

Supplementary Information for Characterization of Diarylethene-Based Photoswitches with Core and Valence Photoabsorption and Photoemission Spectroscopies

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Mass spectra

1

2 Only the mass spectrum of BTF6 was recorded: it is plotted in Figure S1, while Table 1 summarizes the tentative assignment of the most prominent peaks.

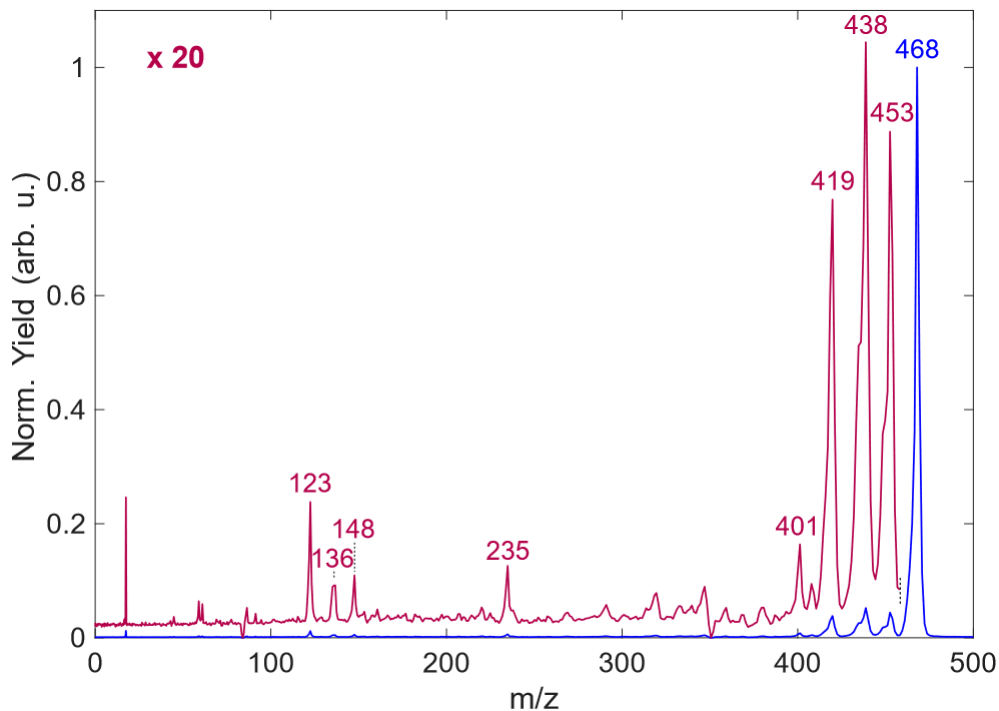


Figure S1: The recorded mass spectrum for BTF6 (blue), overlaid with its 20x magnification (red), where the parent ion ($m/z = 468$) is not shown to highlight the lighter fragments. The m/z of the most prominent peaks is reported above them.

3

BTF6 - Mass Spectrum

| m/z | Assignment |
|-------|----------------------|
| 468 | (parent ion) M^+ |
| 453 | $(M - CH_3)^+$ |
| 438 | $(M - 2CH_3)^+$ |
| 419 | $(M - 2CH_3 - F)^+$ |
| 401 | $(M - 2CH_3 - 2F)^+$ |
| 235 | (dication) M^{2+} |
| 148 | $C_8H_5SCH_3^+$ |
| 136 | $C_8H_5S^+$ |
| 123 | unidentified |

Table 1: Tentative assignment for the most prominent peaks of BTF6 mass spectrum.

4

Geometries

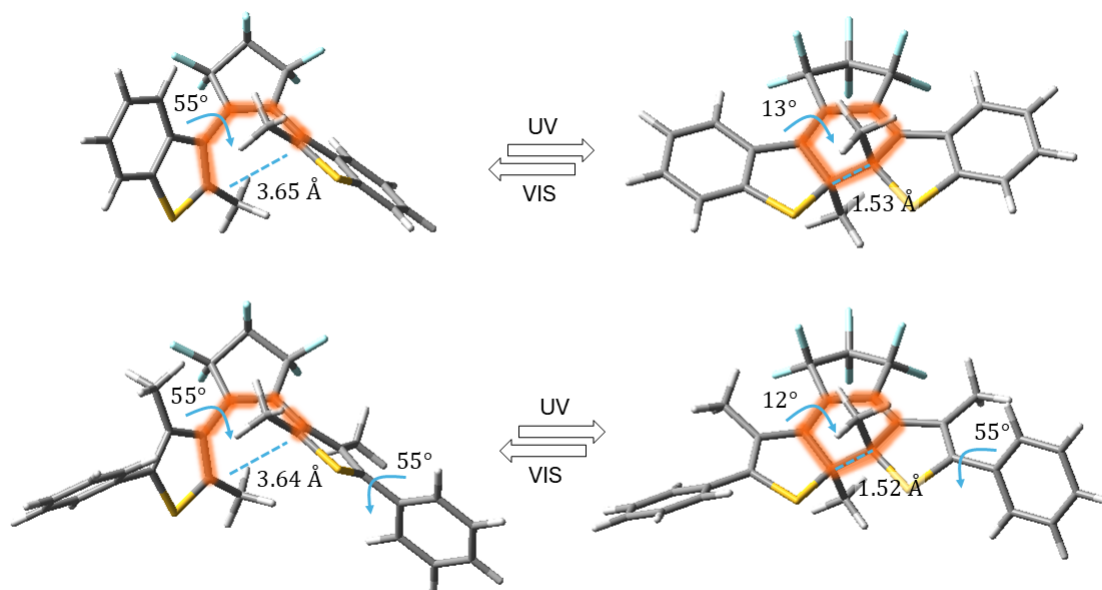


Figure S2: 3D structures of the target molecules (top) BTF6 (parallel conformation) and (bottom) PTF6 in their open-ring (OR) and closed-ring (CR) forms, highlighting hexatriene/cyclohexadiene units involved in the electrocyclic reaction as well as the values of several representative degrees of freedom.

5 Geometries

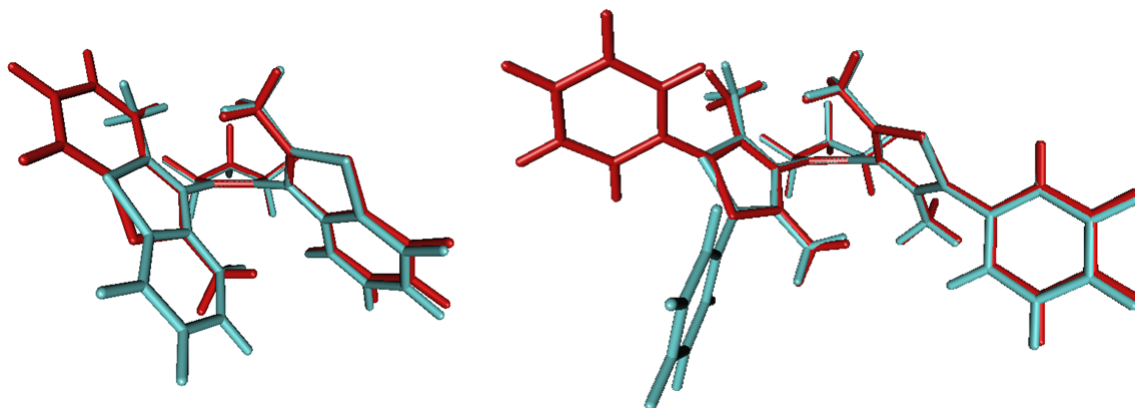


Figure S3: Comparison of the structures of the parallel (cyan) and anti-parallel (red) OR forms of BTF6 (left) and PTF6 (right).

6 Computational details

7 Active space definition for XPS and NEXAFS

8 Core-ionizations and -excitations are realized within the core-valence separation (CVS) approx-
 9 imations. By means of a projection technique all configuration state functions with a maximum

10 occupation of the core orbitals are eliminated from the configurational space. [1]. The AS in-
11 cludes the core orbital(s) of a single atom (1s orbital for carbons, a set of three 2p orbitals for
12 sulfurs) placed in the sub-space RAS1 and kept frozen. A Cholesky localization scheme is used to
13 localize core orbitals on single atomic sites [2]. The use of single atomic sites in the calculations
14 is a common strategy adopted with other levels of theory [3] which streamlines post-processing
15 and analysis. Whereas it neglects possible mixing of core-states involving different atoms, it is
16 justified by the generally excellent agreement between theory and experiment indicating that
17 the mixing is weak. The active space is completed by a set of valence occupied and virtual
18 orbitals - usually of π -nature - distributed between sub-spaces RAS2 and RAS3, respectively,
19 aimed at improving the description of electron correlation. In order to limit the computational
20 cost we allow for up to four intra-valence excitations.

21 For BTF6 we use RAS(15,1,4;1,7,7) and RAS(19,1,4;3,7,7) for XPS simulations at the carbon
22 K-edge and at the sulfur L_{2,3}-edge, respectively. In analogy, for PTF6 we use RAS(19,1,4;1,9,9)
23 and RAS(23,1,4;3,9,9). Thereby, the valence orbital sets are comprised entirely of π molecular
24 orbitals. Only a single core-ionized state was computed for each carbon atom whereas state-
25 averaging over three core-ionized states was performed for the sulfur. Tests with reduced valence
26 sub-spaces, containing only 10 valence orbitals (i.e. 5 occupied and 5 virtual) were carried out
27 at the sulfur L_{2,3}-edge, demonstrating that the computational cost can be significantly lowered
28 without sacrificing accuracy.

29 Specific to NEXAFS spectroscopy, the composition of the sub-space of virtual orbitals (i.e.
30 RAS3) determines the nature of the core-excited states (i.e. core-to- π^* , core-to- σ^* , core-to-
31 Rydberg) and the completeness of the spectrum. Different strategies were pursued for the
32 carbon K-edge and the sulfur L_{2,3}-edge, dictated by the angular momentum selection rule in
33 atomic spectroscopy $\Delta l = \pm 1$ which (loosely) determines signal intensities in core-excited spec-
34 tra of molecules.

35 Specifically, the most intense contributions in the carbon pre-edge spectrum arise from transi-
36 tions to π^* orbitals (exhibiting large MO coefficients from 2p atomic orbitals). Low intensity
37 contributions arise due to transitions in virtual orbitals of (pseudo) Rydberg- and σ^* -type.
38 Thus, a two-step strategy was adopted. In the first step, SA-10-RAS(16,1,4;1,7,7) and SA-10-
39 RAS(20,1,4;1,9,9) for BTF6 and PTF6, respectively, comprised entirely of π molecular orbitals
40 were used to describe accurately core-to- π^* transitions and shake-up satellites, i.e., less intense
41 features in the spectra arising from mixed double excitations in which a core-excitation is accom-
42 panied by a valence π -to- π^* transition. Subsequently, the π and π^* orbitals were rotated out of
43 the AS and replaced with a set of 10 virtual orbitals, which results in a SA-10-RAS(2,1,1;1,0,10)
44 for both BTF6 and PTF6 and targets a qualitative description of core-to-Rydberg and core-to-
45 σ^* transitions.

46 Instead, according to the selection rule the most intense contributions in the sulfur L_{2,3} pre-edge
47 spectrum arise from transitions to orbitals with 4s- and 3d-character on the sulfur. Core-to- π^*
48 transitions are, indeed, present in the pre-edge region but are an order of magnitude less intense.
49 Since a full- π AS extended with additional orbitals of s/d -type is unfeasible, trials were carried
50 out to determine the optimal reduced AS which allows to capture the main contributions. The
51 optimal AS compositions were found to be SA-25-CAS(6,11) and SA-19-CAS(6,9) for BTF6 and
52 PTF6, respectively, containing four virtual orbitals with dominant s/d -character and additional
53 four (BTF6) / two (PTF6) virtual orbitals of π^* -character, which have a negligible contribution
54 to the spectra but confer active space stability. The number of states in the state-averaging was
55 determined considering the electronic ground state and an excitation out of each of the three 2p
56 core orbitals into each virtual orbital. The minimal AS (CAS(6,11) and CAS(6,9)) calculations
57 presented convergence problems in the CASPT2 correction of the electronic ground state, pre-
58 sumably due to near-degenerate configurations involving occupied π -orbitals left outside of the
59 minimal AS. For this reason, no reliable theoretical ground state reference could be obtained for
60 calculating the relative energies of the core-excited states. Hence, for both systems the NEX-
61 AFS spectra were simulated considering only the (absolute) core-excited states energies and the
62 resulting spectra were shifted to make the first peak coincide with its experimental counterpart.

63 **Active space definition for PES**

64 The active space comprises of valence occupied and virtual orbitals of π -nature, distributed
65 between sub-spaces RAS1 and RAS3 (thus leaving RAS2 empty) and allowing for up to four
66 holes and four excitations, respectively. A total of 8 and 10 valence-ionized states were computed
67 in the case of BTF6 and PTF6, respectively, giving rise to SA-8-RASSCF(13,4,4;7,0,7) and
68 SA-10-RASSCF(17,4,4;9,0,9). A different strategy was followed to account for ionization from
69 σ -orbitals. The AS comprised of 20 orbitals with π and σ character, with no virtual orbitals
70 included, leading to a SA-20-CAS(39,20) (i.e. 39 electrons in 20 orbitals). This minimal AS
71 allows to optimize the form of the lowest 20 occupied orbitals subject to ionization but lacks the
72 flexibility to describe electron correlation effects. As a consequence the error with respect to the
73 experimental data (which translates into a rigid shift in the simulated spectrum) is in the order
74 of 1 eV (compared to 0.5 eV correction applied in the case of full- π active spaces). Despite the
75 larger absolute error, the results with CAS(39,20) are presented in the main text as they allow
76 for a more complete description of the PES. The spectra obtained with the restricted AS are
77 presented in Fig. S17.

78 **Sulfur L_{2,3}-edge**

79 **XPS**

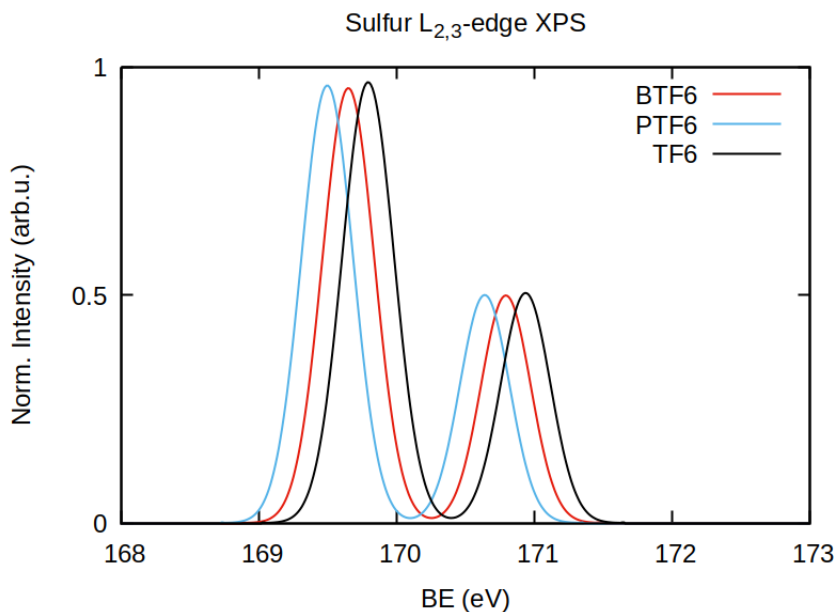


Figure S4: Comparison between the XPS spectra at the L_{2,3} edge of sulfur for BTF6 (red), PTF6 (blue) and TF6 (black). As evident stripping the phenyl groups for PTF6 leads to a blue-shift of the signal due to the reduced extension of the conjugation.

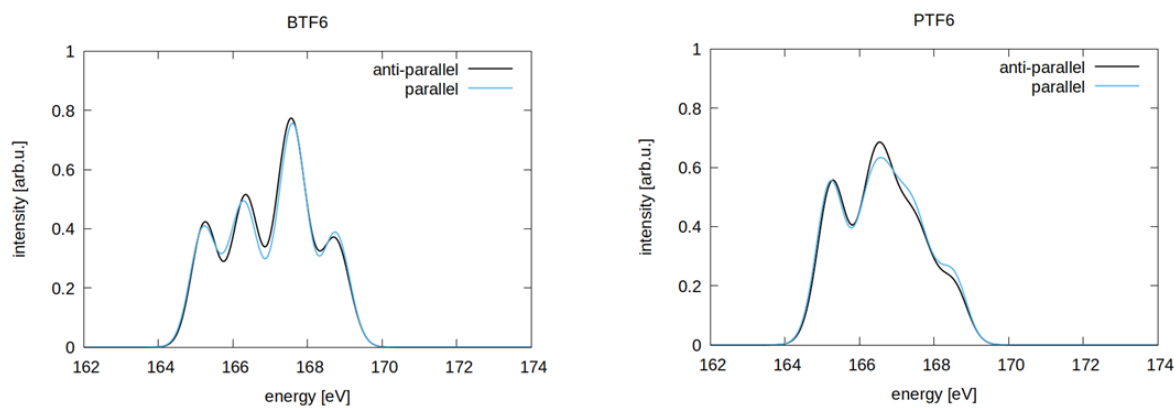


Figure S5: Comparison of the NEXAFS spectra at the sulfur $L_{2,3}$ -edge of the parallel (blue) and anti-parallel (black) OR forms of BTF6 (left) and PTF6 (right) computed at the SS-PT2/SA-25-CAS(6,11) and SS-PT2/SA-19-CAS(6,9) levels of theory, respectively.

Carbon K-edge

82 NEXAFS

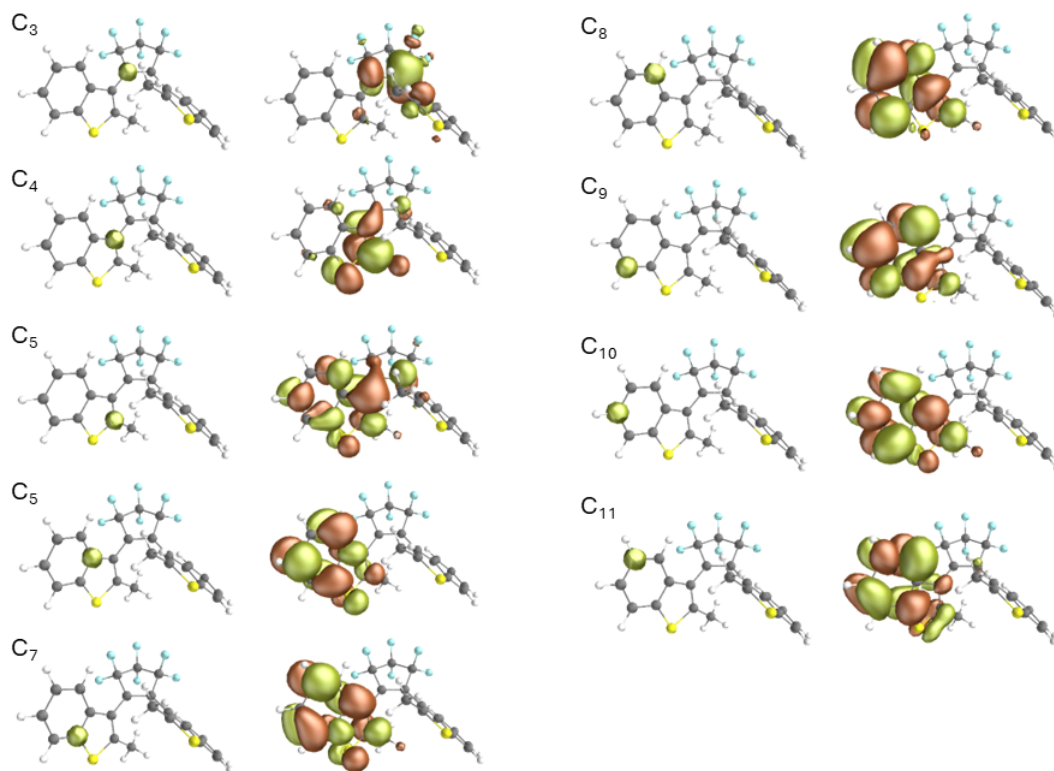


Figure S6: Pairs of 1s core and corresponding relaxed LUMO giving rise to 284.0-286.0 eV signal at the carbon K-edge for the anti-parallel conformation of BTf6. C₁, C₂ and the methyl carbon are omitted as they do not contribute a bright peak in the NEXAFS.

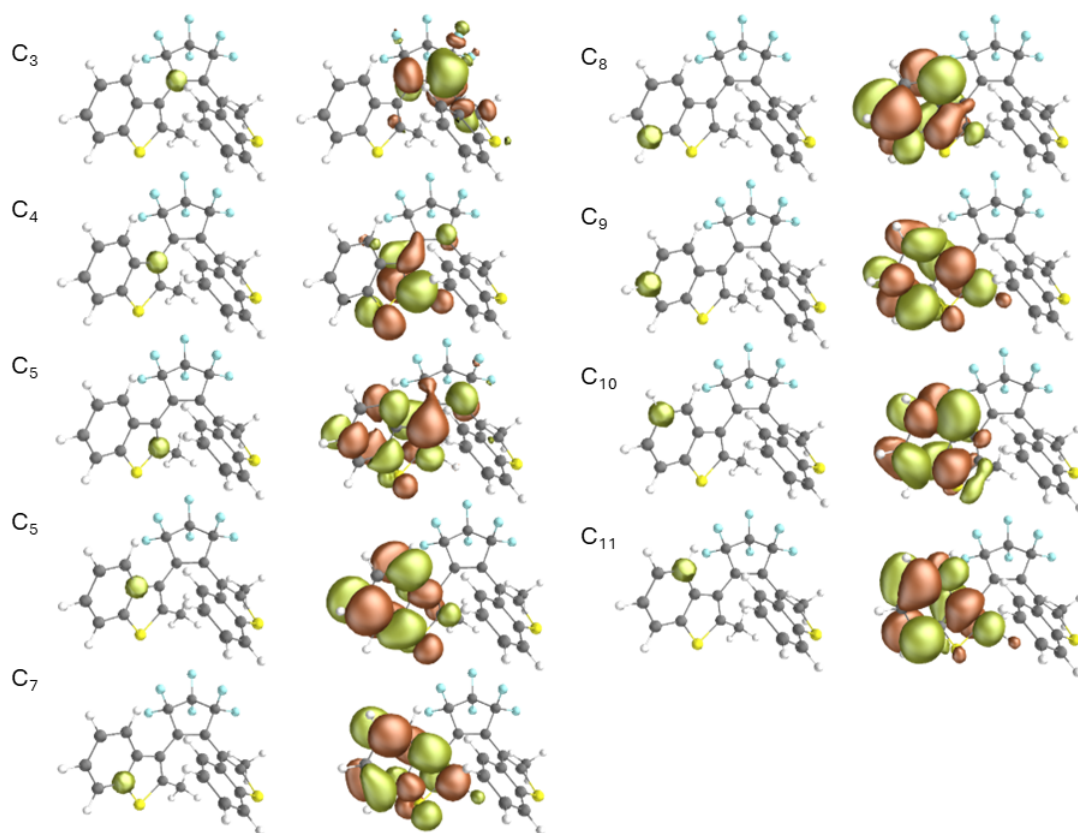


Figure S7: Pairs of 1s core and corresponding relaxed LUMO giving rise to 284.0-286.0 eV signal at the carbon K-edge for the parallel conformation of BTf6. C₁, C₂ and the methyl carbon are omitted as they do not contribute a bright peak in the NEXAFS.

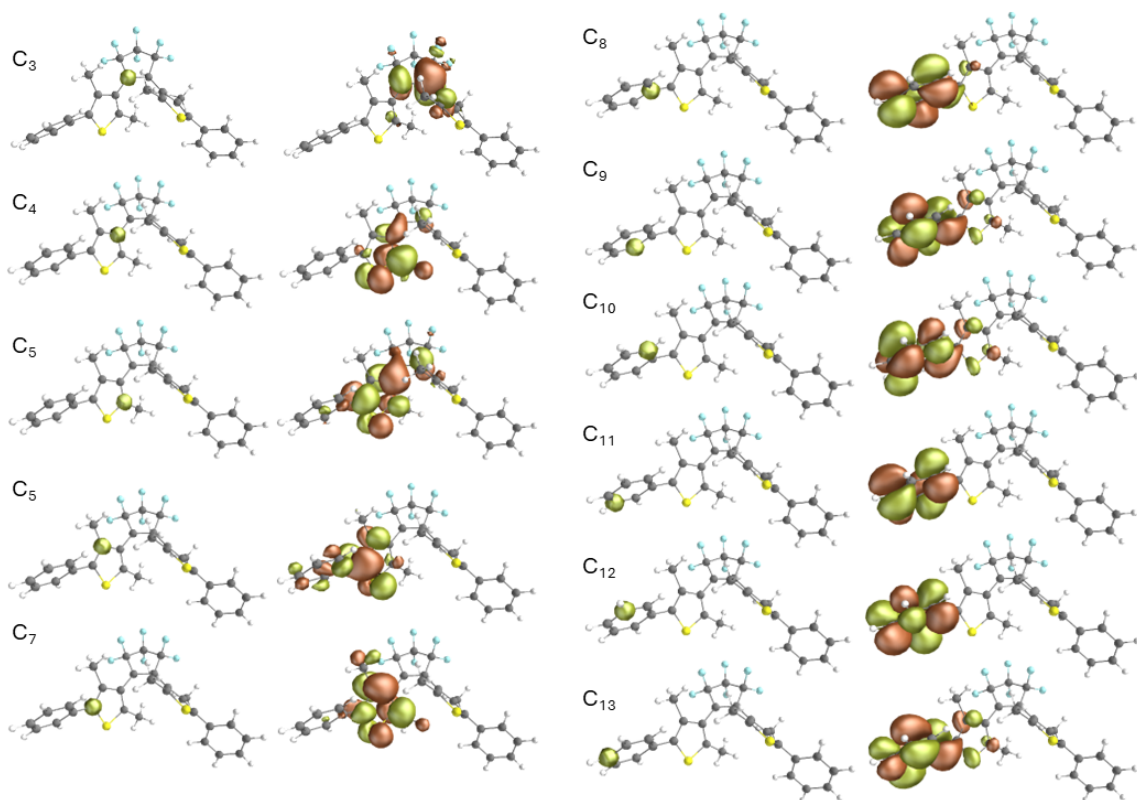


Figure S8: Pairs of 1s core and corresponding relaxed LUMO giving rise to 284.0-286.0 eV signal at the carbon K-edge for the anti-parallel conformation of PTF6. C₁, C₂ and the methyl carbons are omitted as they do not contribute a bright peak in the NEXAFS.

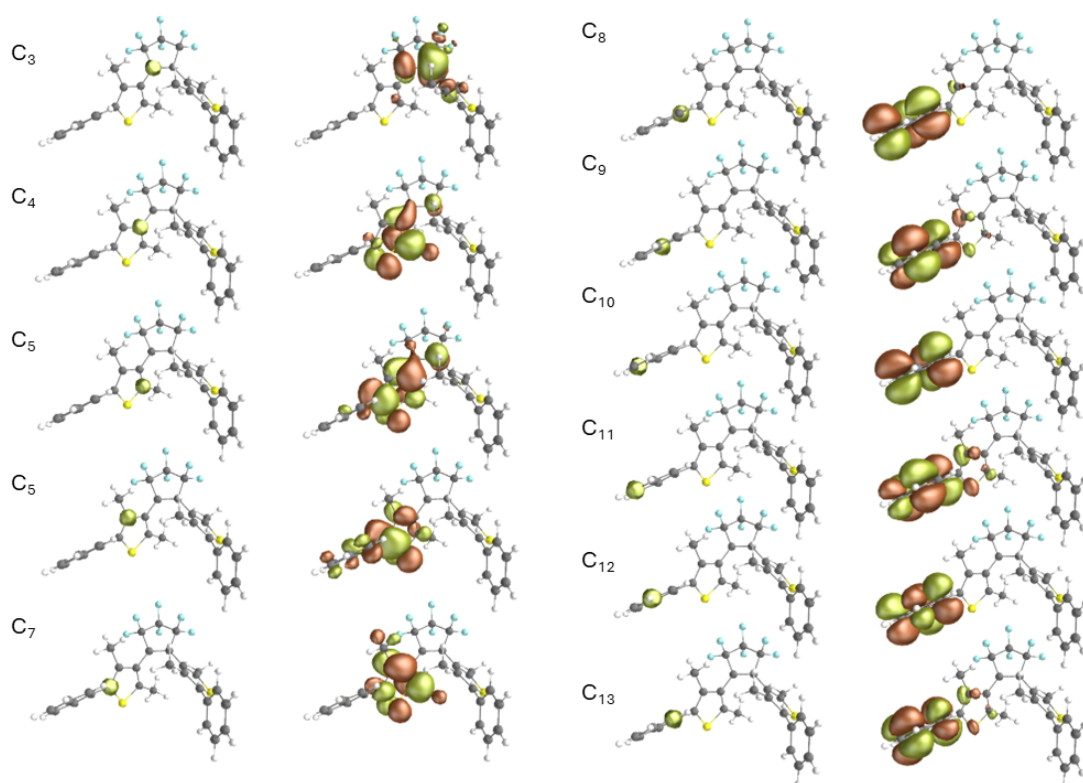


Figure S9: Pairs of 1s core and corresponding relaxed LUMO giving rise to 284.0-286.0 eV signal at the carbon K-edge for the parallel conformation of PTF6. C₁, C₂ and the methyl carbons are omitted as they do not contribute a bright peak in the NEXAFS.

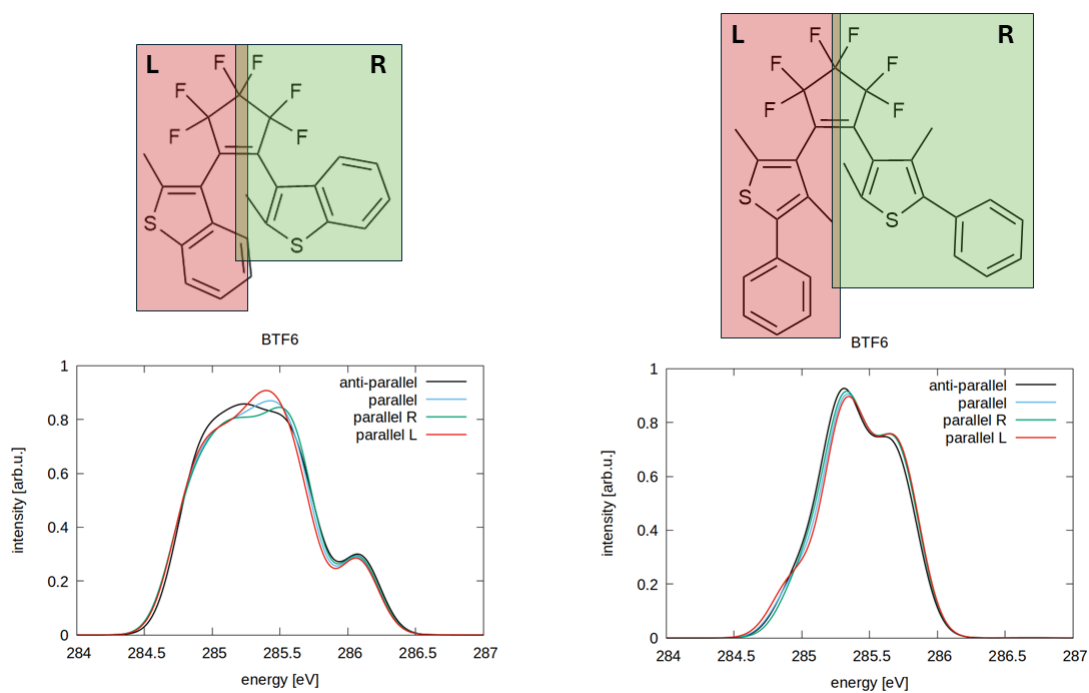


Figure S10: Comparison of the NEXAFS spectra at the carbon K-edge of the parallel (blue) and anti-parallel (black) OR forms of BTF6 (left) and PTF6 (right) computed at the SS-PT2/SA-10-RAS(16,1,4;1,7,7) and SS-PT2/SA-10-RAS(20,1,4;1,9,9) levels of theory, respectively. Only contribution to the first transition $1s \rightarrow \pi^*$ is shown. For the parallel form spectra obtained by considering only the left (L, red) or right (R, green) part of the molecule are also shown.

| carbon # | BTF6 | | | | PTF6 | | | |
|-------------------|-------------------------|--------------------|--------------------|------------------|-------------------------|--------------------|--------------------|------------------|
| | anti-parallel R (eV) | parallel R (eV) | parallel L (eV) | Δ (eV) | anti-parallel R (eV) | parallel R (eV) | parallel L (eV) | Δ (eV) |
| 1 | 291.49 | 291.48 | 291.48 | - | 291.59 | 291.60 | 291.60 | - |
| 2 | 290.70 | 290.68 | 290.69 | -0.01 | 290.78 | 290.80 | 290.75 | 0.05 |
| 3 | 284.83 | 284.79 | 284.80 | -0.01 | 284.97 | 284.99 | 284.93 | 0.06 |
| 4 | 285.48 | 285.48 | 285.43 | 0.05 | 285.67 | 285.70 | 285.70 | 0.00 |
| 5 | 285.57 | 285.59 | 285.55 | 0.04 | 285.64 | 285.67 | 285.69 | -0.02 |
| 6 | 285.71 | 285.68 | 285.68 | 0.00 | 285.75 | 285.74 | 285.76 | -0.02 |
| 7 | 286.09 | 286.08 | 286.07 | 0.01 | 285.78 | 285.79 | 285.78 | 0.01 |
| 8 | 285.21 | 285.20 | 285.21 | -0.01 | 285.56 | 285.56 | 285.57 | -0.01 |
| 9 | 285.10 | 285.09 | 285.09 | 0.00 | 285.28 | 285.31 | 285.32 | -0.01 |
| 10 | 285.32 | 285.34 | 285.35 | -0.01 | 285.33 | 285.33 | 285.36 | -0.03 |
| 11 | 284.96 | 284.97 | 284.96 | 0.01 | 285.22 | 285.23 | 285.24 | -0.01 |
| 12 | - | - | - | - | 285.35 | 285.37 | 285.37 | 0.00 |
| 13 | - | - | - | - | 285.27 | 285.29 | 285.29 | 0.00 |
| CH ₃ 1 | 286.49 | 286.57 | 286.64 | -0.07 | 286.62 | 286.64 | 286.82 | -0.18 |
| CH ₃ 2 | - | - | - | - | 286.71 | 286.72 | 286.53 | 0.19 |

Table 2: Core-excitation energies of the parallel and anti-parallel OR forms of BTF6 and PTF6 and relative difference Δ computed at the SS-PT2/SA-10-RAS(16,1,4;1,7,7) and SS-PT2/SA-10-RAS(20,1,4;1,9,9) levels of theory, respectively. For the parallel form values for "equivalent" atoms from the right (R) and left (L) part of the molecule are compared.

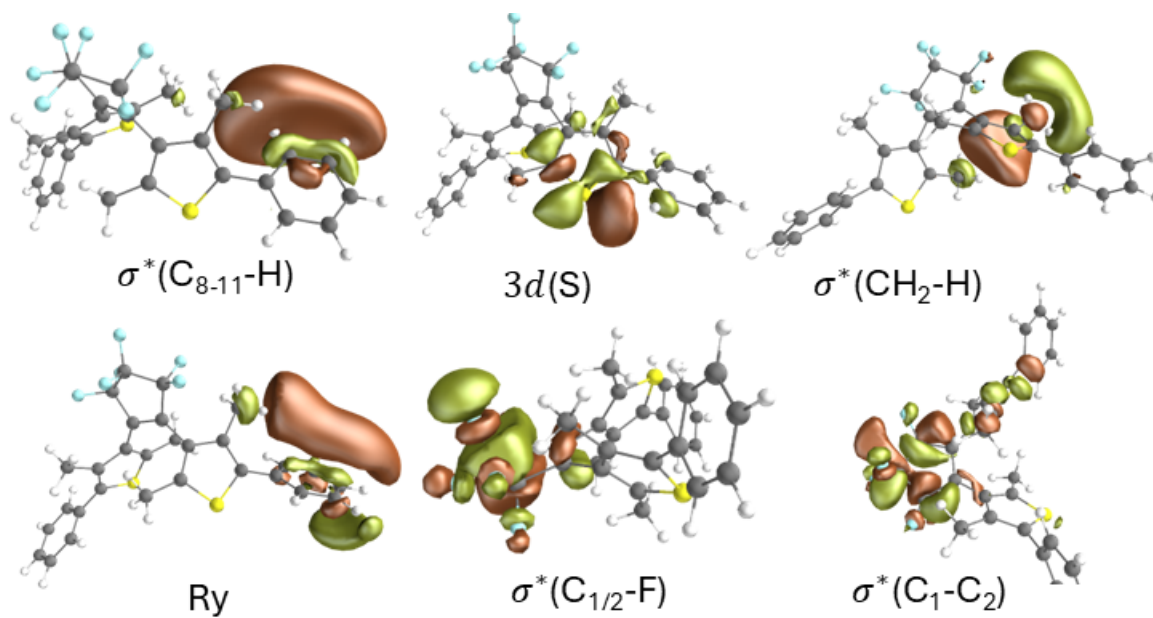


Figure S11: Representative arrival orbitals giving rise to signal above 286.0 eV at the carbon K-edge for the anti-parallel conformation of PTF6.

PES

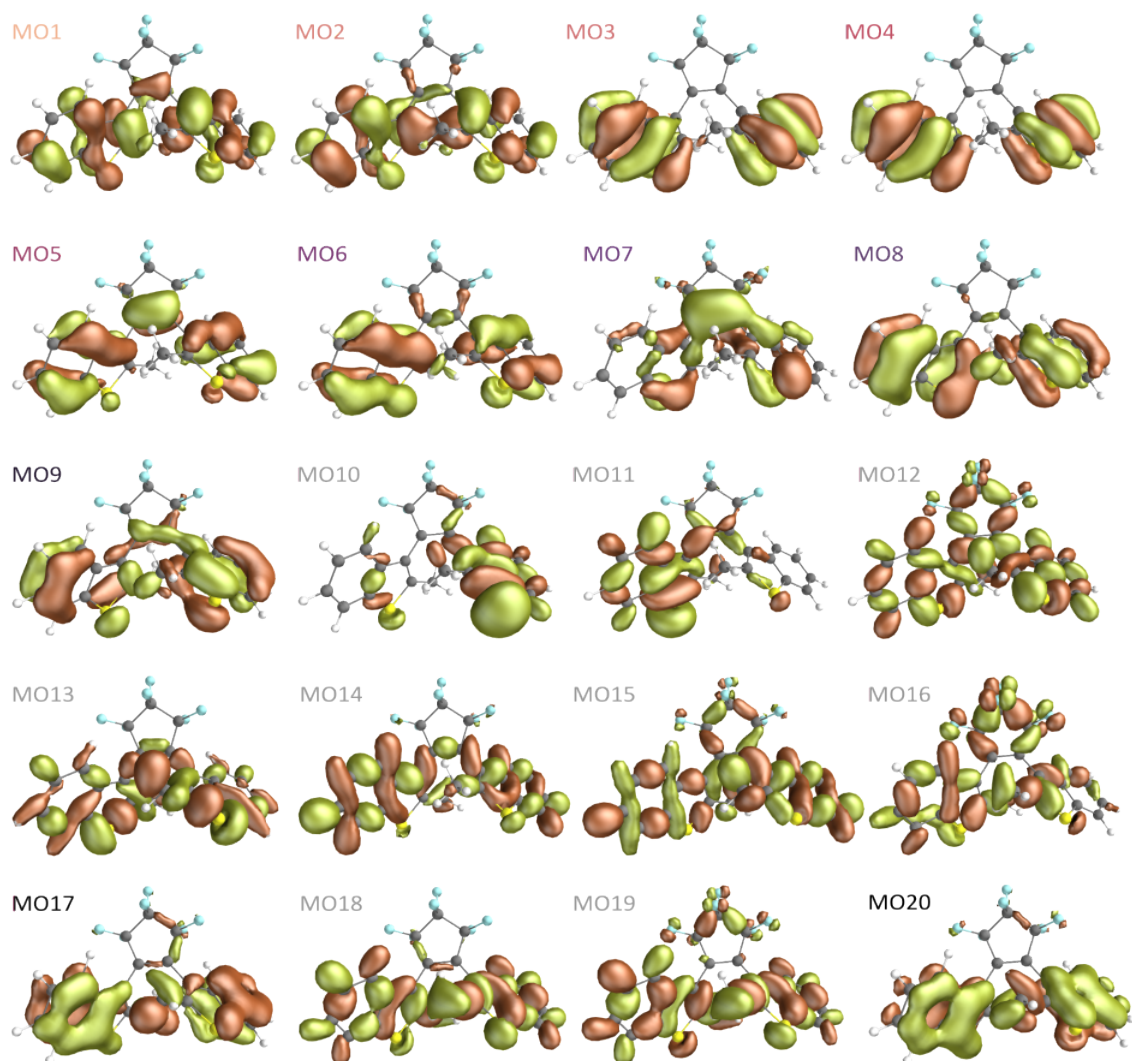


Figure S12: Active space occupied valence orbitals for anti-parallel conformation of BTF6. The color of each label is the same of the correspondent stick representing each contribution in the PES spectrum of the molecular system in panel (a) of Fig. 4 of the main text. The gray labels indicate orbitals with σ character, while colored labels describe orbitals having a π character.

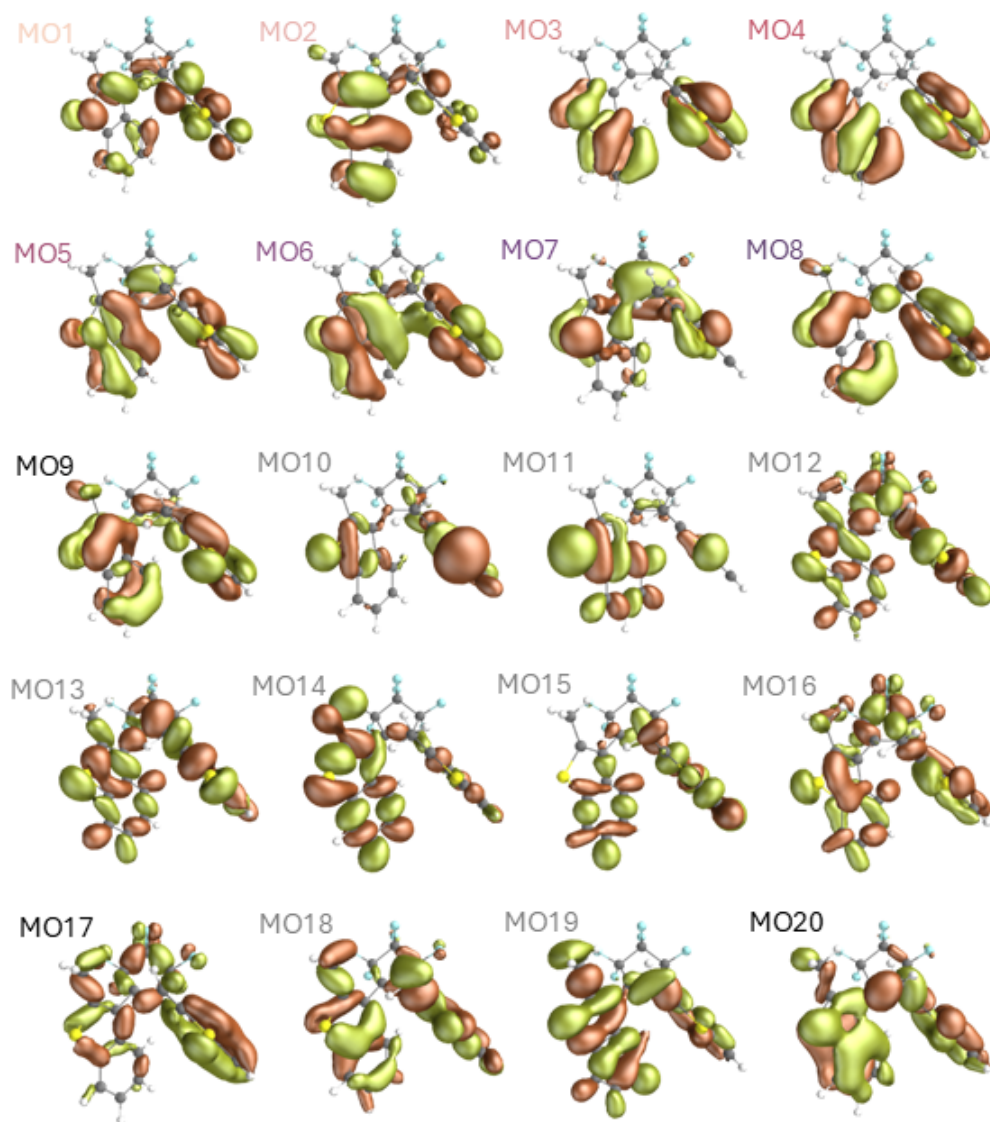


Figure S13: Active space occupied valence orbitals for parallel conformation of BTF6. The color of each label is the same of the correspondent stick representing each contribution in the PES spectrum of the molecular system in panel (a) of Fig. 4 of the main text. The gray labels indicate orbitals with σ character, while colored labels describe orbitals having a π character.

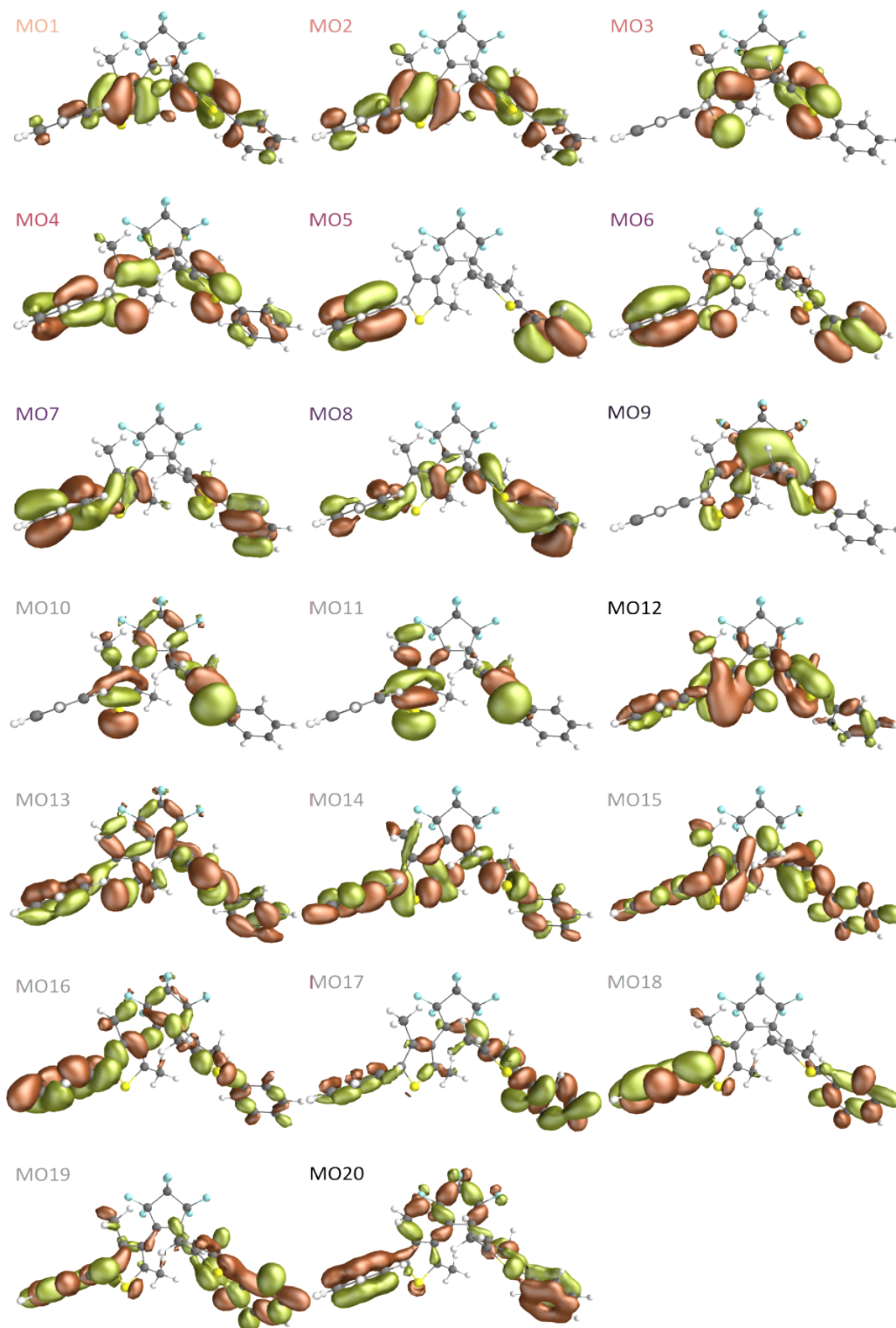


Figure S14: Active space occupied valence orbitals for anti-parallel conformation of PTF6. The color of each label is the same of the correspondent stick representing each contribution in the PES spectrum of the molecular system in panel (b) of Fig. 4 of the main text. The gray labels indicate orbitals with σ character, while colored labels describe orbitals having a π character.

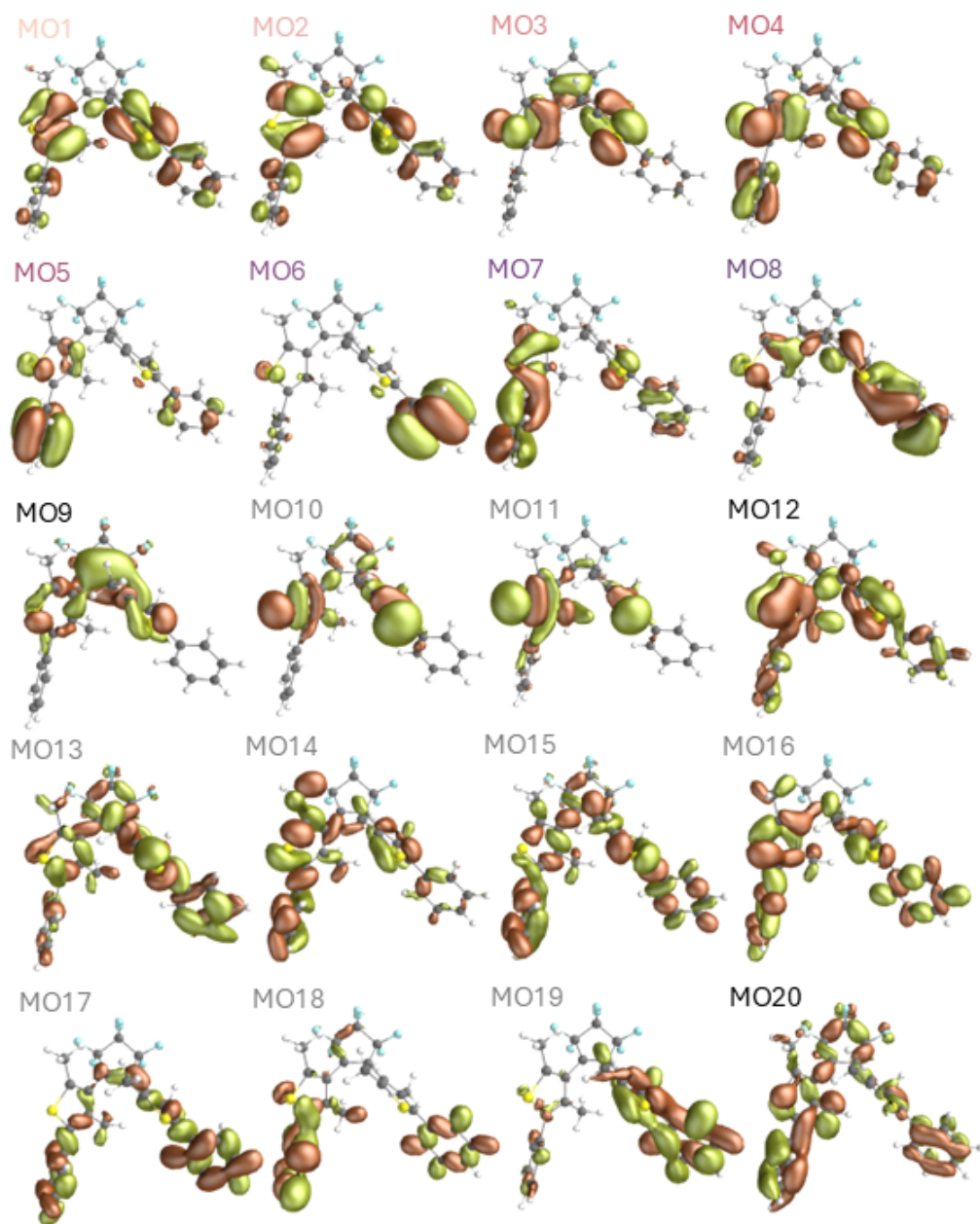


Figure S15: Active space occupied valence orbitals for parallel conformation of PTF6. The color of each label is the same of the correspondent stick representing each contribution in the PES spectrum of the molecular system in panel (b) of Fig. 4 of the main text. The gray labels indicate orbitals with σ character, while colored labels describe orbitals having a π character.

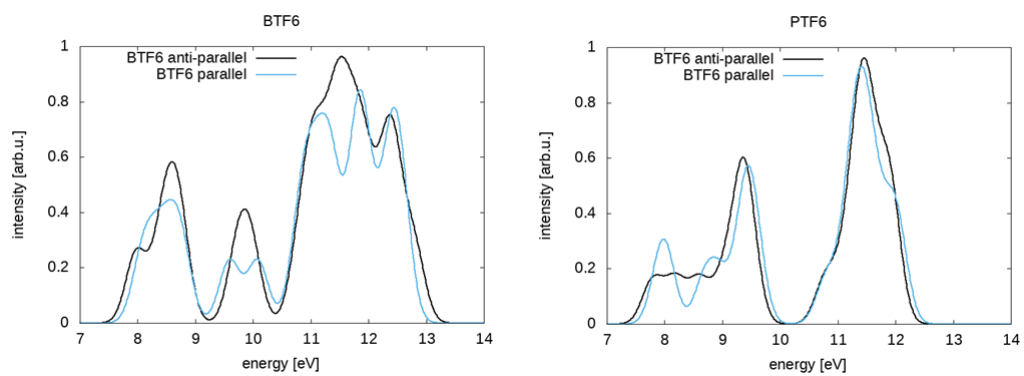


Figure S16: Comparison of the PES spectra of the parallel (blue) and anti-parallel (black) OR forms of BTF6 (left) and PTF6 (right) computed at the SS-PT2/SA-20-CAS(39,20) level of theory.

| transition # | BTF6 | | | PTF6 | | |
|-----------------|-----------------------|------------------|------------------|-----------------------|------------------|------------------|
| | anti-parallel (eV) | parallel (eV) | Δ (eV) | anti-parallel (eV) | parallel (eV) | Δ (eV) |
| 1 | 7.98 | 8.08 | -0.10 | 7.80 | 7.94 | -0.14 |
| 2 | 8.41 | 8.33 | 0.08 | 8.19 | 8.02 | 0.17 |
| 3 | 8.60 | 8.58 | 0.02 | 8.60 | 8.70 | -0.10 |
| 4 | 8.74 | 8.77 | -0.03 | 9.01 | 8.96 | 0.05 |
| 5 | 9.74 | 9.59 | 0.15 | 9.49 | 9.47 | 0.02 |
| 6 | 9.97 | 10.08 | -0.11 | 9.33 | 9.59 | 0.26 |
| 7 | 10.81 | 10.79 | 0.02 | 9.30 | 9.40 | -0.10 |
| 8 | 11.06 | 10.93 | 0.13 | 9.39 | 9.38 | 0.01 |
| 9 | 11.05 | 11.04 | 0.01 | 10.79 | 10.80 | -0.01 |
| 10 | 11.45 | 11.34 | 0.11 | 11.10 | 11.13 | -0.03 |
| 11 | 11.49 | 11.40 | 0.09 | 11.36 | 11.33 | 0.03 |
| 12 | 11.28 | 11.21 | 0.07 | 11.33 | 11.30 | 0.03 |
| 13 | 12.02 | 11.86 | 0.16 | 11.37 | 11.31 | 0.06 |
| 14 | 11.62 | 11.82 | -0.20 | 11.50 | 11.52 | -0.02 |
| 15 | 11.77 | 11.80 | -0.03 | 11.49 | 11.45 | 0.04 |
| 16 | 11.88 | 12.00 | -0.12 | 11.54 | 11.52 | 0.02 |
| 17 | 12.50 | 12.36 | 0.14 | 11.69 | 11.67 | 0.02 |
| 18 | 12.33 | 12.35 | -0.02 | 11.94 | 11.86 | 0.08 |
| 19 | 12.35 | 12.55 | -0.20 | 11.91 | 12.04 | -0.13 |
| 20 | 12.80 | 12.56 | 0.24 | 11.88 | 12.02 | -0.14 |

Table 3: Ionization energies of the parallel and anti-parallel OR forms of BTF6 and PTF6 and relative difference Δ computed at the SS-PT2/SA-20-CAS(39,20) level of theory.

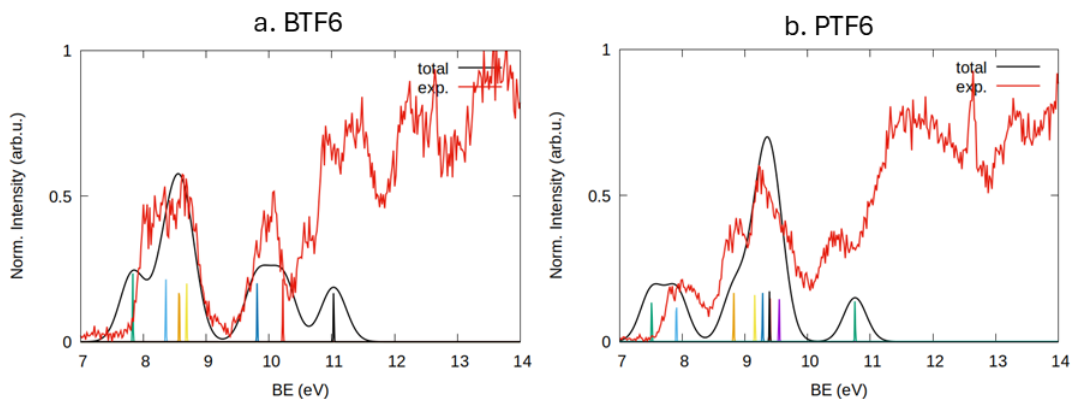


Figure S17: PES spectra of the OR(ap) form of BTF6 (a) and PTF6 (b) obtained with the full- π active spaces SA-8-RASSCF(13,4,4;7,0,7) and SA-10-RASSCF(17,4,4;9,0,9), respectively. Computed signals have been blue-shifted uniformly by 0.5 eV.

84

BTF6 Closed-ring form

XPS - Sulfur $L_{2,3}$ -edge

| | BE (eV) | Leading transition |
|-------|---------|--------------------|
| S_1 | 169.05 | |
| S_2 | 169.15 | 2p ionization |
| S_3 | 170.24 | |

Table 4: Theoretically calculated sulfur $L_{2,3}$ -edge binding energies of the CR form of BTF6, with the inclusion of spin-orbit-coupling. The reported binding energy values have been shifted by +0.05 eV as done for the OR form. The Dyson norm for all the ionization processes in the two molecules is equal to 1.50.

XPS - Carbon K-edge

| Transition | BE (eV) | Dyson norm |
|--------------------|------------|---------------|
| C ₁ | 296.08 | 0.72 |
| C ₂ | 295.93 | 0.72 |
| C ₃ | 290.65 | 0.57 |
| C ₄ | 290.56 | 0.60 |
| C ₅ | 291.53 | 0.72 |
| C ₆ | 290.58 | 0.65 |
| C ₇ | 290.80 | 0.63 |
| C ₈ | 290.28 | 0.64 |
| C ₉ | 290.23 | 0.59 |
| C ₁₀ | 290.15 | 0.57 |
| C ₁₁ | 290.21 | 0.61 |
| C(H ₃) | 290.44 | 0.73 |

Table 5: Theoretically calculated carbon K-edge binding energies for the CR form of BTF6. BEs have been shifted by +0.25 eV as done for the OR form. The different signals contributing to the spectra are due to ionization from the core orbitals 1s of the C_i carbon atoms, with numeration reported in the first row of Fig. 3 of the main paper.

NEXAFS - Carbon K-edge

| Peak | Site | Energy (eV) | arrival MO |
|------|---------------------------------|---------------|--------------|
| A | C ₃ | 284.42 | π^* LUMO |
| | C ₄ | 284.18 | |
| | C ₆ | 284.73 | |
| | C ₇ | 285.16 | |
| | C ₈ -C ₁₁ | 284.38-284.85 | |

Table 6: Theoretically calculated core excited states at carbon K-edge of BTF6 and PTF6. For each carbon contributing to the NEXAFS signal and the transition energies are reported. Only transitions from carbons having oscillator strength larger than 10^{-4} are considered. No shift has been applied.

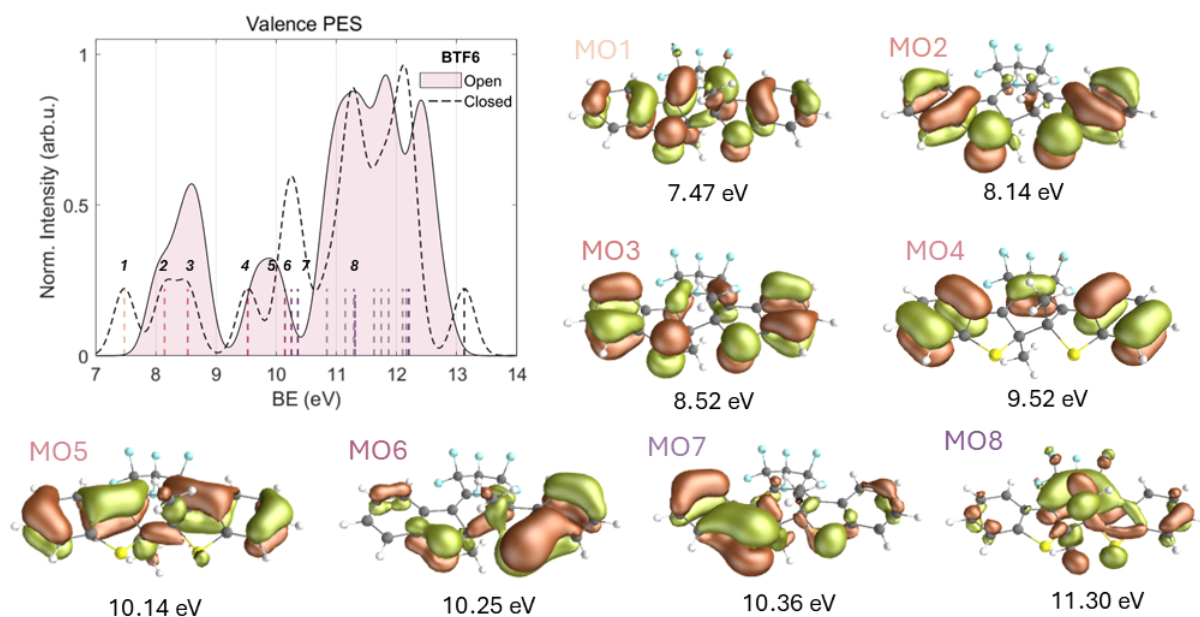


Figure S18: Simulated valence-ionization photoelectron spectra of the OR (shaded area) and CR (dashed line) form of BTF6, both shifted by +0.9 eV. Colored and gray lines stand for ionization from orbitals with and π and σ character, respectively. π MOs giving rise to the signal below 11 eV are shown.

References

- 85
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