

AN ELLIPTIC BOUNDARY VALUE PROBLEM WITH FRACTIONAL NONLINEARITY*

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Abstract. We investigate the existence and uniqueness of solutions to second-order elliptic boundary value problems containing a power nonlinearity applied to a fractional Laplacian. We detect the critical power separating the existence from the nonexistence regimes. For the existence results, we make use of a particular class of weighted Sobolev spaces to compensate for boundary singularities which are naturally built in the problem.

Key words. Dirichlet problem, operators of mixed order, comparison principles, fixed-point arguments, large solutions

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1. Introduction. Given a bounded domain $\Omega \subseteq \mathbb{R}^N$ with $\partial\Omega \in C^2$, $\sigma \in (0, 1)$, and $p \in [1, \infty)$, we study problems of the form

$$(1.1) \quad \begin{cases} -\Delta u + |(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u = f & \text{in } \Omega, \\ u = g & \text{on } \partial\Omega, \\ u = h & \text{in } \mathbb{R}^N \setminus \bar{\Omega}. \end{cases}$$

Here Δ denotes the (classical) Laplace operator, whereas $(-\Delta)^\sigma$ is the fractional Laplacian

$$(1.2) \quad \begin{aligned} (-\Delta)^\sigma u(x) &:= c_{N,\sigma} \text{p.v.} \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{N+2\sigma}} dy \\ &= c_{N,\sigma} \lim_{\varepsilon \downarrow 0} \int_{\mathbb{R}^N \setminus B_\varepsilon(x)} \frac{u(x) - u(y)}{|x - y|^{N+2\sigma}} dy, \end{aligned}$$

a *nonlocal* positive operator of *fractional* order $2\sigma \in (0, 2)$. The positive constant $c_{N,\sigma}$ is a normalization in order to have that the Fourier symbol of the operator is $|\xi|^{2\sigma}$, i.e.,

$$\mathcal{F}[(-\Delta)^\sigma u](\xi) = |\xi|^{2\sigma} \mathcal{F}u(\xi), \quad u \in C_c^\infty(\mathbb{R}^N).$$

We refer to [4, 15] for an introduction to this operator. Let us here simply remark that definition (1.2) only makes pointwise sense for functions which are defined in the whole Euclidean space \mathbb{R}^N . For this reason, the prototypical well-posed boundary

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value problem driven by the fractional Laplacian takes the form

$$(1.3) \quad \begin{cases} (-\Delta)^\sigma u = f & \text{in } \Omega, \\ u = h & \text{in } \mathbb{R}^N \setminus \bar{\Omega}. \end{cases}$$

In this setting, prescribing the values of the solution u on $\partial\Omega$ is immaterial, as the integral operator $(-\Delta)^\sigma$ does not see negligible sets. Nevertheless, it is reasonable to investigate the boundary regularity of solutions and, in particular, their continuity across $\partial\Omega$, like, for example, in [8] or more recently in [7]—in both cases in a more general setting than the one in (1.3).

Problem (1.1) is motivated by an understanding of the interaction and the overlapping of the different boundary conditions required by the Laplacian and the fractional Laplacian. Indeed, although the term $|(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u$ can be interpreted as a mere nonlinear perturbation of the leading one Δu , one needs to know the values of u also outside Ω in order to make it meaningful. Note also that different definitions of u outside Ω might drastically affect the operator on the interior, a peculiarity of its nonlocality.

1.1. Main results. The exponent p has a prominent role in the solvability of (1.1). The range of p splits at $p = 1/\sigma$ into two regimes: below this value, which we call the *subcritical regime*, it is possible to solve (1.1) if the other prescribed data f and g are somewhat well-behaved; above that value, in the *supercritical regime*, a switch in the lead of the equation takes place and the nonlocal term becomes the dominant one, making it impossible (in general) to attain the desired values at the boundary $\partial\Omega$. See Theorems 1.1 and 1.3 for the precise statements.

In all the following, Ω is supposed to be a bounded domain with C^2 boundary.

THEOREM 1.1 (existence in the subcritical regime). *Let $g \in C^0(\partial\Omega)$, $h \in L^\infty(\mathbb{R}^N \setminus \bar{\Omega})$, and f be a locally Hölder continuous function in Ω satisfying*

$$(1.4) \quad \text{dist}(\cdot, \partial\Omega)^{2-\alpha} f \in L^\infty(\Omega)$$

for some $\alpha > 0$. If

$$(1.5) \quad 1 \leq p < \frac{1}{\sigma},$$

then problem (1.1) has a unique solution $u \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$.

We underline how data g and h are completely unrelated in the above statement: in particular, we do not need to assume that the values of g on $\partial\Omega$ match with those of a suitable extension of h to $\partial\Omega$ (which, in fact, needs not be continuous nor continuously extensible up to the boundary).

In the particular case $p = 1$, the equation becomes linear and it admits a Green function for which sharp two-sided estimates are available, see Chen et al. [12, 13]: although we do not make use of these, let us just mention here that Theorem 1.1 (for $p = 1$) can also be deduced by means of such Green representation. Very recently, Biagi et al. [10] have studied the existence, maximum principles, and regularity for (1.1) with $p = 1$ and $g, h = 0$: their techniques and goals are quite different from our approach, but their main result [10, Theorem 1.7] is, nonetheless, closely related to Theorem 1.1.

The core of the proof of Theorem 1.1 relies on a fixed-point argument for the nonlinear operator

$$(1.6) \quad u \mapsto -\Delta u + |(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u.$$

The (possible) jump discontinuity of u , inherited by the prescription of g and h in (1.1), entails a singularity in the nonlocal part of (1.6) at the boundary of Ω . This represents a major challenge in solving (1.1), which turns out to be not only a nonlinear problem but also a singular one. To overcome this issue, we consider an approximating family of regularized problems, run the fixed-point argument to solve these problems, and pass to the limit via uniform estimates in Sobolev spaces of fractional order with boundary weights (cf. (2.2)–(2.3)). Their definition is due to Lototsky [24]: we briefly outline their construction and prove some new results (cf. Lemma 2.3) in section 2. The full proof of Theorem 1.1 is contained in section 5.

An important feature we will need in the proofs is a comparison principle.

PROPOSITION 1.2 (weak comparison principle). *Let $\Omega \subseteq \mathbb{R}^N$ be a bounded open set and $\Phi : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ be a Carathéodory function with $\Phi(x, \cdot)$ nondecreasing for a.e. $x \in \Omega$. Let $\underline{w}, \bar{w} \in L^1_\sigma(\mathbb{R}^N) \cap C^2(\Omega)$ be two functions satisfying*

$$(1.7) \quad \begin{cases} -\Delta \underline{w} + \Phi(\cdot, (-\Delta)^\sigma \underline{w}) \leq -\Delta \bar{w} + \Phi(\cdot, (-\Delta)^\sigma \bar{w}) & \text{in } \Omega, \\ \underline{w} \leq \bar{w} & \text{in } \mathbb{R}^N \setminus \bar{\Omega}. \end{cases}$$

For any $x_0 \in \partial\Omega$, suppose in addition that

$$(1.8) \quad [-\infty, +\infty) \ni \limsup_{\Omega \ni x \rightarrow x_0} \underline{w}(x) \leq \liminf_{\Omega \ni x \rightarrow x_0} \bar{w}(x) \in (-\infty, +\infty].$$

Then, $\underline{w} \leq \bar{w}$ also in Ω .

Here, $L^1_\sigma(\mathbb{R}^N)$ denotes the space of functions $v \in L^1_{\text{loc}}(\mathbb{R}^N)$ for which $(1+|\cdot|)^{-N-2\sigma} v \in L^1(\mathbb{R}^N)$. As is known, this assumption on v , along with its local C^2 (or $C^{2\sigma+\varepsilon}$) regularity, is enough to have $(-\Delta)^\sigma v$ well-defined in the pointwise sense as the integral operator (1.2). We will mostly apply Proposition 1.2 to functions \underline{w} and \bar{w} continuous up to the boundary of Ω . In this case, assumption (1.8) simply boils down to

$$\underline{w} \leq \bar{w} \quad \text{on } \partial\Omega.$$

Notice, however, that (1.8) makes sense even when $\underline{w}|_\Omega$ and $\bar{w}|_\Omega$ cannot be continuously extended up to $\partial\Omega$. This feature will be crucial in order to deal with solutions of (1.1) which blow up at the boundary. Section 4 deals with the proof of Proposition 1.2 and its consequences.

Passing to nonexistence results, we have the following.

THEOREM 1.3 (nonexistence in the critical and supercritical regimes). *Let $g \in C^0(\partial\Omega)$ and $h \in L^\infty(\mathbb{R}^N \setminus \bar{\Omega})$ be such that $g \not\leq 0$ on $\partial\Omega$ and $h \leq 0$ in $\mathbb{R}^N \setminus \bar{\Omega}$. If $\sigma p \geq 1$, then problem*

$$(1.9) \quad \begin{cases} -\Delta u + |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u = 0 & \text{in } \Omega, \\ u = g & \text{on } \partial\Omega, \\ u = h & \text{in } \mathbb{R}^N \setminus \bar{\Omega} \end{cases}$$

has no solution $u \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$.

Note that the threshold $p = 1/\sigma$ is indeed quite natural. The analogue local problem

$$(1.10) \quad \begin{cases} -\Delta u + |\nabla u|^p = 0 & \text{in } \Omega, \\ u = g & \text{on } \partial\Omega \end{cases}$$

is known to always have a solution for $p \in [1, 2]$ [31, Hilfssatz 3] (or also [30, section 11] for a more general statement), whereas existence is lost in general for $p > 2$ (see [30, section 16, Theorem 1]). Interestingly, in this case the critical power $p = 2$ is included in the existence regime, whereas in (1.1) $p = 1/\sigma$ falls into the nonexistence one.

The next theorem shows that condition (1.4) on the right-hand side is almost sharp.

THEOREM 1.4 (nonexistence for large sources). *Let $f \in L^\infty_{\text{loc}}(\Omega)$ be such that*

$$\text{dist}(\cdot, \partial\Omega)^2 f \geq \kappa \quad \text{in } \Omega$$

for some $\kappa > 0$. Then, for all p as in (1.5), problem

$$(1.11) \quad \begin{cases} -\Delta u + |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \bar{\Omega} \end{cases}$$

has no solution $u \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$.

Theorems 1.3 and 1.4 are proved in section 6.

We finally show, in section 7, how the nonlinear character of (1.1) allows for solutions that become singular¹ at $\partial\Omega$ and are therefore called *boundary blow-up solutions* or, simply, *large solutions*.

THEOREM 1.5 (large solutions). *For any*

$$p \in \left(\frac{3 - \sigma}{1 + \sigma}, \frac{1}{\sigma} \right),$$

problem

$$(1.12) \quad \begin{cases} -\Delta u + |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u = 0 & \text{in } \Omega, \\ u = +\infty & \text{on } \partial\Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \bar{\Omega} \end{cases}$$

admits a solution $u \in L^1(\mathbb{R}^N) \cap C^2(\Omega)$. Moreover, there exists a $C > 0$ such that this solution satisfies

$$(1.13) \quad 0 < u \leq C \text{dist}(\cdot, \partial\Omega)^{-2(1-\sigma p)/(p-1)} \quad \text{in } \Omega.$$

1.2. Notation. We will use the following notation without further notice.

We set $\delta = \text{dist}(\cdot, \mathbb{R}^N \setminus \Omega)$ in \mathbb{R}^N . By D^j we denote the collection of partial derivatives of order $j \in \mathbb{N}$. The letters C and \bar{C} will be used to indicate constants that have values larger than 1 and that may change from line to line. Constants denoted by C will depend only on the structural quantities involved in the problem, i.e., N , Ω , p , σ , and α , whereas \bar{C} may also depend on the norms of the data— $\|\delta^{2-\alpha} f\|_{L^\infty(\Omega)}$, $\|g\|_{L^\infty(\partial\Omega)}$, and $\|h\|_{L^\infty(\mathbb{R}^N \setminus \Omega)}$. Whenever a constant also depends on an additional quantity, we will emphasize it by using subscripts—for instance, C_j will denote a constant that also depends on j on top of the previously specified parameters.

¹These solutions appear in classical semilinear problems when the nonlinearity fulfills the so-called *Keller–Osserman condition* [21, 29]. They also show up in problems of the same type as (1.10) as noted, for example, in [22]. For fractional-order equations the situation is more involved; the interested reader might want to check [11, 18, 1, 2].

2. Fractional Sobolev spaces with weights. Following [24], we introduce weighted spaces $L_\theta^{s,p}$ modelled upon Bessel potential spaces, from which we borrow the usual notation. For $s \geq 0$ and $p \geq 1$, we let $L^{s,p}(\mathbb{R}^N)$ denote the Bessel potential space, obtained as the completion of $C_c^\infty(\mathbb{R}^N)$ with respect to the norm

$$(2.1) \quad \|u\|_{L^{s,p}(\mathbb{R}^N)} := \|(1 - \Delta)^{s/2}u\|_{L^p(\mathbb{R}^N)}.$$

We recall that $(1 - \Delta)^{s/2}$ is the operator defined by

$$(1 - \Delta)^{s/2}u := \mathcal{F}^{-1}\left(\left(1 + |\cdot|^2\right)^{s/2}\mathcal{F}u\right) \quad \text{on any } u \in C_c^\infty(\mathbb{R}^N),$$

where \mathcal{F} denotes the Fourier transform. Note that $L^{0,p}(\mathbb{R}^N)$ reduces to the standard Lebesgue space $L^p(\mathbb{R}^N)$.

For any $k \in \mathbb{Z}$, we define

$$A_k := \left\{x \in \Omega : 2^{-k-1} < \delta(x) < 2^{-k+1}\right\}.$$

Let $(\zeta_k)_{k \in \mathbb{Z}}$ be a smooth partition of unity, namely, a family of nonnegative functions $\zeta_k \in C_c^\infty(\mathbb{R}^N)$ such that $\text{supp}(\zeta_k) \subseteq A_k$,

$$|D^j \zeta_k(x)| \leq C_j 2^{jk} \quad \text{for any } x \in \mathbb{R}^N \text{ and } j \in \mathbb{N}_0,$$

for some constant $C_{|\alpha|} > 0$ and with $|\alpha| = \alpha_1 + \dots + \alpha_N$, and

$$\sum_{k \in \mathbb{Z}} \zeta_k(x) = 1 \quad \text{for any } x \in \Omega.$$

Given another parameter $\theta \in \mathbb{R}$, introduce the space

$$(2.2) \quad L_\theta^{s,p}(\Omega) := \left\{u \in L^p(\Omega) : \|u\|_{L_\theta^{s,p}(\Omega)} < +\infty\right\},$$

where

$$(2.3) \quad \|u\|_{L_\theta^{s,p}(\Omega)}^p := \sum_{k \in \mathbb{Z}} 2^{-k\theta} \|\zeta_k(2^{-k}\cdot)u(2^{-k}\cdot)\|_{L^{s,p}(\mathbb{R}^N)}^p.$$

For simplicity, we just write $L_\theta^p(\Omega)$ instead of $L_\theta^{0,p}(\Omega)$.

The space $L_\theta^{s,p}(\Omega)$ is a Sobolev space with a weight at the boundary $\partial\Omega$, introduced by Lototsky in [23, 24]. In the next statement we collect some of its basic properties that will be used in the rest of the paper—see [24, Proposition 2.2] for their proofs.

PROPOSITION 2.1 ([24, Proposition 2.2]). *The following statements hold true.*

- (i) *The space $C_c^\infty(\Omega)$ is dense in $L_\theta^{s,p}(\Omega)$.*
- (ii) *The space $L_\theta^{s,p}(\Omega)$ is independent of the choice of partition of unity $(\zeta_k)_{k \in \mathbb{Z}}$, and different partitions lead to equivalent norms.*
- (iii) *The quantity*

$$(2.4) \quad \|u\|_{L_\theta^p(\Omega)}^* := \left(\int_\Omega |u(x)|^p \delta(x)^{\theta-N} dx\right)^{1/p}$$

defines an equivalent norm for the space $L_\theta^p(\Omega)$.

(iv) If $s = k$ is a nonnegative integer, then

$$L_{\theta}^{k,p}(\Omega) = \left\{ u \in L_{\theta}^p(\Omega) : \delta^j D^j u \in L_{\theta}^p(\Omega) \text{ for all } j = 1, \dots, k \right\}$$

and the norm

$$\|u\|_{L_{\theta}^{k,p}(\Omega)}^* := \left(\sum_{j=0}^k (\|\delta^j D^j u\|_{L_{\theta}^p(\Omega)}^*)^p \right)^{1/p}$$

is equivalent to the one defined in (2.3).

(v) For $s_i \geq 0$, $p_i > 1$, $\theta_i \in \mathbb{R}$ ($i = 0, 1$), and $\nu \in (0, 1)$, it holds that

$$L_{\theta}^{s,p}(\Omega) = \left[L_{\theta_0}^{s_0,p_0}(\Omega), L_{\theta_1}^{s_1,p_1}(\Omega) \right]_{\nu}$$

with

$$s := (1 - \nu)s_0 + \nu s_1, \quad \frac{1}{p} := \frac{1 - \nu}{p_0} + \frac{\nu}{p_1}, \quad \theta := (1 - \nu)\theta_0 + \nu\theta_1,$$

and where $[X, Y]_{\nu}$ denotes the complex interpolation space of X and Y .

In view of (2.4), one has $L_N^p(\Omega) = L^p(\Omega)$ —note, however, that $L_N^{s,p}(\Omega)$ differs from the unweighted Bessel potential space $L^{s,p}(\Omega)$ when $s > 0$.

One of the most important outcomes of the analysis carried out in [24] is a weighted Calderón–Zygmund-type estimate for solutions of a class of degenerate elliptic second-order equations. In section 5, we will take advantage of a very particular case of this result, which we state here below for the reader's convenience and for further reference—see [24, section 5] for its proof.

PROPOSITION 2.2 ([24, Lemma 5.2]). *Let $p > 1$, $f \in L_{\theta}^p(\Omega)$, and $u \in C^2(\Omega) \cap L_{\theta}^p(\Omega)$ be a solution of $-\Delta u = f$ in Ω . Then, $u \in L_{\theta}^{2,p}(\Omega)$ and*

$$\|u\|_{L_{\theta}^{2,p}(\Omega)} \leq C \left(\|f\|_{L_{\theta+2p}^p(\Omega)} + \|u\|_{L_{\theta}^p(\Omega)} \right)$$

for some constant $C > 0$ depending only on N , p , θ , and Ω .

One last element of this theory that we will need is an estimate on the $L_{\theta}^p(\Omega)$ norm of the fractional Laplacian of a regular function. Before heading to its statement, we note that, for $p > 1$, the norm $\|\cdot\|_{L^{s,p}(\mathbb{R}^N)}$ introduced in (2.1) for $L^{s,p}(\mathbb{R}^N)$ is equivalent to

$$\|u\|_{L^{s,p}(\mathbb{R}^N)} := \|u\|_{L^p(\mathbb{R}^N)} + \|(-\Delta)^{s/2} u\|_{L^p(\mathbb{R}^N)}.$$

Indeed, for $s \in \mathbb{N}$ this follows from [27, Corollary 10.1.2/1] and [5, Corollary 4.16 and Theorem 7.63(f)]. In this case, $L^{s,p}(\mathbb{R}^N)$ coincides with the usual Sobolev space $W^{s,p}(\mathbb{R}^N)$. When s is not an integer, the equivalence can be deduced from [27, Theorem 10.1.2/4].

Thanks to this observation, we can prove the following estimate. We will use it in sections 5 and 7, respectively, with $r = p$ and $r = 1$.

LEMMA 2.3. *Let $p > 1$, $r \in [1, p]$, $\sigma \in (0, 1)$, and $\theta > (N/r + 2\sigma)p$. Then, for every $\varepsilon > 0$, there exists a constant $C > 0$, depending only on N , p , r , σ , θ , Ω , and ε , such that*

$$(2.5) \quad \|(-\Delta)^{\sigma} u\|_{L_{\theta}^p(\Omega)} \leq C \left(\|u\|_{L_{\theta-2\sigma p-\varepsilon}^{2\sigma,p}(\Omega)} + \|u\|_{L^r(\Omega)} \right)$$

for every function $u : \mathbb{R}^N \rightarrow \mathbb{R}$ that vanishes outside Ω and is of class C^2 in Ω .

Proof. By the properties of the ζ_k 's, the fact that $u = 0$ in $\mathbb{R}^N \setminus \Omega$, and a suitable change of variables, we have

$$\begin{aligned}
 \|(-\Delta)^\sigma u\|_{L^p_\theta(\Omega)} &\leq \sum_{j \in \mathbb{Z}} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p_\theta(\Omega)} \\
 (2.6) \qquad &= \sum_{j \in \mathbb{Z}} \left(\sum_{k \in \mathbb{Z}} 2^{-k\theta} \left\| \zeta_k(2^{-k}\cdot)(-\Delta)^\sigma(\zeta_j u)(2^{-k}\cdot) \right\|_{L^p(\mathbb{R}^N)}^p \right)^{1/p} \\
 &\leq \sum_{j \in \mathbb{Z}} \left(\sum_{k \in \mathbb{Z}} 2^{-(\theta-N)k} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)}^p \right)^{1/p}.
 \end{aligned}$$

Now, we consider separately the cases $k \geq j - 2$ and $k \leq j - 3$.

In the first situation, we estimate

$$\begin{aligned}
 (2.7) \qquad &\sum_{k=j-2}^{+\infty} 2^{-(\theta-N)k} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)}^p \\
 &\leq 4^{\theta-N} 2^{-(\theta-N)j} \sum_{k=j-2}^{+\infty} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)}^p \\
 &\leq 3 \cdot 4^{\theta-N} 2^{-(\theta-N)j} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(\mathbb{R}^N)}^p,
 \end{aligned}$$

where for the last inequality we used that the A_k 's intersect at most three times.

On the other hand, take $k \leq j-3$ and observe that $|x-y| \geq 2^{-k-2}$ for every $x \in A_k$ and $y \in A_j$. We compute

$$\begin{aligned}
 \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)} &= \left(\int_{A_k} \left| \int_{\mathbb{R}^N} \frac{\zeta_j(y)u(y)}{|x-y|^{N+2\sigma}} dy \right|^p dx \right)^{1/p} \\
 &\leq \left(\int_{\mathbb{R}^N} \left(\int_{\mathbb{R}^N} \zeta_j(y)|u(y)| \frac{\chi_{\mathbb{R}^N \setminus B_{2^{-k-2}}(x-y)}}{|x-y|^{N+2\sigma}} dy \right)^p dx \right)^{1/p}.
 \end{aligned}$$

From this and Young's convolution inequality, we deduce that

$$\begin{aligned}
 \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)} &\leq \|\zeta_j u\|_{L^r(\mathbb{R}^N)} \left(\int_{\mathbb{R}^N \setminus B_{2^{-k-2}}} |z|^{-\frac{(N+2\sigma)pr}{pr-p+r}} dz \right)^{\frac{pr-p+r}{pr}} \\
 &\leq 2^{\left(\frac{N}{r} - \frac{N}{p} + 2\sigma\right)k} C \|\zeta_j u\|_{L^r(\mathbb{R}^N)}.
 \end{aligned}$$

Also, since Ω is bounded, there exists $k_0 \in \mathbb{Z}$ such that $A_k = \emptyset$ for all $k < k_0$. In light of these facts, we have

$$\begin{aligned}
 \sum_{k=-\infty}^{j-3} 2^{-(\theta-N)k} \|(-\Delta)^\sigma(\zeta_j u)\|_{L^p(A_k)}^p &\leq C \|\zeta_j u\|_{L^r(\mathbb{R}^N)}^p \sum_{k=k_0}^{j-3} 2^{-k(\theta - \frac{Np}{r} - 2\sigma p)} \\
 &\leq C \|u\|_{L^r(A_j)}^p,
 \end{aligned}$$

where the last inequality holds since $\theta > (N/r + 2\sigma)p$.

By combining the above estimate with (2.7), changing coordinates, and using the scaling property

$$(-\Delta)^\sigma [v(\lambda \cdot)] = \lambda^{2\sigma} (-\Delta)^\sigma v(\lambda \cdot) \quad \text{for } \lambda > 0,$$

we deduce from (2.6) that

$$\begin{aligned} \|(-\Delta)^\sigma u\|_{L^p_\theta(\Omega)} &\leq C \sum_{j=k_0}^{+\infty} \left(2^{-\frac{\theta-N}{p}j} \|(-\Delta)^\sigma (\zeta_j u)\|_{L^p(\mathbb{R}^N)} + \|u\|_{L^r(A_j)} \right) \\ &\leq C \left(\sum_{j=k_0}^{+\infty} 2^{-\frac{\theta-2\sigma p}{p}j} \left\| (-\Delta)^\sigma [\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)] \right\|_{L^p(\mathbb{R}^N)} + \|u\|_{L^r(\Omega)} \right). \end{aligned}$$

Since, by the definition of $\|\cdot\|_{L^{2\sigma,p}(\mathbb{R}^N)}$ and its equivalence to $\|\cdot\|_{L^{2\sigma,p}(\mathbb{R}^N)}$,

$$\begin{aligned} \left\| (-\Delta)^\sigma [\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)] \right\|_{L^p(\mathbb{R}^N)} &\leq \|\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)\|_{L^{2\sigma,p}(\mathbb{R}^N)} \\ &\leq C \|\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)\|_{L^{2\sigma,p}(\mathbb{R}^N)}, \end{aligned}$$

we conclude that

$$\|(-\Delta)^\sigma u\|_{L^p_\theta(\Omega)} \leq C \left(\sum_{j=k_0}^{+\infty} 2^{-\frac{\theta-2\sigma p}{p}j} \|\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)\|_{L^{2\sigma,p}(\mathbb{R}^N)} + \|u\|_{L^r(\Omega)} \right).$$

To obtain (2.5), we are left to show that the above series can be controlled by the $L^{2\sigma,p}_{\theta-2\sigma p-\varepsilon}(\Omega)$ norm of u for every $\varepsilon > 0$. This is a consequence of a simple inequality for numerical series. Indeed, given any sequence $(a_j)_{j \in \mathbb{N}}$ of nonnegative numbers, Hölder’s inequality gives that

$$\begin{aligned} \sum_{j=k_0}^{+\infty} 2^{-\frac{\theta-2\sigma p}{p}j} a_j &= \sum_{j=k_0}^{+\infty} \left(2^{-\frac{\theta-2\sigma p-\varepsilon}{p}j} a_j \right) 2^{-\frac{\varepsilon}{p}j} \\ &\leq \left(\sum_{j=k_0}^{+\infty} \left(2^{-\frac{\theta-2\sigma p-\varepsilon}{p}j} a_j \right)^p \right)^{1/p} \left(\sum_{j=k_0}^{+\infty} 2^{-\frac{\varepsilon}{p-1}j} \right)^{(p-1)/p} \\ &\leq C \left(\sum_{j \in \mathbb{Z}} 2^{-(\theta-2\sigma p-\varepsilon)j} a_j^p \right)^{1/p}. \end{aligned}$$

Applying this with $a_j = \|\zeta_j(2^{-j}\cdot)u(2^{-j}\cdot)\|_{L^{2\sigma,p}(\mathbb{R}^N)}$ and recalling (2.3), we infer that (2.5) holds true. \square

3. Barriers. In this section, we present the construction of positive supersolutions for both the classical and the fractional Laplace operator in bounded domains. Ultimately, they will be used as barriers for the nonlinear operator

$$u \mapsto -\Delta u + |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u.$$

We begin with a simple barrier that will be used for equations involving bounded right-hand sides. In this case, we may restrict ourselves to consider balls as underlying domains, and thus the construction is rather standard. As usual, for $x, y > 0$ we denote by $B(x, y) = \int_0^1 t^{x-1} (1-t)^{y-1} dt$ the beta function.

LEMMA 3.1. *For $R > 0$, the function $u_\sigma(x) := (R^2 - |x|^2)_+^\sigma$, $x \in \mathbb{R}^N$, satisfies*

$$(3.1) \quad -\Delta u_\sigma(x) = 2\sigma \frac{N(R^2 - |x|^2) + 2(1-\sigma)|x|^2}{(R^2 - |x|^2)^{2-\sigma}} \geq \frac{2\sigma N}{R^{2-2\sigma}} \quad \text{for every } x \in B_R$$

and

$$(3.2) \quad (-\Delta)^\sigma u_\sigma(x) = \frac{c_{N,\sigma} B(\sigma, 1 - \sigma) |\partial B_1|}{2} > 0 \quad \text{for every } x \in B_R.$$

Proof. A straightforward computation gives (3.1). Identity (3.2) is due to [19]—see [17] for an elementary proof and for more general relations. \square

In order to deal with right-hand sides that blow up at the boundary, we can no longer limit ourselves to balls, and instead we need to construct barriers tailored to each specific domain. We do this in the next lemma, by considering powers of the so-called torsion function, i.e., of the solution of the Dirichlet problem

$$(3.3) \quad \begin{cases} -\Delta \tau = 1 & \text{in } \Omega, \\ \tau = 0 & \text{on } \partial\Omega. \end{cases}$$

The existence and uniqueness of the solution τ of the torsion problem (3.3) is classical. Furthermore, $\tau > 0$ in Ω thanks to the strong maximum principle.

LEMMA 3.2. *Let $\Omega \subseteq \mathbb{R}^N$ be a bounded domain with boundary of class C^2 and $\tau \in C^\infty(\Omega) \cap C^1(\bar{\Omega})$ be the solution of (3.3). For $\alpha \geq 0$, set*

$$v_\alpha := \tau^\alpha \chi_{\bar{\Omega}} \quad \text{in } \mathbb{R}^N.$$

Then, the following statements hold true.

- (i) *There exists a constant $C_1 \geq 1$, depending only on Ω , such that*

$$(3.4) \quad C_1^{-1} \delta^\alpha \leq v_\alpha \leq C_1 \delta^\alpha \quad \text{in } \Omega$$

for all $\alpha \in (0, 1]$.

- (ii) *There exists a constant $C_2 \geq 1$, depending only on N , Ω , and σ , such that*

$$(3.5) \quad C_2^{-1} \alpha(1 - \alpha) \delta^{\alpha-2} \leq -\Delta v_\alpha \leq C_2 \alpha \delta^{\alpha-2} \quad \text{in } \Omega,$$

$$(3.6) \quad -C_2 \text{diam}(\Omega)^\alpha \delta^{-2\sigma} \leq (-\Delta)^\sigma v_\alpha \leq C_2 \delta^{\alpha-2\sigma} \quad \text{in } \Omega$$

for all $\alpha \in (0, 1)$.

- (iii) *There exists a constant $C_3 \geq 1$, depending only on N and Ω , such that*

$$(3.7) \quad C_3^{-1} \delta^{-2\sigma} \leq (-\Delta)^\sigma v_0 \leq C_3 \delta^{-2\sigma} \quad \text{in } \Omega.$$

Proof. First, we notice that

$$(3.8) \quad C_1^{-1} \delta \leq \tau \leq C_1 \delta \quad \text{in } \Omega$$

for some constant $C_1 \geq 1$ depending only on Ω . The upper bound on τ follows from its $C^1(\bar{\Omega})$ regularity, whereas the lower bound is a consequence of its positivity inside Ω and of Hopf’s lemma.

Knowing this, we address the validity of the claims made in the statement. Of course, (3.4) follows from (3.8) right away. In order to prove (3.5), a straightforward computation gives that

$$(3.9) \quad \partial_{ij} v_\alpha = \alpha(\alpha - 1) \tau^{\alpha-2} \partial_i \tau \partial_j \tau + \alpha \tau^{\alpha-1} \partial_{ij} \tau \quad \text{in } \Omega$$

for every $i, j = 1, \dots, N$, and thus, by (3.3)

$$-\Delta v_\alpha = \alpha(1 - \alpha) \tau^{\alpha-2} |\nabla \tau|^2 + \alpha \tau^{\alpha-1} \quad \text{in } \Omega.$$

As $\tau \in C^1(\bar{\Omega})$, we infer that $-\Delta v_\alpha \leq C\alpha\tau^{\alpha-2}$ in Ω , for some constant $C \geq 1$ depending only on Ω . On the other hand, using again Hopf's lemma, we get that $|\nabla\tau| \geq C^{-1}$ in $\Gamma_\varepsilon = \{x \in \Omega : \delta(x) < \varepsilon\}$ for some $\varepsilon > 0$ small enough. Since we also clearly have that $\tau \geq C^{-1}$ in $\Omega \setminus \Gamma_\varepsilon$, we deduce that $-\Delta v_\alpha \geq C^{-1}\alpha(1-\alpha)\tau^{\alpha-2}$ in Ω . These facts combined with (3.8) immediately lead to (3.5).

We now proceed to verify (3.6). To this aim, we claim that

$$(3.10) \quad \|D^2v_\alpha\|_{L^\infty(B_{\delta(x)/4}(x))} \leq C_*\delta(x)^{\alpha-2} \quad \text{for every } x \in \Omega$$

for some constant $C_* > 0$ depending only on N and Ω . This estimate is a simple consequence of the Schauder theory. Indeed, by, e.g., [20, Theorem 4.6], (3.3), and the upper bound in (3.8), we have

$$\|D^2\tau\|_{L^\infty(B_{\delta(x)/4}(x))} \leq C\left(\delta(x)^{-2}\|\tau\|_{L^\infty(B_{\delta(x)/2}(x))} + \|1\|_{L^\infty(B_{\delta(x)/2}(x))}\right) \leq C\delta(x)^{-1}$$

for some constant $C > 0$ depending only on N and Ω . Taking advantage of this, the $C^1(\bar{\Omega})$ regularity of τ , and (3.9), one easily deduces (3.10). To establish (3.6), we take a point $x \in \Omega$ and write

$$c_{N,\sigma}^{-1}(-\Delta)^\sigma v_\alpha(x) = \text{p.v.} \int_{B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y)}{|x - y|^{N+2\sigma}} dy + \int_{\mathbb{R}^N \setminus B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y)}{|x - y|^{N+2\sigma}} dy.$$

On the one hand, by (3.10), we have

$$\begin{aligned} & \left| \text{p.v.} \int_{B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y)}{|x - y|^{N+2\sigma}} dy \right| \\ &= \left| \int_{B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y) + \nabla v_\alpha(x) \cdot (y - x)}{|x - y|^{N+2\sigma}} dy \right| \\ &\leq \|D^2v_\alpha\|_{L^\infty(B_{\delta(x)/4}(x))} |\partial B_1| \int_0^{\delta(x)/4} t^{1-2\sigma} dt \leq C\delta(x)^{\alpha-2\sigma}. \end{aligned}$$

On the other hand, using (3.4), we find that

$$\int_{\mathbb{R}^N \setminus B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y)}{|x - y|^{N+2\sigma}} dy \leq |\partial B_1| v_\alpha(x) \int_{\delta(x)/4}^{+\infty} \frac{dt}{t^{1+2\sigma}} \leq C\delta(x)^{\alpha-2\sigma}$$

and

$$\begin{aligned} \int_{\mathbb{R}^N \setminus B_{\delta(x)/4}(x)} \frac{v_\alpha(x) - v_\alpha(y)}{|x - y|^{N+2\sigma}} dy &\geq -C \int_{\Omega \setminus B_{\delta(x)/4}(x)} \frac{\delta(y)^\alpha}{|x - y|^{N+2\sigma}} dy \\ &\geq -C \text{diam}(\Omega)^\alpha \delta(x)^{-2\sigma}. \end{aligned}$$

Putting together the last four formulas, we arrive at (3.6).

We are left to prove (3.7). The right-hand inequality is straightforward, as

$$c_{N,\sigma}^{-1}(-\Delta)^\sigma v_0(x) = \int_{\mathbb{R}^N \setminus \Omega} \frac{dy}{|x - y|^{N+2\sigma}} \leq \int_{\mathbb{R}^N \setminus B_{\delta(x)}(x)} \frac{dy}{|x - y|^{N+2\sigma}} \leq C\delta(x)^{-2\sigma}$$

for every $x \in \Omega$. To check that the left-hand one holds as well, we first observe that it suffices to establish it at points in $\Gamma_\varepsilon = \{x \in \Omega : \delta(x) < \varepsilon\}$ for some $\varepsilon > 0$ arbitrarily small but depending at most on N , Ω , and σ . Let $x \in \Gamma_\varepsilon$ and denote by $z_x \in \partial\Omega$ a

point for which $\delta(x) = |x - z_x|$. By the C^2 regularity of $\partial\Omega$, there exists an exterior tangent ball $B_{\delta(x)}(w_x)$ to Ω at z_x , provided ε is small enough (in dependence of Ω only). By virtue of this, we compute

$$c_{N,\sigma}^{-1}(-\Delta)^\sigma v_0(x) = \int_{\mathbb{R}^N \setminus \Omega} \frac{dy}{|x - y|^{N+2\sigma}} \geq \int_{B_{\delta(x)}(w_x)} \frac{dy}{|x - y|^{N+2\sigma}} \geq C^{-1} \delta(x)^{-2\sigma},$$

and the lower bound in (3.7) follows. The proof of the lemma is thus complete. \square

A simple application of the barriers constructed in Lemma 3.2 is the following result, which provides estimates on how fast solutions of the Poisson equation attain their data in the presence of a right-hand side that blows up at the boundary. Half of this result is included in [20, Theorem 4.9] when the domain is a ball and in [20, Problem 4.6] for the general case.

LEMMA 3.3. *Let $\alpha \in (0, 1)$, $\Omega \subseteq \mathbb{R}^N$ be a bounded domain with boundary of class C^2 , and f be such that $\delta^{2-\alpha} f \in L^\infty(\Omega)$. Let $u \in C^2(\Omega) \cap C^0(\bar{\Omega})$ be a solution of*

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Then,

$$(3.11) \quad \|\delta^{-\alpha} u\|_{L^\infty(\Omega)} \leq C \alpha^{-1} (1 - \alpha)^{-1} \|\delta^{2-\alpha} f\|_{L^\infty(\Omega)}$$

for some constant $C \geq 1$ depending only on Ω . Furthermore, if $f \geq 0$ in Ω , then

$$(3.12) \quad \inf_{\Omega} (\delta^{-\alpha} u) \geq C^{-1} \alpha^{-1} \inf_{\Omega} (\delta^{2-\alpha} f).$$

Proof. Let v_α and C_2 be as in Lemma 3.2. By the left-hand inequality in (3.5), the function

$$\bar{v} := \alpha^{-1} (1 - \alpha)^{-1} C_2 \|\delta^{2-\alpha} f\|_{L^\infty(\Omega)} v_\alpha$$

satisfies

$$\begin{cases} -\Delta \bar{v} \geq \|\delta^{2-\alpha} f\|_{L^\infty(\Omega)} \delta^{\alpha-2} & \text{in } \Omega, \\ \bar{v} = 0 & \text{on } \partial\Omega. \end{cases}$$

Thus, by the weak maximum principle (applied with \bar{v} and $-\bar{v}$, respectively, as super- and subsolution), we have that $|u| \leq \bar{v} = C_2 \alpha^{-1} (1 - \alpha)^{-1} \|\delta^{2-\alpha} f\|_{L^\infty(\Omega)} v_\alpha$ in Ω . This and (3.4) give (3.11).

When $f \geq 0$, estimate (3.12) can be established again via the maximum principle, this time taking advantage of the right-hand inequality in (3.5) and using $\underline{v} := \alpha^{-1} C_2^{-1} \inf_{\Omega} (\delta^{2-\alpha} f) v_\alpha$ as a subsolution. \square

To deal with solutions that blow up at the boundary, we need a different class of barriers. They are provided by the next lemma, which is essentially due to [11]. Following [11, section 3], we define, for $\beta \in (-1, 0)$,

$$(3.13) \quad V_\beta(x) := \begin{cases} \eta(x) & \text{for } x \in \Omega \setminus \Gamma_{\delta_0}, \\ \delta(x)^\beta & \text{for } x \in \Gamma_{\delta_0}, \\ 0 & \text{for } x \in \mathbb{R}^N \setminus \Omega, \end{cases}$$

where $\Gamma_t = \{x \in \Omega : \delta(x) < t\}$, the parameter $\delta_0 > 0$ is sufficiently small to have that $\delta \in C^2(\bar{\Gamma}_{\delta_0})$, and η is any positive function for which V_β is of class C^2 in Ω .

LEMMA 3.4. *Let $\beta \in (-1 + \sigma, 0)$ and $\Omega \subseteq \mathbb{R}^N$ be a bounded domain with boundary of class C^2 . Then, there exist two constants $\delta_1 \in (0, \delta_0]$ and $C_{\frac{3}{4}} \geq 1$, depending only*

on N , Ω , σ , and β , such that

$$(3.14) \quad C_{\sharp}^{-1} \delta^{\beta-2} \leq \Delta V_{\beta} \leq C_{\sharp} \delta^{\beta-2} \quad \text{in } \Gamma_{\delta_1},$$

$$(3.15) \quad C_{\sharp}^{-1} \delta^{\beta-2\sigma} \leq (-\Delta)^{\sigma} V_{\beta} \leq C_{\sharp} \delta^{\beta-2\sigma} \quad \text{in } \Gamma_{\delta_1}.$$

Proof. The inequalities in (3.14) are straightforward. Indeed, a simple computation using that $|\nabla\delta| = 1$ yields

$$\Delta V_{\beta} = |\beta| \delta^{\beta-2} (1 - \beta - \delta \Delta \delta) \quad \text{in } \Gamma_{\delta_0}.$$

Hence, (3.14) follows from the C^2 regularity of δ in Γ_{δ_0} and taking δ_1 suitably small.

On the other hand, (3.15) is the content of [11, Proposition 3.2(ii)], once one realizes that the quantity labeled as $\tau_0(\alpha)$ in [11] is equal to $-1 + \alpha$. This has already been observed in [2, Remark 3.1] and is a consequence of the fact that the function $x_{+}^{-1+\alpha}$ is α -harmonic in $(0, +\infty)$ —the function appearing in [11, formula (1.13)] is equal, up to an irrelevant factor, to the α -Laplacian of x_{+}^{τ} evaluated at $x = 1$. The α -harmonicity of $x_{+}^{-1+\alpha}$ in $(0, +\infty)$ can be verified in several ways—see, e.g., [3, Lemma 4.1] for a proof based on the computations of [17]. \square

4. Comparison principles. In this section, we prove the weak comparison principle of Proposition 1.2 and deduce from it some estimates on the supremum of subsolutions of (1.1).

Proof of Proposition 1.2. Assume first that both inequalities in (1.8) and on the first line of (1.7) are strict, i.e., that \underline{w} and \bar{w} satisfy

$$(4.1) \quad \begin{cases} -\Delta \underline{w} + \Phi(\cdot, (-\Delta)^{\sigma} \underline{w}) < -\Delta \bar{w} + \Phi(\cdot, (-\Delta)^{\sigma} \bar{w}) & \text{in } \Omega, \\ \underline{w} \leq \bar{w} & \text{in } \mathbb{R}^N \setminus \bar{\Omega}, \\ \limsup_{\Omega \ni x \rightarrow x_0} \underline{w}(x) < \liminf_{\Omega \ni x \rightarrow x_0} \bar{w}(x) & \text{for all } x_0 \in \partial\Omega. \end{cases}$$

Let $w := \underline{w} - \bar{w}$ and

$$M := \sup_{\Omega} w.$$

We claim that

$$(4.2) \quad M \leq 0.$$

Of course, if (4.2) is valid, then we are done. Therefore, we argue by contradiction and suppose that $M > 0$.

By the continuity of \underline{w} and \bar{w} inside Ω and the strict inequality on the third line of (4.1), there exists a point $x_M \in \Omega$ at which $w(x_M) = M$. As $w \leq 0$ outside of $\bar{\Omega}$, we infer that

$$w(x_M) = M = \max_{\mathbb{R}^N} w.$$

Accordingly,

$$-\Delta w(x_M) \geq 0 \quad \text{and} \quad (-\Delta)^{\sigma} w(x_M) \geq 0,$$

that is,

$$-\Delta \underline{w}(x_M) \geq -\Delta \bar{w}(x_M) \quad \text{and} \quad (-\Delta)^{\sigma} \underline{w}(x_M) \geq (-\Delta)^{\sigma} \bar{w}(x_M).$$

In view of this and the monotonicity of $\Phi(x_M, \cdot)$, we obtain that

$$-\Delta \underline{w}(x_M) + \Phi(x_M, (-\Delta)^{\sigma} \underline{w}(x_M)) \geq -\Delta \bar{w}(x_M) + \Phi(x_M, (-\Delta)^{\sigma} \bar{w}(x_M)),$$

in contradiction with the first inequality in (4.1). Thus, (4.2) holds true and the lemma is proved when (4.1) is in force.

Suppose now that w and \bar{w} satisfy the weaker hypotheses (1.7) and (1.8). Let $R > 0$ be large enough to have $\bar{\Omega} \subseteq B_R$ and consider the function $u_\sigma(x) := (R^2 - |x|^2)_+^\sigma$. By Lemma 3.1, we know that

$$(4.3) \quad (-\Delta)^\sigma u_\sigma > 0 \quad \text{and} \quad -\Delta u_\sigma \geq 2\sigma NR^{2\sigma-2} > 0 \quad \text{in } B_R.$$

Consequently, letting $\bar{w}_\varepsilon := \bar{w} + \varepsilon u_\sigma$ for any small $\varepsilon > 0$ and using again the monotonicity of Φ with respect to the second variable, we see that

$$\left\{ \begin{array}{ll} -\Delta w + \Phi(\cdot, (-\Delta)^\sigma w) < -\Delta \bar{w}_\varepsilon + \Phi(\cdot, (-\Delta)^\sigma \bar{w}_\varepsilon) & \text{in } \Omega, \\ w \leq \bar{w}_\varepsilon & \text{in } \mathbb{R}^N \setminus \Omega, \\ \limsup_{\Omega \ni x \rightarrow x_0} w(x) < \liminf_{\Omega \ni x \rightarrow x_0} \bar{w}_\varepsilon(x) & \text{for all } x_0 \in \partial\Omega. \end{array} \right.$$

By what we established before, $w \leq \bar{w}_\varepsilon \leq \bar{w} + \varepsilon R^{2\sigma}$ in the whole \mathbb{R}^N . The conclusion of the lemma now follows by letting $\varepsilon \downarrow 0$. \square

As applications of Proposition 1.2, we have two results providing upper bounds on the supremum of subsolutions of (1.1). Of course, from these one may easily deduce the corresponding lower bounds for supersolutions and two-sided bounds for solutions.

First, we suppose the right-hand side f in (1.1) to be a bounded function. In this case, it suffices to apply Proposition 1.2 in conjunction with the barrier of Lemma 3.1.

COROLLARY 4.1. *Let $\Omega \subseteq \mathbb{R}^N$ be a bounded open set and $\Phi : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ be a Carathéodory function with $\Phi(x, \cdot)$ nondecreasing and $\Phi(x, 0) \geq 0$ for a.e. $x \in \Omega$. Let $f \in L^\infty(\Omega)$, $g \in L^\infty(\partial\Omega)$, $h \in L^\infty(\mathbb{R}^N \setminus \bar{\Omega})$, and $w \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$ be such that*

$$(4.4) \quad \left\{ \begin{array}{ll} -\Delta w + \Phi(\cdot, (-\Delta)^\sigma w) \leq f & \text{in } \Omega, \\ w \leq g & \text{on } \partial\Omega, \\ w \leq h & \text{in } \mathbb{R}^N \setminus \bar{\Omega}. \end{array} \right.$$

Then,

$$\sup_\Omega w \leq \sup_{\partial\Omega} g_+ + \sup_{\mathbb{R}^N \setminus \bar{\Omega}} h_+ + C \text{diam}(\Omega)^2 \sup_\Omega f_+$$

for some constant $C > 0$ depending only on N and σ .

Proof. Write $R := \text{diam}(\Omega)$ and pick $x_0 \in \mathbb{R}^N$ in such a way that $\bar{\Omega} \subseteq B_R(x_0)$. Without loss of generality, we may assume that $x_0 = 0$. Similarly to what we did at the end of the proof of Proposition 1.2, we consider the function $u_\sigma(x) := (R^2 - |x|^2)_+^\sigma$, which satisfies (4.3). Hence, the function

$$\bar{w}(x) := \sup_\Omega g_+ + \sup_{\mathbb{R}^N \setminus \bar{\Omega}} h_+ + \frac{R^{2(1-\sigma)} \sup_\Omega f_+}{2N\sigma} u_\sigma(x), \quad x \in \mathbb{R}^N,$$

is such that

$$(4.5) \quad \left\{ \begin{array}{ll} -\Delta \bar{w} + \Phi(\cdot, (-\Delta)^\sigma \bar{w}) \geq f & \text{in } \Omega, \\ \bar{w} \geq g & \text{on } \partial\Omega, \\ \bar{w} \geq h & \text{in } \mathbb{R}^N \setminus \Omega. \end{array} \right.$$

The conclusion now follows by Proposition 1.2. \square

By combining Proposition 1.2 with Lemma 3.2, we may tackle the case when f is merely in $L_{\text{loc}}^\infty(\Omega)$ and blows up at the boundary of Ω at a strictly slower rate than the square of the inverse distance function.

COROLLARY 4.2. *Let $\Omega \subseteq \mathbb{R}^N$ be a bounded open set with boundary of class C^2 and $\Phi : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ be a Carathéodory function with $\Phi(x, \cdot)$ nondecreasing and $\Phi(x, 0) \geq 0$ for a.e. $x \in \Omega$. Let f be such that $\delta^{2-\alpha} f \in L^\infty(\Omega)$ for some $0 < \alpha \leq \bar{\alpha} < 1$, $g \in L^\infty(\partial\Omega)$, and $h \in L^\infty(\mathbb{R}^N \setminus \bar{\Omega})$. Let $\underline{w} \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$ be such that (4.4) holds true. Then,*

$$(4.6) \quad \sup_{\Omega} \underline{w} \leq \sup_{\partial\Omega} g_+ + \sup_{\mathbb{R}^N \setminus \bar{\Omega}} h_+ + C\alpha^{-1} \sup_{\Omega} (\delta^{2-\alpha} f_+)$$

for some constant $C > 0$ depending only on N , Ω , σ , and $\bar{\alpha}$.

Proof. For $M > 0$, define the nonnegative function

$$w(x) := \alpha^{-1} \sup_{\Omega} (\delta^{2-\alpha} f_+) (Mv_0(x) + C_2(1-\alpha)^{-1}v_\alpha(x)), \quad x \in \mathbb{R}^N,$$

where v_0 , v_α , and C_2 are as in Lemma 3.2. By estimates (3.5)–(3.7), in Ω we then have

$$\begin{aligned} -\Delta w &\geq \sup_{\Omega} (\delta^{2-\alpha} f_+) \delta^{\alpha-2}, \\ (-\Delta)^\sigma w &\geq \alpha^{-1} \sup_{\Omega} (\delta^{2-\alpha} f_+) (C_3^{-1}M - C_2^2(1-\alpha)^{-1} \text{diam}(\Omega)^\alpha) \delta^{-2\sigma} \geq 0, \end{aligned}$$

provided M is large enough, in dependence of N , Ω , σ , and $\bar{\alpha}$ only. Thus,

$$-\Delta w + \Phi(\cdot, (-\Delta)^\sigma w) \geq f \quad \text{in } \Omega.$$

This yields that $\bar{w} := w + \sup_{\partial\Omega} g_+ + \sup_{\mathbb{R}^N \setminus \bar{\Omega}} h_+$ satisfies (4.5). From Proposition 1.2 it follows that $\underline{w} \leq \bar{w}$. Estimate (4.6) is then a consequence of Lemma 3.2. \square

5. Existence. Proof of Theorem 1.1. We present here the proof of Theorem 1.1 under the notational conventions explained in subsection 1.2. For readability purposes, we split the proof into four intermediate steps:

- Step (1) First, we reduce (1.1) to an equivalent problem having vanishing boundary and exterior data.
- Step (2) The so-obtained Dirichlet problem will contain singular terms originating from the lack of smoothness of the fractional Laplacian at the boundary and we circumvent this issue by solving a family of regularized problems by cutting the singularities off.
- Step (3) We then obtain uniform estimates on the solutions to these regularized problems; to get stronger estimates, we use the family of weighted Sobolev spaces introduced in section 2, wherein all necessary notation can be found.
- Step (4) Finally, the estimates enable us to conclude that the solutions to the regularized problems accumulate at a solution of the original one. Its uniqueness then immediately follows from Proposition 1.2.

5.1. Reduction to homogeneous data. Let \bar{g} be the harmonic extension of g inside Ω , i.e., $\bar{g} \in C^2(\Omega) \cap C^0(\bar{\Omega})$ is the unique solution of the Dirichlet problem

$$\begin{cases} -\Delta \bar{g} = 0 & \text{in } \Omega, \\ \bar{g} = g & \text{on } \partial\Omega. \end{cases}$$

We define

$$\psi := \chi_{\bar{\Omega}}\bar{g} + \chi_{\mathbb{R}^N \setminus \bar{\Omega}}h = \begin{cases} \bar{g} & \text{in } \bar{\Omega}, \\ h & \text{in } \mathbb{R}^N \setminus \bar{\Omega}. \end{cases}$$

Notice that $\psi \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$.

Letting $v := u - \psi$, it is clear that (1.1) is equivalent to the problem²

$$(5.1) \quad \begin{cases} -\Delta v + P[v] = f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^N \setminus \Omega \end{cases}$$

with

$$P[v] := \left\{ |(-\Delta)^\sigma v + (-\Delta)^\sigma \psi|^{p-1} ((-\Delta)^\sigma v + (-\Delta)^\sigma \psi) \right\} \Big|_\Omega.$$

Observe that both $(-\Delta)^\sigma \psi$ and f are locally bounded and Hölder continuous functions in Ω . However, they may in general blow up at the boundary of Ω . Indeed, we have that

$$(5.2) \quad |(-\Delta)^\sigma \psi(x)| \leq C \left(\|g\|_{L^\infty(\partial\Omega)} + \|h\|_{L^\infty(\mathbb{R}^N \setminus \Omega)} \right) \delta(x)^{-2\sigma} \quad \text{for all } x \in \Omega.$$

To see this, on the one hand, by the maximum principle and the classical Schauder theory (e.g., [20, Theorem 4.6]), one gets that

$$\|\bar{g}\|_{L^\infty(\Omega)} + \delta(x)^2 \|D^2 \bar{g}\|_{L^\infty(B_{\delta(x)/4}(x))} \leq C \|g\|_{L^\infty(\partial\Omega)} \quad \text{for all } x \in \Omega.$$

Consequently, arguing as we did to get estimate (3.6) in Lemma 3.2, we find that

$$|(-\Delta)^\sigma (\chi_{\bar{\Omega}}\bar{g})(x)| \leq C \|g\|_{L^\infty(\partial\Omega)} \delta(x)^{-2\sigma}$$

for all $x \in \Omega$. On the other hand, by computing directly,

$$\begin{aligned} |(-\Delta)^\sigma (\chi_{\mathbb{R}^N \setminus \bar{\Omega}}h)(x)| &\leq c_{N,\sigma} \int_{\mathbb{R}^N \setminus B_{\delta(x)}(x)} \frac{\chi_{\mathbb{R}^N \setminus \bar{\Omega}}(y) |h(y)|}{|x-y|^{N+2\sigma}} dy \\ &\leq C \|h\|_{L^\infty(\mathbb{R}^N \setminus \bar{\Omega})} \delta(x)^{-2\sigma} \end{aligned}$$

for all $x \in \Omega$. The combination of the last two estimates leads to (5.2).

In light of these diverging behaviors, to solve (5.1) it is convenient to consider a family of suitably regularized problems. This will be the content of the next subsection.

5.2. Approximating problems. For any large integer j , consider the open set

$$\Omega_j := \{x \in \Omega : \delta(x) > 2^{-j}\}.$$

Then, let $\eta_j \in C_c^\infty(\mathbb{R}^N)$ be a cut-off function satisfying $0 \leq \eta_j \leq 1$ in \mathbb{R}^N , $\text{supp}(\eta_j) \subseteq \Omega_j$, $\eta_j \equiv 1$ in Ω_{j-1} , and $|\nabla \eta_j| \leq C_j$ in \mathbb{R}^N . We take into account the auxiliary problem

$$(5.3) \quad \begin{cases} -\Delta v + P_j[v] = \eta_j f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^N \setminus \Omega \end{cases}$$

²Let us stress here that $\mathbb{R}^N \setminus \Omega = \partial\Omega \cup (\mathbb{R}^N \setminus \bar{\Omega})$.

with $P_j[v] := \eta_j P[v]$. To find a solution of (5.3), we will look at it as a fixed-point problem.

Let $\beta \in (2\sigma, 2) \setminus \{1\}$ to be chosen later, in dependence of σ and p only, and consider the Banach space

$$\mathcal{X} := \left\{ w \in C^0(\mathbb{R}^N) \cap C^\beta(\overline{\Omega}) : w = 0 \text{ in } \mathbb{R}^N \setminus \Omega \right\},$$

endowed with the norm $\|w\|_{\mathcal{X}} := \|w\|_{C^\beta(\overline{\Omega})}$.

First, we claim that $P_j : \mathcal{X} \rightarrow L^\infty(\Omega)$ is a continuous mapping and that

$$(5.4) \quad \|P_j[w]\|_{L^\infty(\Omega)} \leq \overline{C}_j (1 + \|w\|_{\mathcal{X}}^p) \quad \text{for every } w \in \mathcal{X}.$$

The continuity easily follows from the fact that $\text{supp}(\eta_j) \subseteq \Omega_j$ and the estimate

$$\begin{aligned} & |(-\Delta)^\sigma w(x)| \\ &= \frac{c_{N,\sigma}}{2} \left| \int_{\mathbb{R}^N} \frac{2w(x) - w(x+z) - w(x-z)}{|z|^{N+2\sigma}} dz \right| \\ &\leq C \left(\int_{B_{2^{-j-1}}} \frac{[w]_{C^\beta(B_{2^{-j-1}}(x))}}{|z|^{N+2\sigma-\beta}} dz + \int_{\mathbb{R}^N \setminus B_{2^{-j-1}}} \frac{\|w\|_{L^\infty(\mathbb{R}^N \setminus B_{2^{-j-1}}(x))}}{|z|^{N+2\sigma}} dz \right) \\ &\leq C_j \|w\|_{\mathcal{X}}, \end{aligned}$$

which holds true for every $w \in \mathcal{X}$ and $x \in \Omega_j$. From this and (5.2), we also infer that

$$\begin{aligned} \|P_j[w]\|_{L^\infty(\Omega)} &\leq C \left(\|(-\Delta)^\sigma w\|_{L^\infty(\Omega_j)}^p + \|(-\Delta)^\sigma \psi\|_{L^\infty(\Omega_j)}^p \right) \\ &\leq C_j \left(\|w\|_{\mathcal{X}}^p + \|g\|_{L^\infty(\partial\Omega)}^p + \|h\|_{L^\infty(\mathbb{R}^N \setminus \Omega)}^p \right), \end{aligned}$$

which gives (5.4).

Denote now by $(-\Delta)^{-1}$ the inverse of the Dirichlet Laplacian in Ω , i.e., let $(-\Delta)^{-1}F$ be the only solution ϕ of the problem

$$\begin{cases} -\Delta\phi = F & \text{in } \Omega, \\ \phi = 0 & \text{on } \partial\Omega. \end{cases}$$

By the classical Calderón–Zygmund theory and the Sobolev embedding, the operator $(-\Delta)^{-1}$ maps $L^\infty(\Omega)$ into $W_0^{1,q}(\Omega) \cap W^{2,q}(\Omega) \cap C^\gamma(\overline{\Omega})$ for every $q \in (1, +\infty)$ and $\gamma \in (0, 2)$. Also,

$$(5.5) \quad \|(-\Delta)^{-1}F\|_{C^\gamma(\overline{\Omega})} \leq C_\gamma \|F\|_{L^\infty(\Omega)}.$$

Pick now any $q \in [1, +\infty)$ and $\gamma \in (\beta, 2)$. Then, the standard inclusion $\iota : C^\gamma(\overline{\Omega}) \rightarrow C^\beta(\overline{\Omega})$ is compact. Hence, the mapping $T_j : \mathcal{X} \rightarrow \mathcal{X}$ defined by

$$T_j[w] := \begin{cases} \iota((-\Delta)^{-1}(\eta_j f - P_j[w])) & \text{in } \Omega, \\ 0 & \text{in } \mathbb{R}^N \setminus \Omega \end{cases}$$

for all $w \in \mathcal{X}$ is also compact. We stress that $T_j[w]$ defines a continuous function in \mathbb{R}^N —and is thus an element of \mathcal{X} —since its restriction to Ω belongs to $W_0^{1,q}(\Omega) \cap C^\beta(\overline{\Omega})$.

Notice then that $v \in \mathcal{X} \cap W^{2,q}(\Omega)$ is a solution of (5.3) if and only if it is a fixed point of the map T_j . Since T_j is compact, we can show the existence of a fixed point using the Leray–Schauder theorem (see, e.g., [20, Theorem 11.3]), provided we check that

$$(5.6) \quad \|v\|_{\mathcal{X}} \leq \bar{C}_j \quad \text{for every } v \in \mathcal{X} \text{ such that } v = \lambda T_j[v] \text{ for some } \lambda \in [0, 1].$$

To see this, note that if $v \in \mathcal{X}$ satisfies $v = \lambda T_j[v]$, then v is a $C^0(\mathbb{R}^N) \cap W^{2,q}(\Omega)$ solution of

$$\begin{cases} -\Delta v + \lambda P_j[v] = \lambda \eta_j f & \text{in } \Omega, \\ v = 0 & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$

Then, by standard elliptic regularity, v is actually of class C^2 in Ω , and therefore the function $u := v + \psi$ is a $L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$ solution of

$$\begin{cases} -\Delta u + \lambda \eta_j |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u = \lambda \eta_j f & \text{in } \Omega, \\ u = g & \text{on } \partial\Omega, \\ u = h & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$

Hence, by the comparison principle of Corollary 4.2, we infer that $\|u\|_{L^\infty(\Omega)}$ is universally bounded, and thus

$$(5.7) \quad \|v\|_{L^\infty(\Omega)} \leq \bar{C}.$$

Knowing this, we may proceed to show the validity of (5.6). First, we remark that

$$\|\varphi\|_{C^{\alpha_1}(\bar{\Omega})} \leq C \|\varphi\|_{L^\infty(\Omega)}^{1-\alpha_1/\alpha_2} \|\varphi\|_{C^{\alpha_2}(\bar{\Omega})}^{\alpha_1/\alpha_2} \quad \text{for all } \varphi \in C^{\alpha_2}(\bar{\Omega}) \text{ and } 0 < \alpha_1 < \alpha_2 < 2.$$

See, e.g., [25, Proposition 1.1.3(iii)]. Thanks to this, (5.7), (5.5), and (5.4), we compute

$$\begin{aligned} \|v\|_{\mathcal{X}} &= \|v\|_{C^\beta(\bar{\Omega})} \leq C \|v\|_{L^\infty(\Omega)}^{1-\beta/\gamma} \|v\|_{C^\gamma(\bar{\Omega})}^{\beta/\gamma} \leq \bar{C} \lambda^{\beta/\gamma} \|T_j[v]\|_{C^\gamma(\bar{\Omega})}^{\beta/\gamma} \\ &\leq \bar{C}_\gamma \|\eta_j f - P_j[v]\|_{L^\infty(\Omega)}^{\beta/\gamma} \leq \bar{C}_\gamma \left(\|\eta_j f\|_{L^\infty(\Omega)}^{\beta/\gamma} + \|P_j[v]\|_{L^\infty(\Omega)}^{\beta/\gamma} \right) \\ &\leq \bar{C}_{\gamma,j} \left(1 + \|v\|_{\mathcal{X}}^{\beta p/\gamma} \right). \end{aligned}$$

Notice that, since $\sigma p < 1$, we can choose $\beta \in (2\sigma, 2) \setminus \{1\}$ (close to 2σ) and $\gamma \in (\beta, 2)$ (close to 2) in a way that $\beta p/\gamma < 1$. By doing this and applying the weighted Young’s inequality, claim (5.6) easily follows from the above estimate.

Accordingly, we can apply the Leray–Schauder theorem and conclude that there exists a fixed point $v_j \in \mathcal{X}$ for the map T_j , i.e., a solution $v_j \in C^0(\mathbb{R}^N) \cap C^2(\Omega)$ of problem (5.3). Also,

$$(5.8) \quad \|v_j\|_{L^\infty(\Omega)} \leq \bar{C},$$

as a consequence of (5.7).

5.3. Uniform estimates. We now want to let $j \uparrow \infty$ and show that the limit of the v_j ’s is a solution of (5.1). In order to do this, we need estimates for v_j that do not depend on j . Note that we already know that each v_j satisfies the uniform L^∞ bound (5.8).

Let $q \in (1, +\infty)$. As v_j is a solution of (5.3), by Proposition 2.2, Lemma 2.3 (applied here with pq in place of p , $r = pq$, $\theta = N + 2q > N + 2\sigma pq$, and $\varepsilon = 2(1 - \sigma p)q > 0$), and estimate (5.8), we have

$$\begin{aligned}
 (5.9) \quad & \|v_j\|_{L_N^{2,q}(\Omega)} \\
 & \leq C_q \left(\|\eta_j f - P_j[v_j]\|_{L_{N+2q}^q(\Omega)} + \|v_j\|_{L_N^q(\Omega)} \right) \\
 & \leq C_q \left(\|(-\Delta)^\sigma v_j\|_{L_{N+2q}^{pq}(\Omega)}^p + \|(-\Delta)^\sigma \psi\|_{L_{N+2q}^{pq}(\Omega)}^p + \|f\|_{L_{N+2q}^q(\Omega)} + \|v_j\|_{L^q(\Omega)} \right) \\
 & \leq \bar{C}_q \left(1 + \|v_j\|_{L_N^{2\sigma,pq}(\Omega)}^p \right).
 \end{aligned}$$

Notice that, to get the last inequality, we also took advantage of the fact that, thanks to (5.2) and (1.4),

$$\begin{aligned}
 \|(-\Delta)^\sigma \psi\|_{L_{N+2q}^{pq}(\Omega)}^p + \|f\|_{L_{N+2q}^q(\Omega)}^q & \leq C_q \int_{\Omega} \left(|(-\Delta)^\sigma \psi(x)|^{pq} + |f(x)|^q \right) \delta(x)^{2q} dx \\
 & \leq \bar{C}_q \int_{\Omega} \left(\delta(x)^{2q(1-\sigma p)} + \delta(x)^{\alpha q} \right) dx \leq \bar{C}_q.
 \end{aligned}$$

The interpolation inequality of, say, [26, Corollary 2.1.8], along with the representation of Proposition 2.1(v) for the space $L_N^{2\sigma,pq}(\Omega)$ and again (5.8) then give that

$$\|v_j\|_{L_N^{2\sigma,pq}(\Omega)} \leq C_q \|v_j\|_{L_N^{\frac{1-\sigma}{1-\sigma p}}(\Omega)}^{1-\sigma} \|v_j\|_{L_N^{2,q}(\Omega)}^\sigma \leq \bar{C}_q \|v_j\|_{L_N^{2,q}(\Omega)}.$$

By plugging this into (5.9) and taking advantage of the weighted Young’s inequality, we conclude that

$$(5.10) \quad \|v_j\|_{L_N^{2,q}(\Omega)} \leq \bar{C}_q \quad \text{for every } j \in \mathbb{N}.$$

Note that, once again, we used in a crucial way that $\sigma p < 1$.

Next, we claim that, for any small $\varepsilon \in (0, 1)$,

$$(5.11) \quad |v_j(x)| \leq \bar{C}_\varepsilon \delta(x)^{\min\{1-\varepsilon, 2(1-\sigma p)-\varepsilon, \alpha\}} \quad \text{for every } x \in \Omega \text{ and every } j \in \mathbb{N}.$$

To check this, let $q > N$ and notice that, using (5.10) and the standard Morrey’s inequality,

$$\begin{aligned}
 [\nabla v_j]_{C^{1-\frac{N}{q}}(B_{\delta(x)/2}(x))} & \leq C_q \left(\delta(x)^{-q} \|\nabla v_j\|_{L^q(B_{\delta(x)/2}(x))}^q + \|D^2 v_j\|_{L^q(B_{\delta(x)/2}(x))}^q \right)^{\frac{1}{q}} \\
 & \leq C_q \delta(x)^{-2} \left(\int_{\Omega} |\nabla v_j(y)|^q \delta(y)^q dy + \int_{\Omega} |D^2 v_j(y)|^q \delta(y)^{2q} dy \right)^{\frac{1}{q}} \\
 & \leq C_q \delta(x)^{-2} \|v_j\|_{L_N^{2,q}(\Omega)} \leq \bar{C}_q \delta(x)^{-2}.
 \end{aligned}$$

Hence, from this and (5.8), it easily follows that

$$\begin{aligned}
 |(-\Delta)^\sigma v_j(x)| & \leq C \left(\int_{B_{\delta(x)/2}} \frac{[\nabla v_j]_{C^{1-N/q}(B_{\delta(x)/2}(x))}}{|z|^{N+2\sigma-2+N/q}} dz + \int_{\mathbb{R}^N \setminus B_{\delta(x)/2}} \frac{\|v_j\|_{L^\infty(\Omega)}}{|z|^{N+2\sigma}} dz \right) \\
 & \leq \bar{C}_q \delta(x)^{-2\sigma-N/q}
 \end{aligned}$$

for every $x \in \Omega$ and q large enough. Estimate (5.11) is then a consequence of this inequality, (5.2), and Lemma 3.3, recalling that v_j is a solution of (5.3) and taking $q = Np/\varepsilon$, with $\varepsilon > 0$ sufficiently small.

Thanks to the uniform bounds (5.10) and (5.11), we are now able to get a limit for v_j as $j \uparrow \infty$ and obtain a solution of (5.1). Bear in mind that estimate (5.10) gives in particular that

$$\|v_j\|_{W^{2,q}(\Omega_k)} \leq \bar{C}_{q,k}$$

for any $q > N$. To see this, it is convenient to recall the equivalent representation for $L^{2,q}_N(\Omega)$ given in Proposition 2.1(iv). Thus, by Morrey’s inequality,

$$(5.12) \quad \|v_j\|_{C^\gamma(\bar{\Omega}_k)} \leq \bar{C}_{\gamma,k}$$

for any fixed $\gamma \in (\max\{2\sigma, 1\}, 2)$ and every large integer j and k .

5.4. Passage to the limit. By the compact embedding of Hölder spaces, bound (5.12), and a standard diagonal procedure, $(v_j)_{j \in \mathbb{N}}$ converges (up to a subsequence) to a function v in $C^\gamma_{loc}(\Omega)$. Letting $j \uparrow \infty$ in (5.11), we obtain that the extension of v to 0 outside Ω (that we still call v) defines a continuous function on the whole \mathbb{R}^N . By the dominated convergence theorem and the uniform L^∞ bound (5.8), we also have that $v_j \rightarrow v$ in $L^1(\mathbb{R}^N)$.

Passing (5.3) to the limit, one easily obtains that v is a weak solution of $-\Delta v + P[v] = f$ in every open set compactly contained in Ω , that is,

$$\int_\Omega (\nabla v \cdot \nabla \varphi + P[v]\varphi) = \int_\Omega f\varphi \quad \text{for all } \varphi \in C_c^\infty(\Omega).$$

Notice that this can be done since $P_j[v_j]$ converges to $P[v]$ in $L^\infty_{loc}(\Omega)$, as a consequence of the convergence of v_j to v in $C^\gamma_{loc}(\Omega)$ and $L^1(\mathbb{R}^N)$. Now, since f and $P[v]$ are both locally Hölder continuous functions in Ω (as $v \in C^\gamma_{loc}(\Omega)$ with $\gamma > 2\sigma$), by elliptic regularity we conclude that v belongs to $C^0(\mathbb{R}^N) \cap C^2(\Omega)$ and solves (5.1) pointwise. The proof of Theorem 1.1 is then complete.

6. Nonexistence. Proofs of Theorems 1.3 and 1.4. In this section, we establish our two nonexistence results, which are valid, respectively, when $\sigma p \geq 1$ or when the right-hand side f blows up too rapidly at the boundary of Ω .

First, we deal with the case of vanishing right-hand side and critical or supercritical regime.

Proof of Theorem 1.3. Letting $\alpha, \varepsilon > 0$, $m := \max_{\partial\Omega} g > 0$, $v_0 = \chi_{\bar{\Omega}}$, and v_α be as in Lemma 3.2, we consider the function

$$\bar{w}_{\alpha,\varepsilon} := m(v_0 - \varepsilon v_\alpha).$$

We claim that there exists an $\varepsilon_0 \in (0, 1)$ small enough such that $\bar{w}_{\alpha,\varepsilon}$ is a supersolution to problem (1.1) for any $\alpha \in (0, \sigma)$ and $\varepsilon \in (0, \varepsilon_0]$.

Clearly, $\bar{w}_{\alpha,\varepsilon} \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$. Moreover,

$$\bar{w}_{\alpha,\varepsilon} \geq g \text{ on } \partial\Omega \quad \text{and} \quad \bar{w}_{\alpha,\varepsilon} \geq h \text{ in } \mathbb{R}^N \setminus \Omega,$$

thanks to the definition of m and the fact that h is nonpositive. Therefore, in order

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to prove that $\bar{w}_{\alpha,\varepsilon}$ is a supersolution to (1.9), we only need to check that

$$(6.1) \quad -\Delta \bar{w}_{\alpha,\varepsilon} + |(-\Delta)^\sigma \bar{w}_{\alpha,\varepsilon}|^{p-1} (-\Delta)^\sigma \bar{w}_{\alpha,\varepsilon} \geq 0 \quad \text{in } \Omega,$$

provided ε is sufficiently small, uniformly with respect to $\alpha \in (0, \sigma)$.

To do this, we take into account formulas (3.5)–(3.7) of Lemma 3.2, which give

$$-\Delta v_\alpha \leq C_* \delta^{\alpha-2} \quad \text{and} \quad (-\Delta)^\sigma v_\alpha \leq C_* \delta^{-2\sigma} \quad \text{in } \Omega$$

and

$$-\Delta v_0 = 0 \quad \text{and} \quad (-\Delta)^\sigma v_0 \geq C_*^{-1} \delta^{-2\sigma} \quad \text{in } \Omega$$

for some constant $C_* \geq 1$ depending only on N , σ , and Ω . In particular,

$$-\Delta \bar{w}_{\alpha,\varepsilon} = \varepsilon m \Delta v_\alpha \geq -C_* m \varepsilon \delta^{\alpha-2} \quad \text{in } \Omega$$

and

$$(-\Delta)^\sigma \bar{w}_{\alpha,\varepsilon} = m \left((-\Delta)^\sigma v_0 - \varepsilon (-\Delta)^\sigma v_\alpha \right) \geq m (C_*^{-1} - C_* \varepsilon) \delta^{-2\sigma} \geq \frac{m}{2C_*} \delta^{-2\sigma} \quad \text{in } \Omega,$$

provided $\varepsilon \leq (2C_*^2)^{-1}$. In light of these two relations, we obtain that, in Ω ,

$$\begin{aligned} & -\Delta \bar{w}_{\alpha,\varepsilon} + |(-\Delta)^\sigma \bar{w}_{\alpha,\varepsilon}|^{p-1} (-\Delta)^\sigma \bar{w}_{\alpha,\varepsilon} \\ & \geq -C_* m \varepsilon \delta^{\alpha-2} + \left(\frac{m}{2C_*} \right)^p \delta^{-2\sigma p} \\ & \geq \left(\frac{m}{2C_*} \right)^p \delta^{-2\sigma p} \left(1 - \frac{2^p C_*^{p+1} \text{diam}(\Omega)^{2(\sigma p-1)+\alpha}}{m^{p-1}} \varepsilon \right) \geq 0 \end{aligned}$$

if ε is small enough, depending on N , σ , p , Ω , and m only. Note that the second inequality holds since, by assumption, $\sigma p \geq 1$.

We have, therefore, proved the validity of (6.1), and thus that $\bar{w}_{\alpha,\varepsilon}$ is a supersolution to problem (1.9) for any $\alpha \in (0, \sigma)$ and any $\varepsilon \in (0, \varepsilon_0]$ with $\varepsilon_0 \in (0, 1)$ independent of α . Suppose now that there exists a solution $u \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$ of (1.9). By the weak comparison principle of Lemma 1.2, we then deduce that

$$u(x) \leq \bar{w}_{\alpha,\varepsilon}(x) \quad \text{for any } x \in \Omega, \alpha \in (0, \sigma), \text{ and } \varepsilon \in (0, \varepsilon_0].$$

By taking the limit as $\alpha \downarrow 0$, this in turn implies that

$$u(x) \leq m(1 - \varepsilon_0) \quad \text{for any } x \in \Omega.$$

In particular, since $m > 0$, we infer that

$$\sup_\Omega u \leq m(1 - \varepsilon_0) = (1 - \varepsilon_0) \max_{\partial\Omega} g < \max_{\partial\Omega} g,$$

which contradicts the fact that u attains continuously the boundary datum g . \square

Next, we establish the nonexistence of solutions also in the case when the right-hand side is too singular at the boundary.

Proof of Theorem 1.4. For $\alpha \in (0, \sigma)$, let v_α be as in Lemma 3.2, and define

$$\underline{w}_{\alpha,\varepsilon} := \varepsilon v_\alpha$$

for any $\varepsilon \in (0, 1)$. With the help of (3.5), (3.6), and the fact that $\sigma p < 1$, we compute

$$\begin{aligned} & -\Delta \underline{w}_{\alpha,\varepsilon} + |(-\Delta)^\sigma \underline{w}_{\alpha,\varepsilon}|^{p-1} (-\Delta)^\sigma \underline{w}_{\alpha,\varepsilon} \\ & \leq C_2 \varepsilon \delta^{\alpha-2} + C_2^p \varepsilon^p \text{diam}(\Omega)^{\alpha p} \delta^{-2\sigma p} \\ & \leq \varepsilon C_2^p \left(\text{diam}(\Omega)^\alpha + \text{diam}(\Omega)^{\alpha p + 2(1-\sigma p)} \right) \delta^{-2} \leq \kappa \delta^{-2} \leq f, \end{aligned}$$

provided ε is chosen small enough, depending on N, σ, p, Ω , and κ only. Note that ε can be chosen uniformly with respect to $\alpha \in (0, \sigma)$. Accordingly, $\underline{w}_{\alpha,\varepsilon}$ is a subsolution of problem (1.11) for any $\alpha \in (0, \sigma)$. By the comparison principle of Proposition 1.2, we then have that any solution $u \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\overline{\Omega})$ of (1.11) must satisfy $u \geq \underline{w}_{\alpha,\varepsilon} = \varepsilon v_\alpha$ in Ω . Hence,

$$u(x) \geq \varepsilon \lim_{\alpha \downarrow 0} \tau(x)^\alpha = \varepsilon \quad \text{for any } x \in \Omega,$$

in contradiction with the fact that $u \in C^0(\overline{\Omega})$ and the homogeneous boundary condition in (1.11). □

7. Boundary blow-up solutions. Proof of Theorem 1.5. Here, we construct solutions of problem (1.12) which blow up at the boundary of Ω , thus establishing Theorem 1.5. We will do this by first solving approximating Dirichlet problems with larger and larger data on $\partial\Omega$ and then passing to the limit. This last step will be possible thanks to the barriers provided by the following preliminary result.

LEMMA 7.1. *Let $\Omega \subseteq \mathbb{R}^N$ be a bounded domain with boundary of class C^2 . For*

$$p \in \left(\frac{3-\sigma}{1+\sigma}, \frac{1}{\sigma} \right),$$

let

$$(7.1) \quad \gamma = \gamma(\sigma, p) := -\frac{2(1-\sigma p)}{p-1} \in (-1+\sigma, 0)$$

and V_γ be defined as in (3.13). Then, there exist two constants $A, B \geq 1$, depending only on N, Ω, σ , and p , such that the $L^1(\mathbb{R}^N) \cap C^2(\Omega)$ function

$$\bar{u}(x) := AV_\gamma(x) + B\chi_\Omega(x), \quad x \in \mathbb{R}^N,$$

satisfies

$$-\Delta \bar{u} + |(-\Delta)^\sigma \bar{u}|^{p-1} (-\Delta)^\sigma \bar{u} \geq 0 \quad \text{in } \Omega.$$

Proof. In light of estimates (3.14)–(3.15) of Lemma 3.4 and (3.7) of Lemma 3.2 (recall that $v_0 = \chi_\Omega$ a.e. in \mathbb{R}^N), we have that

$$-\Delta \bar{u} \geq -C_\sharp A \delta^{\gamma-2}$$

and

$$(-\Delta)^\sigma \bar{u} \geq C_\sharp^{-1} A \delta^{\gamma-2\sigma} + C_3^{-1} B \delta^{-2\sigma} \geq C_\sharp^{-1} A \delta^{\gamma-2\sigma} \quad \text{in } \Gamma = \{x \in \Omega : \delta(x) < \delta_1\}.$$

Since, by (7.1), we have $\gamma - 2 = (\gamma - 2\sigma)p$, the above two inequalities give that

$$-\Delta \bar{u} + |(-\Delta)^\sigma \bar{u}|^{p-1} (-\Delta)^\sigma \bar{u} \geq C_\sharp^{-p} A^p \delta^{(\gamma-2\sigma)p} \left(1 - C_\sharp^{p+1} A^{1-p} \right) \geq 0 \quad \text{in } \Gamma,$$

provided A is large enough. The fact that, for B large, the same inequality also holds in $\Omega \setminus \Gamma$ is a simple consequence of the left-hand bound in (3.7) and of the $C^2(\Omega) \cap L^1(\mathbb{R}^N)$ regularity of V_γ —which yields in particular that $-\Delta V_\gamma$ and $(-\Delta)^\sigma V_\gamma$ are both bounded in $\Omega \setminus \Gamma$. \square

Proof of Theorem 1.5. For any $j \in \mathbb{N}$, consider the solution $u_j \in L^\infty(\mathbb{R}^N) \cap C^2(\Omega) \cap C^0(\bar{\Omega})$ of problem (1.1) associated to $g \equiv j$ on $\partial\Omega$, $f \equiv 0$ in Ω , and $h \equiv 0$ in $\mathbb{R}^N \setminus \bar{\Omega}$ —its existence and uniqueness is guaranteed by Theorem 1.1. By the comparison principle of Proposition 1.2, $(u_j)_{j \in \mathbb{N}}$ is a nondecreasing sequence bounded above by the function \bar{u} of Lemma 7.1. Let now $j \uparrow \infty$ to get that $(u_j)_{j \in \mathbb{N}}$ converges monotonically to some $u \leq \bar{u}$. We will show that this pointwise limit is the sought solution.

Our argument is similar to the one displayed in subsections 5.3–5.4. Let $q > 1$ and $\theta \geq Npq$ be chosen later. By Proposition 2.2 and Lemma 2.3 (applied with pq in place of p , $\theta + 2q$ in place of θ , $r = 1$, and $\varepsilon = 2(1 - \sigma p)q > 0$), we have

$$\begin{aligned} \|u_j\|_{L_\theta^{2,q}(\Omega)} &\leq C \left(\|(-\Delta)^\sigma u_j\|_{L_{\theta+2q}^{pq}(\Omega)}^p + \|u_j\|_{L_\theta^q(\Omega)} \right) \\ &\leq C \left(\|u_j\|_{L_\theta^{2\sigma,pq}(\Omega)}^p + \|u_j\|_{L^1(\Omega)}^p + \|u_j\|_{L_\theta^q(\Omega)} \right) \end{aligned}$$

for some constant $C > 0$ depending only on N, p, q, θ, σ , and Ω . In view of Proposition 2.1(v) and [26, Corollary 2.1.8], we estimate

$$\|u_j\|_{L_\theta^{2\sigma,pq}(\Omega)} \leq C \|u_j\|_{L_\theta^{2,q}(\Omega)}^\sigma \|u_j\|_{L_\theta^{\frac{pq(1-\sigma)}{1-\sigma p}}(\Omega)}^{1-\sigma}.$$

Thus, using the weighted Young's inequality along with the facts that $\sigma p < 1$, $p \geq 1$, and $0 \leq u_j \leq \bar{u}$,

$$\begin{aligned} \|u_j\|_{L_\theta^{2,q}(\Omega)} &\leq C \left(\|u_j\|_{L_\theta^{\frac{pq(1-\sigma)}{1-\sigma p}}(\Omega)}^{\frac{(1-\sigma)p}{1-\sigma p}} + \|u_j\|_{L^1(\Omega)}^p + \|u_j\|_{L_\theta^q(\Omega)} \right) \\ &\leq C \left(\|\bar{u}\|_{L_\theta^{\frac{pq(1-\sigma)}{1-\sigma p}}(\Omega)}^{\frac{(1-\sigma)p}{1-\sigma p}} + \|\bar{u}\|_{L^1(\Omega)}^p + 1 \right). \end{aligned}$$

Notice that the first term involving \bar{u} on the right-hand side is finite, provided θ is taken sufficiently large in dependence of N, p, q , and σ only, whereas the second term is always finite, as $\bar{u} \in L^1(\Omega)$.

As the last estimate holds for every $q > 1$, by compactness we deduce that $(u_j)_{j \in \mathbb{N}}$ actually converges to u in $C_{\text{loc}}^{1,\alpha}(\Omega)$ for every $\alpha \in (0, 1)$. Using this, it is easy to see that u satisfies

$$\int_{\Omega} \left(\nabla u \cdot \nabla \varphi + |(-\Delta)^\sigma u|^{p-1} (-\Delta)^\sigma u \varphi \right) = 0 \quad \text{for all } \varphi \in C_c^\infty(\Omega).$$

By standard elliptic regularity, we then get that $u \in C^2(\Omega)$ and solves the equation in the pointwise sense.

Estimate (1.13) is an immediate consequence of the pointwise inequalities $0 \leq u \leq \bar{u}$. The fact that $u > 0$ in Ω follows from a simple strong maximum principle. Finally, for all $x_0 \in \partial\Omega$ we have

$$\liminf_{\Omega \ni x \rightarrow x_0} u(x) \geq \sup_{j \in \mathbb{N}} \lim_{\Omega \ni x \rightarrow x_0} u_j(x) = \sup_{j \in \mathbb{N}} j = +\infty,$$

and the proof is complete. \square

8. Comments, open questions, and motivations. We conclude the paper with a few remarks on possible extensions of our results, points left open by our analysis, and possible applications.

- (i) Though stated for the specific operator $u \mapsto -\Delta u + |(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u$, the main results of this paper can be extended to a larger class of operators having p -growth in $(-\Delta)^\sigma u$ and satisfying the comparison principle of Proposition 1.2. For instance, Theorem 1.1 is also valid for operators of the form

$$u \mapsto -\Delta u + \Phi(\cdot, (-\Delta)^\sigma u),$$

where $\Phi : \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is a uniformly Hölder continuous function satisfying $\Phi(\cdot, 0) \geq 0$ in Ω and the growth condition

$$|\Phi(x, t)| \leq C(1 + |t|^p) \quad \text{for all } x \in \Omega, t \in \mathbb{R}$$

for some constant $C > 0$, in addition to the assumptions of Proposition 1.2.

- (ii) While, in light of the existence/nonexistence dichotomy provided by Theorems 1.1 and 1.3, the p -growth structure clearly cannot be fully abandoned, it would be nice to understand whether our results could be extended to operators which do not satisfy the hypotheses of Proposition 1.2, such as $u \mapsto -\Delta u + |(-\Delta)^\sigma u|^p$ or $u \mapsto -\Delta u - |(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u$. Indeed, in these cases the absence of comparison principles prevents one from using barrier arguments: these are at the core of a priori estimates (see Corollary 4.2), which are in turn essential to deduce uniform estimates (see subsection 5.3) making our approximating strategy feasible (see subsection 5.4).

For some more details regarding this issue in the case $p = 1$, see also [10, Appendix A].

- (iii) Theorem 1.5 gives the existence of a solution to $-\Delta u + |(-\Delta)^\sigma u|^{p-1}(-\Delta)^\sigma u = 0$ in Ω which vanishes a.e. outside of Ω and blows up at its boundary, from the inside. As a byproduct of the method of construction, we obtain the upper bound (1.13) on its blow-up rate. Unfortunately, we are not able to determine either a corresponding lower bound or the uniqueness of the solution. We believe it would be interesting to investigate both these issues.
- (iv) As a matter of fact, (1.1) is related to the fractional Lane–Emden equation (with sign-changing nonlinearity)

$$(-\Delta)^s v + |v|^{p-1}v = f \quad \text{in } \Omega, \quad s = 1 - \sigma \in (0, 1).$$

The two can be bridged by simply relabeling, at least formally, $(-\Delta)^\sigma u = v$ in \mathbb{R}^N . In doing so, one has to pay attention to what happens to the boundary conditions.

For example, if we perform this change of variable on a nonnegative solution u of (1.1) with $g \geq 0$ and $h \leq 0$, then v solves

$$\begin{cases} (-\Delta)^s v + |v|^{p-1}v = f & \text{in } \Omega, \\ v \leq 0 & \text{in } \mathbb{R}^N \setminus \bar{\Omega} \end{cases}$$

and may therefore act as a subsolution to the fractional problem. We believe that this could contribute to the study of *very large solutions* of the fractional

Laplacian, i.e., solutions of

$$\begin{cases} (-\Delta)^s v + |v|^{p-1} v = f & \text{in } \Omega, \\ \frac{v}{\delta^{s-1}} = +\infty & \text{on } \partial\Omega, \\ v = 0 & \text{in } \mathbb{R}^N \setminus \bar{\Omega} \end{cases}$$

for which there are still a number open problems, such as uniqueness and clear boundary asymptotics (see [2, 11]).

- (v) For $p = 1$, the operator $-\Delta + (-\Delta)^\sigma$, $\sigma \in (0, 1)$, is the infinitesimal generator of a discontinuous Markov process which has both diffusion and jump components: roughly, this means that trajectories look like disconnected portions of Brownian motions. This process belongs to the class of *Lévy processes* (stochastic processes with stationary and independent increments): these are, in particular, uniquely characterized by their characteristic function via the *Lévy–Khintchine formula*; see, for example, [6]. The interested reader can also check the self-contained presentation in [16, Appendix B] on how to recover this process as a limit of discrete random walks.

Lévy processes are widely used in models from mathematical finance [14], especially dealing with *option pricing*, and are related to optimization [9] and stochastic control [28] problems described by a Hamilton–Jacobi–Bellman fully nonlinear integro-differential equation. The classical theory of Hamilton–Jacobi–Bellman equations, which is making use only of the continuous diffusion entailed by the Laplacian, gives rise to boundary value problems of the form (1.10)—possibly with $g \equiv +\infty$; see [22]—for the optimal *cost function*.

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