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This is the final peer-reviewed author's accepted manuscript (postprint) of the following publication:

Published Version:

Farhad Zeighami, Leonardo Sandoval, Alberto Guadagnini, Vittorio Di Federico (2023). Uncertainty quantification and global sensitivity analysis of seismic metabarriers. *ENGINEERING STRUCTURES*, 277, 1-13 [10.1016/j.engstruct.2022.115415].

Availability:

This version is available at: <https://hdl.handle.net/11585/909731> since: 2024-07-03

Published:

DOI: <http://doi.org/10.1016/j.engstruct.2022.115415>

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Uncertainty quantification and global sensitivity analysis of seismic metabarriers

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Abstract

Seismic metabarriers consist of an array of locally resonant elements (i.e., mechanical resonators) installed over the soil surface, whose design is rationally engineered to reduce ground-induced vibrations and shield vulnerable structures against seismic surface waves. Successful design and implementation of seismic metabarriers require a comprehensive knowledge and characterization of the role played by the model parameters (and their associated uncertainty) governing soil-barrier dynamic interaction. In this context, sensitivity analysis techniques allow assessing the response of a given model through the quantification of the influence and action of model inputs (and model input uncertainties) concerning a target model output. This study relies on global sensitivity analysis techniques to investigate the influence that the uncertainty associated with three key mechanical parameters of a metabarrier (i.e., soil density, soil shear modulus, and mass of mechanical resonators) has on its seismic isolation performance. The latter is measured in terms of transmission coefficient (TC). We do so by employing a two-dimensional wave finite element model developed under the plane-strain conditions to evaluate the dispersion relation and transmission coefficient of a metabarrier interacting with Rayleigh waves in the low-frequency regime (i.e., frequencies between 2Hz and 7Hz). Our results suggest that the shear modulus is the uncertain parameter with the most significant influence on the transmission coefficient of the metabarrier across the entire frequency range of interest. Otherwise, the resonator mass plays a substantial role in the frequency range close to the metabarrier resonant frequency.

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Keywords: Seismic metamaterials, Metabarrier, Seismic surface waves, Global sensitivity analysis, Surrogate modeling, Polynomial chaos expansion.

1. Introduction

Elastic metamaterials are artificial composite structures designed to possess unconventional mechanical properties (Deymier et al., 2013). Since their emergence around two decades ago, elastic metamaterials have found several applications in mechanical and civil engineering fields. In the context of earthquake engineering, seismic metamaterials (SMs) are proposed as an innovative isolation technique to safeguard critical and vulnerable structures such as high-rise buildings, urban areas, historical places, and heritage sites (Brûlé et al., 2014; Krödel et al., 2015; Colombi et al., 2016c; Miniaci et al., 2016). SMs are generally classified according to their application, regulation mechanisms, and arrangement patterns (Mu et al., 2020). From an application perspective, SMs are categorized as periodic (Cheng and Shi, 2013; Cheng et al., 2020) and resonant foundations (Basone et al., 2019; Sun et al., 2019) that interact with seismic waves to shield unprotected residential buildings and industrial facilities, as well as periodic (Huang and Shi, 2013; Ni and Shi, 2022) and resonant barriers (Palermo et al., 2016; Colombi et al., 2017) to protect critical infrastructures or structures against incident surface waves. Among these systems, locally resonant barriers, namely metabarriers, incorporate wave-impeding devices with feasible dimensions from an engineering attitude that do not require any structural intervention to the target infrastructure.

A seismic metabarrier is a passive resonant barrier organized as an arbitrary arrangement of resonant structures/units with dimensions much smaller than the wavelength of seismic waves (Palermo et al., 2016). Metabarriers are typically installed in the vicinity of target structures and activated by the motion of incoming waves. The operating resonant frequency of a metabarrier is usually tuned in the low-frequency regime (i.e. below 10 Hz), where most of the elastic energy of seismic waves exists (Colombi et al., 2016c; Palermo et al., 2016, 2018b). The dynamic interaction between metabarrier and seismic waves is described analytically through the effective medium approach (Boechler et al., 2013) and the multiple scattering theory (Pu et al., 2021) and further assessed through finite element (FE) numerical analyses (Palermo et al., 2018b; Zeighami et al., 2019). The attenuation performance of a metabarrier was verified in small-scale laboratory tests for shear vertical (Palermo et al., 2016; Zaccherini et al., 2020a) and shear horizontal surface waves (Zaccherini et al., 2020b). Since then, metabarriers have been associated with various applications in waveguiding (Maznev and Gusev, 2015), wave filtering (Colombi

35 et al., 2016c), wave focusing (Colombi et al., 2016b), and energy harvesting (De Ponti
36 et al., 2020).

37 After the introduction of the metabarrier concept in civil and material engineering
38 communities, various design strategies have been proposed to enhance their efficiency
39 in terms of seismic wave attenuation. These include the exploitation of mechanical
40 oscillators (Palermo et al., 2016), resonant pillars (Colombi et al., 2016a), or locally
41 resonant inclusions (Zeighami et al., 2021a) as the fundamental unit of metabarriers.
42 Several studies are then keyed to the assessment of optimal spatial arrangement of
43 mechanical resonators. These resonant elements can be either installed at the soil
44 surface (Boutin and Roussillon, 2006) or buried inside soil layers (Zaccherini et al.,
45 2020a; Zeighami et al., 2021a). To the best of our knowledge, previous literature
46 studies evaluate the seismic wave attenuation associated with the metabarriers by
47 considering their design parameters as deterministic quantities. Otherwise, in the
48 context of geophysical sciences, various types of uncertainties are associated with
49 the soil system and mechanical resonators. Henceforth, the aim of this research is
50 two-fold: (i) to identify the uncertain parameters of the coupled soil-barrier dynamic
51 system, and (ii) to quantify the influence of these uncertain parameters on the sur-
52 face wave attenuation performance of metabarriers measured as the transmission
53 coefficient (TC).

54 In this work, we employ global sensitivity analysis (GSA) techniques to quan-
55 tify the influence of the uncertainties associated with the mechanical properties of
56 metabarrier components on the seismic wave attenuation of the latter. We rely on
57 (i) the classical variance-based Sobol indices, which quantify the expected reduction
58 of a model output variance due to the knowledge of (or conditioning on) a parameter
59 value, and (ii) the moment-based AMA indices, which quantify the normalized ex-
60 pected deviation of the statistical moment of a model output due to the knowledge of
61 (or conditioning on) a parameter value. Sobol indices are broadly used in structural
62 engineering problems, including for instance load-carrying capacity analysis of steel
63 plane frames (Kala, 2011) and axial load evaluation of tie-rod elements (De Falco
64 et al., 2021). AMA indices (termed after the initials of the authors) have been re-
65 cently proposed by Dell’Oca et al. (2017). These global sensitivity metrics recognize
66 that the uncertainty of a model parameter can be imprinted onto diverse statisti-
67 cal moments of model outputs. They have been applied to quantify uncertainty in
68 different civil engineering scenarios, including degradation of contaminants in soils
69 (la Cecilia et al., 2020), groundwater flow (Bianchi Janetti et al., 2019), and gas flow
70 in low permeable materials (Sandoval et al., 2022). In the context of geotechnical
71 and earthquake engineering, other global sensitivity analysis techniques have been
72 employed for uncertainty quantification of layered periodic foundations (Liu et al.,

73 2020), offshore wind turbine foundations (Velarde et al., 2019), and ground motion
74 modeling in seismic risk assessment (Vetter and Taflanidis, 2012).

75 Successful application of GSA techniques typically requires multiple evaluations
76 of the model to be analyzed. In cases where such evaluation is associated with a
77 heavy computational burden, applying GSA may become unfeasible (Sudret, 2008).
78 In such a case, using a reduced complexity model minimizes the computational bur-
79 den associated with the evaluation of the model whilst preserving the relationships
80 between the inputs and outputs of the model (Dell’Oca et al., 2017). In this study,
81 we rely on polynomial chaos expansion (PCE) to construct models of reduced com-
82 plexity, our approach being otherwise fully compatible with other model reduction
83 techniques. PCE-based techniques have been widely employed in studies to quantify
84 uncertainty in diverse civil engineering areas such as dam engineering (Hariri-Ardebili
85 and Sudret, 2020), hydraulic fracturing operations (Gläser et al., 2016), and CO₂ se-
86 questration (Zhang and Sahinidis, 2013).

87 The paper is structured as follows: first, the analytical expression underlying the
88 design of seismic metabarriers is given in Sec. 2.1. A two-dimensional Finite Element
89 (FE) model of the metabarrier basic module is developed in Sec. 2.2 to demonstrate
90 the dispersive features of a single resonator. Then, a full numerical model of the
91 entire barrier is realized to assess its seismic isolation efficiency in Sec. 2.3. Next,
92 uncertain model parameters influencing the attenuation efficiency of a metabarrier
93 are introduced in Sec. 2.4. Secs. 2.5 and 2.6 describe the GSA and surrogate model-
94 ing techniques employed in the study. The accuracy of the surrogate models, as well
95 as the results of the GSA, are presented in Sec. 3. Finally, conclusions and future
96 research directions are addressed in Sec. 4.

98 2. Methodology and materials

99 In this Section, we explain the design methodology of a locally resonant metabar-
100 rier rationally engineered to impede the propagation of seismic surface waves. Metabar-
101 riers are composed of a series of passive mechanical resonators organized in a regular
102 grid with identical spacing to be directly placed over the soil surface. The metabarrier
103 is installed in the proximity of a target structure/infrastructure and activated by the
104 motion of incoming waves. The dynamic coupling between a seismic metabarrier and
105 vertically-polarized seismic waves (also called Rayleigh waves) is studied analytically
106 via a closed-form dispersion law proposed by Palermo et al. (2016), where resonators
107 are assumed to have a linear elastic behavior and the soil is modeled as an isotropic
108 and homogeneous half-space. In this dynamic system, the exchange of stress between

109 resonators and the soil generates a low-frequency band gap (BG) in the frequency
110 spectrum of surface waves, where a significant ground-motion attenuation is expected
111 (Colquitt et al., 2017). The band gap width depends on the operating frequency and
112 mass per unit area of the resonators (Palermo et al., 2016). Since no surface mode
113 can propagate within the band gap frequency region, the elastic energy of the seismic
114 waves diverges from the soil surface trajectory to the bulk media (Colquitt et al.,
115 2017).

116 Recently, more complex theoretical models have been developed to account for the
117 soil heterogeneity (Zeng et al., 2022) and its non-linear behavior (Kanellopoulos et al.,
118 2022), as well as exploiting non-linear resonators inside the metabarrier arrangement
119 (Palermo et al., 2022). For the case of stratified soil, the emergence of the band
120 gap in the dispersion relation is discarded by the propagation of higher-order surface
121 modes. Thus, surface-to-shear wave conversion is hindered. However, a substantial
122 ground motion reduction is observed around the collective resonant frequencies of
123 the resonators (Palermo et al., 2018a; Zeng et al., 2022). These additional complex-
124 ities are usually ignored in the long-wavelength (low-frequency) regime to develop
125 theoretical frameworks that capture the fundamental physics of wave propagation
126 settings. Hence, we resort to a linear elastic resonator and a linear, homogeneous,
127 and isotropic soil model in this study.

128 The overall workflow and research methodology are depicted in Fig. 1 and dis-
129 cussed extensively in the following. We first present the analytical dispersion law
130 for the seismic metabarrier. We then illustrate the seismic isolation assessment of
131 the barrier through a surrogate model and by a global sensitivity analysis of the
132 uncertain parameters.

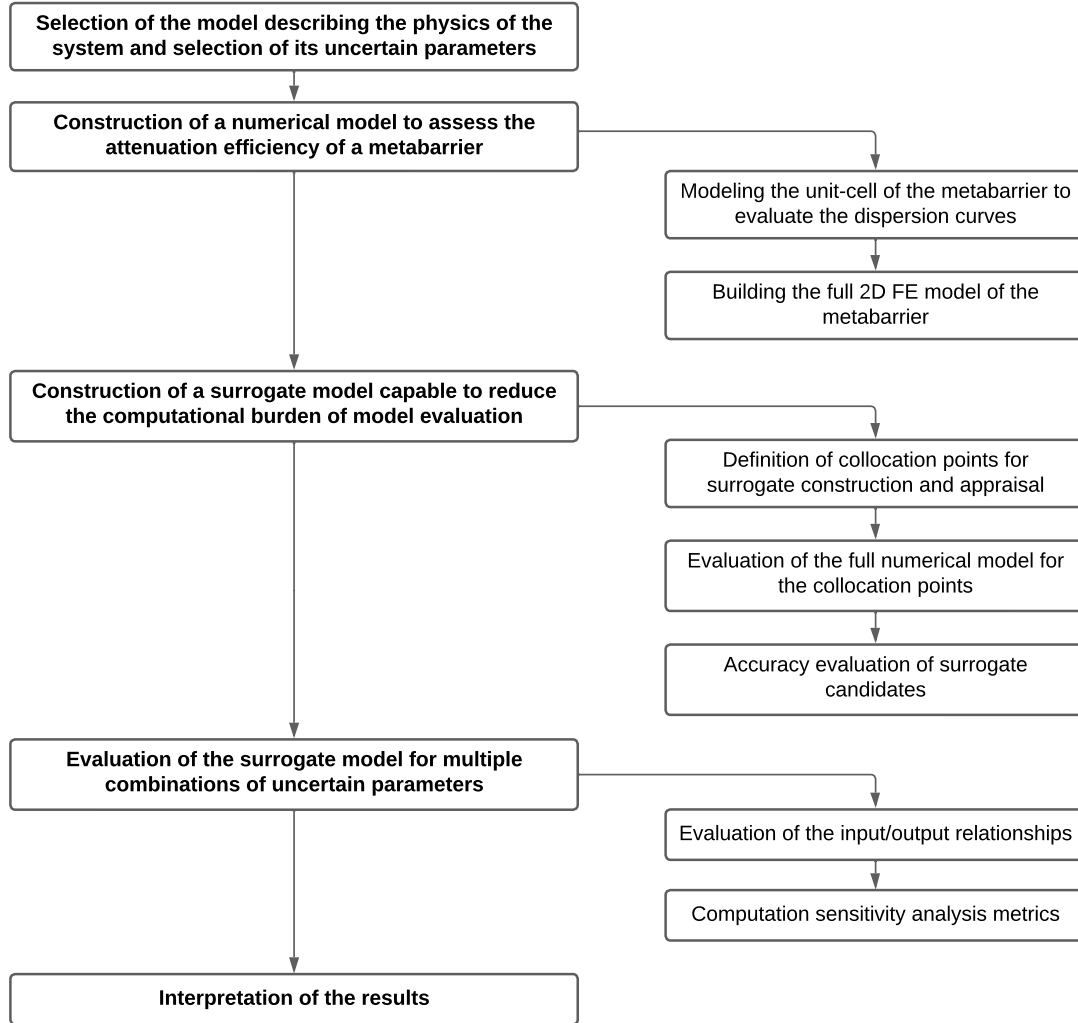


Figure 1: Global sensitivity analysis framework used to assess seismic isolation performance of a compact seismic metabarrier.

133 *2.1. Analytical dispersion relation of seismic metabarriers*

134 We study Rayleigh wave propagation through an elastic half-space equipped with
 135 a finite-size array of surface resonators, as shown in Fig. 2. Rayleigh waves are
 136 propagating in the x -direction, while they are polarizing in the z -direction. We note
 137 that parameters related to resonator and soil are denoted with subscription r and
 138 s , respectively. The resonators have a length of a_r and an out-of-plane depth of l_r

139 that assumes a unit value in a 2D problem. As such, the resonators have dimensions
 140 much smaller than the wavelength of surface Rayleigh waves, i.e. $a_r \ll \lambda_{RW}$, and
 141 they are distributed in a regular arrangement with an equal spacing of a_r , identical
 142 to the resonator length. Hence, the resonator influence area is $A_r = a_r \times l_r$.

143 The half-space is constituted by an isotropic and homogeneous soil with density
 144 ρ_s , Poisson's ratio ν_s , Young's modulus E_s , and shear modulus $G_s = E_s/(2(1 + \nu_s))$.
 145 Each resonator of the metabarrier has a mass m_r linked to the soil surface via linear
 146 elastic springs with axial stiffness K_r . Relying on a 2D instead of a 3D model can
 147 enable one to capture the main physics of resonator-soil coupling scenarios (see, e.g.
 148 (Palermo et al., 2016; Colquitt et al., 2017; Palermo et al., 2018b)).

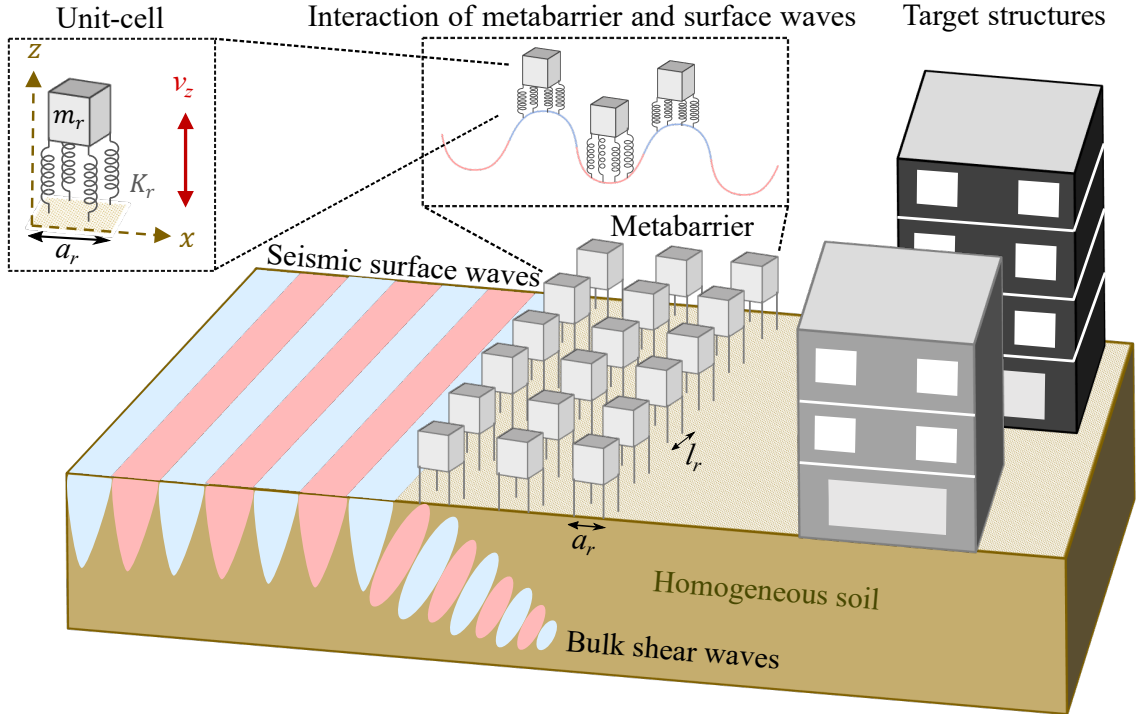


Figure 2: Schematic of a seismic metabarrier composed of surface resonators interacting with incoming seismic surface waves. The inset shows a unit-cell of the metabarrier consisting of a rigid mass and four elastic springs.

149 The interaction between the sub-wavelength resonators of a metabarrier, the elastic
 150 soil domain, and surface Rayleigh waves is defined via the effective medium approach
 151 (Maznev and Gusev, 2015). According to the latter, a metabarrier can be
 152 seen as a thin resonant boundary layer that exerts uniform vertical stress on the soil

153 surface. This approximation allows for deriving a dispersion law (i.e., the relation-
 154 ship between the wavenumber and frequency space) that predicts the fundamental
 155 dispersive features of metabarriers and guides their design procedure. We exploit the
 156 dispersion relation originally developed by Boechler et al. (2013), and update it in
 157 terms of the elastic modulus of a homogeneous soil and of the mechanical parameters
 158 of the resonators as:

$$\left(\frac{m_r \omega^2}{K_r} - 1\right) \left[\left(2 - \frac{\rho_s \omega^2}{G_s k^2}\right)^2 - 4 \sqrt{1 - \frac{\rho_s C^2 \omega^2}{G_s k^2}} \sqrt{1 - \frac{\rho_s \omega^2}{G_s k^2}} \right] = \frac{m_r \rho_s \omega^4}{a_r l_r G_s^2 k^3} \sqrt{1 - \frac{\rho_s C^2 \omega^2}{G_s k^2}}, \quad (1)$$

159 where $\omega = 2\pi f$ [rad/s] represents angular frequency, f [Hz] is frequency, k [rad/m] is
 160 a wavenumber vector that can vary from 0 to the edge of the first Brillouin zone (i.e.,
 161 π/a_r), and $C = \sqrt{(1 - 2\nu_s)/(2 - 2\nu_s)} = c_T/c_L$ is a dimensionless quantity expressing
 162 the ratio between the shear and longitudinal wave velocities. The dispersion relation
 163 can be formulated either in terms of Lamé parameters by exploiting the expressions
 164 $\mu = G_s$ and $\lambda = 2G_s/(1 - 2\nu_s)$ or via the longitudinal and shear wave speeds, whose
 165 expressions are given in Eqs. (2).

$$c_L = \sqrt{\frac{2G_s(1 - \nu_s)}{\rho_s(1 - 2\nu_s)}}, \quad c_T = \sqrt{\frac{G_s}{\rho_s}}. \quad (2a-b)$$

166 Eq. (1) shows that the wavenumber (k) is a function of the frequency (f),
 167 the mechanical parameters of the resonator (m_r, a_r, l_r, K_r), and the soil param-
 168 eters (ρ_s, G_s, ν_s). For the sake of completeness, the dispersion relation (Eq. (1)) is
 169 recast in a dimensionless format in Appendix A.

170 2.2. Numerical assessment of the metabarrier through a Wave Finite Element Method

171 Our study relies on the Wave Finite Element Method (WFEM), initially proposed
 172 by Mace and Manconi (2008) and further developed by Palermo et al. (2018b) to
 173 assess the ground vibration attenuation capability of locally resonant metabarriers.
 174 The WFEM rests on the conventional finite element model of a small portion of the
 175 composite waveguide (namely the unit-cell) to quantify its dynamic response against
 176 incoming waves. So far, this method has been applied to various composite structures
 177 such as beams, pipes, laminated plates, sandwich panels, and thin-walled structures.
 178 The efficiency of the WFEM for the case of seismic surface wave propagation through
 179 a finite-length barrier placed on top of a soil column has been assessed in recent

180 studies (Zeighami et al., 2019, 2021b,a; Palermo et al., 2022). According to WFEM,
 181 a numerical model of a single resonator linked to a 2D soil column can be envisaged to
 182 satisfy the analytical dispersion law of Eq. (1). Such an approach allows modeling the
 183 entire barrier to evaluate its seismic isolation performance. The coupled resonator-
 184 soil column represents a unit-cell (fundamental module) of the barrier. The dynamic
 185 response of a finite-size chain of unit-cells obtained from frequency domain analysis
 186 yields the numerical dispersion curves. The accuracy of the developed FE model will
 187 then be verified against analytical solutions of Eq. (1).

188 We develop a two-dimensional FE model of a metabarrier unit-cell according to
 189 a realistic engineering design of a metabarrier. The unit-cell (see Fig. 3a) consists
 190 of a rigid mass (m_r) attached to a soil column via two vertical elastic springs, each
 191 having an axial stiffness of $K_r/2$. The length of the soil column is being taken to
 192 coincide with the resonators' spacing, a_r . The soil depth is considered large enough
 193 to mimic a semi-infinite soil media as $3\lambda_0$, where $\lambda_0 = c_R/f_r$, c_R is the Rayleigh wave
 194 velocity, and $f_r = (1/2\pi)\sqrt{K_r/m_r}$ is the resonant frequency of the resonators. The
 195 vertical and horizontal displacements of the soil bottom are restricted to avoid any
 196 undesirable motion. The horizontal displacement of the resonator mass is suppressed
 197 while the springs are allowed to elongate and compress along their vertical axis.
 198 Bloch periodicity conditions (Brillouin, 1946) are imposed on the lateral edges of the
 199 soil column to construct the numerical dispersion curve. The resonator mass and
 200 soil domain are discretized via triangular mesh elements with a minimum dimension
 201 $L_m = \lambda_0/10$. Each spring is represented by a single truss element with one node at
 202 each joint. We seek the eigenfrequency solutions of the unit-cell in the wavenumber
 203 interval of $k = [0, \pi/a_r]$.

204 *2.3. Application to a finite-size metabarrier attached to different homogeneous soil* 205 *layer*

206 We implement WFEM on a compact metabarrier to assess its seismic surface
 207 waves isolation efficiency. The proposed metabarrier is composed of a finite number
 208 of equidistant resonators distributed over a total length λ_0 which is equivalent to an
 209 array of 20 unit-cells with an equal spacing $a_r = 1$ m (see Fig. 3b). We note that this
 210 is the minimum length of the barrier enabling one to detect the attenuation effect.
 211 Increasing the length of the barrier (or, conversely, the number of resonators) would
 212 result in an enhanced attenuation (Pu et al., 2021). The full 2D model of the barrier
 213 is developed on the basis of the reduced numerical model of the unit-cell (see Fig.
 214 3a); its height is equal to the depth of unit-cell ($3\lambda_0$) and has a total length of $18\lambda_0$, as
 215 shown in Fig. 3c. A harmonic input source excites the entire domain. The left part
 216 of the domain represents a reference soil (i.e., soil without metabarrier), whereas the

217 right side of the domain includes the metabarrier zone. The metabarrier is placed
 218 at a distance of $5\lambda_0$ from the input source to eliminate near-source effects. Both
 219 bottom corners of the soil domain are fixed to maintain the static stability of the
 220 model during simulations. In addition, Low-Reflective Boundary (LRB) conditions
 221 are imposed on the lateral and bottom edges of the model to minimize the back-
 222 reflection of surface waves from the boundaries. The barrier response is extracted
 223 from an output region with a length $L_{out} = \lambda_0$ placed after the metabarrier zone,
 224 namely barrier output. The same quantity is measured for the reference soil from
 225 the left part of the model (see soil output in Fig. 3c) to compare the soil response
 226 equipped with resonators with bare soil condition.

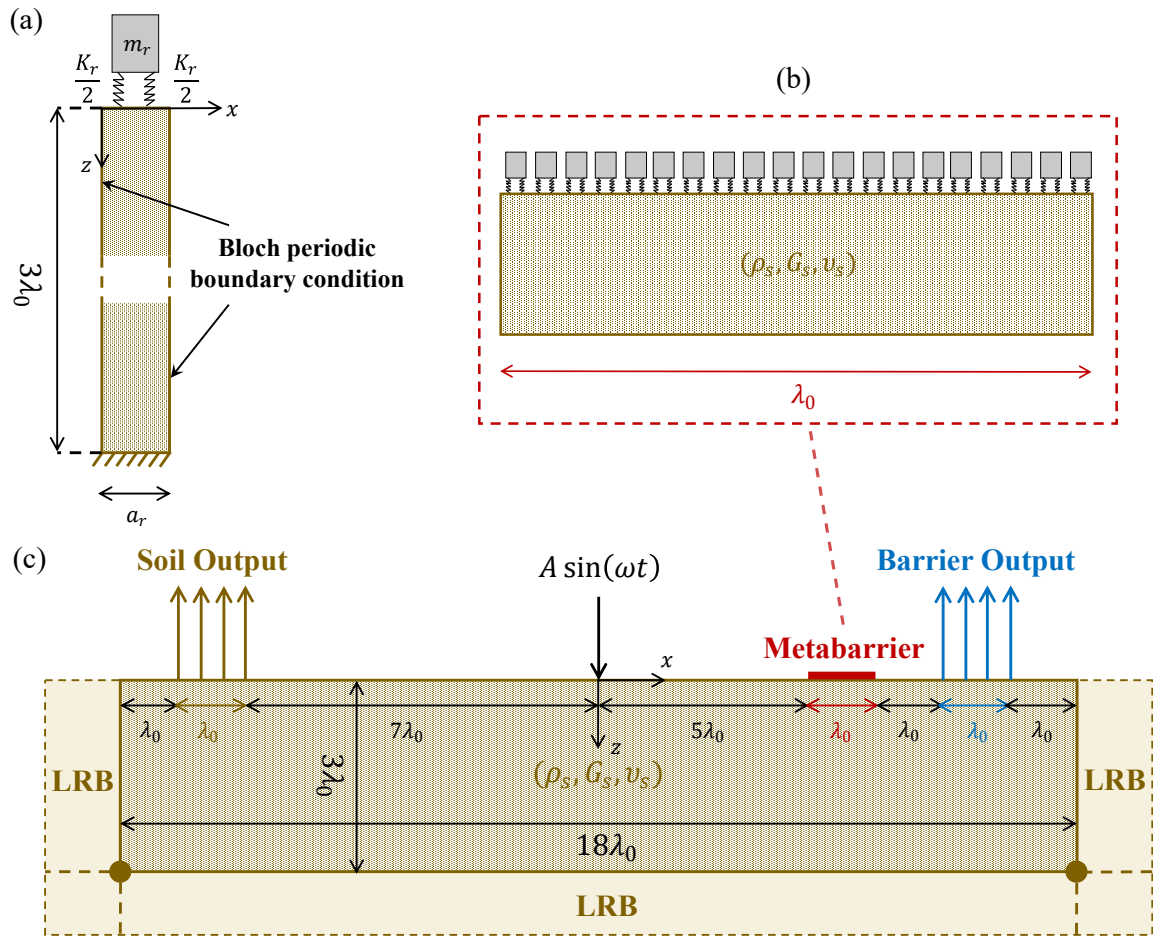


Figure 3: Schematics of the seismic metabarrier FE model. (a) A representative unit-cell of the metabarrier. (b) An array of 20 resonators forms the metabarrier. (c) The 2-D wavefield model used to calculate the Transmission Coefficient (TC).

227 To quantify the seismic isolation performance of the barrier, one can either per-
 228 form a time-history analysis and measure the soil response via Fourier transform
 229 (Zeighami et al., 2021b) or perform harmonic analysis to explicitly obtain the re-
 230 sult in the frequency domain (Palermo et al., 2018b). In this study, we employ
 231 the latter approach (*i*) to understand the physics of the problem by evaluating the
 232 transmission/attenuation performance of the metabarrier at different frequencies and
 233 comparing the results with their counterparts obtained from the dispersion analy-
 234 sis and (*ii*) to minimize the computational cost of simulations. In this context, a
 235 harmonic displacement excites the model from the center to evaluate the Transmis-
 236 sion Coefficient (TC) and Attenuation Coefficient (AC) of the metabarrier (Palermo
 237 et al., 2018b) as:

$$TC(f) = \frac{\int_0^{L_{out}} |v_b| dx}{\int_0^{L_{out}} |v_s| dx}, \quad AC(f) = 1 - TC(f), \quad (3a-b)$$

238 where v_b is the vertical nodal displacement of the soil measured from barrier
 239 output (see Fig. 3c), to be averaged over the output length L_{out} , and v_s is the same
 240 quantity evaluated from the reference soil in the absence of the metabarrier.

241 2.4. Uncertain model parameters

242 During the design and implementation phases of a seismic metabarrier, the out-of-
 243 plane depth (l_r) and the resonators spacing (a_r) are typically considered as determi-
 244 nistic parameters. The mass (m_r) and stiffness (K_r) of the resonator can be regarded
 245 as uncertain parameters due to manufacturing imperfections in their geometries or
 246 weight differences. To streamline the analysis, we consider mass as the only uncer-
 247 tain parameter of the resonator, since the mass per unit area ratio ($m_r/(a_r l_r)$) in
 248 Eq. (1) controls the hybridization of the fundamental surface mode (Palermo et al.,
 249 2016). Regarding the mechanical parameters of the soil, these data are generally
 250 obtained from experiments. Thus, their estimates are typically uncertain, even for
 251 homogeneous and isotropic geomaterials. Since the variation of the Poisson ratio (ν_s)
 252 is generally less pronounced than variations of mass density (ρ_s), and given that the
 253 shear modulus (G_s) is related to ν_s via the Young modulus (E_s), we consider ρ_s and
 254 G_s as the uncertain parameters of the soil and set ν_s as a deterministic parameter.
 255 In summary, the uncertain model parameters of the present study are ρ_s , G_s , and
 256 m_r .

257 To rigorously assess the significance of uncertain parameters on the isolation
 258 properties of metabarriers, we consider three common soil scenarios found in nature.

259 We assume that such soils have homogeneous mechanical properties throughout the
 260 depth of the elastic waveguide. The soil types are (i) sedimentary soil (S1), (ii) com-
 261 pletely weathered granite (S2), and (iii) silty-clay soil (S3). These soil types are not
 262 ideal for construction purposes and require engineering intervention to increase their
 263 bearing capacities. Sedimentary soils are loose sediments usually found near river
 264 basins. Weathered granite soils are found in mountainous areas where infrastructures
 265 (dams, mountain roads, and railways) are built. Silty clay soil is an intermediate be-
 266 tween sandy and clay soils that tends to shift due to moisturizing/drying; therefore,
 267 they require deep foundations to protect the infrastructures against seismic action.

268 In this study, the uncertain model parameters are considered independent and
 269 identically distributed random variables, each characterized by a uniform distribution
 270 within the intervals listed in Table 1. The choice of the latter distribution enables
 271 one to give the same weight to all parameter values across their support. We assume
 272 that the ranges of the variability of ρ_s and G_s are centered around the values ex-
 273 perimentally found by Cai et al. (2021), who documented Rayleigh waves velocities
 274 $c_R = [94.91, 64.70, 55.80]$ m/s for S1, S2, and S3, respectively. An effective coupling
 275 between the metabarrier and the soil is seen for sedimentary soils with $c_R < 1000$
 276 m/s (Colombi et al., 2016c). Even as values of the soil parameters of the current
 277 study (see Table 1) are representative of sedimentary-basins-like (soft) soils with low
 278 Rayleigh wave speeds, the sensitivity of the results discussed in Sec. 3 are not strictly
 279 limited to these wave velocity ranges, and can be extended in future studies to the
 280 soils with average Rayleigh celerity ($300 \leq c_R \leq 500$ m/s).

281 We consider the uniform distributions of ρ_s and G_s to be characterized by a
 282 coefficient of variation of 10% to encompass a range of values typical of common
 283 engineering applications. The lower and upper bound of the support associated with
 284 the distribution of the resonator mass are defined upon considering a 5% coefficient
 285 of variation. The mean value of the resonator mass is assumed to be 1500 kg. We
 286 consider a target resonant frequency of $f_r = 5$ Hz for the incoming waves. This, in
 287 turn, leads to an axial stiffness of the resonator $K_r = 1480$ kN/m and $\lambda_0 = 19$ m,
 288 13 m, and 11 m for S1, S2, and S3, respectively. Table 1 lists, for each Scenario,
 289 the range of variability of the model’s uncertain parameters and the value of the
 290 deterministic parameters.

291 2.5. Global sensitivity analysis (GSA)

292 We employ GSA techniques to diagnose the influence that the uncertainties on
 293 model parameters (ρ_s , G_s , m_r) have on the attenuation performance of the metabar-
 294 rier (measured as TC and evaluated via Eq. 3a). As stated in Section 1, we rely on (i)
 295 the classical variance-based approach grounded on the evaluation of the well-known

Model parameters			Scenario 1	Scenario 2	Scenario 3
Parameter	Units	CV [%]	Range/Value		
G_s	MPa	10	13.5 - 16.5	9.36 - 11.44	6.14 - 7.50
ρ_s	kg/m ³	10	1350 - 1650	1890 - 2310	1665 - 2035
m_r	kg	5	1425 - 1575	1425 - 1575	1425 - 1575
ν_s	-	-	0.45	0.25	0.32
K_r	kN/m	-	1480		
a_r	m	-	1		
l_r	m	-	1		

Table 1: Ranges of variability for the model uncertain parameters considered in the GSA and values of deterministic model parameters considered in this study. Values of the coefficient of variation of the uncertain model parameters are also listed.

296 Sobol indices (Saltelli and Sobol', 1995) and (ii) the moment-based GSA framework
 297 introduced by Dell'Oca et al. (2017).

298 2.5.1. Variance-based Sobol Indices

299 Sobol indices (Saltelli and Sobol', 1995) quantify the relative expected reduction
 300 of the variance of a target model output due to knowledge of (or conditioning on)
 301 an uncertain model parameter. In this context, considering a model output ζ , which
 302 depends on P random parameters collected in vector $\boldsymbol{\theta} = (\theta_1, \theta_2, \dots, \theta_P)$ and defined
 303 within the space $\Gamma = \Gamma_1 \times \Gamma_2 \times \dots \times \Gamma_P$ ($\Gamma_p = [\theta_{p,min}, \theta_{p,max}]$ corresponding to the
 304 support of the p -th parameter, θ_p), the principal Sobol' index S_{θ_p} associated with a
 305 given model parameter θ_p is evaluated as

$$S_{\theta_p} = \frac{V[E[\zeta|\theta_p]]}{V[\zeta]}. \quad (4)$$

306 Here, $E[\cdot]$ and $V[\cdot]$ represent expectation and variance operators, respectively; the
 307 notation $\zeta|\theta_p$ denotes conditioning of ζ on a value of θ_p . Note that S_{θ_p} describes
 308 the relative contribution to $V[\zeta]$ due to variability of only θ_p . Joint contributions of
 309 θ_p with other model parameters included in $\boldsymbol{\theta}$ to the variance of ζ are embedded in
 310 the total Sobol indices (details not shown). Note that by relying on Sobol indices
 311 to diagnose the relative importance of uncertain model parameters to model outputs
 312 one assumes that the uncertainty of a model output is completely characterized by its
 313 variance. Thus, even though Sobol indices are characterized by conceptual simplicity
 314 and straightforward implementation and use, they provide only limited information
 315 about the way variations of model parameters can influence the complete probability

316 density function (pdf) of model outputs.

317 2.5.2. Moment-Based AMA Indices

318 The moment-based GSA approach proposed by Dell’Oca et al. (2017, 2020) rests
319 on the idea that quantifying the effects of model parameter uncertainty on various
320 statistical moments of the ensuing pdf of model outputs can provide an enhanced
321 understanding of model functioning. Dell’Oca et al. (2017) introduce Moment-Based
322 sensitivity metrics (termed AMA indices) according to which one can evaluate the
323 influence of uncertain model parameters on key elements of the model output pdf,
324 as embedded in its associated statistical moments. The AMA indices are defined as
325 follows (Dell’Oca et al. (2017)):

$$\text{AMAM}_{\theta_p} = \frac{1}{|M[\zeta]|} E [|M[\zeta] - M[\zeta|\theta_p]|]. \quad (5)$$

326 Here, AMAM_{θ_p} represents the indices associated with a model parameter θ_p and
327 a given statistical moment M of the pdf of model output ζ . For the purpose of our
328 study we focus on the first two moments (i.e., the mean ($M = E$) and the variance
329 ($M = V$)) of the model output pdf. The AMA indices are intended to quantify
330 the relative importance of parameter θ_p on a given statistical moment of ζ . Large
331 values of these indices indicate that $\zeta|\theta_p$ strongly deviates from its unconditional
332 counterpart.

333 2.6. Surrogate model

334 To employ the previously described GSA techniques, several evaluations of TC
335 under diverse combinations of ρ_s , G_s , and m_r are required. Such a procedure is
336 impractical in our scenario due to the heavy computational burden associated with
337 the evaluation of TC . One single simulation takes approximately 72 seconds on
338 an Intel Core i7-116G7 @ 2.80GHz with 32GB of Memory. Thus, here we rely on
339 a generalized Polynomial Chaos Expansion (gPCE) surrogate of the full numerical
340 model that allows for reducing the computational time associated with the execution
341 of the GSA technique (Dell’Oca et al., 2017; Sudret, 2008).

342 In the context of gPCE, a model $g(\boldsymbol{\theta})$ can be expressed as a linear combination
343 of the multivariate polynomials, $\psi_i(\boldsymbol{\theta})$, i.e.,

$$\begin{aligned}
g(\boldsymbol{\theta}) &\approx \sum_{i \in \Lambda^{P,D}} \beta_i \psi_i(\boldsymbol{\theta}), \\
\psi_i(\boldsymbol{\theta}) &= \prod_{p=1}^P \psi_p^d(\theta_p).
\end{aligned} \tag{6}$$

344 Here, β_i is the coefficient of the i -th term of the model surrogate; $\psi_p^d(\theta_p)$ is
345 a univariate polynomial of order d of the parameter θ_p ; and $\Lambda^{P,D}$ is a multi-index
346 containing the indices of all the multivariate polynomials ($\psi_i(\boldsymbol{\theta})$) with degree equal or
347 smaller than the surrogate degree, D , (i.e., multivariate polynomials where $\sum_{p=1}^P d \leq$
348 D).

349 Note that in the context of gPCE the univariate polynomials must satisfy the
350 orthonormality condition, i.e., $E[\psi_p^j \psi_p^k] = \delta_{jk}$, where δ_{jk} is the Kronecker-delta func-
351 tion, $\delta_{jk} = 1$ if $j = k$ and zero otherwise. Multiple families of polynomials satisfy this
352 condition; however, the selection of the suitable family of polynomials is made based
353 on the pdf of the model parameters, which in this study are considered uniform.
354 Thus, the Legendre polynomial family is employed to construct the surrogates.

355 The construction of a gPCE surrogate requires the evaluation of the surrogate
356 model coefficients, $\beta = \{\beta_i, \forall i \in \Lambda^{P,D}\}$, and the selection of the surrogate model
357 degree, D (Sudret, 2008). Regarding the evaluation of β , we rely on least-square
358 minimization (also termed as regression approach). According to this technique,
359 the surrogate coefficients β are those that minimize the mean square error between
360 TC values computed with the full numerical model, $y(\boldsymbol{\theta})$, and the corresponding
361 outputs of the surrogate model. Thus, several full numerical simulations need to
362 be performed in order to estimate the coefficients β . Generally, as the number of
363 full numerical simulations employed for the construction of the surrogate increases,
364 also the accuracy of the surrogate increases. In this study, the maximum admissible
365 computational burden allows us to perform 1233 full model evaluations encompassing
366 an equal number of randomly selected sets of parameters, such sets of parameters are
367 randomly sampled employing a Quasi-Monte Carlo approach which guarantees that
368 the parameter space is sampled uniformly. The evaluation of β is then performed by
369 minimizing

$$\sum_{s=1}^{1233} \left[g(\boldsymbol{\theta}_s) - \sum_{i \in \Lambda^{P,D}} \beta_i \psi_i(\boldsymbol{\theta}_s) \right]^2, \tag{7}$$

370 where $\boldsymbol{\theta}_s$ is the s -th randomly selected set of the uncertain model parameters.

371 In our analyses, the selection of D is performed on the basis of an accuracy test
 372 of surrogates with degrees varying between 4 and 13. In such a test, the TC of 50
 373 randomly selected sets of parameters (different from the sets of parameters employed
 374 for the estimation of β) is evaluated with the numerical model and the surrogate.
 375 Then, the mean absolute error between these two quantities is evaluated, and the
 376 surrogate associated with the smallest error is selected and employed for the GSA.

377 3. Results and Discussion

378 This Section provides the resulting dispersion curves of each soil Scenario analyzed
 379 for a single resonator and the transmission coefficients of the entire metabarrier. The
 380 accuracy of the surrogate models generated for the GSA is then discussed. Finally,
 381 the GSA results are presented, and some conclusions about the impact of uncertain
 382 input parameters on the output results are drawn.

383 3.1. Dispersion analysis results

384 Analytical dispersion curves for each Scenario from Eq. (1) are depicted in Fig. 4a
 385 (solid curves). The dispersion curve highlights the hybridization of the fundamental
 386 surface mode around the local resonance of the resonators in two avoided-crossing
 387 branches observed in previous studies (Boechler et al., 2013; Palermo et al., 2016;
 388 Colquitt et al., 2017). The split of the fundamental mode results in the generation
 389 of a low-frequency band gap typical of the local resonance mechanism, where the
 390 propagation of seismic surface waves is impeded within this frequency range. Having
 391 an identical mass and stiffness of the resonators, a comparison between different soil
 392 Scenarios reveals that the dynamic coupling between surface waves and metabarrier
 393 is stronger for the silty clay soil (S3). This is due to the lower relative density between
 394 resonator and soil, noting that Rayleigh and shear wave velocities are slower for silty
 395 clay soil in comparison with other Scenarios. Under this rationale, soil type 3 presents
 396 a flattened in-phase branch (lower branch in Fig. 4a) and propagates with higher
 397 velocity in the frequency ranges above the band gap, as evidenced by the larger slope
 398 of its out-of-phase branch (upper branch in Fig. 4a). The band gap width falls in the
 399 interval $f_{BG} = [f_r, f_r(\alpha + \sqrt{\alpha^2 + 1})]$, where $\alpha = (\pi m_r f_r) / (a_r l_r) \sqrt{(1 - C^2) / (\rho_s G_s)}$
 400 is a dimensionless number that relates resonator mass to the soil mass excited by
 401 Rayleigh waves at the resonance (Palermo et al., 2016). Consequently, the band gap
 402 width is $\Delta f_{BG} = \alpha - 1 + \sqrt{\alpha^2 + 1}$ which takes the values 0.80, 0.70, and 0.97 Hz for
 403 S1, S2, and S3, respectively. The largest band gap is associated with the silty clay
 404 soil, having the lowest shear modulus and Rayleigh velocity among all the cases.

405 We perform a frequency domain analysis of the FE unit-cell model (see Fig. 3a)
 406 in the frequency range of 0 to 10 Hz corresponding to Rayleigh wavelengths between

407 20 to 240 m in Comsol Multiphysics (COMSOL Multiphysics®), 2022) under plane-
 408 strain conditions. The simulations are performed for 1233 realization of collocation
 409 points as specified in Sec. 3.3, each containing a set of uncertain parameters. The
 410 simulation outcomes are averaged for each Scenario and depicted in Fig. 4a. The
 411 eigenfrequency solutions of the FE analysis are marked with dots and super-imposed
 412 on the analytical dispersion curves. There is a good agreement between the analyti-
 413 cal and numerical solutions. Otherwise, it can be noted that the numerical resonant
 414 frequencies are shifted toward the lower values. This frequency shift stems from
 415 uncertainties associated with the resonator mass or, equivalently, with the resonant
 416 frequency, together with a soft soil mechanism that emerged from the inertia dif-
 417 ference between the resonator and soil. For heavy resonating masses, the soil acts
 418 as a soft spring with the longitudinal stiffness of K_s working in serial configuration
 419 with resonators' springs K_r . The equivalent stiffness of the coupled system would
 420 be $K_{eq} = K_r K_s / (K_r + K_s) < K_r$. Thus, the numerical resonant frequency will be
 421 $f_{r,FE} = 2\pi K_{eq}^{1/2} m_r^{-1/2}$. Since $K_{S3} < K_{S2} < K_{S1}$, the frequency shift is less pro-
 422 nounced for the soft soil S1. The opposite can be observed for S3. The resulting
 423 numerical resonant frequencies are $f_{r,FE} = [4.5, 4.4, 4.3]$ Hz for S1, S2, and S3, while
 424 $f_{r,AN} = [4.91, 4.78, 4.84]$ Hz correspond to their analytically evaluated counterparts.

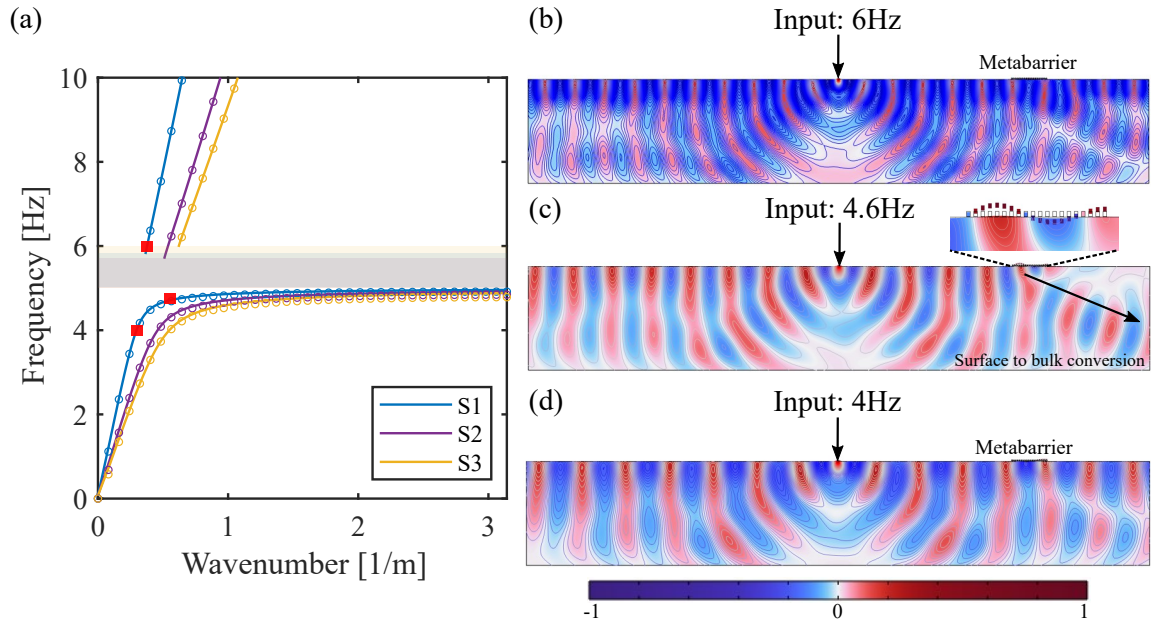


Figure 4: (a) Averaged dispersion curves obtained from the unit-cells of seismic metabARRIER for all three Scenarios S1-S3. Analytical surface modes obtained from Eq. (1) are shown with solid lines and FE eigensolutions super-imposed with circles. Highlighted boxes show the band gap region associated with each scenario. The full wavefields are corresponding to (b) out-of-phase surface mode (upper branch the dispersion curve), (c) resonant frequency, and (d) in-phase surface mode (lower branch) of S1. Excitation frequencies are marked with red squares on the dispersion curve.

425 The full wavefield of the plane-strain FE model (see Fig. 3c) is shown for the
 426 out-of-phase and in-phase surface modes of S1 in Fig. 4b and d, respectively. The
 427 introduction of the metabARRIER (right side of the individual figures) modifies the
 428 surface wave propagation compared to the reference soil condition (left side of the
 429 figures). The excitation frequencies are denoted with red squares in the dispersion
 430 curve of Fig. 4a. For a vertical harmonic excitation with a carrier frequency close to
 431 the operational frequency of the metabARRIER, the surface-to-shear wave conversion
 432 due to the local resonance of the resonators is observed in Fig. 4c. This phenomenon
 433 originates from the dynamic interaction between the harmonic motion of seismic
 434 surface waves (Rayleigh-like waves) and vertical displacement of resonators inside
 435 the barrier (see inset of Fig. 4c). The outcomes of the dispersion relation anticipate
 436 a considerable surface wave amplitude reduction around the band gap region. Similar
 437 results are obtained for S2 and S3 with minor differences in their numerical resonant
 438 frequencies.

439 3.2. The surface wave attenuation coefficient

440 We perform a harmonic analysis of the full FE model results of Fig. 3c in the
 441 frequency interval of 2 Hz to 7 Hz using a frequency resolution of 0.1 Hz. The
 442 triangular mesh elements with identical sizes of the unit-cell model are incorporated
 443 in all Scenarios to discretize the models. We obtained TC and AC (Eq. (3)) from the
 444 vertical nodal displacements of the barrier output averaged over its length ($L_{out} = \lambda_0$)
 445 divided by the same quantity measured from the soil output, as described in Sec.
 446 2.3. Similar to the unit-cell case, we execute the frequency domain simulations for
 447 1233 realization of collocation points and average the outputs of TC and AC for
 448 each Scenario. Fig. 5 summarizes the results of the harmonic analysis of the seismic
 449 metabarrier.

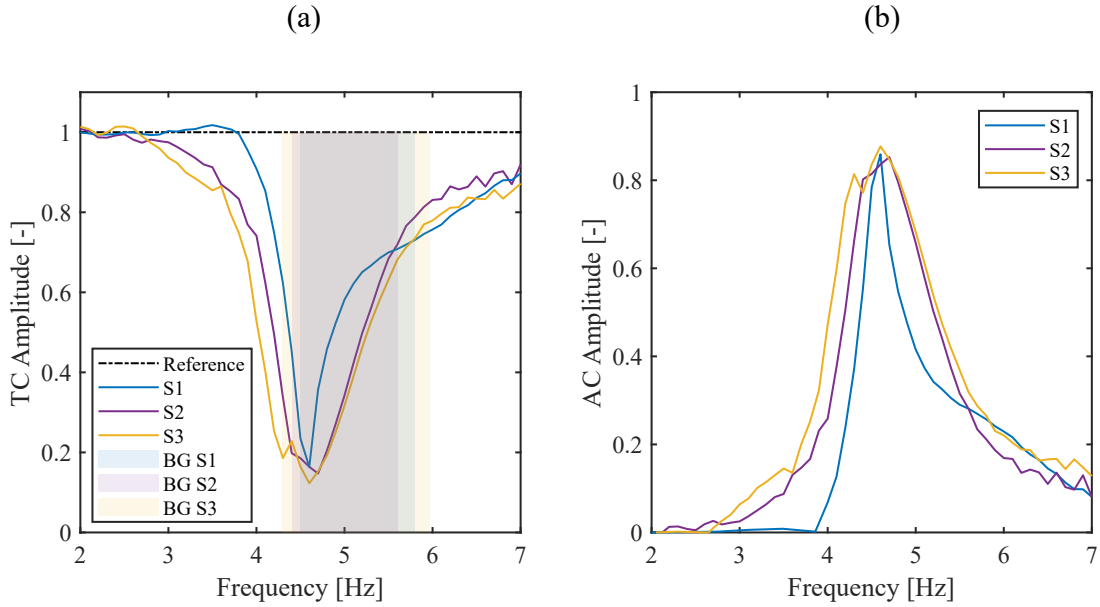


Figure 5: (a) Mean transmission and (b) mean attenuation coefficient of the proposed seismic metabarrier obtained from Eq. (3) in the low-frequency range for all realizations of each Scenario. Shaded areas in Figure 5a represent the band gap (BG) for each Scenario.

450 The responses of the soils equipped with the resonators are shown with solid lines,
 451 while the bare soil response (reference case) is marked with a solid-dashed line. The
 452 associated FE band gap regions of each Scenario are highlighted with rectangular
 453 boxes. The lower edges of BG zones coincide with the first peak attenuation of each
 454 soil model. In all Scenarios, the surface wave attenuation starts around 3 Hz. It
 455 becomes more indicative within the BG zone, while the peak attenuation (4.6, 4.7,

456 and 4.6 Hz for S1, S2, and S3) appears in the proximity of the collective resonant
 457 frequencies of the oscillators. In the frequency ranges above the resonance, the
 458 transmission/attenuation coefficients present an almost linear increasing/decreasing
 459 trend. In the high-frequency regime, the surface wave attenuation in the presence
 460 of the metabarrier approaches the reference soil case. The silty-clay soil (S3) is
 461 characterized by the broadest attenuation frequency range with the most significant
 462 peak analogous to the predictions of the dispersion curve previously discussed in Sec.
 463 3.1. Instead, the soft sedimentary soil (S1) has the narrowest attenuation zone as a
 464 result of higher relative density difference and weaker dynamic coupling between the
 465 metabarrier and the soil.

466 3.3. Surrogate models

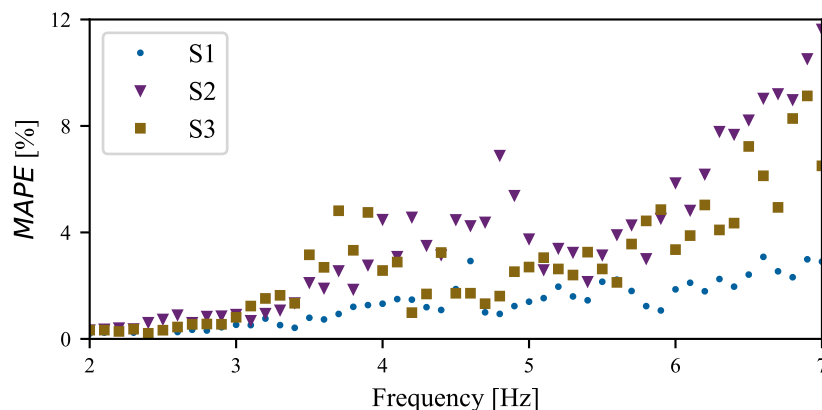


Figure 6: Accuracy of the best surrogate models for diverse frequency values and Scenarios considered in this study.

467 Fig. 6 presents the accuracy values of the surrogate models constructed for the
 468 Scenarios considered in this study. Accuracy is assessed by comparing the results
 469 of the surrogate predictions of TC with the results of full numerical simulations;
 470 the latter is different from the simulations employed for the surrogate construction.
 471 Accuracy is measured with $MAPE$ metric, which is defined as

$$MAPE = \sum_{s=1}^{50} \left[\frac{TC_{surr} - TC_{full}}{TC_{full}} \right]. \quad (8)$$

472 Note that we construct one surrogate per Scenario and per frequency. Therefore,
 473 Fig. 6 presents the accuracy of 153 surrogates. These results show that the accuracy

474 of the models is generally good ($MAPE \leq 5\%$), except for frequency values larger
 475 than 6 Hz. Regarding the different Scenarios, S1 is characterized by the lowest values
 476 of $MAPE$, whereas S2 is associated with the largest values.

477 We construct polynomials of order 12 for the 3 Scenarios. Surrogates of lower
 478 polynomial order are associated with reduced accuracy while relying on a higher order
 479 was not possible due to the ill-conditioning of the minimization problem required for
 480 the evaluation of the surrogate coefficients (details not shown).

481 3.4. Global sensitivity analysis

482 The surrogate models obtained in Section 3.3 are evaluated 10^6 times with ran-
 483 dom combinations of the model parameters. While performing such a task would
 484 have required a computational effort of more than two years by employing the full
 485 numerical model, relying on the surrogate model considered involves an overall com-
 486 putational time of about 60 seconds on a laptop with an Intel Core i7-8550 @ 1.8GHz
 487 with 8GB of Memory. Note that the number of model evaluations is defined to ensure
 488 the stability of the sensitivity metrics considered (details not shown). The results of
 489 these model evaluations are then employed for the global sensitivity analysis (GSA).

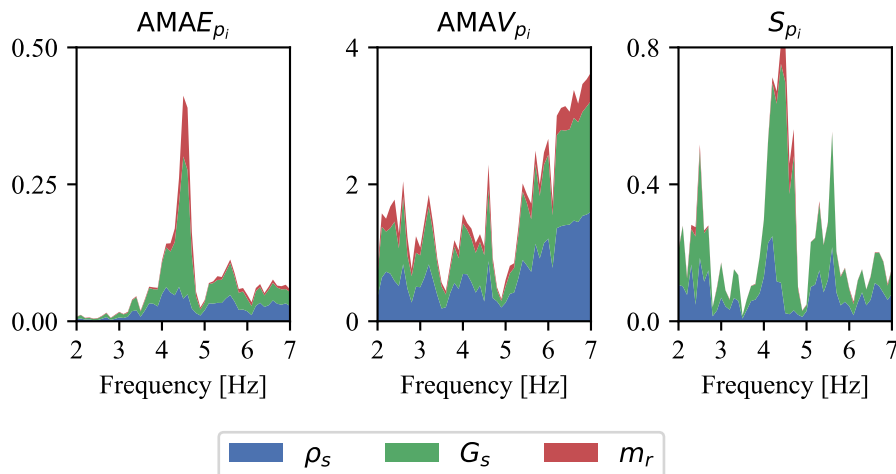


Figure 7: Stacked diagram of the first two moment-based AMA indices ($AMAE_{p_i}$ and $AMAV_{p_i}$), and Sobol principal indices (S_{p_i}) of the model uncertain parameters considered in this study for Scenario 1.

490 Fig. 7 depicts the evolution of the first two moment-based AMA indices and the
 491 variance-based Sobol indices of the uncertain model parameters of Scenario 1. The
 492 results are presented for frequency values in the interval between 2 and 7 Hz, for

493 which a unique GSA was performed every 0.1 Hz. GSA metrics of ρ_s , G_s , and m_r
494 are depicted in blue, green, and red, respectively. Values of AMAE index indicate
495 that for frequencies below 4 Hz and over 5 Hz, the mean of TC is practically inde-
496 pendent of the model parameter values. Conversely, in the interval between 4 Hz and
497 5 Hz (i.e., the frequency range near the local resonance of the resonators in which
498 the TC/AC assumes minimum/maximum values), knowledge of (or conditioning on)
499 model parameters may significantly modify the mean of TC . In this interval, all un-
500 certain model parameters play an important role in determining the mean of TC , the
501 parameters playing the most and least essential roles being G_s and ρ_s , respectively.
502 This suggests that for $4 \text{ Hz} \leq f < 5 \text{ Hz}$, the shear modulus of the soil, which controls
503 the bulk waves velocities and the resonator mass, plays an indispensable role in the
504 definition of the TC mean. For frequency values close to the resonance and due to
505 the effective coupling between resonators and soil, the influence of resonator mass
506 becomes more prominent. Within the band gap frequency range ($5 \text{ Hz} \leq f \leq 5.8 \text{ Hz}$),
507 the impact of all uncertain model parameters is less pronounced, and resonator mass
508 has the minimum contribution.

509 Values of the AMAV index suggest that the variance of the model output can be
510 significantly modified by the knowledge of (or conditioning on) a model parameter,
511 such variability being large for high-frequency values. In general, TC variance is
512 governed by G_s and ρ_s . The results of Sobol indices are consistent with those of
513 AMAE and AMAV indices. Sobol indices suggest that the variance of TC can
514 be reduced significantly by the knowledge of (or conditioning on) G_s and ρ_s . The
515 resonator mass (m_r) may play a non-negligible role also for $4 \text{ Hz} \leq f < 5 \text{ Hz}$.

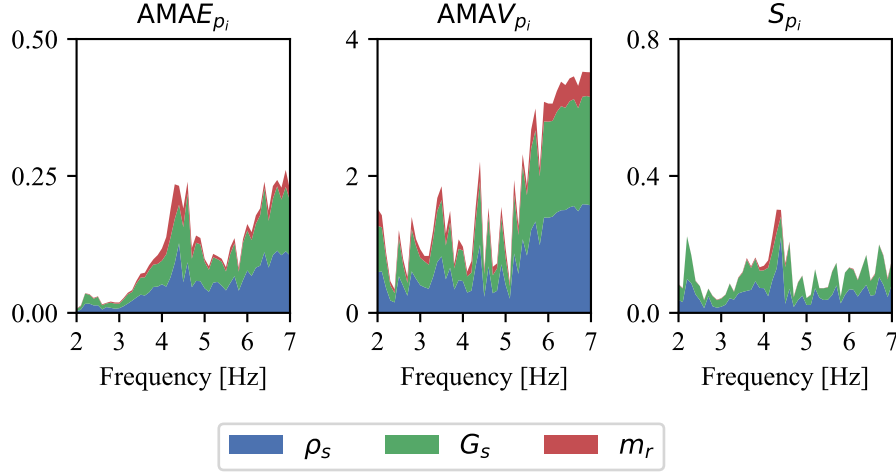


Figure 8: Stacked diagram of the first two moment-based AMA indices ($AMAE_{p_i}$ and $AMAV_{p_i}$), and Sobol principal indices (S_{p_i}) of the model uncertain parameters considered in this study for Scenario 2.

516 Fig. 8 presents the evolution of the first two moment-based AMA indices and the
517 variance-based Sobol indices of the uncertain model parameters of Scenario 2. GSA
518 metrics of ρ_s , G_s , and m_r are depicted in blue, green, and red, respectively. Values of
519 $AMAE$ index document that for frequencies below 3 Hz the mean of TC is practically
520 independent of the uncertain model parameter values. Differently from Scenario 1,
521 the knowledge of (or conditioning on) model parameters significantly influences the
522 value of the TC mean for almost all the considered frequencies. The most and least
523 influential parameters are ρ_s and m_r , respectively. Note that the relative importance
524 of ρ_s is close to the relative importance of G_s . m_r has a marginal contribution only
525 in the frequency range between 3.5 Hz and 4.5 Hz, where resonators start moving
526 in phase with the soil domain and ultimately reach the resonance condition. The
527 resonator mass does not have a remarkable impact on the determination of the TC
528 mean outside this frequency interval, and soil parameters govern the overall dynamic
529 behavior of the system. With reference to the $AMAV$ index, the results are similar
530 to those of Scenario 1. This suggests that the variance of the model output can be
531 significantly modified by the knowledge of (or conditioning on) a model parameter,
532 such variability being large for high-frequency values. In general, TC variance is
533 governed by G_s and ρ_s . The results of Sobol indices confirm the outcomes of $AMAE$
534 and $AMAV$ indices. Sobol indices indicate that the variance of TC can be reduced
535 significantly by the knowledge of (or conditioning on) G_s and ρ_s . As expected, the
536 resonator mass m_r comes into play in the frequency range of $4 \text{ Hz} \leq f < 5 \text{ Hz}$.

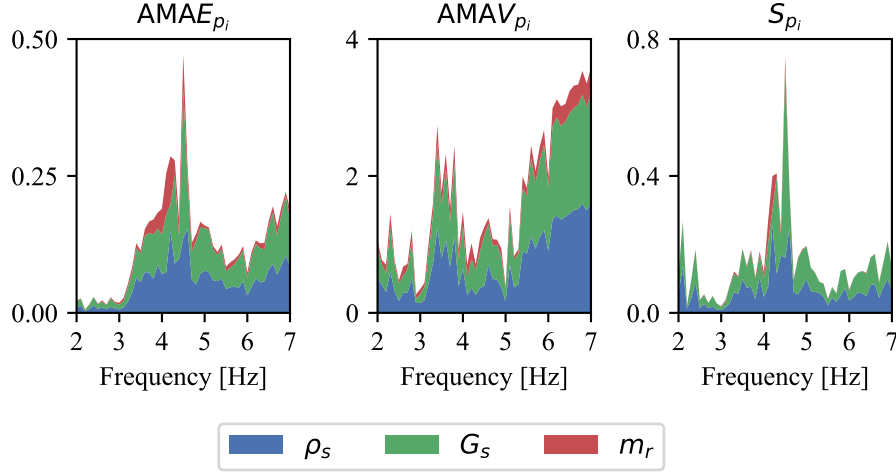


Figure 9: Stacked diagram of the first two moment-based AMA indices ($AMAE_{p_i}$ and $AMAV_{p_i}$), and Sobol principal indices (S_{p_i}) of the model uncertain parameters considered in this study for Scenario 3.

537 Fig. 9 presents the same indices as previous Scenarios. GSA metrics of ρ_s , G_s , and
 538 m_r are depicted in blue, green, and red, respectively. Similar to the previous cases,
 539 the $AMAE$ index values do not depend on the uncertain model parameters in the
 540 low-frequency range below 3 Hz. Unlike Scenario 1, the knowledge of (or conditioning
 541 on) model parameters considerably influences the values of the TC mean for almost
 542 all the considered frequencies. The most and least influential parameters are ρ_s and
 543 m_r , respectively. m_r plays a substantial role in the determination of TC mean for
 544 $3.5 \text{ Hz} \leq f \leq 4.5 \text{ Hz}$, close to the frequency range where the maximum ground-
 545 motion attenuation is predicted (see Fig. 5a). Regarding $AMAV$ index, the results
 546 are similar to Scenario 1, indicating that the variance of the model output can be
 547 significantly modified by the knowledge of (or conditioning on) a model parameter,
 548 such variability being large for high-frequency values. In general, TC variance is
 549 governed by G_s and ρ_s . According to the outcomes of Sobol indices, the variance of
 550 TC can be reduced remarkably by the knowledge of (or conditioning on) uncertain
 551 soil parameters. The influence of the resonator mass is negligible for the frequency
 552 ranges far from resonance. The results of Sobol indices are consistent with those of
 553 $AMAE$ and $AMAV$ indices.

554 The values of the $AMAE$ metric associated with different Scenarios indicate that
 555 the mechanical parameters of the soil are more influential than the resonator mass
 556 in producing a wider attenuation frequency range for stiff soils such as the weath-
 557 ered granite soil considered in Scenario 3. Such a phenomenon arises from a more

558 extensive transfer of stress from heavy resonators to the stiff soil and is consistent
 559 with the outcomes of the full numerical model (see Fig. 5). Conversely, results of the
 560 sensitivity analysis performed for a softer soil like the one considered in Scenario 1
 561 suggest that the resonator mass becomes more dominant only in a narrow frequency
 562 range between $4.3 \text{ Hz} \leq f \leq 4.7 \text{ Hz}$. These results are also consistent with those of
 563 the WFEM numerical simulations. Soils of moderate stiffness like the one consid-
 564 ered in Scenario 2 present an intermediate case, similar to the results of dispersion
 565 analysis and attenuation/transmission coefficients.

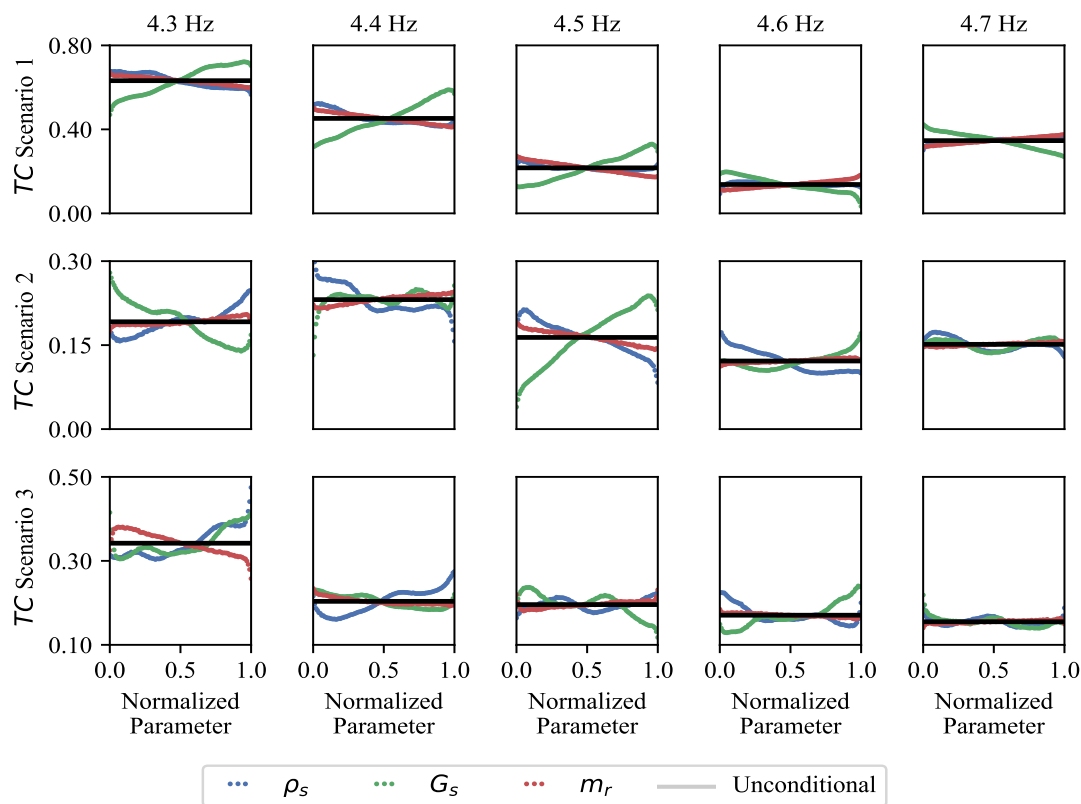


Figure 10: TC conditional on values of the uncertain model parameters considered in this study (ρ_s in blue, G_s in green, and m_r in red). The corresponding unconditional mean is also depicted (black bold horizontal lines). Intervals of variation of the uncertain model parameters are scaled within the unit interval for graphical representation purposes. Results are presented for Scenario 1 (first row), Scenario 2 (second row) and Scenario 3 (third row).

566 We extend our investigations by analyzing the effect of each uncertain parameter
 567 on TC at specific frequency values. Since most of the elastic energy of seismic surface

568 waves is trapped in the frequency range of $4 \text{ Hz} \leq f \leq 6 \text{ Hz}$ (see Fig. 5a), we calculate
 569 TC values conditioned to values of uncertain model parameters for diverse values of
 570 $f = [4.3, 4.4, 4.5, 4.6, \text{ and } 4.7] \text{ Hz}$ where resonant frequencies and frequencies of peak
 571 surface wave attenuation can be found for all Scenarios. Fig. 10 presents TC values
 572 conditioned to values of ρ_s (in blue), G_s (in green), and m_r (in red). The range of
 573 variability of the uncertain parameters is scaled to the unit interval for graphical
 574 representation purposes. The unconditional value of TC is also depicted (in black)
 575 as a reference.

576 Fig. 10 indicates that across all 3 Scenarios, at the frequency of maximum at-
 577 tenuation (4.6 Hz, 4.7 Hz, and 4.6 Hz for Scenarios 1, 2, and 3, respectively) the
 578 influence of uncertain model parameters on TC is less marked than what can be
 579 documented at resonance frequencies (4.5, 4.4, and 4.3 Hz for Scenarios 1, 2, and 3
 580 respectively). At the frequency of maximum attenuation, the influence of m_r on the
 581 mean of TC is minute compared to the influence that ρ_s and G_s have on the mean of
 582 TC . Instead, the resonator mass becomes more influential in the resonant frequency
 583 of all Scenarios. In Scenario 1 the resonator mass has a directly proportional effect
 584 on the TC mean (i.e., the larger the values of m_r , the larger the values of TC).
 585 Whereas in Scenarios 2 and 3, the influence of m_r on the values of TC is virtually
 586 negligible.

587 4. Conclusions

588 A seismic metabarrier is a passive barrier composed of locally resonant elements
 589 devised to control the propagation of seismic surface waves in the long-wavelength
 590 regime. Metabarriers have found application in safeguarding unprotected structures
 591 and infrastructures by reducing the ground motion generated from seismic surface
 592 waves during an earthquake. The engineering implementation of metabarriers de-
 593 mands detailed knowledge of the uncertain parameters affecting their seismic atten-
 594 uation capability during their design phase. Within the context of seismic isolation,
 595 the current study proposes a rigorous methodology to quantify the impact of the
 596 uncertainty associated with system parameters on the seismic attenuation efficiency
 597 of a metabarrier.

598 In this study, we analyze the effect that the uncertainty associated with three pa-
 599 rameters driving the physical behavior of a metabarrier (i.e., the mass of resonators
 600 m_r , soil density ρ_s , and soil shear modulus G_s) has on its attenuation efficiency.
 601 A numerical model is developed to obtain the dispersion relation of the proposed
 602 metabarrier; its results coincide with those derived via analytical dispersion curves.
 603 A low-frequency band gap typical of the local resonance mechanism appears in the

604 dispersion curve, where a significant seismic surface wave attenuation is anticipated.
605 The narrowest and widest band gap regions are associated with granite (S2) and
606 silty-clay (S3) soils, respectively. In the latter case, stronger dynamic coupling oc-
607 curs between the resonators and soil due to their higher relative inertia. A companion
608 full numerical model is developed, and frequency domain analysis is performed to
609 measure the transmission/attenuation coefficients of the metabarrier. The stiffest
610 soil, silty-clay soil of S3, presents the largest peak attenuation with the widest atten-
611 uation frequency range. On the contrary, the sedimentary soil (S1) has the narrowest
612 global attenuation with the smallest peak attenuation due to its weak dynamic in-
613 teraction with heavy resonating masses of resonators; this stems from the soft soil
614 mechanism.

615 The uncertainty of the transmission coefficient TC (i.e., a measure of the attenu-
616 ation efficiency of a metabarrier) is governed by the uncertain parameters associated
617 with the mechanical properties of the soil (ρ_s and G_s). The influence that such
618 parameters have on TC varies with the frequency at which Rayleigh waves propa-
619 gate. In general, for Rayleigh waves oscillating at frequencies close to the resonant
620 frequency of the metabarrier, G_s is the parameter with the largest influence on TC ,
621 followed by ρ_s and by m_r , a parameter whose influence is not negligible only for
622 frequencies close to the resonance.

623 The influence of the resonator mass (m_r) on the attenuation efficiency of the
624 metabarrier depends on the type of soil analyzed. For soft soils (e.g., the soils of
625 Scenario 1 and Scenario 2), an increase in the mass of the resonator enhances the
626 attenuation efficiency of the metabarrier (i.e., decreases TC values) by increasing
627 the relative density of resonators with respect to the soil; whereas for the stiffest soil
628 analyzed in Scenario 3 variations of m_r do not significantly modify the values of TC .

629 Overall, our approach provides new insights into the design and analysis of locally
630 resonant devices to extend current knowledge of metabarriers in different application
631 areas by including the uncertainties associated with the design parameters. The
632 presented methodology does not include some complexities which might arise during
633 the practical implementation of metabarriers. These include, e.g., soil stratigraphy,
634 the presence of groundwater, and soil bearing capacity failure under heavy resonating
635 masses. Otherwise, our approach, in spite of its simplified assumptions, can still be
636 used as a benchmark for advanced numerical models typically adopted in the context
637 of actual geophysical scenarios. The results of this study can impact several segments
638 of engineering, including earthquake engineering, geotechnical engineering, road and
639 railway traffic, and acoustics. Future extensions of the proposed methodology can
640 be tailored to include dampening effects of soil and resonators, the presence of a
641 water table and (possibly) its dynamics, non-linear effects of the soil, and scenarios

642 encompassing sedimentary soils characterized by large Rayleigh velocities to assess
643 the shielding efficiency of seismic metabarriers.

644 **CRedit authorship contribution statement**

645 **Farhad Zeighami:** Conceptualization, Methodology, Software, Data curation,
646 Writing - original draft, Writing - review and editing. **Leonardo Sandoval:** Method-
647 ology, Software, Data curation, Writing - original draft, Writing - review and editing.
648 **Alberto Guadagnini:** Conceptualization, Methodology, Writing - review and edit-
649 ing, Supervision. **Vittorio Di Federico:** Conceptualization, Methodology, Writing
650 - review and editing, Supervision, Funding acquisition.

651 **Conflict of interest**

652 The authors declare that they have no conflict of interest. There are no data
653 sharing issues since all numerical information is provided in the figures produced by
654 solving the equations in the paper.

655 **Acknowledgments**

656 V. Di Federico acknowledges support from the University of Bologna through the
657 Ricerca Fondamentale Orientata (RFO) Grant 2021.

658 **Appendix A. Dimensionless analysis of the dispersion relation**

659 This Section provides the dimensionless form of the dispersion law of Eq. (1).
660 We exploit the resonator parameters (m_r , K_r , and a_r) as scales. Hence, we introduce
661 a set of dimensionless parameters as follows:

$$\omega' = \frac{\omega}{K_r^{1/2} m_r^{-1/2}} = \frac{\omega}{\omega_r}, \quad k' = \frac{k}{k_{max}} = \frac{k}{\frac{\pi}{a_r}}, \quad G' = \frac{G_s}{\frac{K_r}{a_r}}, \quad l' = \frac{l_r}{a_r}, \quad \rho' = \frac{\rho_s}{\frac{m_r}{a_r^3}} \approx \frac{\rho_s}{\rho_r}, \quad (\text{A.1})$$

662 where, ω' is the angular frequency normalized by the angular resonant frequency
663 of the resonator, k' the wavenumber normalized by the maximum wavenumber at
664 the edge of Brillouin zone, G' the shear modulus normalized by the approximate
665 longitudinal modulus of the resonator, l' the dimensionless shape parameter of the
666 resonator, and ρ' the ratio between the mass density of soil and resonator. The

667 dimensionless dispersion law is derived by substituting the dimensionless parameters
 668 of Eq. (A.1) into Eq. (1) to obtain

$$\begin{aligned} & \left(\omega'^2 - 1 \right) \left[\left(2 - \frac{\rho'}{G'} \left(\frac{\omega'}{\pi k'} \right)^2 \right)^2 - 4 \sqrt{1 - \frac{\rho'}{G'} \left(\frac{C\omega'}{\pi k'} \right)^2} \sqrt{1 - \frac{\rho'}{G'} \left(\frac{\omega'}{\pi k'} \right)^2} \right] \\ & = \frac{\rho' \omega'}{l' G'^2} \left(\frac{\omega'}{\pi k'} \right)^3 \sqrt{1 - \frac{\rho'}{G'} \left(\frac{C\omega'}{\pi k'} \right)^2}. \end{aligned} \quad (\text{A.2})$$

669 In the original dispersion equation of Eq. (1) and taking into account Eq.
 670 (2), the angular frequency is an implicit function of eight parameters, i.e. $\omega =$
 671 $fun(k, K_r, m_r, a_r, l_r, \rho_s, G_s, \nu_s)$, while in the dimensionless equation, the dimension-
 672 less angular frequency is an implicit function of five parameters, including i.e. $\omega' =$
 673 $fun(k', l', \rho', G', \nu')$. This is in agreement with the Buckingham π theorem, stating
 674 that the original dispersion law with $n = 8$ dimensional physical variables can be
 675 written in a dimensionless format using $p = n - m = 5$ pure numbers, where $m = 3$
 676 is the number of dimensionally independent scales, chosen to coincide with the vari-
 677 ables describing the resonator. As an alternative to the procedure adopted in the
 678 main body of the paper, GSA could be performed on Eq. (A.2), where the number
 679 of model parameters is reduced.

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