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Investigating the (Poly)Radicaloid Nature of Real-World Organic Compounds with DFT-Based Methods

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11 **Investigating the (poly)radicaloid nature**  
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13 **DFT-based methods**  
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## Abstract

Recent advances in the synthesis of stable organic (open-shell) polyradicaloids have opened their application as active compounds for emerging technologies. These systems typically exhibit small energy differences between states with different spin multiplicities, which are intrinsically difficult to calculate by theoretical methods. We thus apply here some DFT-based variants (FT-DFT, SF-DFT, and SF-TDDFT) on a test set of large and real-world molecules, as test systems for which such energy differences are experimentally available, also comparing systematically with RAS-SF results to infer if shortcomings of previous DFT applications are corrected. Additionally, we explore the spin-spin contribution to the ZFS tensor, of high interest for EPR spectroscopy, and derive the spatial extent of the corresponding (photoexcited) triplet state.

*Key words:* organic (poly)radicals, low- and high-spin states, finite-temperature DFT, spin-flip (TD)DFT, ZFS tensor.

# 1 Introduction

Studying the (poly)radical character of organic molecules is a long-standing field of research due to the many envisioned applications of these compounds.<sup>1,2</sup> Recent experimental developments, in ultrahigh-vacuum surfaces or using non-standard synthetic routes, have propelled the synthesis of highly challenging species including classically studied carbon-rich radicals<sup>3</sup> like triangulenes,<sup>4</sup> graphene nanoribbons with zigzag edges,<sup>5,6</sup> kekulenes,<sup>7</sup> long acenes,<sup>8,9</sup> cyclic nanobelts,<sup>10,11</sup> etc. All these systems share a complicated electronic structure, with (near-)degenerated orbitals lying within the gap between occupied and virtual ones, leading to small exchange interactions and thus close in energy low- (e.g. singlet or doublet) and high-spin (e.g. triplet or quartet) many-body states.<sup>12</sup> Furthermore, C-based magnetism is gaining attention for nanographene fragments since long time ago,<sup>13</sup> and it has been recently demonstrated for well-defined geometrical C-based structures, like those arising from planar conjugated hydrocarbons, how to anticipate the spin multiplicity and energy ordering of the corresponding states,<sup>14</sup> thus complementing the Ovchinnikov's rule<sup>15</sup> and the Lieb theorem.<sup>16</sup> However, for more general situations, one should rely on robust, accurate, and cost-effective theoretical methods, which is still a difficult task not exempted from computational limitations, especially for large systems.<sup>17</sup>

On the other hand, the application of standard Density Functional Theory (DFT) methods to these (poly)radical systems is known to be affected by some pitfalls and/or artifacts: the intrinsic one-determinantal nature of Kohn-Sham (KS) DFT precludes to deal with orbital degeneracies, thus neglecting non-dynamical or static correlation effects, and the use of a Broken-

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8 Symmetry (BS) solution for open-shell systems introduces spin-contamination  
9 (also scaling with size<sup>18</sup>) issues mostly affecting the energy of the low-spin  
10 solution.<sup>19</sup> This situation has historically prompted the development of non-  
11 standard methods able to cope with these subtle electronic effects, namely  
12 based on the two-body on-top pair density<sup>20-29</sup> with a revisited interest  
13 nowadays,<sup>30-33</sup> the balanced coupling of *ab initio* and density functional  
14 expressions,<sup>34-37</sup> the use of natural orbitals<sup>38-40</sup> or the specific *ensemble*  
15 of pure spin states,<sup>41-43</sup> to name just a few of the existing non-standard  
16 methods. Another possible route is the use of fractional spin<sup>44</sup> or orbital  
17 occupation,<sup>45,46</sup> mimicking the situation when multiconfigurational *ab ini-*  
18 *tio* methods are instead employed, or spin-flip techniques,<sup>47-49</sup> describing  
19 target states from a high-spin reference state.  
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31 To further explore (*vide infra*) the applicability and accuracy of modern  
32 DFT variants, in the search for the best trade-off between accuracy and  
33 computational cost, we have chosen a set of large (and real-world) organic  
34 radical compounds recently synthesized and crystallized with diverse struc-  
35 tural motifs (see Figure 1). Note that for all of the systems selected, their  
36 stability has allowed the original authors to perform experimental measure-  
37 ments such as Electron Spin Resonance (ESR) or Superconducting QUan-  
38 tum Interference Device (SQUID), among others, to extract e.g., the energy  
39 difference between low- and high-spin states, thus allowing to bracket the  
40 accuracy of the theoretical methods employed after the comparison with ex-  
41 perimental results. The systems selected here (and their short names used  
42 in the following) are: (i) substituted Blatter-like radicals<sup>50,51</sup> (Diradical  
43 I and II); (ii) [6]cyclo-para-phenylmethine<sup>52</sup> (6CPPM-Mes); (iii) [n]cyclo-  
44 para-biphenylmethines<sup>53</sup> ([n]CPBM-Ant) with  $n = 3 - 6$ ; (iv) ethynylene-  
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bridged fluorenyl macrocycle<sup>54</sup> (MC-F3A3); and (v) cyclopenta-ring-fused oligo(m-phenylene) macrocyclic<sup>55</sup> (8MC). Note that the DFT-based results will also be compared with those from the Restricted-Active-Space Spin-Flip (RAS-SF) method,<sup>56</sup> to bracket their accuracy also for magnitudes for which experimental results are not available.

## 2 Theoretical Methods

### 2.1 The FT-DFT method

The Finite-Temperature DFT (FT-DFT) method relies on the fractional occupation of molecular orbitals induced by (near-)degeneracy effects, with the associated density written as:

$$\rho(\mathbf{r}) = \sum_i^{\infty} f_i |\phi_i(\mathbf{r})|^2, \quad (1)$$

where  $\phi_i(\mathbf{r})$  is a molecular spin-orbital and  $f_i$  its fractional electron occupation numbers ( $0 \leq f_i \leq 1$ ). The self-consistency of the procedure is achieved by minimizing the Gibbs electronic free energy ( $G_{el} = E_{el} - T_{el}S_{el}$ ) of the system at a fictitious pseudo-temperature (i.e., electronic) called  $T_{el}$ , with the  $f_i$  values obtained by a Fermi-Dirac distribution around the Fermi level  $E_F$ :

$$f_i = \frac{1}{1 + e^{(\epsilon_i - E_F)/\theta}}, \quad (2)$$

depending on  $\theta = k_B T_{el}$ . The corresponding energy difference between the low-spin (LS) and high-spin (HS) solutions can be calculated after imposing the desired spin multiplicity,  $\Delta E(\text{LS} - \text{HS}) = E(\text{LS}) - E(\text{HS})$ , with  $\Delta E(\text{LS} - \text{HS}) < 0$  indicating a favoured low-spin ground-state (antiferromagnetic). Note the similarities between this method and the Thermally-Assisted-Occupation (TAO) DFT method of Chai *et al.*<sup>46,57</sup>



### 2.1.1 Characterization of the radical character

Additionally, the set  $\{f_i, \phi_i(\mathbf{r})\}$  can be used to define a density of unpaired electrons, that is the Fractional Orbital Density (FOD)<sup>58,59</sup> as:

$$\rho^{FOD}(\mathbf{r}) = \sum_i^M (\delta_1 - \delta_2 f_i) |\phi_i(\mathbf{r})|^2, \quad (3)$$

where  $\delta_1$  and  $\delta_2$  are chosen to become (1, 1) if the single-particle energy level ( $\epsilon_i$ ) associated with the orbital  $\phi_i$  is lower than the energy of the Fermi level,  $E_F$ , or (0, -1) otherwise. This density also leads upon integration to a measure of the number of strongly correlated electrons,  $N_{FOD} = \int \rho^{FOD}(\mathbf{r}) d\mathbf{r}$ , which is a concept equivalent to the (linear) metrics introduced by Head-Gordon,<sup>60</sup> typically labelled as  $N_U$  and obtained from natural orbital occupation numbers (NOONs), i.e., the eigenvalues of the one-electron reduced density matrix.

Complementarily, the radical character of electronic states can be quantified by means of the radical indices  $0 \leq y_i \leq 1$  ( $i = 0, 1, 2, 3$ ). Within the FT-DFT methodology, they can be directly assigned to the electronic occupations above the Fermi level as  $y_i = f_{LUMO+i}^\alpha + f_{LUMO+i}^\beta$ , with  $f_{LUMO+i}^\sigma$  the fractional occupation number of the lowest unoccupied LUMO+ $i$  spin-orbital (since approximately  $f_{LUMO+i}^\sigma + f_{HOMO-i}^\sigma = 1$ ). For systems with a significant (poly)radical nature, the indices  $y_i$  can be used to estimate their di- or tri- ( $y_0$ ), tetra- or penta- ( $y_1$ ), hexa- or hepta- ( $y_2$ ), and octa- or nonradical ( $y_3$ ) character, respectively. Large indices ( $y_i \approx 1$ ) indicate high radical character, while intermediate values are indicative of moderate (poly)radicaloid character. The similarity of these fractional occupation numbers with the NOONs has been recently confirmed for polycyclic aromatic hydrocarbons,<sup>61</sup> as well as the trend between  $N_{FOD}$  and global

biradical values arising from experimental measurements.<sup>62,63</sup>

## 2.2 The SF-DFT method

The Spin-Flip DFT (SF-DFT) method relies on the exchange of the  $\alpha$  and  $\beta$  spin blocks of the density on certain user-defined centers, thus generating a Broken-Symmetry (BS) solution after converging the high-spin wavefunction. The energy difference between both considered configurations is given by  $\Delta E(\text{BS} - \text{HS}) = E(\text{BS}) - E(\text{HS})$ , which can be used as a first approximation to the energy difference between LS and HS solutions. Energy differences between pure spin states can be estimated by the Yamaguchi's procedure:<sup>64-66</sup>

$$\Delta E(\text{LS} - \text{HS}) = \frac{n_S}{\langle \hat{S}^2 \rangle_{\text{HS}} - \langle \hat{S}^2 \rangle_{\text{BS}}} \Delta E(\text{BS} - \text{HS}), \quad (4)$$

where  $n_S$  corresponds to the  $\langle \hat{S}^2 \rangle$  difference between the ideal spin multiplicities, e.g.,  $n_S = 2$  for a LS singlet and a HS triplet,  $n_S = 3$  for a LS doublet and a HS quartet, etc.

## 2.3 The SF-TDDFT method

The Spin-Flip Time-Dependent DFT (SF-TDDFT) method is recognized to uniformly describe excited states of single, double, and mixed excitation character in molecular systems,<sup>67</sup> and more specifically in conjugated molecules featuring diradical or (poly)radical character,<sup>68,69</sup> starting from a high-spin (e.g., triplet) reference state. The formalism is based on the (linear response) TDDFT equations for excitation energies:

$$\begin{bmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{B}^* & \mathbf{A}^* \end{bmatrix} \begin{bmatrix} \mathbf{X} \\ \mathbf{Y} \end{bmatrix} = \Omega \begin{bmatrix} \mathbf{1} & \mathbf{0} \\ \mathbf{0} & -\mathbf{1} \end{bmatrix} \begin{bmatrix} \mathbf{X} \\ \mathbf{Y} \end{bmatrix}, \quad (5)$$

with  $\mathbf{X}$  ( $\mathbf{Y}$ ) the set of (de-)excitation amplitudes,  $\mathbf{A}$  and  $\mathbf{B}$  the linear-response matrices, and  $\Omega$  the excitation energies. In SF-TDDFT, Eq. (5) is solved for the subspace of spin-flip ( $\alpha \rightarrow \beta$ ) operators.<sup>47</sup> For this case, the general expression for  $\mathbf{A}$  and  $\mathbf{B}$  elements simplify to  $A_{i\bar{a},j\bar{b}} = \delta_{ij}\delta_{\bar{a},\bar{b}}(\epsilon_{\bar{a}} - \epsilon_i) - C_x(ij|\bar{a}\bar{b})$  and  $B_{i\bar{a},j\bar{b}} = -C_x(ib|\bar{a}\bar{j})$ , with  $C_x$  the weight of exact exchange of the density functional used,  $i, j$  ( $a, b$ ) refer to occupied (virtual) orbitals (the overbar on orbital indices indicates  $\beta$ -spin),  $\epsilon_p$  is the energy associated to the KS  $p$ -spin-orbital, and  $(pq|st)$  is the two electron interaction integral given in Mülliken's notation:

$$(pq|st) = \int \phi_p^*(\mathbf{x})\phi_q(\mathbf{x})\frac{1}{|\mathbf{x} - \mathbf{x}'|}\phi_t^*(\mathbf{x}')\phi_s(\mathbf{x}')d\mathbf{x}d\mathbf{x}' \quad (6)$$

## 2.4 The RAS-SF method

Spin-flip methods with single spin-flip excitations<sup>70–72</sup> are not capable to properly describe low-spin states of molecular systems with more than three unpaired electrons, e.g., tetraradicals. This limitation can be overcome through the generalization of the excitation operator to multiple spin-flip excitations, as in the RAS-SF method. In RAS-SF the orbital space of the high-spin reference is split in three subspaces: doubly occupied (RAS1), singly occupied (RAS2), and virtual (RAS3). The eigenstates of the RAS-SF Hamiltonian are obtained as  $n$ -spin-flip excitations expanded in terms of the number of holes (electrons) in the doubly (virtual) spaces:

$$\hat{R}^{n\text{SF}} = \hat{r}_0^{n\text{SF}} + \hat{r}_h^{n\text{SF}} + \hat{r}_p^{n\text{SF}} + \hat{r}_{hp}^{n\text{SF}} + \hat{r}_{2h}^{n\text{SF}} + \hat{r}_{2p}^{n\text{SF}} + \dots, \quad (7)$$

where  $\hat{r}_0^{n\text{SF}}$  performs all possible spin-flip excitations within RAS2 and  $h$  and  $p$  subindices indicate the number of holes and electrons in RAS1 and

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8 RAS3, respectively.  
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## 10 11 12 **2.5 Computational details** 13

14 We use a semi-local (TPSS<sup>73</sup>) and a pair of hybrid<sup>74</sup> (TPSS0 and TPSSHH)  
15 exchange-correlation meta-GGA functionals differing in the weight ( $C_x$ ) of  
16 the EXact-eXchange (EXX) introduced (i.e., 0% for TPSS, 25% for TPSS0,  
17 and 50% for TPSSHH) for the FT-DFT and SF-DFT calculations reported  
18 here. Note that the original FT-DFT method employed the TPSS functional,  
19 which will be respected here, but we will also complementarily explored the  
20 dependence of the results with respect to the EXX weight. The electronic  
21 temperature ( $T_{el}$ ) was fixed for the FT-DFT calculations following the rec-  
22 ommended expression  $T_{el}/K = 5000 + 20000 C_x$  as a function of the EXX  
23 weight  $C_x$ .  
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33 We use the cost-effective 6-31G\*\* (SF-DFT) and the large def2-TZVP<sup>75</sup>  
34 (FT-TPSS) basis sets for those calculations, together with the RIJCOSX  
35 technique<sup>76</sup> (with the def2/JK auxiliary basis sets<sup>77</sup>) to reduce the increase  
36 in computational cost associated to the TPSS0 and TPSSHH functionals.  
37 The plots of the  $\rho^{\text{FOD}}(\mathbf{r})$  density were generated by the UCSF Chimera<sup>78</sup>  
38 (version 1.12) package. The FT-DFT and SF-DFT calculations were done  
39 with the ORCA (version 4.0.1.2) quantum-chemical package<sup>79</sup> employing  
40 ultrafine numerical integration grids (i.e., Grid6, NoFinalGrid) in all cases.  
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50 The SF-TDDFT calculations employed the collinear approximation as  
51 implemented in the GAMESS package,<sup>80</sup> together with the BHHLYP func-  
52 tional<sup>81</sup> and the cost-effective 6-31G\* basis set. Note that the use of a  
53 functional with a high  $C_x = 0.50$  value is recommended for this kind of  
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8 calculations,<sup>47,82</sup> and that the accuracy is not expected to vary using an-  
9 other exchange-correlation functional like TPSSHH of PBEHH (both with  
10 the same  $C_x = 0.50$  value than BHHLYP).<sup>83</sup>  
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15 The RAS-SF calculations have been done within the hole and electron  
16 approximation, that is including the three first terms in the rhs of Eq. (5)  
17 using a ROHF (Restricted Open-Shell) high-spin reference: triplet (Blatter-  
18 like diradicals), quartet ([3]CPBM), quintet ([4]CPBM), sextet ([5]CPBM),  
19 septet ([6]CPPM and [6]CPBM), and nonet (8MC). Further details can be  
20 found at the Supporting Information and elsewhere.<sup>52-55</sup> These calculations  
21 have been done with the Q-Chem (version 5.2) program<sup>84</sup> and the 6-31G\*\*  
22 basis set.  
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31 Finally, the Zero-Field-Splitting (ZFS) calculations were performed with  
32 the  $\omega$ B97X-D functional<sup>85</sup> and the IGLO-II basis set,<sup>86</sup> intended for com-  
33 puting magnetic properties with high accuracy,<sup>87</sup> together with the 'Au-  
34 toAux' generation procedure for auxiliary basis sets.<sup>88</sup> The ZFS tensor was  
35 self-consistently calculated on the basis on spin-Unrestricted Natural Or-  
36 bitals (UNO)<sup>89</sup> as recommended.<sup>90</sup> The ZFS calculations were done with  
37 the ORCA (version 4.0.1.2) quantum-chemical package<sup>79</sup> employing a tight  
38 threshold for convergence (i.e., TightSCF) and ultrafine numerical integra-  
39 tion grids (i.e., Grid6, NoFinalGrid) in all cases.  
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### 3 Results and discussion

#### 3.1 Quantification of the (poly)radicaloid character.

First, we aim to evaluate the extent of the radical character, i.e., the number of unpaired electrons ( $N_{FOD}$ ), of the considered molecular species by means of FT-DFT calculations. The  $N_{FOD}$  values obtained by FT-TPSS, FT-TPSS0, and FT-TPSSHH methods are presented in Table 1 for both the low- and high-spin states of all studied compounds. Complementarily, Figure 2 compares the calculated  $N_{FOD}$  values for the low-spin state of all compounds, from which we can recognize a close agreement between RAS-SF and FT-TPSS results. The results discussed along this section will thus limit to those obtained at the FT-TPSS level, with FT-TPSS0 slightly (FT-TPSSHH largely) overestimating the RAS-SF results. Moreover, all molecules present significant  $N_{FOD}$  values, indicating their open-shell (poly)radical character. The radical character is also preserved for the high-spin states, i.e. qualitatively similar  $N_{FOD}(\text{LS})$  and  $N_{FOD}(\text{HS})$  values are found except for Diradical I and II systems. Because fractional occupation is induced by near degeneracy, the smaller values of  $N_{FOD}$  for the HS state of Diradical I and II can be rationalized by their HOMO-1/HOMO and LUMO/LUMO+1 gaps, considerably larger than those of the other systems investigated. Interestingly, the series of [n]CPBM-Ant ( $n = 3 - 6$ ) compounds is predicted to increase their radical character as a function of their increasing size, in perfect agreement with experimental and RAS-SF results.<sup>53</sup>

Following the agreement found between  $N_{FOD}$  values at the FT-TPSS and RAS-SF levels, see also Table S1, we represent in Figure 3 the topology (real-space distribution) of the corresponding density,  $\rho^{FOD}(\mathbf{r})$ , at the FT-TPSS level and using the recommended threshold<sup>58,59</sup> for the isocon-

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8 tour values ( $\sigma = 0.005 \text{ e/bohr}^3$ ). For the Blatter-like radicals, the FOD  
9 density concentrates on the N atoms of the conjugated backbone, and on  
10 the nitrosyl substituents, in agreement with what one would expect from the  
11 resonance Lewis structures of the molecules. For 6CPPM-Mes, [n]CPBM-  
12 Ant ( $n = 3 - 6$ ), and MC-F3A3 compounds we can observe how the FOD  
13 density locates mainly at those C atoms bringing the mesityl and anthracene  
14 substituents, respectively, acting effectively as protective synthons. For the  
15 8MC compound we observe a delocalization of the FOD density on the non-  
16 bridging C atoms, resembling the results found for other systems with cyclic  
17 topologies as cyclacenes (i.e. cyclic oligoacenes<sup>91</sup>).  
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### 28 **3.1.1 Radical(oid) indices.**

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30 In order to get a deeper insight into the radical nature of these com-  
31 pounds, in the following we explore them by means of their  $\{y_i\}$  indices.  
32 Table 2 presents the  $y_0$ ,  $y_1$ ,  $y_2$ , and  $y_3$  values for all the systems studied at  
33 the FT-TPSS level. The Blatter-like diradicals exhibit nearly ideal diradical  
34 character, with  $y_0 \simeq 1.0$  and  $y_{i>0} \simeq 0$  for the low and high-spin  $T_0$  states.  
35 Note that this is in agreement with the smaller  $N_{FOD}$  values discussed in  
36 the previous section for the HS state of these two systems. The 6CPPM-Mes  
37 molecule holds a sizable tetraradicaloid character, with moderate  $y_0$  and  $y_1$   
38 values for the ground-state singlet. For the [n]CPBM-Ant ( $n = 3 - 6$ ) sys-  
39 tems, we observe an increase of the number of strongly correlated electrons  
40 as a function of their size, in agreement with the trend found for the  $N_{FOD}$   
41 values. Inspection of their  $y_i$  values allows to classify them as tri-, tetra-  
42 , penta- and hexaradicaloid molecules, respectively. Finally, for the 8MC  
43 molecule we obtain moderate values for all the  $y_{0-3}$  indices, indicating a  
44 moderate octaradicaloid behaviour. Tables S2-S3 present the  $y_i$  values ob-  
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8 tained at the FT-TPSS0 and FT-TPSSHH levels, respectively, which follow  
9 the same trend found for FT-TPSS, but with  $\{y_i\}$  indices being systemati-  
10 cally larger, like for  $N_{FOD}$  values.  
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### 15 **3.2 Energy difference between low- and high-spin states.**

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18 Table 3 presents the energy difference  $\Delta E(\text{LS} - \text{HS})$  between the low-  
19 and high-spin states of all the systems considered, calculated at the FT-  
20 TPSS, FT-TPSS0, and FT-TPSSHH levels. Note that, except for the Blatter-  
21 like diradicals considered, the electronic ground-state of these systems is al-  
22 ways the one with the lowest spin-multiplicity, thus denoted as  $S_0$  (singlet) or  
23  $D_0$  (doublet). Therefore,  $\Delta E(\text{LS} - \text{HS})$  refers to the  $S_0$ - $T_1$  or  $D_0$ - $Q_1$  energy  
24 difference, respectively, and will hold a negative sign:  $\Delta E(\text{LS} - \text{HS}) < 0$ .  
25 For the Blatter-like diradicals, the triplet electronic ground-state is instead  
26 favoured, and in these cases it should be  $\Delta E(\text{LS} - \text{HS}) > 0$  accordingly.  
27 First of all, inspecting the evolution of values in Table 3, we can see how the  
28 relative stabilization of the HS state with respect to the LS solution increases  
29 with the amount of Hartree-Fock exchange, i.e., upon going from FT-TPSS  
30 to FT-TPSS0, and to FT-TPSSHH, with the latter being systematically  
31 closer to experimental energy gaps. However, the agreement with experi-  
32 mental results largely differs among the set of compounds, even looking at  
33 the FT-TPSSHH results (i.e. best estimates) providing the lowest MSE and  
34 MUE values. For [6]CPBM or 8MC, employing any of the FT-DFT variants  
35 will lead to an error close or even less than 1 kcal/mol, commonly known  
36 as the chemical accuracy threshold. On the other hand, for [3]CPBM and  
37 MC-F3A3 compounds the computed gaps are a few kcal/mol too negative,  
38 even with the FT-TPSSH method. Figure 4 compares the FT-DFT cal-  
39 culated values with the experimental results, for which we can also easily  
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8 observe a different behaviour for odd and even  $[n]$ CPBM compounds. These  
9 facts, together with the spread of the results for MC-F3A3, allow us to con-  
10 clude that the FT-DFT (with the default electronic temperatures) tends to  
11 overestimate the relative stability of low-spin (singlet or doublet) state with  
12 respect to the next higher spin state (triplet or quartet). Energy differences  
13 can be systematically improved, to some extent, by increasing the amount  
14 of exact exchange.  
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22 We compare next the SF-DFT and the experimental results in Table  
23 4, also using the TPSS, TPSS0, and TPSSHH functionals to disclose the  
24 effect of linearly increasing the exact-exchange weight. First of all, we con-  
25 sider the FOD density as the criteria to select those atoms to flip, with the  
26 highest density localized on them, which could also be roughly estimated by  
27 inspecting the corresponding spin density. In this case, spin contamination  
28 becomes a key factor and results progressively deteriorates upon increas-  
29 ing the exact-exchange weight, contrarily to what happened with FT-DFT  
30 methods. The (spin-corrected) energy gaps  $\Delta E(\text{LS} - \text{HS})$  keep an accuracy  
31 similar to that obtained for the uncorrected  $\Delta E(\text{BS} - \text{HS})$  values, still with  
32 the SF-TPSS or SF-TPSS0 methods providing the closest agreement with  
33 experimental results (e.g., MUEs of 6.0 and 4.0–5.0 kcal/mol, respectively).  
34 Remarkably, the SF-TPSS method provides the correct lowest-energy spin-  
35 state for all the molecules considered, contrarily to SF-TPSS0 and especially  
36 SF-TPSSHH. Inspecting now the SF-TDDFT results in Table 5, done with  
37 the BHHLYP and thus comparable with TPSSHH in terms of having a sim-  
38 ilar exact exchange proportion, we observe larger averaged errors than for  
39 previous FT-DFT or SF-DFT methods, with a reverse state ordering for  
40 Diradicals I and II. The method yields too large energy differences for the  
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set of  $[n]$ CPBM compounds, but it keeps the correct trend of decreasing the  $\Delta E$  values with increasing size.

For the sake of comparison of all these results, RAS-SF gives a MSE (MUE) of  $-0.55$  ( $0.85$ ) kcal/mol with respect to experimental results, with a maximum deviation of  $3.2$  kcal/mol and producing thus more accurate relative energies than the investigated DFT-based methods. Actually, this method is able to provide the chemical accuracy sought for the whole set of compounds. Discarding the case of  $[6]$ CPPM-Mes, the MSE (MUE) would decrease to  $-0.22$  ( $0.56$ ) kcal/mol, and thus being considerably low.

### 3.3 Zero-field splitting interactions

The magnetic dipole-dipole (i.e., spin-spin) interaction leads to the splitting of the triplet sublevels ( $M_s = 0, \pm 1$ ) even in the absence of any external field; a physical effect described by the Zero-Field Splitting (ZFS) Hamiltonian:

$$\hat{H}_{ZFS} = \hat{\mathbf{S}} \cdot \hat{\mathbf{D}} \cdot \hat{\mathbf{S}} = D_{xx} \hat{S}_x^2 + D_{yy} \hat{S}_y^2 + D_{zz} \hat{S}_z^2 = D \left( \hat{S}_z^2 - \frac{1}{3} \hat{\mathbf{S}}^2 \right) + E \left( \hat{S}_x^2 - \hat{S}_y^2 \right), \quad (8)$$

with  $D_{ii}$  the principal values of the ZFS diagonal tensor  $\hat{\mathbf{D}}$ , which by convention are recasted as  $D = D_{zz} - \frac{1}{2} (D_{xx} + D_{yy})$  and  $E = \frac{1}{2} (D_{xx} - D_{yy})$ . For systems having  $S > 1/2$ , the ZFS usually dominates the spectral shape of the Electron Paramagnetic Resonance (EPR) spectra, and thus the absolute values of  $D$  and the  $E/D$  ratio determine the energies of the three magnetic sublevels.<sup>92</sup> Additionally, provided that the point-dipole approximation holds,  $D$  also relates with the averaged distance ( $\Delta r$ ) between ideally localized spin densities, and can be thus used to estimate the size of the photoexcited triplet exciton<sup>93</sup> (see the Supporting Information for further

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8 details). Note that in the following we will restrict the study to those sys-  
9 tems possessing ground-state or low-lying triplet states, i.e., with an even  
10 number of electrons.  
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15 First of all, we have thoroughly assessed the accuracy of DFT methods  
16 to calculate the  $D$  and  $E$  parameters for the pair of systems (Diradicals  
17 I and II) for which experimental measurements are available.<sup>51</sup> For both  
18 compounds it is clear that  $D/hc < 0$  from the experimental EPR spec-  
19 tra, thus indicating a prolate-like distribution of the spin density for the  
20 triplet state. The sign of  $D$  indicates whether the  $M_s = 0$  ( $D > 0$ ) or  
21  $M_s = \pm 1$  ( $D < 0$ ) spin substrates are the lowest energy states at zero ex-  
22 ternal fields. However, previous results at the B3LYP/EPR-II level,<sup>51</sup> and  
23 with different exchange-correlation functionals and basis sets (see Tables S4-  
24 S5), predicted the wrong sign for Diradical II ( $D/hc > 0$ ), which is properly  
25 characterized only by certain range-separated functionals (i.e.,  $\omega$ B97X-D<sup>85</sup>  
26 and LC-BLYP<sup>94</sup>) together with basis sets suited for electric and magnetic  
27 properties, i.e. IGLO-II or EPR-II, previously applied<sup>95</sup> to the study of  
28 spin-spin contributions to the ZFS tensor in organic radicals too.  
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42 This deficiency of DFT methods has also been documented before<sup>96,97</sup>  
43 and prompted us to apply in the following the  $\omega$ B97X-D functional for con-  
44 sistency. Note also that the range-separated CAM-B3LYP functional<sup>99</sup> was  
45 also used but did not bring the correct sign of  $D$  for Diradical II. The main  
46 difference between the  $\omega$ B97X-D/LC-BLYP and CAM-B3LYP schemes is a  
47 relatively large (35%) fixed DFT exchange contribution in the latter, and  
48 thus a maximum screened exact exchange of 65%, which seems to corrob-  
49 orate the importance of that variable part (80% and 100% for  $\omega$ B97X-D  
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8 and LC-BLYP, respectively). On the other hand, looking again at Tables  
9 S4-S5, the relative error for the calculation of  $E$  was found larger than for  
10  $D$ , in agreement with previous applications to heavy-atoms coordination  
11 complexes.<sup>98</sup>  
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17 Table 6 presents the  $D$ ,  $E$ , and  $\Delta r$  calculated values (at the  $\omega$ B97X-  
18 D/IGLO-II level) for the lowest triplet state of the set of compounds studied.  
19 Interestingly, for the Diradicals I and II, a high-spin ground-state together  
20 with a negative  $D$  could lead to effective molecular magnets.<sup>13</sup> In the case  
21 of 8MC, we can see how  $E = 0$  due to the perfect axial symmetry of this  
22 compound, with  $E$  being considerably lower than  $D$  as expected in all other  
23 cases. Inspecting the  $\Delta r$  values, i.e., the mean inter-spin distance in a dipole-  
24 dipole approximation, see the Supporting Information for further details  
25 about the explicit derivation, we can see how it decreases with the system  
26 size; a fact also documented before for linear polyenes and polyacenes.<sup>100</sup>  
27 This is rationalized by the dependence  $D \propto r^{-3}$  with  $r$  the distance between  
28 the spins of the unpaired electrons. We can also compare these results with  
29 the estimated exciton size ( $\Delta r$ ) for the triplet ground-state of 2,6,10-Tri-  
30 *tert*-Butyltriangulene,<sup>101</sup> around 5.6 Å, or for the photoexcited triplet state  
31 of tetracene and pentacene,<sup>102</sup> around 3.8 Å, or for the photoexcited triplet  
32 state of B- and N-doped nanographenes,<sup>103</sup> around 4.4-5.2 Å depending on  
33 their size.  
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## 50 4 Conclusions

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53 We report here a benchmark study of a set of real-world (poly)radicaloids,  
54 focusing on the extent of the radical character, spatial distribution of the  
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8 unpaired electrons, and singlet-triplet (or doublet-quartet) energy difference  
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10 obtained with different electronic structure methods. Current research on  
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12 organic (poly)radicaloid character and its applications has prompted the ap-  
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14 plication here of both (cost-effective) DFT-based and RAS-SF methods, with  
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16 the latter method behaving more accurately than the others as compared  
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18 with reference experimental results. Complementarily, we have systemati-  
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20 cally compared finite-temperature (FT-DFT) and spin-flip approaches (SF-  
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22 DFT and SF-TDDFT) with various exchange-correlation functionals, mostly  
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24 differing in their exact exchange weight, to disentangle the effect of the un-  
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26 derlying expression as well as the effect of the spin-contamination intro-  
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28 duced. Evidently, the use of any of these approaches with a GGA form (i.e.,  
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30 TPSS) is less costly than using a hybrid expression. Finally, we have also  
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32 calculated the ZFS parameters for the triplet states of the compounds, as  
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34 well as their exciton size. Overall, we have shown how the cost-effective char-  
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36 acterization of (poly)radicaloid nature in conjugated organic compounds is  
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38 still a challenging issue, precluding the blind application of DFT variants.

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## 48 Supplementary Material

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51 The Supplementary Material contains in this order: (i) the metrics er-  
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53 ror used to compare the performance of the different methods; (ii)  $N_U$  and  
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55  $\Delta E(\text{LS} - \text{HS})$  values obtained at the RAS-SF level for all the compounds;  
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(iii) calculated radical indices ( $y_i$ ) at the FT-TPSS0 and FT-TPSSHH levels for all the compounds; (iv) comparison between calculated and experimental EPR parameters for Diradicals I and II; (v) notes on the theoretical estimates of the exciton size and the sign of the D-tensor; (vi) cartesian coordinates of all the compounds.

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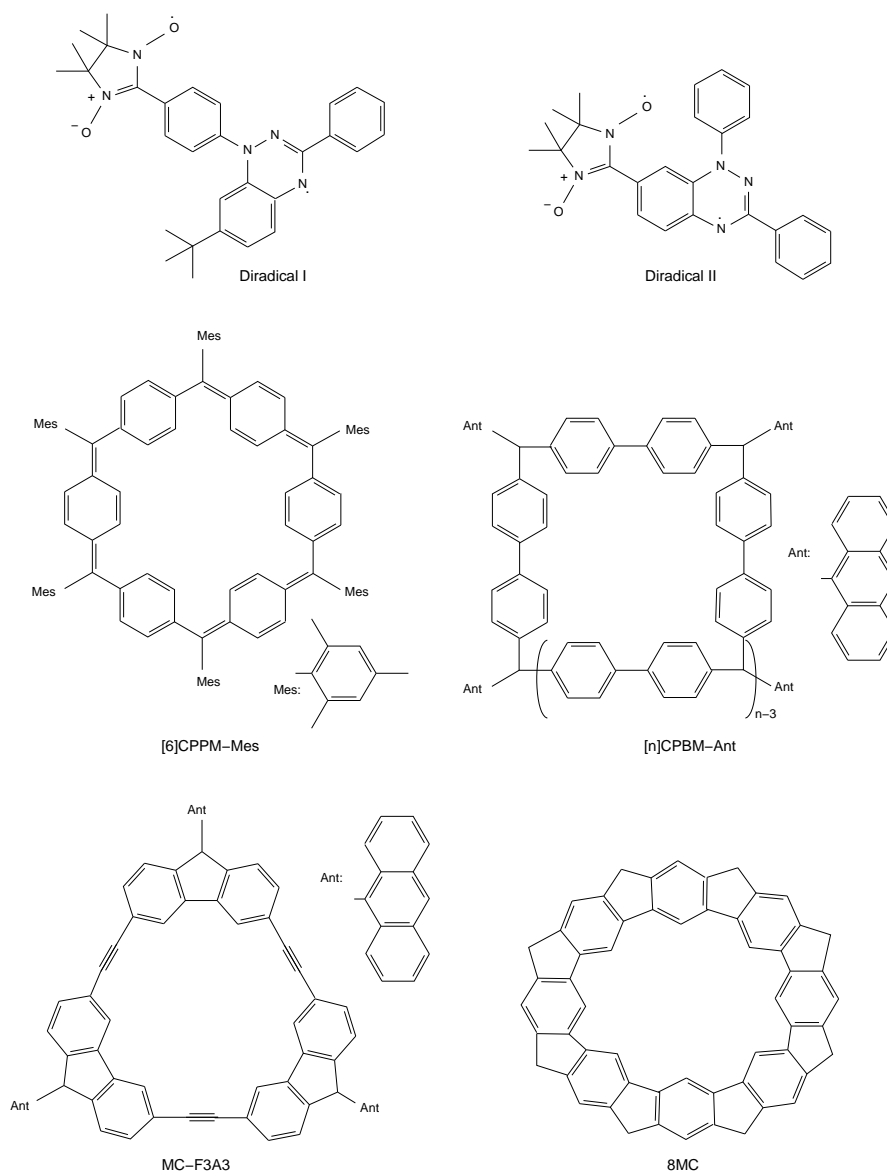


Figure 1: Chemical structures of the investigated compounds

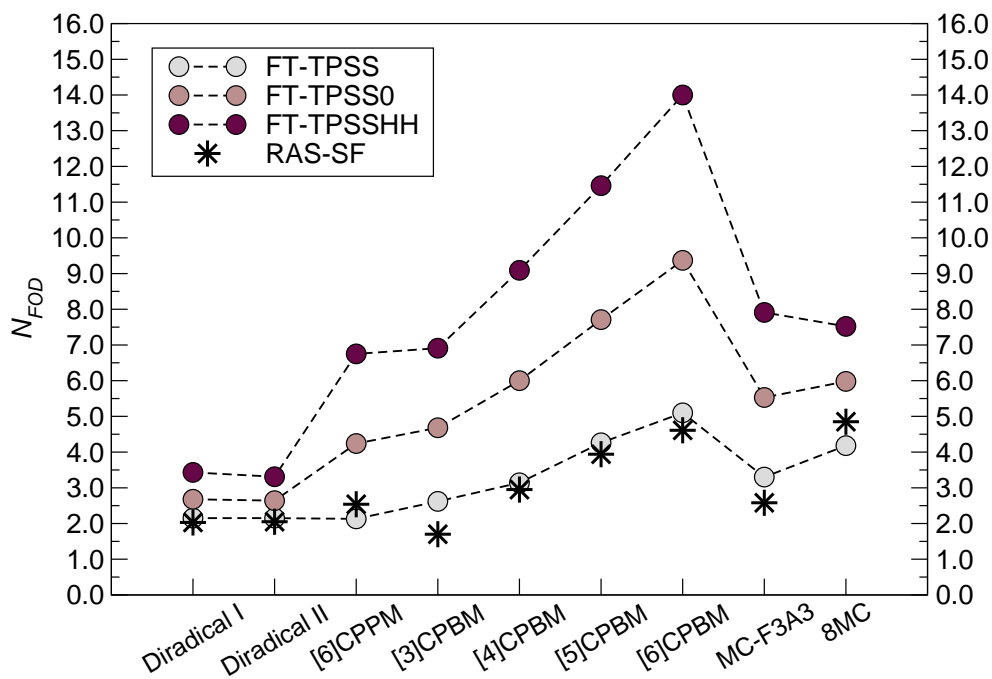


Figure 2: Comparison between FT-DFT and RAS-SF  $N_{FOD}$  values for the low-spin state of the set of studied compounds.

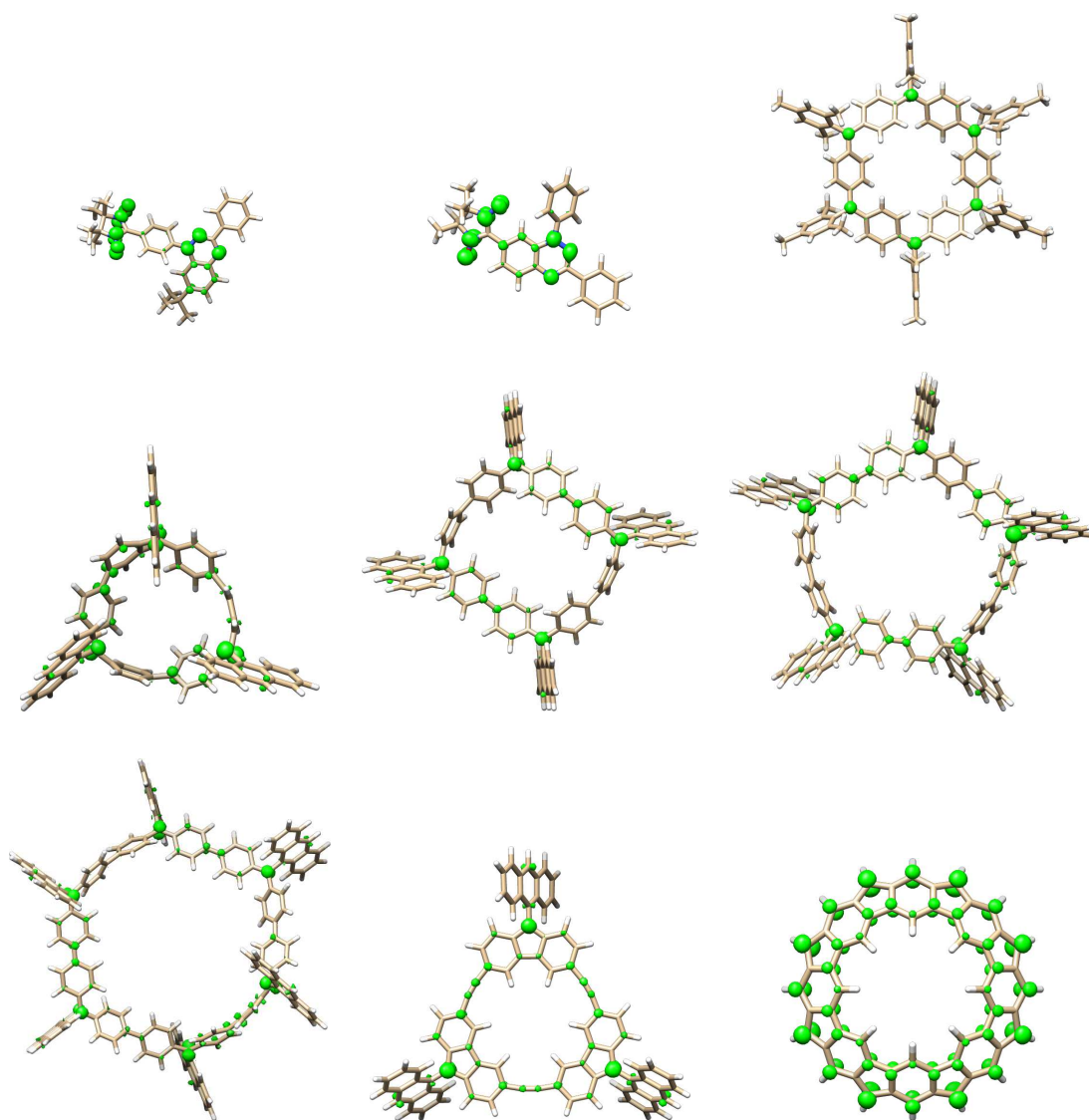


Figure 3: FOD density plots ( $\sigma = 0.005 \text{ e/bohr}^3$ ) obtained from the FT-TPSS/def2-TZVP method for the set of studied compounds.

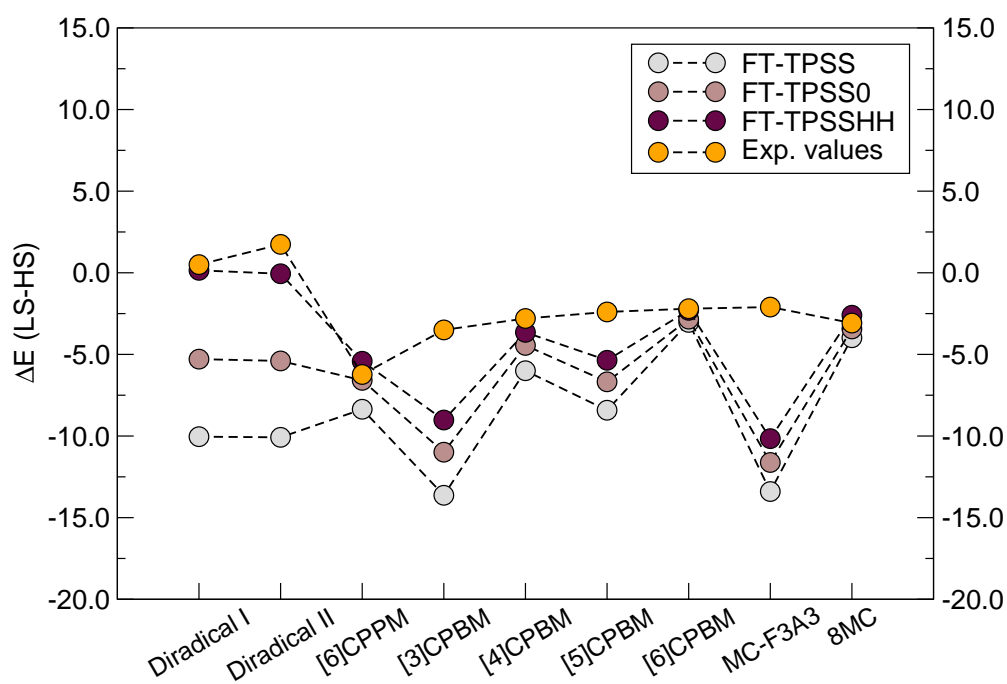


Figure 4: Comparison between  $\Delta E(\text{LS} - \text{HS})$  (kcal/mol) computed (FT-DFT) and experimental values for the set of studied compounds.



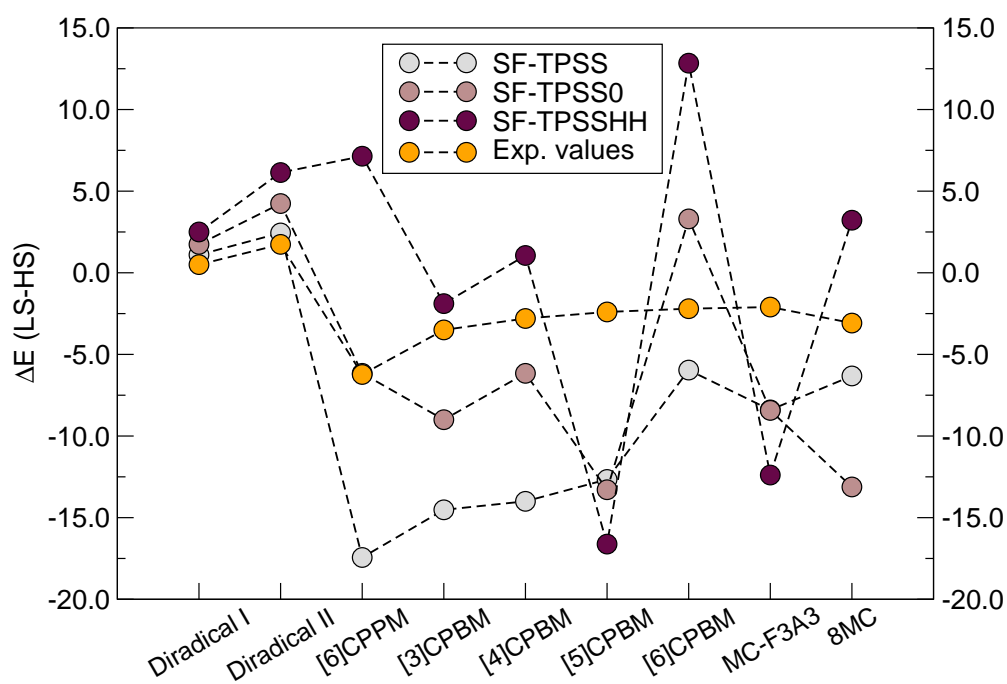


Figure 5: Comparison between  $\Delta E(\text{LS} - \text{HS})$  (kcal/mol) computed (SF-DFT) and experimental values for the set of studied compounds.

Table 1:  $N_{FOD}$  values obtained at different theoretical levels.

Compound	GS	FT-TPSS		FT-TPSS0		FT-TPSSH	
		$N_{FOD}(\text{LS})$	$N_{FOD}(\text{HS})$	$N_{FOD}(\text{LS})$	$N_{FOD}(\text{HS})$	$N_{FOD}(\text{LS})$	$N_{FOD}(\text{HS})$
Diradical I <sup>a</sup>	T <sub>0</sub>	2.15	0.81	2.68	1.44	3.43	2.23
Diradical II <sup>a</sup>	T <sub>0</sub>	2.15	0.79	2.64	1.42	3.31	2.10
[6]CPPM-Mes	S <sub>0</sub>	2.13	2.90	4.24	4.82	6.75	7.24
[3]CPBM-Ant	D <sub>0</sub>	2.62	2.78	4.68	4.78	6.91	6.99
[4]CPBM-Ant	S <sub>0</sub>	3.14	3.70	6.00	6.43	9.09	9.44
[5]CPBM-Ant	D <sub>0</sub>	4.26	4.61	7.71	7.95	11.46	11.66
[6]CPBM-Ant	S <sub>0</sub>	5.10	5.50	9.37	9.63	14.00	14.17
MC-F3A3	D <sub>0</sub>	3.30	3.13	5.53	5.37	7.91	7.73
8MC	S <sub>0</sub>	4.18	4.34	5.98	5.93	7.52	7.36

<sup>a</sup> Note that for these systems the header classification do not apply, since the ground-state is already the T<sub>0</sub> and thus the HS state.

Table 2: Calculated radical indices ( $y_i^\alpha$ ) by the FT-DFT method at the TPSS/def2-TZVP level.

Compound	$y_0^\alpha$	$y_1^\alpha$	$y_2^\alpha$	$y_3^\alpha$
Diradical I	0.49	0.03	0.02	0.00
Diradical II	0.49	0.03	0.00	0.00
[6]CPPM-Mes	0.24	0.22	0.04	0.00
[3]CPBM-Ant	0.35	0.08	0.08	0.08
[4]CPBM-Ant	0.28	0.21	0.06	0.06
[5]CPBM-Ant	0.31	0.29	0.08	0.08
[6]CPBM-Ant	0.34	0.28	0.21	0.06
MC-F3A3	0.46	0.07	0.07	0.06
8MC	0.30	0.30	0.21	0.21

Table 3: Energy difference (kcal/mol) between the low-spin (LS) and high-spin (HS) states,  $\Delta E(\text{LS} - \text{HS})$ , obtained at the FT-DFT level.

Compound	GS	FT-TPSS	FT-TPSS0	FT-TPSSHH	Exp.
Diradical I	T <sub>0</sub>	-10.04	-5.29	0.15	0.50±0.02
Diradical II	T <sub>0</sub>	-10.08	-5.40	-0.06	1.74±0.07
[6]CPPM-Mes	S <sub>0</sub>	-8.36	-6.50	-5.42	-6.23±0.78
[3]CPBM-Ant	D <sub>0</sub>	-13.63	-10.99	-9.02	-3.5
[4]CPBM-Ant	S <sub>0</sub>	-6.00	-4.45	-3.65	-2.8
[5]CPBM-Ant	D <sub>0</sub>	-8.42	-6.68	-5.36	-2.4
[6]CPBM-Ant	S <sub>0</sub>	-3.54	-2.84	-2.30	-2.2
MC-F3A3	D <sub>0</sub>	-13.40	-11.62	-10.16	-2.10
8MC	S <sub>0</sub>	-3.98	-3.44	-2.60	-3.08
	MSE	-6.4	-4.1	-2.0	
	MUE	6.4	4.1	2.3	
	MIN	0.9	0.4	0.4	
	MAX	11.8	9.5	8.1	

Table 4: Energy difference (kcal/mol) between the Broken-Symmetry (BS) and high-spin (HS) states,  $\Delta E(\text{BS} - \text{HS})$ , and the corresponding  $\Delta E(\text{LS} - \text{HS})$  corrected, obtained at the SF-DFT level.

Compound	GS	SF-TPSS		SF-TPSS0		SF-TPSSHH		Exp.
		$\Delta E(\text{BS} - \text{HS})$	$\Delta E(\text{LS} - \text{HS})$	$\Delta E(\text{BS} - \text{HS})$	$\Delta E(\text{LS} - \text{HS})$	$\Delta E(\text{BS} - \text{HS})$	$\Delta E(\text{LS} - \text{HS})$	
Diradical I	T <sub>0</sub>	0.55	1.10	0.88	1.74	1.32	2.50	0.50±0.02
Diradical II	T <sub>0</sub>	1.23	2.44	2.18	4.24	3.34	6.14	1.74±0.07
[6]CPPM-Mes	S <sub>0</sub>	-18.44	-17.44	-8.61	-6.16	12.94	7.14	-6.23±0.78
[3]CPBM-Ant	D <sub>0</sub>	-14.61	-14.52	-9.21	-9.00	-2.05	-1.89	-3.5
[4]CPBM-Ant	S <sub>0</sub>	-14.07	-14.00	-8.62	-6.16	1.27	1.06	-2.8
[5]CPBM-Ant	D <sub>0</sub>	-12.60	-12.66	-10.16	-13.29	-11.41	-16.62	-2.4
[6]CPBM-Ant	S <sub>0</sub>	-6.29	-5.96	3.48	3.30	15.85	12.84	-2.2
MC-F3A3	D <sub>0</sub>	-7.89	-8.40	-5.97	-8.43	-7.47	-12.39	-2.10
8MC	S <sub>0</sub>	-6.72	-6.32	-8.06	-13.12	1.93	3.22	-3.08
	MSE	-6.5	-6.2	-2.7	-3.0	4.0	2.4	
	MUE	6.5	6.5	4.1	5.0	7.2	7.9	
	MIN	0.0	0.6	0.4	0.1	0.8	1.6	
	MAX	12.2	11.2	7.8	10.9	19.2	15.0	

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Table 5: Energy difference (kcal/mol) between the low-spin (LS) and high-spin (HS) states,  $\Delta E(\text{LS} - \text{HS})$ , obtained at the SF-TDDFT level.

Compound	GS	SF-TDBHLYP	Exp.
Diradical I	T <sub>0</sub>	-9.34	0.50±0.02
Diradical II	T <sub>0</sub>	-9.94	1.74±0.07
[6]CPPM-Mes	S <sub>0</sub>	-3.37	-6.23±0.78
[3]CPBM-Ant	D <sub>0</sub>	-17.99	-3.5
[4]CPBM-Ant	S <sub>0</sub>	-17.80	-2.8
[5]CPBM-Ant	D <sub>0</sub>	-15.86	-2.4
[6]CPBM-Ant	S <sub>0</sub>	-9.34	-2.2
MC-F3A3	D <sub>0</sub>	-6.69	-2.10
8MC	S <sub>0</sub>	-0.57	-3.08
	MSE	-7.9	
	MUE	9.1	
	MIN	2.5	
	MAX	15.0	

Table 6: Calculated  $D$  and  $E$  EPR parameters ( $D/hc$  and  $E/hc$  in  $10^3\text{cm}^{-1}$ ) and exciton size ( $\Delta r$ , in  $\text{\AA}$ ) at the  $\omega\text{B97X-D/IGLO-II}$  level, of the lowest triplet state of the selected compounds.

Compound	$D$	$E$	$\Delta r$
Diradical I	-5.52	-0.80	7.8
Diradical II	-11.13	-3.66	6.2
[6]CPPM-Mes	-12.60	-0.56	5.9
[4]CPBM-Ant	-13.04	-1.74	5.8
[6]CPBM-Ant	-10.05	-0.04	6.4
8MC	4.29	0.00	8.5