ⁹⁸*Fysiska institutionen, Lunds universitet, Lund, Sweden*
- ⁹⁹*Departamento de Física Teórica C-15 and CIAFF, Universidad Autónoma de Madrid, Madrid, Spain*
- ¹⁰⁰*Institut für Physik, Universität Mainz, Mainz, Germany*
- ¹⁰¹*School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom*

- ¹⁰²CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
- ¹⁰³Department of Physics, University of Massachusetts, Amherst, Massachusetts, USA
- ¹⁰⁴Department of Physics, McGill University, Montreal, Quebec, Canada
- ¹⁰⁵School of Physics, University of Melbourne, Victoria, Australia
- ¹⁰⁶Department of Physics, University of Michigan, Ann Arbor, Michigan, USA
- ¹⁰⁷Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA
- ¹⁰⁸Group of Particle Physics, University of Montreal, Montreal, Quebec, Canada
- ¹⁰⁹Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- ¹¹⁰Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- ¹¹¹Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
- ¹¹²Department of Physics and Astronomy, University of New Mexico, Albuquerque, New Mexico, USA
- ¹¹³Institute for Mathematics, Astrophysics and Particle Physics, Radboud University/Nikhef, Nijmegen, Netherlands
- ¹¹⁴Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- ¹¹⁵Department of Physics, Northern Illinois University, DeKalb, Illinois, USA
- ^{116a}New York University Abu Dhabi, Abu Dhabi, United Arab Emirates
- ^{116b}University of Sharjah, Sharjah, United Arab Emirates
- ¹¹⁷Department of Physics, New York University, New York, New York, USA
- ¹¹⁸Ochanomizu University, Otsuka, Bunkyo-ku, Tokyo, Japan
- ¹¹⁹Ohio State University, Columbus, Ohio, USA
- ¹²⁰Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, Oklahoma, USA
- ¹²¹Department of Physics, Oklahoma State University, Stillwater, Oklahoma, USA
- ¹²²Palacký University, Joint Laboratory of Optics, Olomouc, Czech Republic
- ¹²³Institute for Fundamental Science, University of Oregon, Eugene, Oregon, USA
- ¹²⁴Graduate School of Science, Osaka University, Osaka, Japan
- ¹²⁵Department of Physics, University of Oslo, Oslo, Norway
- ¹²⁶Department of Physics, Oxford University, Oxford, United Kingdom
- ¹²⁷LPNHE, Sorbonne Université, Université Paris Cité, CNRS/IN2P3, Paris, France
- ¹²⁸Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania, USA
- ¹²⁹Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, Pennsylvania, USA
- ^{130a}Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa, Portugal
- ^{130b}Departamento de Física, Faculdade de Ciências, Universidade de Lisboa, Lisboa, Portugal
- ^{130c}Departamento de Física, Universidade de Coimbra, Coimbra, Portugal
- ^{130d}Centro de Física Nuclear da Universidade de Lisboa, Lisboa, Portugal
- ^{130e}Departamento de Física, Universidade do Minho, Braga, Portugal
- ^{130f}Departamento de Física Teórica y del Cosmos, Universidad de Granada, Granada (Spain), Spain
- ^{130g}Departamento de Física, Instituto Superior Técnico, Universidade de Lisboa, Lisboa, Portugal
- ¹³¹Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
- ¹³²Czech Technical University in Prague, Prague, Czech Republic
- ¹³³Charles University, Faculty of Mathematics and Physics, Prague, Czech Republic
- ¹³⁴Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- ¹³⁵IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France
- ¹³⁶Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, California, USA
- ^{137a}Departamento de Física, Pontificia Universidad Católica de Chile, Santiago, Chile
- ^{137b}Millennium Institute for Subatomic physics at high energy frontier (SAPHIR), Santiago, Chile
- ^{137c}Instituto de Investigación Multidisciplinario en Ciencia y Tecnología, y Departamento de Física, Universidad de La Serena, La Serena, Chile
- ^{137d}Universidad Andres Bello, Department of Physics, Santiago, Chile
- ^{137e}Instituto de Alta Investigación, Universidad de Tarapacá, Arica, Chile
- ^{137f}Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- ¹³⁸Department of Physics, University of Washington, Seattle, Washington, USA
- ¹³⁹Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- ¹⁴⁰Department of Physics, Shinshu University, Nagano, Japan
- ¹⁴¹Department Physik, Universität Siegen, Siegen, Germany
- ¹⁴²Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada
- ¹⁴³SLAC National Accelerator Laboratory, Stanford, California, USA
- ¹⁴⁴Department of Physics, Royal Institute of Technology, Stockholm, Sweden
- ¹⁴⁵Departments of Physics and Astronomy, Stony Brook University, Stony Brook, New York, USA
- ¹⁴⁶Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
- ¹⁴⁷School of Physics, University of Sydney, Sydney, Australia
- ¹⁴⁸Institute of Physics, Academia Sinica, Taipei, Taiwan

- ^{149a}*E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia*
^{149b}*High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
^{149c}*University of Georgia, Tbilisi, Georgia*
¹⁵⁰*Department of Physics, Technion, Israel Institute of Technology, Haifa, Israel*
¹⁵¹*Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
¹⁵²*Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
¹⁵³*International Center for Elementary Particle Physics and Department of Physics, University of Tokyo, Tokyo, Japan*
¹⁵⁴*Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
¹⁵⁵*Department of Physics, University of Toronto, Toronto, Ontario, Canada*
^{156a}*TRIUMF, Vancouver, British Columbia, Canada*
^{156b}*Department of Physics and Astronomy, York University, Toronto, Ontario, Canada*
¹⁵⁷*Division of Physics and Tomonaga Center for the History of the Universe, Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
¹⁵⁸*Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA*
¹⁵⁹*United Arab Emirates University, Al Ain, United Arab Emirates*
¹⁶⁰*Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA*
¹⁶¹*Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
¹⁶²*Department of Physics, University of Illinois, Urbana, Illinois, USA*
¹⁶³*Instituto de Física Corpuscular (IFIC), Centro Mixto Universidad de Valencia - CSIC, Valencia, Spain*
¹⁶⁴*Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada*
¹⁶⁵*Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada*
¹⁶⁶*Fakultät für Physik und Astronomie, Julius-Maximilians-Universität Würzburg, Würzburg, Germany*
¹⁶⁷*Department of Physics, University of Warwick, Coventry, United Kingdom*
¹⁶⁸*Waseda University, Tokyo, Japan*
¹⁶⁹*Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot, Israel*
¹⁷⁰*Department of Physics, University of Wisconsin, Madison, Wisconsin, USA*
¹⁷¹*Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany*
¹⁷²*Department of Physics, Yale University, New Haven, Connecticut, USA*

^aDeceased.

^bAlso at Department of Physics, King's College London, London, United Kingdom.

^cAlso at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^dAlso at Lawrence Livermore National Laboratory, Livermore, California, USA.

^eAlso at TRIUMF, Vancouver, British Columbia, Canada.

^fAlso at Department of Physics, University of Thessaly, Greece.

^gAlso at An-Najah National University, Nablus, Palestine.

^hAlso at Department of Physics, University of Fribourg, Fribourg, Switzerland.

ⁱAlso at University of Colorado Boulder, Department of Physics, Boulder, Colorado, USA.

^jAlso at Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada.

^kAlso at Department of Physics, Westmont College, Santa Barbara, USA.

^lAlso at Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain.

^mAlso at Affiliated with an institute covered by a cooperation agreement with CERN.

ⁿAlso at The Collaborative Innovation Center of Quantum Matter (CICQM), Beijing, China.

^oAlso at Department of Physics, Ben Gurion University of the Negev, Beer Sheva, Israel.

^pAlso at Università di Napoli Parthenope, Napoli, Italy.

^qAlso at Institute of Particle Physics (IPP), Victoria, British Columbia, Canada.

^rAlso at Borough of Manhattan Community College, City University of New York, New York, New York, USA.

^sAlso at National Institute of Physics, University of the Philippines Diliman (Philippines), Philippines.

^tAlso at Department of Financial and Management Engineering, University of the Aegean, Chios, Greece.

^uAlso at Department of Physics, Stanford University, Stanford, California, USA.

^vAlso at Centro Studi e Ricerche Enrico Fermi, Italy.

^wAlso at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

^xAlso at Technical University of Munich, Munich, Germany.

^yAlso at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.

^zAlso at CERN, Geneva, Switzerland.

^{aa}Also at Center for Interdisciplinary Research and Innovation (CIRI-AUTH), Thessaloniki, Greece.

^{bb}Also at Hellenic Open University, Patras, Greece.

^{cc}Also at Center for High Energy Physics, Peking University, Beijing, China.

^{dd}Also at APC, Université Paris Cité, CNRS/IN2P3, Paris, France.

^{ee}Also at Department of Physics, Royal Holloway University of London, Egham, United Kingdom.

^{ff} Also at L2IT, Université de Toulouse, CNRS/IN2P3, UPS, Toulouse, France.

^{gg} Also at Department of Physics, California State University, Sacramento, California, USA.

^{hh} Also at Département de Physique Nucléaire et Corpusculaire, Université de Genève, Genève, Switzerland.

ⁱⁱ Also at Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany.

^{jj} Also at Washington College, Chestertown, Maryland, USA.

^{kk} Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

^{ll} Also at Institute of Applied Physics, Mohammed VI Polytechnic University, Ben Guerir, Morocco.

^{mm} Also at Institute of Physics and Technology, Ulaanbaatar, Mongolia.

ⁿⁿ Also at Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan, USA.