

Alma Mater Studiorum Università di Bologna Archivio istituzionale della ricerca

Lipschitz regularity of almost minimizers in one-phase problems driven by the p-Laplace operator

This is the final peer-reviewed author's accepted manuscript (postprint) of the following publication:

Published Version:

Dipierro, S., Ferrari, F., Forcillo, N., Valdinoci, E. (2024). Lipschitz regularity of almost minimizers in one-phase problems driven by the p-Laplace operator. INDIANA UNIVERSITY MATHEMATICS JOURNAL, 73(3), 813-854 [10.1512/iumj.2024.73.9926].

Availability:

This version is available at: https://hdl.handle.net/11585/975974 since: 2024-07-29

Published:

DOI: http://doi.org/10.1512/iumj.2024.73.9926

Terms of use:

Some rights reserved. The terms and conditions for the reuse of this version of the manuscript are specified in the publishing policy. For all terms of use and more information see the publisher's website.

This item was downloaded from IRIS Università di Bologna (https://cris.unibo.it/). When citing, please refer to the published version.

(Article begins on next page)

LIPSCHITZ REGULARITY OF ALMOST MINIMIZERS IN ONE-PHASE PROBLEMS DRIVEN BY THE p-LAPLACE OPERATOR

SERENA DIPIERRO, FAUSTO FERRARI, NICOLÒ FORCILLO, AND ENRICO VALDINOCI

ABSTRACT. We prove that, given $p > \max\left\{\frac{2n}{n+2}, 1\right\}$, the nonnegative almost minimizers of the nonlinear free boundary functional

$$J_p(u,\Omega) := \int_{\Omega} \left(|\nabla u(x)|^p + \chi_{\{u>0\}}(x) \right) dx$$

are Lipschitz continuous.

1. Introduction

In this article we consider a nonlinear free boundary problem and we establish the Lipschitz continuity of its almost minimizers. The classical motivations for free boundary problems of these types stem from flows with jets and cavities (see e.g. Section 1.1 in [CS05]). In this context, the conditions arising from a free boundary problem can be seen as the variational counterpart of Bernoulli's law according to which pressure is prescribed on the free streamline as a balance with the velocity (or the kinetic energy) of the fluid.

The nonlinear feature of the corresponding differential operator aims at modeling possibly non-Newtonian fluids, in which linear relations between physical quantities are replaced by power-laws.

Similar models also appear in the study of electrical impedance tomography (see [AP98]), optimal heat flows (see [Ack77]), electrochemical machining (see [LS87]) and high activation energy in combustion theory (see [BCN90,DPS03,Kar06]). Also, free boundary problems can be used as a sharp-interface approximation of phase coexistence models (see e.g. [PV05a, PV05b, Val06], roughly speaking as a replacing for phases modeled by a smooth state parameter with a merely continuous one).

Date: June 8, 2022.

S.D. and E.V. are members of AustMS. F.F. and N.F. are members of INdAM. S.D. has been supported by the Australian Research Council DECRA DE180100957 "PDEs, free boundaries and applications". F.F. and N.F. have been supported by 2022- INDAM- GNAMPA project Regolarità locale e globale per problemi completamente nonlineari. N.F. has been supported by GHAIA Horizon 2020 MCSA RISE programme grant No 777822. E.V. has been supported by the Australian Laureate Fellowship FL190100081 "Minimal surfaces, free boundaries and partial differential equations". Part of this work has been completed during a very pleasant visit of N.F. to the University of Western Australia, that we thank for the warm hospitality. The authors wish to thank Daniela De Silva and Ovidiu Savin for useful discussions.

The setting of almost minimizers is also classical in the calculus of variations (dating back, at least up to a certain extent, to a famous sentence in Leibniz's *Specimen Geometriae Luciferae*, probably written in the mid-1690s, "pro minimis adhiberi possunt quasi minima", that is "the almost minimizers can be exploited in place of minimizers").

More specifically, the mathematical setting that we consider here goes as follows. Let $\Omega \subset \mathbb{R}^n$ be a given domain and $p > \max\left\{\frac{2n}{n+2}, 1\right\}$. We consider the energy functional

(1.1)
$$J_p(u,\Omega) := \int_{\Omega} \left(|\nabla u(x)|^p + \chi_{\{u>0\}}(x) \right) dx$$

for all $u \in W^{1,p}(\Omega)$ with $u \geqslant 0$.

The condition that u is nonnegative corresponds, in the framework of free boundary problems, to considering "one-phase" solutions (solutions which may change sign being related to "two-phase" problems).

The precise notion of almost minimizers that we use in this paper is the following one:

Definition 1.1. Let $\kappa \geqslant 0$ and $\beta > 0$. We say that $u \in W^{1,p}(\Omega)$ is an almost minimizer for J_p in Ω , with constant κ and exponent β , if $u \geqslant 0$ a.e. in Ω and

(1.2)
$$J_p(u, B_{\varrho}(x)) \leqslant (1 + \kappa \varrho^{\beta}) J_p(v, B_{\varrho}(x)),$$

for every ball $B_{\varrho}(x)$ such that $\overline{B_{\varrho}(x)} \subset \Omega$ and for every $v \in W^{1,p}(B_{\varrho}(x))$ such that v = u on $\partial B_{\varrho}(x)$ in the sense of the trace.

In some sense, Definition 1.1 is one of the possible modern formalizations of Leibniz's initial intuition reported at the beginning of this paper: namely, almost minimizers are natural objects to look at, for instance, to deal with minimizers of "perturbed" functionals. As a concrete example, if we consider

$$\widetilde{J}_p(u,\Omega) := J_p(u,\Omega) + \iint_{\Omega \times \Omega} \Phi(u(y)) \,\Phi(u(z)) \,\Phi(u(y) - u(z)) \,dy \,dz$$

for a function $\Phi: \mathbb{R} \to [0,1]$ with $\Phi=0$ in $(-\infty,0]$, we readily see that $J_p(u,B_r(x)) \leqslant \widetilde{J}_p(u,B_r(x))$ and

$$\widetilde{J}_p(u, B_r(x)) \leqslant J_p(u, B_r(x)) + \iint_{B_r(x) \times B_r(x)} \chi_{\{u > 0\}}(y) \, \chi_{\{u > 0\}}(z) \, dy \, dz$$

$$\leqslant J_p(u, B_r(x)) + |B_r| \int_{B_r(x)} \chi_{\{u > 0\}}(y) \, dy \leqslant (1 + |B_1| \, r^n) J_p(u, B_r(x)).$$

Accordingly, a minimizer for the "complicated" functional \widetilde{J}_p turns out to be an almost minimizer for the "simpler" functional J_p .

As usual, the constants depending only on n and p are called universal. If u is an almost minimizer, the structural constants may depend on κ and β as well.

Our main result establishes the Lipschitz regularity of the one-phase almost minimizers as follows:

Theorem 1.2. Let $p > \max\left\{\frac{2n}{n+2}, 1\right\}$ and u be an almost minimizer for J_p in B_1 with constant κ and exponent β .

Then,

$$\|\nabla u\|_{L^{\infty}(B_{1/2})} \leqslant C(\|\nabla u\|_{L^{p}(B_{1})} + 1),$$

where C > 0 is a constant depending on n, p, κ and β .

In addition, u is uniformly Lipschitz continuous in a neighborhood of $\{u = 0\}$, namely if u(0) = 0 then

$$|\nabla u| \leqslant C \quad in \ B_{r_0},$$

for some C > 0, depending only on n, p, κ and β , and $r_0 \in (0,1)$, depending on n, p, κ , β and $\|\nabla u\|_{L^p(B_1)}$.

We stress that Theorem 1.2 is new for $p \neq 2$, the case p = 2 being treated in [DSS20b].

We recall that, when p=2, minimizers of (1.1) were studied in [AC81], where the Lipschitz regularity of minimizers and the regularity of flat free boundaries were established. In [Caf87, Caf89, Caf88] Caffarelli developed a viscosity approach to the free boundary problem for p=2 in the two-phase setting (see also [DS11] and the references therein for related free boundary regularity properties). For one-phase problems the viscosity approach for operators governed by the p-Laplace or p(x)-Laplace operators has been developed in [LR18, FL21] and in [FL22] as well, where Lipschitz regularity of viscosity solutions of inhomogeneous one-phase problems, governed by the p(x)-Laplace operator and some regularity properties of their free boundaries were established.

The Lipschitz regularity of the minimizers of the functional in (1.1) has been obtained in [DK18a], which has also provided a proof of the Lipschitz regularity when p=2 without using monotonicity formulae. We refer to [DK18b] as well, where a discrete version of the Weiss monotonicity formula for the functional in (1.1) has been established for p close to 2.

Minimizers of the functional in (1.1) have been also considered in [DP05], in which the regularity of the free boundary near flat points was established. See also [DP06, Kar08, MW08, LN10, LN12, MW14, BM14, LdQT15, Kar21] for regularity results on p-Laplacian free boundary problems.

Regarding the setting of almost minimizers, the case p=2 has been investigated in [DT15, DET19, DSS20b]. The case of almost minimizers for $p \neq 2$ was, to the best of our knowledge, not fully investigated, hence Theorem 1.2 aims at starting a research line in this direction, and the results presented in this article are part of the PhD thesis of the third author [For21]. In this paper we take inspiration from the approach used in [DSS20b], see also [DSS20a, DSS21].

We point out that the case $p \neq 2$ provides significant technical complications, especially due to the fact that the sum of two solutions is not a solution any longer, thus making it difficult to develop an exhaustive theory of harmonic replacements in the nonlinear scenario.

We highlight the fact that the regularity results put forth in this paper do not follow from the classical ones in the calculus of variations, since the integrand of the functional that we consider here is a discontinuous function, due to the presence of the term $\chi_{\{u>0\}}$ (instead, the classical cases dealt with require the integrand to be continuous, or even Lipschitz continuous, see e.g. assumption (8.48) in [Giu03]).

The paper is organized as follows. In Section 2 we recall the notion of p-harmonic replacement and put forth some basic energy estimates needed in our main arguments.

Then, in Section 3, we develop a dichotomy theory according to which, roughly speaking, the average of the energy of an almost minimizer decreases in a smaller ball, unless we are arbitrarily close to the case of linear functions.

In Section 4 we show that almost minimizers of J_p are Lipschitz continuous, namely we prove Theorem 1.2.

2. The p-harmonic replacement

One of the classical ingredients in nonlinear partial differential equations is the notion of p-harmonic replacement, which we now recall:

Definition 2.1. Let r > 0, $x_0 \in \mathbb{R}^n$ and $u \in W^{1,p}(B_r(x_0))$. We say that $v \in W^{1,p}(B_r(x_0))$ is the p-harmonic replacement of u in $B_r(x_0)$ if

$$\int_{B_r(x_0)} |\nabla v|^p \ dx = \min_{u - w \in W_0^{1,p}(B_r(x_0))} \int_{B_r(x_0)} |\nabla w|^p \ dx.$$

We remark that if v is the p-harmonic replacement of u in $B_r(x_0)$, then in particular it satisfies

(2.1)
$$\int_{B_r(x_0)} |\nabla v(x)|^{p-2} \nabla v(x) \cdot \nabla \varphi(x) dx = 0$$

for all $\varphi \in W_0^{1,p}(B_r(x_0))$, namely v is a weak solution of $\Delta_p v = 0$ in $B_r(x_0)$.

We now provide some energy estimates of classical flavor for the p-harmonic replacement that we will use in the sequel.

Lemma 2.2. Let $x_0 \in \mathbb{R}^n$, r > 0 and $u \in W^{1,p}(B_r(x_0))$. Let v be the p-harmonic replacement of u in $B_r(x_0)$. Then,

(i) if
$$1 , then$$

(2.2)
$$\int_{B_{r}(x_{0})} \left| \nabla u(x) - \nabla v(x) \right|^{p} dx$$

$$\leq C \left(\int_{B_{r}(x_{0})} \left(\left| \nabla u(x) \right|^{p} - \left| \nabla v(x) \right|^{p} \right) dx \right)^{\frac{p}{2}} \left(\int_{B_{r}(x_{0})} \left(\left| \nabla u(x) \right| + \left| \nabla v(x) \right| \right)^{p} dx \right)^{1-\frac{p}{2}},$$

for some positive universal constant C;

(ii) if $p \ge 2$, then

(2.3)
$$\int_{B_r(x_0)} |\nabla u(x) - \nabla v(x)|^p \, dx \leqslant C \int_{B_r(x_0)} \left(|\nabla u(x)|^p - |\nabla v(x)|^p \right) dx,$$

for some positive universal constant C.

Proof. For all $s \in [0, 1]$, we consider the family of functions $u_s(x) := su(x) + (1 - s)v(x)$. Notice that $u_0 = v$ and $u_1 = u$. As a consequence,

$$\int_{B_{r}(x_{0})} \left(|\nabla u(x)|^{p} - |\nabla v(x)|^{p} \right) dx = \int_{B_{r}(x_{0})} \left(\int_{0}^{1} \frac{d}{ds} |\nabla u_{s}(x)|^{p} ds \right) dx$$

$$= \int_{B_{r}(x_{0})} \left(\int_{0}^{1} p |\nabla u_{s}(x)|^{p-2} |\nabla u_{$$

This and (2.1) (used here with $\varphi := u - v$) lead to

$$\int_{B_{r}(x_{0})} \left(\left| \nabla u(x) \right|^{p} - \left| \nabla v(x) \right|^{p} \right) dx$$

$$= p \left[\int_{B_{r}(x_{0})} \left(\int_{0}^{1} \left| \nabla u_{s}(x) \right|^{p-2} \nabla u_{s}(x) \cdot \nabla (u-v)(x) ds \right) dx$$

$$- \int_{0}^{1} \left(\int_{B_{r}(x_{0})} \left| \nabla v(x) \right|^{p-2} \nabla v(x) \cdot \nabla (u-v)(x) dx \right) ds \right]$$

$$= p \int_{0}^{1} \left(\int_{B_{r}(x_{0})} \left(\left| \nabla u_{s}(x) \right|^{p-2} \nabla u_{s}(x) - \left| \nabla v(x) \right|^{p-2} \nabla v(x) \right) \cdot \nabla (u-v)(x) dx \right) ds.$$

We also point out that

(2.4)
$$u_s(x) - v(x) = s(u(x) - v(x)),$$

and therefore we arrive at

(2.5)
$$\int_{B_r(x_0)} \left(\left| \nabla u(x) \right|^p - \left| \nabla v(x) \right|^p \right) dx$$

$$= p \int_0^1 \frac{1}{s} \left(\int_{B_r(x_0)} \left(\left| \nabla u_s(x) \right|^{p-2} \nabla u_s(x) - \left| \nabla v(x) \right|^{p-2} \nabla v(x) \right) \cdot \nabla (u_s - v)(x) dx \right) ds.$$

Now we recall the inequality

$$(|\xi|^{p-2}\xi - |\zeta|^{p-2}\zeta) \cdot (\xi - \zeta) \ge \gamma \begin{cases} |\xi - \zeta|^2 (|\xi| + |\zeta|)^{p-2} & \text{if } 1$$

for any ξ , $\zeta \in \mathbb{R}^n \setminus \{0\}$, for some positive universal constant γ , see e.g. page 100 in [DP05]. Thus, using this inequality with $\xi := \nabla u_s$ and $\zeta := \nabla v$ into (2.5), we obtain that

$$\int_{B_{r}(x_{0})} \left(|\nabla u(x)|^{p} - |\nabla v(x)|^{p} \right) dx$$

$$\geqslant \begin{cases} p \gamma \int_{0}^{1} \frac{1}{s} \left(\int_{B_{r}(x_{0})} |\nabla u_{s}(x) - \nabla v(x)|^{2} \left(|\nabla u_{s}(x)| + |\nabla v(x)| \right)^{p-2} dx \right) ds, & \text{if } 1$$

Furthermore, recalling (2.4), we conclude that

$$(2.6)$$

$$\int_{B_{r}(x_{0})} \left(\left| \nabla u(x) \right|^{p} - \left| \nabla v(x) \right|^{p} \right) dx$$

$$\geqslant \begin{cases} p \gamma \int_{0}^{1} s \left(\int_{B_{r}(x_{0})} \left| \nabla u(x) - \nabla v(x) \right|^{2} \left(\left| \nabla u_{s}(x) \right| + \left| \nabla v(x) \right| \right)^{p-2} dx \right) ds, & \text{if } 1$$

Now, if $1 , since <math>s \in [0, 1]$ we have that

$$|\nabla u_s| + |\nabla v| \leqslant s |\nabla u| + (1-s)|\nabla v| + |\nabla v| = s |\nabla u| + (2-s)|\nabla v| \leqslant 2(|\nabla u| + |\nabla v|).$$

Hence, plugging this information into (2.6), we find that

$$\int_{B_{r}(x_{0})} \left(\left| \nabla u(x) \right|^{p} - \left| \nabla v(x) \right|^{p} \right) dx$$

$$\geqslant p \gamma 2^{p-2} \int_{0}^{1} s \left(\int_{B_{r}(x_{0})} \left| \nabla u(x) - \nabla v(x) \right|^{2} \left(\left| \nabla u(x) \right| + \left| \nabla v(x) \right| \right)^{p-2} dx \right) ds$$

$$= p \gamma 2^{p-3} \int_{B_{r}(x_{0})} \left| \nabla u(x) - \nabla v(x) \right|^{2} \left(\left| \nabla u(x) \right| + \left| \nabla v(x) \right| \right)^{p-2} dx.$$

We now apply the Hölder's inequality with Hölder exponent 2/p and conjugate exponent

$$\left(\frac{2}{p}\right)' = \frac{2/p}{2/p-1} = \frac{2}{2-p},$$

to see that

$$\int_{B_{r}(x_{0})} |\nabla u(x) - \nabla v(x)|^{p} dx$$

$$= \int_{B_{r}(x_{0})} |\nabla u(x) - \nabla v(x)|^{p} (|\nabla u(x)| + |\nabla v(x)|)^{\frac{p(p-2)}{2}} (|\nabla u(x)| + |\nabla v(x)|)^{-\frac{p(p-2)}{2}} dx$$

$$\leq \left(\int_{B_{r}(x_{0})} |\nabla u(x) - \nabla v(x)|^{2} (|\nabla u(x)| + |\nabla v(x)|)^{p-2} dx \right)^{\frac{p}{2}}$$

$$\times \left(\int_{B_{r}(x_{0})} (|\nabla u(x)| + |\nabla v(x)|)^{p} dx \right)^{1-\frac{p}{2}}.$$

This and (2.7) yield that

$$\int_{B_{r}(x_{0})} \left|\nabla u(x) - \nabla v(x)\right|^{p} dx$$

$$\leq C \left(\int_{B_{r}(x_{0})} \left(\left|\nabla u(x)\right|^{p} - \left|\nabla v(x)\right|^{p}\right) dx\right)^{\frac{p}{2}} \left(\int_{B_{r}(x_{0})} \left(\left|\nabla u(x)\right| + \left|\nabla v(x)\right|\right)^{p} dx\right)^{1-\frac{p}{2}},$$

for some universal constant C > 0, which proves the desired result in (2.2) when 1 . $If instead <math>p \ge 2$, we deduce from (2.6) that

$$\int_{B_r(x_0)} \left(\left| \nabla u(x) \right|^p - \left| \nabla v(x) \right|^p \right) dx \geqslant p \gamma \int_0^1 s^{p-1} \left(\int_{B_r(x_0)} \left| \nabla u(x) - \nabla v(x) \right|^p dx \right) ds$$

$$= \gamma \int_{B_r(x_0)} \left| \nabla u(x) - \nabla v(x) \right|^p dx,$$

which establishes (2.3) when $p \ge 2$ and completes the proof of Lemma 2.2.

3. Some dichotomy results

With the preliminary work carried out so far, we can now provide a dichotomy statement. Roughly speaking, either the average of the energy of an almost minimizer decreases in a smaller ball, or the distance of its gradient and a suitable constant vector becomes as small as we wish (that is, linear functions are the "only ones for which the average does not improve in small balls"). The precise result goes as follows:

Proposition 3.1. There exists $\varepsilon_0 \in (0,1)$ such that for every $\varepsilon \in (0,\varepsilon_0)$ there exist $\eta \in (0,1)$, $M \geqslant 1$ and $\sigma_0 \in (0,1)$, depending on ε , n and p, such that if $\sigma \in [0,\sigma_0]$ and $a \geqslant M$ then the following statement holds true.

Let $u \in W^{1,p}(B_1)$ be such that

(3.1)
$$J_p(u, B_1) \leq (1 + \sigma)J_p(v, B_1)$$

for all $v \in W^{1,p}(B_1)$ such that v = u on ∂B_1 , with

(3.2)
$$a := \left(\int_{B_1} |\nabla u(x)|^p \, dx \right)^{1/p}.$$

Then, either

$$\left(\oint_{B_{\eta}} \left| \nabla u(x) \right|^p \, dx \right)^{1/p} \leqslant \frac{a}{2},$$

or

(3.4)
$$\left(\int_{B_{\eta}} |\nabla u(x) - q|^p \ dx \right)^{1/p} \leqslant \varepsilon a,$$

with $q \in \mathbb{R}^n$ such that

$$(3.5) \qquad \frac{a}{4} < |q| \leqslant C_0 a,$$

and $C_0 > 0$ universal.

Proof. Let v be the p-harmonic replacement of u in B_1 . By Theorem 3.19 in [MZ97] we have that, for every $x \in B_{1/2}$,

$$|\nabla v(x)|^p \le \sup_{B_{1/4}(x)} |\nabla v|^p \le C \int_{B_{1/2}(x)} |\nabla v(y)|^p \ dy \le C \int_{B_1} |\nabla u(y)|^p \ dy.$$

As a consequence, for all $x \in B_{1/2}$,

$$(3.6) |\nabla v(x)| \leqslant C_0 a,$$

for some positive universal constant C_0 .

Accordingly, we denote by $q := \nabla v(0)$ and we deduce from (3.6) that $|q| \leq C_0 a$. This together with Theorem 2 in [Man86], gives that, for all $\eta \in (0, 1/2]$,

(3.7)
$$\int_{B_{\eta}} |\nabla v(x) - q|^{p} dx \leqslant \int_{B_{\eta}} \left(C \left(\frac{\eta}{1/2} \right)^{\alpha} \|\nabla v\|_{L^{\infty}(B_{1/2})} \right)^{p} dx
\leqslant C \eta^{\alpha p} \|\nabla v\|_{L^{\infty}(B_{1/2})}^{p}
\leqslant C_{1} \eta^{\alpha p} a^{p},$$

for some $\alpha \in (0,1)$ and $C_1 > 0$ universal.

Furthermore, using (3.1) we have that

$$\int_{B_{1}} \left(\left| \nabla u(x) \right|^{p} - \left| \nabla v(x) \right|^{p} \right) dx \leqslant J_{p}(u, B_{1}) - \int_{B_{1}} \left| \nabla v(x) \right|^{p} dx
\leqslant (1 + \sigma) J_{p}(v, B_{1}) - \int_{B_{1}} \left| \nabla v(x) \right|^{p} dx
\leqslant C \left(\sigma \int_{B_{1}} \left| \nabla v(x) \right|^{p} dx + 1 \right)
\leqslant C \left(\sigma \int_{B_{1}} \left| \nabla u(x) \right|^{p} dx + 1 \right)$$

We distinguish two cases, $p \ge 2$ and 1 .

If $p \ge 2$, then from (2.3) and (3.8) we deduce that

$$\int_{B_1} |\nabla u(x) - \nabla v(x)|^p dx \le C \left(\sigma \int_{B_1} |\nabla u(x)|^p dx + 1\right).$$

Taking the average over B_1 , we thereby obtain that

$$\int_{B_1} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C(\sigma a^p + 1).$$

Using this and (3.7), we get that

$$\int_{B_{\eta}} |\nabla u(x) - q|^{p} dx \leq 2^{p-1} \int_{B_{\eta}} \left(|\nabla u(x) - \nabla v(x)|^{p} + |\nabla v(x) - q|^{p} \right) dx
\leq 2^{p-1} \left(\frac{|B_{1}|}{|B_{\eta}|} \int_{B_{1}} |\nabla u(x) - \nabla v(x)|^{p} dx + C_{1} \eta^{\alpha p} a^{p} \right)
\leq 2^{p-1} \left(C \eta^{-n} (\sigma a^{p} + 1) + C_{1} a^{p} \eta^{\alpha p} \right),
= 2^{p-1} C \eta^{-n} \sigma a^{p} + 2^{p-1} C \eta^{-n} + 2^{p-1} C_{1} a^{p} \eta^{\alpha p}.$$

This yields that

$$(3.10) \qquad \int_{B_n} |\nabla u(x)|^p \, dx \leqslant 2^{2(p-1)} C \eta^{-n} \sigma a^p + 2^{2(p-1)} C \eta^{-n} + 2^{2(p-1)} C_1 a^p \eta^{\alpha p} + 2^{p-1} |q|^p.$$

If instead 1 , by virtue of (2.2) and (3.8), we obtain that

$$\int_{B_1} |\nabla u(x) - \nabla v(x)|^p dx$$

$$\leq C \left(\sigma \int_{B_1} |\nabla u(x)|^p dx + 1\right)^{\frac{p}{2}} \left(\int_{B_1} \left(|\nabla u(x)| + |\nabla v(x)| \right)^p dx \right)^{1-\frac{p}{2}}$$

$$\leq C \left(\sigma^{\frac{p}{2}} \left(\int_{B_1} |\nabla u(x)|^p dx \right)^{\frac{p}{2}} + 1\right) \left(\int_{B_1} \left(|\nabla u(x)|^p + |\nabla v(x)|^p \right) dx \right)^{1-\frac{p}{2}}.$$

Thus, since v is the p-harmonic replacement of u in B_1 ,

$$\int_{B_1} |\nabla u(x) - \nabla v(x)|^p dx$$

$$\leq C \left(\sigma^{\frac{p}{2}} \left(\int_{B_1} |\nabla u(x)|^p dx \right)^{\frac{p}{2}} + 1 \right) \left(\int_{B_1} |\nabla u(x)|^p dx \right)^{1-\frac{p}{2}}$$

$$= C \left(\sigma^{\frac{p}{2}} \int_{B_1} |\nabla u(x)|^p dx + \left(\int_{B_1} |\nabla u(x)|^p dx \right)^{1-\frac{p}{2}} \right).$$

Consequently, taking the average integral, we have that

$$\int_{B_1} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C \left(\sigma^{\frac{p}{2}} a^p + |B_1|^{-\frac{p}{2}} a^{p(1-\frac{p}{2})} \right) \leqslant C \left(\sigma^{\frac{p}{2}} a^p + a^{p(1-\frac{p}{2})} \right).$$

From this and (3.7), we obtain that

$$\int_{B_{\eta}} |\nabla u(x) - q|^{p} dx \leq 2^{p-1} \int_{B_{\eta}} \left(|\nabla u(x) - \nabla v(x)|^{p} + |\nabla v(x) - q|^{p} \right) dx
\leq 2^{p-1} \left(\frac{|B_{1}|}{|B_{\eta}|} \int_{B_{1}} |\nabla u(x) - \nabla v(x)|^{p} dx + C_{1} \eta^{\alpha p} a^{p} \right)
\leq 2^{p-1} C \eta^{-n} \sigma^{\frac{p}{2}} a^{p} + 2^{p-1} C \eta^{-n} a^{p \left(1 - \frac{p}{2}\right)} + 2^{p-1} C_{1} a^{p} \eta^{\alpha p}.$$

This gives that

$$(3.12) \int_{B_{\eta}} |\nabla u(x)|^p dx \leq 2^{2(p-1)} C \eta^{-n} \sigma^{\frac{p}{2}} a^p + 2^{2(p-1)} C \eta^{-n} a^{p(1-\frac{p}{2})} + 2^{2(p-1)} C_1 a^p \eta^{\alpha p} + 2^{p-1} |q|^p.$$

Now, given $\varepsilon_0 \in (0, 1/4]$, we claim that for every $\varepsilon \in (0, \varepsilon_0)$ there exists η small enough (depending on ε) such that if σ is chosen sufficiently small and a sufficiently large (depending on η , and thus on ε) then

$$\begin{cases}
2^{2(p-1)}C\eta^{-n}\sigma a^{p} + 2^{2(p-1)}C\eta^{-n} + 2^{2(p-1)}C_{1}a^{p}\eta^{\alpha p} \leqslant 2^{p-1}\varepsilon^{p}a^{p} \leqslant \frac{a^{p}}{2^{p+1}} \\
\text{if } p \geqslant 2, \\
2^{2(p-1)}C\eta^{-n}\sigma^{\frac{p}{2}}a^{p} + 2^{2(p-1)}C\eta^{-n}a^{p\left(1-\frac{p}{2}\right)} + 2^{2(p-1)}C_{1}a^{p}\eta^{\alpha p} \leqslant 2^{p-1}\varepsilon^{p}a^{p} \leqslant \frac{a^{p}}{2^{p+1}} \\
\text{if } 1$$

To prove this we distinguish two cases. If $p \ge 2$, we pick $\eta > 0$ sufficiently small such that $\varepsilon^p - 2^{p-1}C\eta > 2^{p-1}C_1\eta^{\alpha p}$. This allows us to define

$$M := \left(\frac{2^{p-1}C\eta^{-n}}{\varepsilon^p - 2^{p-1}C\eta - 2^{p-1}C_1\eta^{\alpha p}}\right)^{1/p}.$$

Note also that we can suppose $M \ge 1$ by taking η small enough. Let also

$$\sigma_0 := \eta^{n+1}$$

With this setting, we obtain that, for every $a \ge M$ and for every $0 < \sigma \le \sigma_0$,

$$\begin{split} & 2^{2(p-1)}C\eta^{-n}\sigma a^p + 2^{2(p-1)}C\eta^{-n} + 2^{2(p-1)}C_1a^p\eta^{\alpha p} \\ \leqslant & a^p \Big(2^{2(p-1)}C\eta + 2^{2(p-1)}C_1\eta^{\alpha p} \Big) + 2^{2(p-1)}C\eta^{-n} \\ = & a^p \Big(2^{2(p-1)}C\eta + 2^{2(p-1)}C_1\eta^{\alpha p} \Big) + 2^{p-1} \, M^p \Big(\varepsilon^p - 2^{p-1}C\eta - 2^{p-1}C_1\eta^{\alpha p} \Big) \\ \leqslant & a^p \Big(2^{2(p-1)}C\eta + 2^{2(p-1)}C_1\eta^{\alpha p} \Big) + 2^{p-1} \, a^p \Big(\varepsilon^p - 2^{p-1}C\eta - 2^{p-1}C_1\eta^{\alpha p} \Big) \\ = & 2^{p-1}\varepsilon^p a^p + a^p \Big(2^{2(p-1)}C\eta + 2^{2(p-1)}C_1\eta^{\alpha p} - 2^{2(p-1)}C\eta - 2^{2(p-1)}C_1\eta^{\alpha p} \Big) \\ = & 2^{p-1}\varepsilon^p a^p, \end{split}$$

which proves (3.13) when $p \ge 2$.

If instead $1 , we pick <math>\eta > 0$ small enough such that $\varepsilon^p > 2^{p-1}C\eta + 2^{p-1}C_1\eta^{\alpha p}$. In this way, we can define

$$M := \left(\frac{C2^{p-1}\eta^{-n}}{\varepsilon^p - 2^{p-1}C\eta - 2^{p-1}C_1\eta^{\alpha p}}\right)^{2/p^2}.$$

Let also $\sigma_0 := \eta^{(n+1)\frac{2}{p}}$. Then, whenever $a \geqslant M$ and $0 < \sigma \leqslant \sigma_0$ it follows that

$$2^{2(p-1)}C\eta^{-n}\sigma^{\frac{p}{2}}a^{p} + 2^{2(p-1)}C\eta^{-n}a^{p\left(1-\frac{p}{2}\right)} + 2^{2(p-1)}C_{1}a^{p}\eta^{\alpha p}$$

$$\leqslant 2^{2(p-1)}C\eta a^{p} + 2^{2(p-1)}C\eta^{-n}a^{p\left(1-\frac{p}{2}\right)} + 2^{2(p-1)}C_{1}a^{p}\eta^{\alpha p}$$

$$= 2^{p-1}a^{p}\left(2^{p-1}C\eta + 2^{p-1}C\eta^{-n}a^{-\frac{p^{2}}{2}} + 2^{p-1}C_{1}\eta^{\alpha p}\right)$$

$$\leqslant 2^{p-1}a^{p}\left(2^{p-1}C\eta + 2^{p-1}C\eta^{-n}M^{-\frac{p^{2}}{2}} + 2^{p-1}C_{1}\eta^{\alpha p}\right)$$

$$= 2^{p-1}a^{p}\left(2^{p-1}C\eta + \varepsilon^{p} - 2^{p-1}C\eta - 2^{p-1}C_{1}\eta^{\alpha p} + 2^{p-1}C_{1}\eta^{\alpha p}\right)$$

$$= 2^{p-1}\varepsilon^{p}a^{p},$$

which establishes (3.13) in the case 1 as well.

In order to complete the proof of Proposition 3.1, we now distinguish two cases according to the size of |q|. More precisely, we first suppose that

$$|q| \leqslant \frac{a}{4}.$$

Then, we use either (3.10) (if $p \ge 2$) or (3.12) (if 1), and (3.13) to conclude that

$$\int_{B_n} |\nabla u(x)|^p dx \leqslant \frac{a^p}{2^{p+1}} + 2^{p-1} \frac{a^p}{2^{2p}} = \frac{a^p}{2^{p+1}} + \frac{a^p}{2^{p+1}} = \frac{a^p}{2^p},$$

and thus

$$\left(\oint_{B_{\eta}} |\nabla u(x)|^p dx \right)^{1/p} \leqslant \frac{a}{2},$$

which is the first alternative in (3.3).

Otherwise, it holds that

$$\frac{a}{4} < |q| \leqslant C_0 a,$$

and therefore, by either (3.9) (if $p \ge 2$) or (3.11) (if 1), and (3.13), we have that

$$\left(\int_{B_{\eta}} |\nabla u(x) - q|^p \ dx \right)^{1/p} \leqslant \varepsilon a,$$

which is the second alternative in (3.4). The proof of Proposition 3.1 is thereby complete. \square

We will now show that the alternative in (3.4) can be "improved" when ε and σ are sufficiently small. This result is the counterpart of Lemma 2.3 in [DSS20b] in the more general setting dealt with in this paper. The main difficulty here with respect to Lemma 2.3 in [DSS20b] relies on the fact that the problem is not linear when $p \neq 2$, and therefore, if v_1 and v_2 are the p-harmonic replacements of u_1 and u_2 in B_1 , then it is not true that $v_1 + v_2$ is the p-harmonic replacement of $u_1 + u_2$, unless p = 2.

We will overcome this difficulty by showing that a "uniformly elliptic" equation is still satisfied from the sum of p-harmonic replacements. More precisely, we will show the following result:

Lemma 3.2. Let $q \in \mathbb{R}^n$ and let $\eta : \mathbb{R}^n \to \mathbb{R}^n$ be such that $|\eta(x)| < \frac{|q|}{2}$. Let $F(z) := |z|^{p-2}z$ and

$$A(x) := \int_0^1 DF(q + t\eta(x)) dt.$$

Then,

$$\lambda |q|^{p-2} \leqslant A\xi \cdot \xi \leqslant \Lambda |q|^{p-2}$$

for all $\xi \in \partial B_1$, for some $\Lambda \geqslant \lambda > 0$, depending on p.

Proof. We notice that, for all $t \in (0, 1)$,

(3.14)
$$|q + t\eta| \le |q| + |\eta| < |q| + \frac{|q|}{2} = \frac{3|q|}{2}$$
 and
$$|q + t\eta| \ge |q| - |\eta| > |q| - \frac{|q|}{2} = \frac{|q|}{2}.$$

Furthermore, we observe that

$$DF(z) = (p-2)|z|^{p-4}z \otimes z + |z|^{p-2}Id,$$

where Id denotes the identity matrix in dimension n, and therefore, for all $\xi \in \partial B_1$,

$$DF(z)\xi \cdot \xi = (p-2)|z|^{p-4}(z \cdot \xi)^2 + |z|^{p-2}|\xi|^2$$

$$\geqslant -(2-p)^+|z|^{p-2} + |z|^{p-2}|\xi|^2$$

$$= \left[1 - (2-p)^+\right]|z|^{p-2}|\xi|^2$$

$$= \left[1 - (2-p)^+\right]|z|^{p-2}.$$
12

Consequently,

$$A\xi \cdot \xi \geqslant \left[1 - (2-p)^{+}\right] \int_{0}^{1} |q + t\eta|^{p-2} dt.$$

Accordingly, we can use (3.14) to find that

$$A\xi \cdot \xi \geqslant \frac{1 - (2 - p)^{+}}{2^{p-2}} |q|^{p-2}$$

if $p \ge 2$, and

$$A\xi \cdot \xi \geqslant \left[1 - (2 - p)^{+}\right] \left(\frac{3}{2}\right)^{p-2} |q|^{p-2}$$

if $p \in (1, 2)$.

Similarly, since

$$DF(z)\xi \cdot \xi \leqslant (p-2)^{+}|z|^{p-2} + |z|^{p-2}|\xi|^{2}$$

$$= [1 + (p-2)^{+}]|z|^{p-2}|\xi|^{2}$$

$$= [1 + (p-2)^{+}]|z|^{p-2},$$

we have that

$$A\xi \cdot \xi \le [1 + (p-2)^+] \int_0^1 |q + t\eta|^{p-2} dt.$$

Hence, making again use of (3.14),

$$A\xi \cdot \xi \le \left[1 + (p-2)^+\right] \left(\frac{3}{2}\right)^{p-2} |q|^{p-2}$$

if $p \ge 2$, and

$$A\xi \cdot \xi \leqslant \frac{1 + (p-2)^+}{2^{p-2}} |q|^{p-2}$$

if $p \in (1, 2)$.

These considerations complete the proof of Lemma 3.2.

With this, we can now state the following result:

Lemma 3.3. Let $a_1 > a_0 > 0$ and

$$(3.15) p > \max\left\{\frac{2n}{n+2}, 1\right\}.$$

There exist $\alpha_0 \in (0,1]$ and $\widetilde{C} > 0$, depending on n and p, such that for every $\alpha \in (0,\alpha_0)$ there exist

- $\rho \in (0,1)$, depending on n, p and α ,
- $\varepsilon_0 \in (0,1)$, depending on n, p, a_0 , a_1 and α ,
- $c_0 > 0$, depending on n, p, a_0 , a_1 and α ,

such that, if $\varepsilon \in (0, \varepsilon_0]$ and $\sigma \in (0, c_0 \varepsilon^P]$, with $P := \max\{p, 2\}$, then the following statement holds true.

Let $u \in W^{1,p}(B_1)$ be such that $u \ge 0$ in B_1 and

$$(3.16) J_p(u, B_1) \leq (1 + \sigma)J_p(v, B_1)$$

for all $v \in W^{1,p}(B_1)$ such that v = u on ∂B_1 .

Let

$$(3.17) a := \left(\int_{B_1} |\nabla u(x)|^p dx \right)^{1/p}$$

and suppose that

$$(3.18) a \in [a_0, a_1].$$

Assume also that

(3.19)
$$\left(\int_{B_1} |\nabla u(x) - q|^p \ dx \right)^{1/p} \leqslant \varepsilon a,$$

for some $q \in \mathbb{R}^n$ such that

$$(3.20) \frac{a}{8} < |q| \leqslant 2C_0 a,$$

where $C_0 > 0$ is the universal constant given by Proposition 3.1. Then

(3.21)
$$\left(\int_{B_{\rho}} |\nabla u(x) - \widetilde{q}|^p dx \right)^{1/p} \leqslant \rho^{\alpha} \varepsilon a,$$

with $\widetilde{q} \in \mathbb{R}^n$ such that

$$(3.22) |q - \widetilde{q}| \leqslant \widetilde{C}\varepsilon a.$$

The quantity α_0 in the statement of Lemma 3.3 measures the oscillation of ∇u , see [Ura68], [DiB83], [Tol84] or, specifically for our case, Theorem 2 in [Man86]. When p=2, we have that $\alpha_0 = 1$, according to the full regularity of harmonic functions, but when $p \neq 2$ in general we only know that $\alpha_0 \in (0,1)$ (but when n=2 a sharp regularity exponent has been determined in [IM89]).

Proof of Lemma 3.3. Let \bar{v} denote the p-harmonic replacement of u in $B_{9/10}$ and let v be defined as

(3.23)
$$v := \begin{cases} \bar{v} & \text{in } B_{9/10}, \\ u & \text{in } B_1 \setminus B_{9/10}. \end{cases}$$

Then, since v = u on ∂B_1 and $v \in W^{1,p}(B_1)$, we can use (3.16) to see that

$$J_p(u, B_1) \leqslant (1 + \sigma)J_p(v, B_1).$$

Consequently,

$$J_{p}(u, B_{9/10}) = J_{p}(u, B_{1}) - J_{p}(u, B_{1} \setminus B_{9/10})$$

$$\leq J_{p}(v, B_{9/10}) + J_{p}(v, B_{1} \setminus B_{9/10}) + \sigma J_{p}(v, B_{1}) - J_{p}(u, B_{1} \setminus B_{9/10})$$

$$= J_{p}(v, B_{9/10}) + J_{p}(u, B_{1} \setminus B_{9/10}) + \sigma J_{p}(v, B_{1}) - J_{p}(u, B_{1} \setminus B_{9/10})$$

$$= J_{p}(v, B_{9/10}) + \sigma J_{p}(v, B_{1}).$$

Hence, recalling the definition of J_p in (1.1),

$$\int_{B_{9/10}} |\nabla u(x)|^p dx + \left| \{u > 0\} \cap B_{9/10} \right| \le \int_{B_{9/10}} |\nabla v(x)|^p dx + \left| B_{9/10} \right| + \sigma J_p(v, B_1),$$

which yields that

(3.24)
$$\int_{B_{9/10}} \left(\left| \nabla u(x) \right|^p - \left| \nabla v(x) \right|^p \right) dx \leqslant \left| \left\{ u = 0 \right\} \cap B_{9/10} \right| + \sigma J_p(v, B_1).$$

Moreover, recalling (3.23), we point out that v is the p-harmonic replacement of u in $B_{9/10}$, and therefore

$$J_{p}(v, B_{1}) = \int_{B_{1}} \left(|\nabla v(x)|^{p} + \chi_{\{v>0\}}(x) \right) dx$$

$$\leq \int_{B_{9/10}} |\nabla v(x)|^{p} dx + \int_{B_{1} \setminus B_{9/10}} |\nabla v(x)|^{p} dx + |B_{1}|$$

$$\leq \int_{B_{1}} |\nabla u|^{p} dx + |B_{1}| \leq a^{p} + |B_{1}|.$$

Furthermore, if $p \ge 2$, by (2.3) and (3.24) we deduce that

$$\int_{B_{9/10}} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C |\{u = 0\} \cap B_{9/10}| + C\sigma J_p(v, B_1),$$

for some positive universal constant C. Consequently, exploiting (3.25), we obtain that

(3.26)
$$\int_{B_{9/10}} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C |\{u = 0\} \cap B_{9/10}| + C\sigma(a^p + 1),$$

up to renaming C.

If instead p < 2, we use (2.2), (3.24) and (3.25) to write that

$$\int_{B_{9/10}} |\nabla u(x) - \nabla v(x)|^{p} dx$$

$$\leq C \left(\int_{B_{9/10}} (|\nabla u(x)|^{p} - |\nabla v(x)|^{p}) dx \right)^{\frac{p}{2}} \left(\int_{B_{9/10}} (|\nabla u(x)| + |\nabla v(x)|)^{p} dx \right)^{1-\frac{p}{2}}$$

$$\leq C \left(\left| \{u = 0\} \cap B_{9/10} \right| + \sigma J_{p}(v, B_{1}) \right)^{\frac{p}{2}} \left(\int_{B_{9/10}} 2^{p} |\nabla u(x)|^{p} dx \right)^{1-\frac{p}{2}}$$

$$\leq C \left(\left| \{u = 0\} \cap B_{9/10} \right| + \sigma (a^{p} + 1) \right)^{\frac{p}{2}} a^{p(1-\frac{p}{2})}.$$

Now we claim that

$$(3.28) |B_{9/10} \cap \{u = 0\}| \leqslant C_1 \varepsilon^{p+\delta},$$

for some $C_1 > 0$ and $\delta > 0$. To prove this, we consider the linear function

(3.29)
$$\ell(x) := b + q \cdot x, \quad \text{with} \quad b := \int_{B_1} u(x) \, dx.$$

We remark that

$$\int_{B_1} (u(x) - \ell(x)) dx = b - \left(b + \int_{B_1} q \cdot x \, dx\right) = -\int_{B_1} q \cdot x \, dx = 0.$$

As a consequence, denoting by

$$(u - \ell)_{B_1} := \int_{B_1} (u(x) - \ell(x)) dx$$

and using the Poincaré inequality, we see that

$$||u - \ell - (u - \ell)_{B_1}||_{L^p(B_1)} = ||u - \ell||_{L^p(B_1)} \leqslant C ||\nabla (u - \ell)||_{L^p(B_1)},$$

for some C > 0 universal.

This and (3.19) entail that

(3.30)
$$\int_{B_1} |u(x) - \ell(x)|^p dx \leqslant C \int_{B_1} |\nabla (u - \ell)(x)|^p dx \leqslant C\varepsilon^p a^p.$$

We also observe that, since $u \ge 0$, it holds that $\ell^- \le |u - \ell|$. Using this information into (3.30), we obtain that

(3.31)
$$\int_{B_1} (\ell^-(x))^p \, dx \leqslant C\varepsilon^p a^p.$$

Now we claim that, if ε is sufficiently small,

(3.32)
$$\ell \geqslant c_1 a \text{ in } B_{9/10},$$

for some $c_1 > 0$. To check this, we argue by contradiction assuming that

$$\min_{x \in \overline{B_{9/10}}} \ell(x) < ca$$

for all c > 0. We notice that, for every $x \in B_{9/10}$,

$$-\frac{9|q|}{10} \leqslant \ell(x) - b \leqslant \frac{9|q|}{10}.$$

and therefore

$$ca > \min_{x \in \overline{B_{1/2}}} \ell(x) \geqslant b - \frac{9|q|}{10}.$$

This leads to

$$(3.33) b \leqslant ca + \frac{9|q|}{10}.$$

Now we define

$$\mathcal{B} := \left\{ x = -\frac{tq}{|q|} + \eta, \text{ for some } t \in \left[\frac{38}{40}, \frac{39}{40} \right] \text{ and } \eta \in B_{1/40} \right\}.$$

Notice that if $x \in \mathcal{B}$ then

$$|x| \le t + |\eta| < \frac{39}{40} + \frac{1}{40} = 1,$$

and this shows that

$$(3.34) \mathcal{B} \subseteq B_1.$$

Furthermore, from (3.20) and (3.33) we see that, if $x \in \mathcal{B}$,

$$\ell(x) = b - t|q| + \eta \cdot q \leqslant ca + \frac{9|q|}{10} - \frac{38|q|}{40} + \frac{|q|}{40} = ca - \left(\frac{38}{40} - \frac{9}{10} - \frac{1}{40}\right)|q|$$
$$= ca - \frac{|q|}{40} < ca - \frac{a}{320}.$$

Hence, taking $c \in (0, \frac{1}{640})$, we infer that

$$\ell(x) \leqslant -\frac{a}{640}.$$

Accordingly, using this and (3.34) into (3.31), we obtain that

$$C|B_1|\varepsilon^p a^p \geqslant \int_{B_1} (\ell^-(x))^p dx \geqslant \int_{\mathcal{B}} (\ell^-(x))^p dx \geqslant \int_{\mathcal{B}} \left(\frac{a}{80 \cdot 2^{(3p+1)/p}}\right)^p dx \geqslant \overline{c} a^p,$$

for some positive universal constant \overline{c} . This establishes the desired contradiction if ε is sufficiently small, and thus the proof of (3.32) is complete.

Now, we distinguish the three following cases: $p \in (1, n)$, p = n and p > n.

Firstly, if p < n, recalling the Poincaré-Sobolev inequality (see e.g. Theorem 2 on page 265 of [Eva98], or Theorem 7.30 in [GT01]) and applying the same argument of formula (7.40) in [GT01], we obtain that

(3.35)
$$\left(\int_{B_1} |u(x) - \ell(x)|^{p^*} dx \right)^{1/p^*} \leq C \left(\int_{B_1} |\nabla (u(x) - \ell(x))|^p dx \right)^{1/p}$$

for some C > 0 universal, where

$$p^* := \frac{np}{n-p}.$$

Then, by virtue of (3.19), (3.29) and (3.35), we get that

$$\left(\int_{B_1} |u(x) - \ell(x)|^{p^*} dx\right)^{1/p^*} \leqslant C \left(\int_{B_1} |\nabla (u(x) - \ell(x))|^p dx\right)^{1/p}$$

$$= C \left(\int_{B_1} |\nabla u(x) - q|^p dx\right)^{1/p} \leqslant C\varepsilon a.$$

This and (3.32) entail that

$$C\varepsilon a \geqslant \left(\int_{B_{9/10}\cap\{u=0\}} |\ell(x)|^{p^*} dx\right)^{1/p^*} \geqslant c_1 a \left|B_{9/10}\cap\{u=0\}\right|^{1/p^*},$$

and thus, up to renaming constants,

$$(3.36) |B_{9/10} \cap \{u=0\}| \leqslant C\varepsilon^{p^*}.$$

Now we notice that

$$p^* = \frac{np + p^2 - p^2}{n - p} = p + \frac{p^2}{n - p}.$$

Therefore, setting

$$\delta := \frac{p^2}{n-p} > 0,$$

we obtain (3.28) from (3.36) in the case p < n.

If instead p > n, we recall Theorem 7.17 as well as (7.12) and (7.16) in [GT01], to see that

$$\sup_{B_1} |u - \ell| \leqslant C \left(\int_{B_1} \left| \nabla \left(u(x) - \ell(x) \right) \right|^p dx \right)^{1/p}$$

for some C>0 universal. From this, and making again use of (3.19) and (3.29), we deduce that

$$\sup_{B_1} |u - \ell| \leqslant C \left(\int_{B_1} |\nabla (u(x) - \ell(x))|^p dx \right)^{1/p} \leqslant C\varepsilon a.$$

As a consequence of this and (3.32), for all $x \in B_{9/10}$,

$$C\varepsilon a \geqslant \sup_{B_1} |u - \ell| \geqslant \sup_{B_{9/10}} (\ell - u) \geqslant \ell(x) - u(x) \geqslant c_1 a - u(x),$$

and thus $u(x) \ge c_1 a - C\varepsilon a > 0$ for all $x \in B_{9/10}$, as long as ε is sufficiently small. Accordingly, it follows that

$$(3.38) |B_{9/10} \cap \{u = 0\}| = 0,$$

and this gives (3.28) in the case p > n as well.

It remains to analyze the case p = n. For this purpose, we point out that $u - \ell \in W^{1,n}(B_1)$, and so we can apply Theorem 7.35 and Theorem 7.16 in [GT01] to obtain that

$$\int_{B_1} \exp\left(\frac{|u(x) - \ell(x)|}{C_2 \|\nabla(u - \ell)\|_{L^n(B_1)}}\right)^{\frac{n}{n-1}} dx \leqslant C_3 |B_1|,$$

where C_2 and C_3 are positive universal constants.

Also, for all $x \in B_1$, we have that $e^x \ge C_4 x^n$ for some $C_4 > 0$, and thus, recalling (3.19) and (3.29),

(3.39)
$$\int_{B_1} |u(x) - \ell(x)|^{\frac{n^2}{n-1}} dx \leqslant C \|\nabla u - q\|_{L^n(B_1)}^{\frac{n^2}{n-1}} \leqslant C\varepsilon^{\frac{n^2}{n-1}} a^{\frac{n^2}{n-1}},$$

for some C > 0, from which we deduce that, using (3.32) and (3.39),

$$C\varepsilon^{\frac{n^2}{n-1}}a^{\frac{n^2}{n-1}} \geqslant \int_{B_1} |u(x) - \ell(x)|^{\frac{n^2}{n-1}} dx \geqslant \int_{B_{9/10} \cap \{u=0\}} |\ell(x)|^{\frac{n^2}{n-1}} dx$$
$$\geqslant c_1^{\frac{n^2}{n-1}}a^{\frac{n^2}{n-1}}|B_{9/10} \cap \{u=0\}|.$$

We point out that

$$\frac{n^2}{n-1} = \frac{n^2 - n + n}{n-1} = n + \frac{n}{n-1},$$

and therefore, choosing

$$\delta := \frac{n}{n-1} > 0,$$

we establish (3.28) when p = n as well.

Gathering together (3.26) and (3.28) we thus conclude that, if $p \ge 2$,

(3.40)
$$\int_{B_{9/10}} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C_1 \varepsilon^{p+\delta} + C\sigma(a^p + 1).$$

Instead, when 1 , we make use of (3.27) and (3.28) and we obtain that

$$(3.41) \qquad \int_{B_{9/10}} |\nabla u(x) - \nabla v(x)|^p dx \leqslant C \left(C_1 \varepsilon^{p+\delta} + \sigma(a^p + 1) \right)^{\frac{p}{2}} a^{p\left(1 - \frac{p}{2}\right)}.$$

Our aim is now to show that, even if $v - q \cdot x$ is not the p-harmonic replacement of $u - q \cdot x$, as in the classical case (see Lemma 2.3 in [DSS20b]), it satisfies a "nice" equation in $B_{9/10}$.

To this end, we observe that if $p \ge 2$, using (3.19) and (3.40),

$$(3.42) \int_{B_{9/10}} |\nabla v(x) - q|^p dx \leq 2^{p-1} \left(\int_{B_{9/10}} |\nabla u(x) - q|^p dx + \int_{B_{9/10}} |\nabla v(x) - \nabla u(x)|^p dx \right)$$

$$\leq C_2 \varepsilon^p a^p + C_1 \varepsilon^{p+\delta} + C \sigma(a^p + 1),$$

up to renaming constants.

If instead 1 , from (3.19) and (3.41) we have that

$$(3.43) \int_{B_{9/10}} |\nabla v(x) - q|^p dx \leq 2^{p-1} \left(\int_{B_{9/10}} |\nabla u(x) - q|^p dx + \int_{B_{9/10}} |\nabla v(x) - \nabla u(x)|^p dx \right)$$

$$\leq C_2 \varepsilon^p a^p + C \left(C_1 \varepsilon^{p+\delta} + \sigma(a^p + 1) \right)^{\frac{p}{2}} a^{p(1 - \frac{p}{2})}.$$

Suppose now that $p \ge 2$ and $\sigma \le c_0 \varepsilon^p$, with c_0 to be made precise later. Recalling (3.18) and (3.42), we infer that

$$(3.44) \qquad \int_{B_{9/10}} |\nabla v(x) - q|^p \ dx \leqslant C_2 \varepsilon^p a^p + C_1 \varepsilon^{p+\delta} + C c_0 \varepsilon^p (a^p + 1) \leqslant C \varepsilon^p a^p,$$

up to relabeling C > 0.

Similarly, if $1 , we take <math>\sigma \leq c_0 \varepsilon^2$, with c_0 to be made precise later. In this case, we deduce from (3.43) that

(3.45)
$$\int_{B_{9/10}} |\nabla v(x) - q|^p dx \leqslant C_2 \varepsilon^p a^p + C \left(C_1 \varepsilon^{p+\delta} + c_0 \varepsilon^2 (a^p + 1) \right)^{\frac{p}{2}} a^{p\left(1 - \frac{p}{2}\right)}$$
$$\leqslant C_1 \varepsilon^{(p+\delta)\frac{p}{2}} a^p + C_2 \varepsilon^p a^p,$$

up to renaming the constants.

We stress that we have not used the assumption on p given in (3.15) so far. We will use it next to reabsorb the term $\varepsilon^{(p+\delta)\frac{p}{2}}a^p$ appearing in (3.45) into the term ε^pa^p . More precisely, we point out that

(3.46) if
$$p \in (\max\{\frac{2n}{n+2}, 1\}, 2)$$
, then $(p+\delta)\frac{p}{2} > p$.

Indeed, when n=1 we are in the situation in which n=1 < p, and therefore we can take $\delta=1$ in (3.28) (recall also (3.38)), thus obtaining that $\frac{p+1}{2} > 1$, which gives (3.46) in this case.

Also, if $n \ge 2$, then we are in the case p < n, and thus, recalling the value of δ given in (3.37), we obtain that

$$p + \delta = p + \frac{p^2}{n - p} = \frac{np}{n - p} > 2,$$

which gives (3.46).

As a consequence of (3.45) and (3.46), we obtain that, if $p \in (\max\{\frac{2n}{n+2}, 1\}, 2)$,

$$\int_{B_{9/10}} |\nabla v(x) - q|^p \ dx \leqslant C\varepsilon^p a^p,$$

for some C > 0.

Putting together this and (3.44), we conclude that, if $p > \max\left\{\frac{2n}{n+2}, 1\right\}$,

(3.47)
$$\int_{B_{9/10}} |\nabla v(x) - q|^p \ dx \leqslant C\varepsilon^p a^p.$$

Furthermore, we claim that, if ε is sufficiently small, for all $x \in B_{1/2}$,

$$(3.48) |\nabla v(x) - q| \leqslant C(\varepsilon a)^{\nu},$$

for some C>0 and $\nu\in(0,1)$. Indeed, suppose by contradiction that

$$\max_{x \in \overline{B_{1/2}}} |\nabla v(x) - q| > C(\varepsilon a)^{\nu}$$

for all C > 0 and $\nu \in (0,1)$. We recall that v is p-harmonic in $B_{9/10}$ and thus the maximum in (3.49) is achieved, see e.g. [DiB83] or [Tol84]. Moreover, let $\overline{x} \in \overline{B_{1/2}}$ be such that

$$(3.50) |\nabla v(\overline{x}) - q| = \max_{x \in \overline{B_{1/2}}} |\nabla v(x) - q| > C(\varepsilon a)^{\nu}.$$

Then, for all $x \in B_{1/8}(\overline{x})$,

$$\left|\nabla v(x) - \nabla v(\overline{x})\right| \leqslant \overline{C}|x - \overline{x}|^{\alpha},$$

for some $\overline{C} > 0$ and $\alpha \in (0,1]$, with α depending only on n and p, and \overline{C} bounded by a constant depending only on n and p times $\|\nabla v\|_{L^{\infty}(B_{1/4}(\overline{x}))}$, see Theorem 2 in [Man86]. Actually, since, by formula (3.44) in [MZ97],

$$\|\nabla v\|_{L^{\infty}(B_{1/4}(\overline{x}))} \leqslant \|\nabla v\|_{L^{\infty}(B_{3/4})} \leqslant \hat{C}\|\nabla v\|_{L^{p}(B_{9/10})} \leqslant \hat{C}\|\nabla u\|_{L^{p}(B_{9/10})} \leqslant \tilde{C}a_{1},$$

for some \hat{C} and $\widetilde{C} > 0$ depending only on n and p, we infer that \overline{C} only depends on n, p and a_1 . As a consequence of this and (3.50), for all $x \in B_{\left(\frac{C(\varepsilon a)^{\nu}}{2\overline{C}}\right)^{1/\alpha}}(\overline{x})$,

$$|\nabla v(x) - q| \geqslant |\nabla v(\overline{x}) - q| - \left| (\nabla v(x) - q) - (\nabla v(\overline{x}) - q) \right|$$
$$\geqslant C(\varepsilon a)^{\nu} - \overline{C}|x - \overline{x}|^{\alpha} \geqslant C(\varepsilon a)^{\nu} - \frac{C(\varepsilon a)^{\nu}}{2} = \frac{C(\varepsilon a)^{\nu}}{2}.$$

Notice also that, if ε is sufficiently small, then $B_{\left(\frac{C(\varepsilon a)^{\nu}}{2\overline{C}}\right)^{1/\alpha}}(\overline{x}) \subseteq B_{9/10}$. Accordingly,

$$\int_{B_{9/10}} |\nabla v(x) - q|^p \, dx \geqslant \int_{B_{\left(\frac{C(\varepsilon a)^{\nu}}{2\overline{C}}\right)^{1/\alpha}(\overline{x})}} |\nabla v(x) - q|^p \, dx \geqslant \left(\frac{C(\varepsilon a)^{\nu}}{2}\right)^p \left|B_{\left(\frac{C(\varepsilon a)^{\nu}}{2\overline{C}}\right)^{1/\alpha}(\overline{x})}\right|$$

$$\left(\frac{C(\varepsilon a)^{\nu}}{2}\right)^p \left(\frac{C(\varepsilon a)^{\nu}}{2\overline{C}}\right)^{n/\alpha} |B_1| = \frac{C^{p + \frac{n}{\alpha}}|B_1|}{2^{p + \frac{n}{\alpha}}\overline{C}^{\frac{n}{\alpha}}} (\varepsilon a)^{\nu(p + \frac{n}{\alpha})}.$$

Hence, exploiting (3.47), we find that

$$C\varepsilon^p a^p \geqslant \frac{C^{p+\frac{n}{\alpha}}|B_1|}{2^{p+\frac{n}{\alpha}}\overline{C}^{\frac{n}{\alpha}}} (\varepsilon a)^{\nu(p+\frac{n}{\alpha})},$$

which yields the desired contradiction as soon as ε is chosen sufficiently small and $\nu \in \left(0, \frac{p}{p + \frac{n}{\alpha}}\right)$. The proof of (3.48) is thus complete.

Let us define now the function $F: \mathbb{R}^n \longrightarrow \mathbb{R}^n$ as $F(z) := |z|^{p-2} z$. We have that

$$|\nabla v|^{p-2} \nabla v - |q|^{p-2} q = F(\nabla v) - F(\nabla (q \cdot x)) = F(\nabla v) - F(q)$$

$$= \int_0^1 \frac{d}{dt} F(t\nabla v + (1-t)q) dt = \int_0^1 DF(t\nabla v + (1-t)q)(\nabla v - q) dt.$$

Hence, taking the divergence of both sides, we obtain that

(3.51)
$$\operatorname{div}(A(\nabla v - q)) = 0 \quad \text{in } B_{9/10},$$

with

$$A(x) := \int_0^1 DF(t\nabla v(x) + (1-t)q) dt.$$

Notice that we are in the setting of Lemma 3.2 with $\eta := \nabla v - q$. Indeed, we point out that, in light of (3.20) and (3.48), for all $x \in B_{1/4}$,

$$|\nabla v(x) - q| \leqslant C(\varepsilon a)^{\nu} \leqslant \frac{a}{16} < \frac{|q|}{2},$$

as long as ε is sufficiently small.

Accordingly, exploiting Lemma 3.2 we find that

$$\lambda |q|^{p-2} \leqslant A\xi \cdot \xi \leqslant \Lambda |q|^{p-2}$$

for all $\xi \in \partial B_1$, for some $\Lambda \geqslant \lambda > 0$, depending only on p. Recalling (3.20), this gives that

$$\lambda a^{p-2} \leqslant A\xi \cdot \xi \leqslant \Lambda a^{p-2}$$

up to renaming λ and Λ , that now may depend on n and p.

We point out that we are in the setting of Section 3.2 in [MZ97] (see in particular formula (3.43) there, with λ and Λ replaced by λa^{p-2} and Λa^{p-2} here, respectively). Hence, we are in the position of applying Theorem 3.19 in [MZ97], and so we obtain that, for every $x \in B_{1/2}$,

$$(3.52) |\nabla v(x) - q|^p \leqslant \sup_{B_{1/8}(x)} |\nabla v - q|^p \leqslant C \int_{B_{1/4}(x)} |\nabla v(y) - q|^p \, dy \leqslant C\varepsilon^p a^p,$$

for some C > 0, depending on n and p, where we have also used (3.47).

Consequently, for all $x \in B_{1/2}$,

$$(3.53) |\nabla v(x) - q| \leqslant C\varepsilon a,$$

up to renaming C, depending on n and p.

Also, denoting $\bar{q} := \nabla v(0) - q$, we have that

$$(3.54) |\bar{q}| = |\nabla v(0) - q| \leqslant C\varepsilon a.$$

Therefore, from this and (3.53), we deduce that, for all $x \in B_{1/2}$,

$$|\nabla v(x) - q - \bar{q}| \leqslant C\varepsilon a.$$

As a consequence, exploiting Theorem 2 in [Man86], we obtain that, for every $\rho \in (0, 1/2)$,

$$(3.55) \quad \int_{B_{\rho}} |\nabla v(x) - q - \bar{q}|^p \ dx \leqslant \int_{B_{\rho}} \left(C \left(\frac{\rho}{1/2} \right)^{\mu} \|\nabla v - q - \bar{q}\|_{L^{\infty}(B_{1/2})} \right)^p \ dx \leqslant C_2 \, \rho^{\mu p} \, \varepsilon^p a^p,$$

for some $\mu \in (0, 1]$ and $C_2 > 0$ depending on n and p.

Then, if $p \ge 2$, putting together (3.40) and (3.55), we arrive at

$$\int_{B_{\rho}} |\nabla u(x) - q - \bar{q}|^{p} dx$$

$$\leq 2^{p-1} \left(\int_{B_{\rho}} |\nabla u(x) - \nabla v(x)|^{p} dx + \int_{B_{\rho}} |\nabla v(x) - q - \bar{q}|^{p} dx \right)$$

$$\leq 2^{p-1} C_{1} \varepsilon^{p+\delta} \rho^{-n} + 2^{p-1} C_{0} \sigma(a^{p} + 1) \rho^{-n} + 2^{p-1} C_{2} \rho^{\mu p} \varepsilon^{p} a^{p}.$$

Thus, setting $\alpha_0 := \mu$ and, for every $\alpha \in (0, \alpha_0)$,

(3.57)
$$\rho := (2^{p+1}C_2)^{\frac{1}{(\alpha-\alpha_0)p}}, \qquad \varepsilon_0 := \left(\frac{\rho^{\alpha p+n}a_0^p}{2^{p+1}C_1}\right)^{\frac{1}{\delta}} \quad \text{and} \quad c_0 := \frac{\rho^{\alpha p+n}a_0^p}{2^{p+1}C(a_1^p+1)},$$

we have that, for every $\varepsilon \in (0, \varepsilon_0]$ and $\sigma \in (0, c_0 \varepsilon^p]$,

$$2^{p-1}C_2\rho^{\mu p} \leqslant \frac{1}{4}\rho^{\alpha p},$$

$$2^{p-1}C_1\varepsilon^{p+\delta}\rho^{-n} \leqslant \frac{1}{4}\rho^{\alpha p}\varepsilon^p a^p$$
and
$$2^{p-1}C\sigma(a^p+1)\rho^{-n} \leqslant \frac{1}{4}\rho^{\alpha p}\varepsilon^p a^p.$$

As a consequence of this and (3.56),

$$\int_{B_{\rho}} |\nabla u(x) - q - \bar{q}|^p dx \leqslant \frac{1}{4} \rho^{\alpha p} \varepsilon^p a^p + \frac{1}{4} \rho^{\alpha p} \varepsilon^p a^p + \frac{1}{4} \rho^{\alpha p} \varepsilon^p a^p \leqslant \rho^{\alpha p} \varepsilon^p a^p,$$

which gives the desired result in (3.21) by setting $\tilde{q} := q + \bar{q}$.

Moreover, from (3.54) we have that $|q - \tilde{q}| = |\bar{q}| \leqslant C\varepsilon a$, which establishes (3.22). This completes the proof of Lemma 3.3 when $p \geqslant 2$.

In the case $\max\left\{\frac{2n}{n+2},1\right\} , we use (3.41) and (3.55) and we see that$

$$\int_{B_{\rho}} |\nabla u(x) - q - \bar{q}|^p \, dx$$

$$\leqslant 2^{p-1} \left(\int_{B_{\rho}} |\nabla u(x) - \nabla v(x)|^{p} dx + \int_{B_{\rho}} |\nabla v(x) - q - \bar{q}|^{p} dx \right)
\leqslant 2^{p-1} C \left(C_{1} \varepsilon^{p+\delta} + \sigma(a^{p} + 1) \right)^{\frac{p}{2}} a^{p\left(1 - \frac{p}{2}\right)} \rho^{-n} + 2^{p-1} C_{2} \rho^{\mu p} \varepsilon^{p} a^{p}
\leqslant 2^{p-1} C_{1} \varepsilon^{\frac{(p+\delta)p}{2}} a^{p\left(1 - \frac{p}{2}\right)} \rho^{-n} + 2^{p-1} C \sigma^{\frac{p}{2}} (a^{p} + 1)^{\frac{p}{2}} a^{p\left(1 - \frac{p}{2}\right)} \rho^{-n} + 2^{p-1} C_{2} \rho^{\mu p} \varepsilon^{p} a^{p}
\leqslant 2^{p-1} C_{1} \varepsilon^{p+\tilde{\delta}} a^{p\left(1 - \frac{p}{2}\right)} \rho^{-n} + 2^{p-1} C \sigma^{\frac{p}{2}} (a^{p} + 1)^{\frac{p}{2}} a^{p\left(1 - \frac{p}{2}\right)} \rho^{-n} + 2^{p-1} C_{2} \rho^{\mu p} \varepsilon^{p} a^{p},$$

where $\widetilde{\delta} := (p+\delta)p/2 - p > 0$, thanks to (3.46) (and up to renaming C and C_1 , depending on n and p).

We set $\alpha_0 := \mu$ and, for all $\alpha \in (0, \alpha_0)$, we take ρ as in (3.57) and

$$\varepsilon_0 := \left(\frac{\rho^{\alpha p + n} a_0^{\frac{p^2}{2}}}{2^{p+1} C_1}\right)^{\frac{1}{\delta}} \quad \text{and} \quad c_0 := \frac{\rho^{2\alpha + \frac{n}{p}} a_0^p}{4^{\frac{p+1}{p}} C^{\frac{2}{p}} (a_1^p + 1)},$$

obtaining that, for all $\varepsilon \in (0, \varepsilon_0]$ and $\sigma \in (0, c_0 \varepsilon^2]$,

$$\int_{B_{\rho}} |\nabla u(x) - q - \bar{q}|^p dx \leqslant \rho^{\alpha p} \varepsilon^p a^p.$$

Hence, setting $\tilde{q} := q + \bar{q}$, the desired results in (3.21) and (3.22) follow from this and (3.54). The proof of Lemma 3.3 is thereby complete.

Iterating Lemma 3.3 we obtain the following estimates:

Corollary 3.4. Let $a_1 > a_0 > 0$ and $p > \max\{\frac{2n}{n+2}, 1\}$. Let u be an almost minimizer for J_p in B_1 (with constant κ and exponent β) and

$$a := \left(\int_{B_1} |\nabla u(x)|^p \ dx \right)^{1/p}.$$

Suppose that

$$(3.58) a \in [a_0, a_1]$$

and that u satisfies (3.5) and (3.19).

Then there exist ε_0 , κ_0 and $\gamma \in (0,1)$, depending on n, p, β , a_0 and a_1 , such that, for every $\varepsilon \in (0,\varepsilon_0]$ and $\kappa \in (0,\kappa_0\varepsilon^P]$, with $P := \max\{p,2\}$, then

(3.59)
$$||u - \ell||_{C^{1,\gamma}(B_{1/2})} \leqslant C\varepsilon a.$$

The positive constant C depends only on n and p, and ℓ is a linear function of slope q.

Moreover,

with $\widetilde{C} > 0$ depending only on n and p.

Remark 3.5. We point out that a consequence of (3.59) in Corollary 3.4 is that, if ε is sufficiently small,

$$(3.61) \nabla u \neq 0 in B_{1/2}.$$

Indeed, by (3.59), we have that, for all $x \in B_{1/2}$,

$$|\nabla u(x) - q| \leqslant C\varepsilon a,$$

which gives that

$$|\nabla u(x)| \geqslant |q| - |\nabla u(x) - q| \geqslant |q| - C\varepsilon a.$$

As a consequence, from (3.5), we get

$$|\nabla u(x)| \geqslant \frac{a}{4} - C\varepsilon a > 0,$$

as soon as ε is sufficiently small, which yields (3.61).

Furthermore,

$$(3.62) u > 0 in B_{1/2}.$$

To check this, we recall that $u \ge 0$ and we suppose by contradiction that there exists a point $x_0 \in B_{1/2}$ such that $u(x_0) = 0$. As a consequence, since $u \in C^{1,\gamma}(B_{1/2})$, we see that $\nabla u(x_0) = 0$, and this contradicts (3.61), thus proving (3.62).

In order to prove Corollary 3.4 we recall the definition of Campanato spaces and a result which we will use in the proof of the corollary, see [Giu03].

Definition 3.6 (Campanato spaces). Let Ω be a bounded open set in \mathbb{R}^n , $1 \leq p < +\infty$ and $\lambda \geq 0$. We denote by $\mathcal{L}^{p,\lambda}(\Omega)$ the space of functions $u \in L^p(\Omega)$ such that

$$[u]_{\mathcal{L}^{p,\lambda}(\Omega)}^{p} := \sup_{\substack{x_0 \in \Omega \\ \rho > 0}} \rho^{-\lambda} \int_{\Omega_{x_0,\rho}} |u(x) - u_{x_0,\rho}|^{p} dx < +\infty,$$

where

$$\Omega_{x_0,\rho} := \Omega \cap B_{\rho}(x_0) \quad and \quad u_{x_0,\rho} := \int_{\Omega_{x_0,\rho}} u(x) \, dx.$$

We point out that the quantity $[u]_{\mathcal{L}^{p,\lambda}(\Omega)}$ is a seminorm in $\mathcal{L}^{p,\lambda}(\Omega)$ and

(3.64) it is equivalent to the quantity
$$\left(\sup_{\substack{x_0 \in \Omega \\ \rho > 0}} \rho^{-\lambda} \inf_{\xi \in \mathbb{R}^n} \int_{\Omega_{x_0,\rho}} |u(x) - \xi|^p \ dx\right)^{1/p}.$$

We also define the norm in $\mathcal{L}^{p,\lambda}(\Omega)$ as

$$||u||_{\mathcal{L}^{p,\lambda}(\Omega)} := ||u||_{L^{p}(\Omega)} + [u]_{\mathcal{L}^{p,\lambda}(\Omega)}.$$

Theorem 3.7. Let $x \in \mathbb{R}^n$ and r > 0. Let $1 \leq p < +\infty$ and $\lambda \geq 0$. Suppose that $n < \lambda \leq n+p$. Then, the space $\mathcal{L}^{p,\lambda}(B_r(x))$ is isomorphic to $C^{0,\alpha}(B_r(x))$, with $\alpha := \frac{\lambda-n}{p}$.

We also point out the following scaling property of almost minimizers:

Lemma 3.8. Let u be an almost minimizer for J_p in B_1 with constant κ and exponent β . For any $r \in (0,1)$, let

$$(3.66) u_r(x) := \frac{u(rx)}{r}.$$

Then, u_r is an almost minimizer for J_p in $B_{1/r}$ with constant κr^{β} and exponent β , namely

$$(3.67) J_p(u_r, B_{\varrho}(x_0)) \leqslant (1 + \kappa r^{\beta} \varrho^{\beta}) J_p(v, B_{\varrho}(x_0)),$$

for every ball $B_{\varrho}(x_0)$ such that $\overline{B_{\varrho}(x_0)} \subset B_{1/r}$ and for every $v \in W^{1,p}(B_{\varrho}(x_0))$ such that $v = u_r$ on $\partial B_{\varrho}(x_0)$ in the sense of the trace.

Proof. By definition of almost minimizers, we know that

$$(3.68) J_p(u, B_{\vartheta}(y_0)) \leqslant (1 + \kappa \vartheta^{\beta}) J_p(w, B_{\vartheta}(y_0))$$

for every ball $B_{\vartheta}(y_0)$ such that $\overline{B_{\vartheta}(y_0)} \subset B_{1/r}$ and for every $w \in W^{1,p}(B_{\vartheta}(y_0))$ such that w = u on $\partial B_{\vartheta}(y_0)$ in the sense of the trace.

Now, given $x_0 \in B_{1/r}$, we take ϱ and v as in the statement of Lemma 3.8 and we define

$$(3.69) w(x) := rv\left(\frac{x}{r}\right).$$

Then, using the notation $y_0 := rx_0$ and $\vartheta := r\varrho$, for all $x \in \partial B_{\vartheta}(y_0)$, we have that $\frac{x}{r} \in \partial B_{\varrho}(x_0)$ and therefore, in the sense of the trace,

$$w(x) = rv\left(\frac{x}{r}\right) = ru_r\left(\frac{x}{r}\right) = u(x).$$

Accordingly, we can use w as a competitor for u in (3.68), thus obtaining that

$$(3.70) \int_{B_{r\varrho}(y_0)} \left(|\nabla u(y)|^p + \chi_{\{u>0\}}(y) \right) dy \leqslant (1 + \kappa r^\beta \varrho^\beta) \int_{B_{r\varrho}(y_0)} \left(|\nabla w(y)|^p + \chi_{\{w>0\}}(y) \right) dy.$$

Furthermore, using, consistently with (3.66), the notation $w_r(x) := \frac{w(rx)}{r}$, with the change of variable $x := \frac{y}{r}$ we see that

(3.71)
$$\int_{B_{r\varrho}(y_0)} \left(\left| \nabla w(y) \right|^p + \chi_{\{w>0\}}(y) \right) dy = r^n \int_{B_{\varrho}(x_0)} \left(\left| \nabla w(rx) \right|^p + \chi_{\{w>0\}}(rx) \right) dx$$
$$= r^n \int_{B_{\varrho}(x_0)} \left(\left| \nabla w_r(x) \right|^p + \chi_{\{w_r>0\}}(x) \right) dx$$

and a similar identity holds true with u and u_r replacing w and w_r .

Also, recalling (3.69), we observe that $v = w_r$. Plugging this information and (3.71) into (3.70), we obtain the desired result in (3.67).

Proof of Corollary 3.4. Up to scaling, we suppose that

(3.72) u is an almost minimizer for J_p in B_2 (with constant κ and exponent β).

We prove that we can iterate Lemma 3.3 indefinitely with $\alpha := \min \left\{ \frac{\alpha_0}{2}, \frac{\beta}{P} \right\}$, where $P := \max\{p, 2\}$ and α_0 is given in Lemma 3.3.

More precisely, we claim that, for all $k \ge 0$, there exists $q_k \in \mathbb{R}^n$ such that

$$|q_{k}| \in \left[\frac{a}{4} - \frac{\widetilde{C}\varepsilon a\left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}}, C_{0}a + \frac{\widetilde{C}\varepsilon a\left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}}\right],$$

$$\left(\int_{B_{\rho^{k}}} |\nabla u(x) - q_{k}|^{p} dx\right)^{1/p} \leq \rho^{k\alpha}\varepsilon a$$
and
$$\left(\int_{B_{\rho^{k}}} |\nabla u(x)|^{p} dx\right)^{1/p} \in \left[\frac{|q|}{2}, 2|q|\right],$$

where $\rho \in (0,1)$, $C_0 > 0$ and $\widetilde{C} > 0$ are universal quantities.

To prove this, we argue by induction. When k = 0, we pick $q_0 := q$. Then, in this case, the desired claims in (3.73) follow from (3.5), (3.17) and (3.19).

Now we perform the inductive step by assuming that (3.73) is satisfied for k and we establish the claim for k+1. To this end, we let $r := \rho^k$ and utilize the rescaling u_r defined in (3.66). In this setting, the inductive assumption gives that

$$\left(\int_{B_1} \left| \nabla u_r(x) - q_k \right|^p dx \right)^{1/p} \leqslant r^{\alpha} \varepsilon a = \varepsilon_k a,$$

with $\varepsilon_k := r^{\alpha} \varepsilon = \rho^{k\alpha} \varepsilon$.

Notice that the inductive assumption also yields (3.20), as soon as ε is chosen conveniently small. Therefore, thanks to Lemma 3.8 as well, we are in position of using Lemma 3.3 on the function u_r with $\sigma := \kappa r^{\beta}$. We stress that the structural condition $\sigma \leqslant c_o \varepsilon^P$ in Lemma 3.3 translates here into $\kappa \leqslant c_0 \varepsilon^P$, which is precisely the requirement in the statement of Corollary 3.4 (by taking κ_0 there less than or equal to c_0). In this way, we deduce from (3.21) and (3.22) that there exists $q_{k+1} \in \mathbb{R}^n$ such that

$$\left(\int_{B_{\rho}} |\nabla u_r(x) - q_{k+1}|^p \ dx \right)^{1/p} \leqslant \rho^{\alpha} \varepsilon_k a \quad \text{and} \quad |q_k - q_{k+1}| \leqslant \widetilde{C} \varepsilon_k a.$$

Scaling back, we find that

$$\left(\int_{B_{\rho^{k+1}}} |\nabla u(x) - q_{k+1}|^p dx \right)^{1/p} \leqslant \rho^{\alpha} \rho^{k\alpha} \varepsilon a = \rho^{(k+1)\alpha} \varepsilon a.$$

Furthermore,

$$|q_{k+1}| \leqslant |q_k - q_{k+1}| + |q_k| \leqslant \widetilde{C}\varepsilon_k a + C_0 a + \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}}$$

$$= C_0 a + \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}} + \widetilde{C}\rho^{k\alpha}\varepsilon a = C_0 a + \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{(k+1)\alpha}\right)}{1 - \rho^{\alpha}}$$

and

$$|q_{k+1}| \geqslant |q_k| - |q_k - q_{k+1}| \geqslant \frac{a}{4} - \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}} - \widetilde{C}\varepsilon_k a$$

$$= \frac{a}{4} - \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}} - \widetilde{C}\rho^{k\alpha}\varepsilon a = \frac{a}{4} - \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{(k+1)\alpha}\right)}{1 - \rho^{\alpha}}.$$

In addition,

$$\left| \left(\int_{B_{\rho^{k+1}}} \left| \nabla u(x) \right|^p dx \right)^{1/p} - q_{k+1} \right| = \frac{1}{|B_{\rho^{k+1}}|^{1/p}} \left| \left(\int_{B_{\rho^{k+1}}} \left| \nabla u(x) \right|^p dx \right)^{1/p} - q_{k+1} |B_{\rho^{k+1}}|^{1/p} \right|$$

$$= \frac{1}{|B_{\rho^{k+1}}|^{1/p}} \left| \left\| \nabla u \right\|_{L^p(B_{\rho^{k+1}})} - \left\| q_{k+1} \right\|_{L^p(B_{\rho^{k+1}})} \right| \leqslant \frac{1}{|B_{\rho^{k+1}}|^{1/p}} \left\| \nabla u - q_{k+1} \right\|_{L^p(B_{\rho^{k+1}})}$$

$$= \left(\int_{B_{\rho^{k+1}}} \left| \nabla u(x) - q_{k+1} \right|^p dx \right)^{1/p} \leqslant \varepsilon a,$$

which yields that

$$\left| \left(\int_{B_{\rho^{k+1}}} |\nabla u(x)|^p \, dx \right)^{1/p} - q \right| = \left| \left(\int_{B_{\rho^{k+1}}} |\nabla u(x)|^p \, dx \right)^{1/p} - q_0 \right|$$

$$\leq \left| \left(\int_{B_{\rho^{k+1}}} |\nabla u(x)|^p \, dx \right)^{1/p} - q_{k+1} \right| + \sum_{j=0}^k |q_{j+1} - q_j|$$

$$\leq \varepsilon a + \widetilde{C} a \sum_{j=0}^k \varepsilon_j \leqslant \varepsilon a + \widetilde{C} \varepsilon a \sum_{j=0}^{+\infty} \rho^{j\alpha} = \left(1 + \frac{\widetilde{C}}{1 - \rho^{\alpha}} \right) \varepsilon a$$

$$\leq \left(1 + \frac{\widetilde{C}}{1 - \rho^{\alpha}} \right) \varepsilon |q| \leqslant \frac{|q|}{2}.$$

These observations conclude the proof of the inductive step and establish (3.73).

By scaling the relations in (3.73) (and recalling (3.72)), we have that, for all $x_0 \in B_{3/4}$,

(3.74)
$$\left(\int_{B_{a^k}(x_0)} |\nabla u(x) - q_{k,x_0}|^p dx \right)^{1/p} \leqslant C \rho^{k\alpha} \varepsilon a,$$

for some positive constant C, depending only on n and p, and suitable points q_{k,x_0} such that

$$|q_{k,x_0}| \in \left[\frac{a}{4} - \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}}, C_0 a + \frac{\widetilde{C}\varepsilon a \left(1 - \rho^{k\alpha}\right)}{1 - \rho^{\alpha}} \right].$$

We now want to exploit the Campanato estimate in Theorem 3.7, here applied to the gradient of u minus q, and used with $\Omega := B_{1/2}$ and $\lambda := n + \alpha p$ (note that, with this choice, $\lambda \in (n, n+p]$ and we are therefore under the structural assumptions of Theorem 3.7). For this purpose, we claim that, for all $x_0 \in B_{3/4}$ and $\tau > 0$,

(3.75)
$$\left(\tau^{-\lambda} \inf_{\xi \in \mathbb{R}^n} \int_{B_{\tau}(x_0) \cap B_{1/2}} |\nabla u(x) - q - \xi|^p dx\right)^{1/p} \leqslant C\varepsilon a,$$

up to renaming C > 0.

To prove this, we distinguish two cases, either $\tau \ge 1$ or $\tau \in (0,1)$.

If $\tau \geqslant 1$, we use (3.19) and we get that

$$\left(\tau^{-\lambda} \inf_{\xi \in \mathbb{R}^n} \int_{B_{\tau}(x_0) \cap B_{1/2}} |\nabla u(x) - q - \xi|^p \, dx\right)^{1/p} \leqslant \left(\int_{B_{1/2}} |\nabla u(x) - q|^p \, dx\right)^{1/p}$$

$$\leqslant \left(|B_1| \oint_{B_1} |\nabla u(x) - q|^p \, dx\right)^{1/p} \leqslant |B_1|^{1/p} \varepsilon a,$$

which gives (3.75) in this case.

If instead $\tau \in (0,1)$, we pick $k \in \mathbb{N}$ such that $\rho^{k+1} < \tau \leqslant \rho^k$ and we make use of (3.74) to infer that

$$\left(\tau^{-\lambda} \inf_{\xi \in \mathbb{R}^n} \int_{B_{\tau}(x_0) \cap B_{1/2}} |\nabla u(x) - q - \xi|^p \, dx\right)^{1/p} \leq \left(\rho^{-\lambda(k+1)} \int_{B_{\rho^k}(x_0)} |\nabla u(x) - q_{k,x_0}|^p \, dx\right)^{1/p} \\
= \left(\rho^{-\lambda(k+1) + kn} |B_1| \int_{B_{\rho^k}(x_0)} |\nabla u(x) - q_{k,x_0}|^p \, dx\right)^{1/p} \\
\leq C|B_1|^{1/p} \rho^{[-(n+\alpha p)(k+1) + kn + k\alpha p]/p} \varepsilon a \\
\leq C|B_1|^{1/p} \rho^{[-n-\alpha p]/p} \varepsilon a,$$

which gives (3.75) up to renaming C. The proof of (3.75) is thereby complete.

By (3.64) and (3.75), we obtain that

$$[\nabla u - q]_{\mathcal{L}^{p,\lambda}(B_{1/2})} \leqslant C\varepsilon a,$$

up to renaming C. Since also

$$\|\nabla u - q\|_{L^p(B_{1/2})} \leqslant C\varepsilon a,$$

thanks to (3.19), we deduce from (3.65) that $\|\nabla u - q\|_{\mathcal{L}^{p,\lambda}(B_{1/2})} \leq C\varepsilon a$, up to renaming C. Then, in view of Theorem 3.7, we find that

with $\gamma = \frac{\lambda - n}{p} = \alpha$.

Now we define $\ell(x) := u(0) + q \cdot x$. For all $x \in B_{1/2}$ we have that

$$|u(x) - \ell(x)| = |u(x) - u(0) - q \cdot x| = \left| \int_0^1 \left(\nabla u(tx) - q \right) \cdot x \, dt \right| \leqslant C \varepsilon a,$$

thanks to (3.76), up to renaming C, and therefore $||u - \ell||_{L^{\infty}(B_{1/2})} \leq C\varepsilon a$. This and (3.76) establish (3.59), as desired.

4. Lipschitz continuity of almost minimizers and proof of Theorem 1.2

We are now in position of establishing the Lipschitz regularity result in Theorem 1.2.

Proof of Theorem 1.2. In the light of Lemma 3.8, up to a rescaling, we can assume that u is an almost minimizer with constant

$$\widetilde{\kappa} := \kappa s^{\beta},$$

which can be made arbitrarily small by an appropriate choice of s > 0.

We let $P := \max\{p, 2\}$. Let also $\alpha_0 \in (0, 1]$ be the structural constant given by Lemma 3.3 and define

$$\alpha := \frac{1}{2} \min \left\{ \alpha_0, \frac{\beta}{P} \right\}.$$

We also consider ε_0 as given by Proposition 3.1 and take $\eta \in (0,1)$ and $M \ge 1$ as in Proposition 3.1 (corresponding here to choice $\varepsilon := \varepsilon_0/2$).

Let us define

(4.2)
$$a(\tau) := \left(\int_{B_{\tau}} |\nabla u(x)|^p dx \right)^{1/p}.$$

We claim that, for every $r \in (0, \eta]$,

(4.3)
$$a(r) \leq C(M, \eta)(1 + a(1)),$$

for some $C(M, \eta) > 0$, possibly depending on n and p as well.

To prove this, we consider the set $\mathcal{K} \subseteq \mathbb{N} = \{0, 1, 2, \dots\}$ containing all the integers $k \in \mathbb{N}$ such that

(4.4)
$$a(\eta^k) \leqslant C(\eta)M + 2^{-k}a(1),$$

where

$$(4.5) C(\eta) := 2\eta^{-n/p}.$$

We stress that for k = 0 formula (4.4) is clearly true, hence

$$(4.6) 0 \in \mathcal{K} \neq \varnothing.$$

We then distinguish two cases, namely whether (4.4) holds for every k (i.e. $\mathcal{K} = \mathbb{N}$) or not (i.e. $\mathcal{K} \subseteq \mathbb{N}$).

To start with, let us assume that we are in the first situation. Thus, for every $r \in (0, \eta]$, we pick $k_0 \in \mathbb{N} = \mathcal{K}$ such that $\eta^{k_0+1} < r \leqslant \eta^{k_0}$. Hence, according to (4.2) and (4.4), we get that

$$a(r) \leqslant \left(\frac{1}{|B_1| \, \eta^{(k_0+1)n}} \int_{B_{\eta^{k_0}}} |\nabla u(x)|^p \, dx\right)^{1/p} = \eta^{-n/p} a(\eta^{k_0}) \leqslant \eta^{-n/p} (C(\eta)M + a(1))$$

$$\leqslant \eta^{-n/p} \max \left\{ C(\eta)M, 1 \right\} (1 + a(1)) \leqslant C(M, \eta)(1 + a(1)),$$

provided that $C(M, \eta) \ge \eta^{-n/p} \max \{C(\eta)M, 1\}$. The proof of (4.3) is thereby complete in this case.

We now consider the second case, that is the one in which $\mathcal{K} \subsetneq \mathbb{N}$. Then, by (4.6), there exists $k_0 \in \mathbb{N}$ such that $\{0, \ldots, k_0\} \in \mathcal{K}$ and

$$(4.7) k_0 + 1 \not\in \mathcal{K}.$$

We notice that, by (4.2) and (4.5),

$$\eta^{-n/p}M < C(\eta)M \le C(\eta)M + 2^{-(k_0+1)}a(1) < a(\eta^{k_0+1}) \le \eta^{-n/p}a(\eta^{k_0})$$

and therefore

$$(4.8) a(\eta^{k_0}) > M.$$

Furthermore, from (4.7),

(4.9)
$$a(\eta^{k_0+1}) > C(\eta)M + 2^{-(k_0+1)}a(1) \geqslant \frac{C(\eta)M + 2^{-k_0}a(1)}{2} \geqslant \frac{a(\eta^{k_0})}{2}.$$

Now, we claim that

(4.10)
$$\left(\int_{B_{\eta^{k_0+1}}} |\nabla u(x) - q|^p \ dx \right)^{1/p} \leqslant \varepsilon a(\eta^{k_0}),$$
 for some $q \in \mathbb{R}^n$ such that $\frac{a(\eta^{k_0})}{4} < |q| \leqslant C_0 a(\eta^{k_0}),$

being C_0 the constant given by Proposition 3.1.

To prove this, we apply the dichotomy result in Proposition 3.1 rescaled in the ball $B_{\eta^{k_0}}$. In this way, we deduce from (3.3), (3.4) and (3.5) that (4.10) holds true, unless

$$a(\eta^{k_0+1}) = \left(\int_{B_{\eta^{k_0+1}}} |\nabla u(x)|^p \ dx \right)^{1/p} \leqslant \frac{a(\eta^{k_0})}{2},$$

in contradiction with (4.9). This ends the proof of (4.10).

Now we apply Corollary 3.4 rescaled in the ball $B_{\eta^{k_0+1}}$: namely, we use here Corollary 3.4 with B_1 replaced by $B_{\eta^{k_0+1}}$ and a replaced by $a(\eta^{k_0+1})$. To this end, we need to verify that the assumptions of Corollary 3.4 are fulfilled in this rescaled situation. Specifically, we note that, by (4.8) and (4.9),

$$a(\eta^{k_0+1}) \geqslant \frac{M}{2}.$$

Also, since $k_0 \in \mathcal{K}$,

$$a(\eta^{k_0+1}) \leqslant \eta^{-n/p} a(\eta^{k_0}) \leqslant \eta^{-n/p} (C(\eta)M + 2^{-k_0} a(1)) \leqslant \eta^{-n/p} (C(\eta)M + a(1)).$$

These observations give that (3.58) is satisfied in this rescaled setting with

(4.11)
$$a_0 := \frac{M}{2}$$
 and $a_1 := \eta^{-n/p} (C(\eta)M + a(1)).$

Moreover, by (4.9) and (4.10),

$$\left(\int_{B_{\eta^{k_0+1}}} |\nabla u(x) - q|^p \ dx \right)^{1/p} \leqslant 2\varepsilon a(\eta^{k_0+1}),$$

showing that (3.19) is satisfied (here with 2ε instead of ε).

We also deduce from (4.9) and (4.10) that

$$\frac{\eta^{n/p}a(\eta^{k_0+1})}{4} \leqslant \frac{a(\eta^{k_0})}{4} < |q| \leqslant C_0 a(\eta^{k_0}) \leqslant 2C_0 a(\eta^{k_0+1}),$$

which gives that (3.5) is satisfied here (though with different structural constants).

We can thereby exploit Corollary 3.4 in a rescaled version to obtain the existence of ε_0 (possibly different from the one given by Proposition 3.1) and κ_0 , depending on n, p, β , a_0 and a_1 , such that, if

$$(4.12) \widetilde{\kappa} \in (0, \kappa_0 \varepsilon_0^P],$$

we have that

(4.13)
$$\|\nabla u\|_{L^{\infty}(B_{\eta^{k_0+1}/2})} \leqslant \bar{C}a(\eta^{k_0}),$$

for some structural constant $\bar{C} > 0$.

We remark that condition (4.12) is satisfied by taking s in (4.1) sufficiently small, namely taking $s := \left(\frac{\kappa_0 \varepsilon_0^P}{2\kappa}\right)^{\frac{1}{\beta}}$. Notice in particular that

(4.14)
$$s$$
 depends on n , p , κ , β and $\|\nabla u\|_{L^p(B_1)}$,

due to (4.11).

As a result of (4.13), for all $r \in \left(0, \frac{\eta^{k_0+1}}{2}\right)$,

(4.15)
$$a(r) = \left(\frac{1}{|B_r|} \int_{B_r} |\nabla u(x)|^p dx\right)^{\frac{1}{p}} \leqslant \bar{C}a(\eta^{k_0}) \leqslant \bar{C}\left(C(\eta)M + 2^{-k_0}a(1)\right)$$
$$\leqslant \bar{C}\left(C(\eta)M + a(1)\right) \leqslant C(M, \eta)(1 + a(1)),$$

as long as $C(M, \eta) \geqslant \bar{C}(C(\eta)M + 1)$.

Moreover, if $r \in \left[\frac{\eta^{k_0+1}}{2}, \eta\right)$ then we take $k_r \in \mathbb{N}$ such that $\eta^{k_r+1} < r \leqslant \eta^{k_r}$. Note that

$$\frac{1}{\eta^{k_r}} \leqslant \frac{1}{r} \leqslant \frac{2}{\eta^{k_0+1}},$$

whence

$$(4.16) k_r \leqslant k_0 + C_{\star},$$

where $C_{\star} := 1 + \frac{\log 2}{\log(1/\eta)}$.

We now distinguish two cases: if $k_r \in \{0, ..., k_0\}$ then $k_r \in \mathcal{K}$ and therefore

$$a(\eta^{k_r}) \leqslant C(\eta)M + 2^{-k_r}a(1).$$

From this we obtain that

(4.17)
$$a(r) \leqslant \left(\frac{1}{|B_{\eta^{k_r+1}}|} \int_{B_{\eta^{k_r}}} |\nabla u(x)|^p dx\right)^{\frac{1}{p}} = \eta^{-n/p} a(\eta^{k_r})$$
$$\leqslant \eta^{-n/p} \left(C(\eta) M + 2^{-k_r} a(1)\right) \leqslant C(M, \eta) (1 + a(1)).$$

If instead $k_r > k_0$, we employ (4.16) to see that

$$a(r) \leqslant \left(\frac{1}{|B_{\eta^{k_0+C_{\star}+1}}|} \int_{B_{\eta^{k_0}}} |\nabla u(x)|^p dx\right)^{\frac{1}{p}} = \eta^{-n(C_{\star}+1)/p} a(\eta^{k_0})$$

$$\leqslant \eta^{-n(C_{\star}+1)/p} \left(C(\eta)M + 2^{-k_0} a(1)\right) \leqslant \eta^{-n(C_{\star}+1)/p} \left(C(\eta)M + a(1)\right)$$

$$\leqslant C(M, \eta)(1 + a(1)),$$

as long as $C(M, \eta)$ is chosen large enough.

This and (4.17) give that for all $r \in \left[\frac{\eta^{k_0+1}}{2}, \eta\right)$, we have that $a(r) \leqslant C(M, \eta)(1 + a(1))$. Combining this with (4.15), we deduce that (4.3) holds true, as desired.

Now, up to scaling and translations, we can extend (4.3) to all balls with center x_0 in $B_{1/2}$ and sufficiently small radius. Namely, we have that, for all $r \in (0, \eta]$,

$$a(r, x_0) \leqslant C(M, \eta)(1 + a(1)),$$

where

$$a(r, x_0) := \left(\int_{B_r(x_0)} |\nabla u(x)|^p dx \right)^{1/p}.$$

Therefore, according to the Lebesgue Differentiation Theorem, recalling that $u \in W^{1,p}(B_1)$, we find that, for all $x_0 \in B_{1/2}$,

$$|\nabla u(x_0)| = \lim_{r \to 0} a(r, x_0) \leqslant C(M, \eta)(1 + a(1)) = C\left(1 + ||\nabla u||_{L^p(B_1)}\right),$$

for some C > 0, which yields

(4.18)
$$\|\nabla u\|_{L^{\infty}(B_{1/2})} \leqslant C \Big(1 + \|\nabla u\|_{L^{p}(B_{1})} \Big).$$

So, the first claim in Theorem 1.2 is proved.

We now show that the second claim in Theorem 1.2 holds true. For this, we can assume that

$$(4.19) {u = 0} \cap B_{s/100} \neq \varnothing$$

and we take a Lebesgue point $\bar{x} \in B_{s/100}$ for ∇u . Up to a translation, we can suppose that $\bar{x} = 0$ and change our assumption (4.19) into

$$(4.20) {u=0} \cap B_{s/50} \neq \varnothing.$$

We claim that, in this situation,

$$(4.21) \mathcal{K} = \mathbb{N}.$$

Indeed, suppose not and let k_0 as above (recall (4.7)). Then, in light of (4.10), we can apply Corollary 3.4 (rescaled as before) and conclude, by (3.62), that u > 0 in $B_{s/2}$, in contradiction with our hypothesis in (4.20). This establishes (4.21).

Therefore, in view of (4.4) and (4.21), for all $k \in \mathbb{N}$,

$$a(\eta^k) \leqslant C(\eta)M + 2^{-k}a(1),$$

and as a result

$$|\nabla u(\bar x)| = |\nabla u(0)| = \lim_{k \to +\infty} a(\eta^k) \leqslant \lim_{k \to +\infty} \left(C(\eta)M + 2^{-k}a(1)\right) = C(\eta)M.$$

Recalling also (4.14), the proof of the second claim in Theorem 1.2 is thereby complete. \Box

References

- [Ack77] Andrew Acker, Heat flow inequalities with applications to heat flow optimization problems, SIAM J. Math. Anal. 8 (1977), no. 4, 604–618, DOI 10.1137/0508048. MR473960
- [AC81] H. W. Alt and L. A. Caffarelli, Existence and regularity for a minimum problem with free boundary, J. Reine Angew. Math. 325 (1981), 105–144. MR618549
- [AP98] N. D. Aparicio and M. K. Pidcock, On a class of free boundary problems for the Laplace equation in two dimensions, Inverse Problems 14 (1998), no. 1, 9–18, DOI 10.1088/0266-5611/14/1/003. MR1607608
- [BCN90] H. Berestycki, L. A. Caffarelli, and L. Nirenberg, Uniform estimates for regularization of free boundary problems, Analysis and partial differential equations, Lecture Notes in Pure and Appl. Math., vol. 122, Dekker, New York, 1990, pp. 567–619. MR1044809
- [BM14] J. Ederson M. Braga and Diego R. Moreira, Uniform Lipschitz regularity for classes of minimizers in two phase free boundary problems in Orlicz spaces with small density on the negative phase, Ann. Inst. H. Poincaré C Anal. Non Linéaire 31 (2014), no. 4, 823–850, DOI 10.1016/j.anihpc.2013.07.006. MR3249814
- [Caf87] Luis A. Caffarelli, A Harnack inequality approach to the regularity of free boundaries. I. Lipschitz free boundaries are C^{1,α}, Rev. Mat. Iberoamericana 3 (1987), no. 2, 139–162, DOI 10.4171/RMI/47. MR990856
- [Caf89] _____, A Harnack inequality approach to the regularity of free boundaries. II. Flat free boundaries are Lipschitz, Comm. Pure Appl. Math. 42 (1989), no. 1, 55–78, DOI 10.1002/cpa.3160420105. MR973745
- [Caf88] ______, A Harnack inequality approach to the regularity of free boundaries. III. Existence theory, compactness, and dependence on X, Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) 15 (1988), no. 4, 583–602 (1989). MR1029856
- [CS05] Luis Caffarelli and Sandro Salsa, A geometric approach to free boundary problems, Graduate Studies in Mathematics, vol. 68, American Mathematical Society, Providence, RI, 2005. MR2145284
- [DP05] Donatella Danielli and Arshak Petrosyan, A minimum problem with free boundary for a degenerate quasilinear operator, Calc. Var. Partial Differential Equations 23 (2005), no. 1, 97–124, DOI 10.1007/s00526-004-0294-5. MR2133664
- [DP06] _____, Full regularity of the free boundary in a Bernoulli-type problem in two dimensions, Math. Res. Lett. 13 (2006), no. 4, 667–681, DOI 10.4310/MRL.2006.v13.n4.a14. MR2250499
- [DPS03] D. Danielli, A. Petrosyan, and H. Shahgholian, A singular perturbation problem for the p-Laplace operator, Indiana Univ. Math. J. **52** (2003), no. 2, 457–476, DOI 10.1512/iumj.2003.52.2163. MR1976085
- [DET19] Guy David, Max Engelstein, and Tatiana Toro, Free boundary regularity for almost-minimizers, Adv. Math. 350 (2019), 1109–1192, DOI 10.1016/j.aim.2019.04.059. MR3948692
- [DT15] G. David and T. Toro, Regularity of almost minimizers with free boundary, Calc. Var. Partial Differential Equations 54 (2015), no. 1, 455–524, DOI 10.1007/s00526-014-0792-z. MR3385167
- [DS11] D. De Silva, Free boundary regularity for a problem with right hand side, Interfaces Free Bound. 13 (2011), no. 2, 223–238, DOI 10.4171/IFB/255. MR2813524
- [DSS20a] D. De Silva and O. Savin, *Thin one-phase almost minimizers*, Nonlinear Anal. **193** (2020), 111507, 23, DOI 10.1016/j.na.2019.04.006. MR4062979
- [DSS20b] ______, Almost minimizers of the one-phase free boundary problem, Comm. Partial Differential Equations 45 (2020), no. 8, 913–930, DOI 10.1080/03605302.2020.1743718. MR4126326
- [DSS21] _____, Quasi-Harnack inequality, Amer. J. Math. **143** (2021), no. 1, 307–331, DOI 10.1353/ajm.2021.0001. MR4201786
- [DiB83] E. DiBenedetto, $C^{1+\alpha}$ local regularity of weak solutions of degenerate elliptic equations, Nonlinear Anal. 7 (1983), no. 8, 827–850, DOI 10.1016/0362-546X(83)90061-5. MR709038

- [DK18a] Serena Dipierro and Aram L. Karakhanyan, Stratification of free boundary points for a two-phase variational problem, Adv. Math. 328 (2018), 40–81, DOI 10.1016/j.aim.2018.01.005. MR3771123
- [DK18b] ______, A new discrete monotonicity formula with application to a two-phase free boundary problem in dimension two, Comm. Partial Differential Equations 43 (2018), no. 7, 1073–1101, DOI 10.1080/03605302.2018.1499776. MR3910196
- [Eva98] Lawrence C. Evans, *Partial differential equations*, Graduate Studies in Mathematics, vol. 19, American Mathematical Society, Providence, RI, 1998. MR1625845
- [FL21] Fausto Ferrari and Claudia Lederman, Regularity of flat free boundaries for a p(x)-Laplacian problem with right hand side, Nonlinear Anal. 212 (2021), Paper No. 112444, 25, DOI 10.1016/j.na.2021.112444. MR4273843
- [FL22] ______, Regularity of Lipschitz free boundaries for a p(x)-Laplacian problem with right hand side, Preprint (2022).
- [For21] Nicolò Forcillo, Regularity in degenerate elliptic and parabolic free boundary problems, Alma Mater Studiorum Università di Bologna. Dottorato di ricerca in Matematica, 34 Ciclo, posted on 2021, 1–227, DOI 10.48676/unibo/amsdottorato/10001.
- [GT01] David Gilbarg and Neil S. Trudinger, Elliptic partial differential equations of second order, Classics in Mathematics, Springer-Verlag, Berlin, 2001. Reprint of the 1998 edition. MR1814364
- [Giu03] Enrico Giusti, Direct methods in the calculus of variations, World Scientific Publishing Co., Inc., River Edge, NJ, 2003. MR1962933
- [IM89] Tadeusz Iwaniec and Juan J. Manfredi, Regularity of p-harmonic functions on the plane, Rev. Mat. Iberoamericana 5 (1989), no. 1-2, 1-19, DOI 10.4171/RMI/82. MR1057335
- [Kar06] Aram L. Karakhanyan, Up-to boundary regularity for a singular perturbation problem of p-Laplacian type, J. Differential Equations 226 (2006), no. 2, 558–571, DOI 10.1016/j.jde.2005.10.014. MR2237691
- [Kar08] _____, On the Lipschitz regularity of solutions of a minimum problem with free boundary, Interfaces Free Bound. 10 (2008), no. 1, 79–86, DOI 10.4171/IFB/180. MR2383537
- [Kar21] Aram Karakhanyan, Full and partial regularity for a class of nonlinear free boundary problems, Ann. Inst. H. Poincaré C Anal. Non Linéaire 38 (2021), no. 4, 981–999, DOI 10.1016/j.anihpc.2020.09.008. MR4266232
- [LS87] A. A. Lacey and M. Shillor, Electrochemical and electro-discharge machining with a threshold current, IMA J. Appl. Math. 39 (1987), no. 2, 121–142, DOI 10.1093/imamat/39.2.121. MR983738
- [LdQT15] Raimundo Leitão, Olivaine S. de Queiroz, and Eduardo V. Teixeira, Regularity for degenerate two-phase free boundary problems, Ann. Inst. H. Poincaré C Anal. Non Linéaire 32 (2015), no. 4, 741–762, DOI 10.1016/j.anihpc.2014.03.004. MR3390082
 - [LR18] Raimundo Leitão and Gleydson Ricarte, Free boundary regularity for a degenerate problem with right hand side, Interfaces Free Bound. 20 (2018), no. 4, 577–595, DOI 10.4171/IFB/413. MR3893420
 - [LN10] John L. Lewis and Kaj Nyström, Regularity of Lipschitz free boundaries in two-phase problems for the p-Laplace operator, Adv. Math. 225 (2010), no. 5, 2565–2597, DOI 10.1016/j.aim.2010.05.005. MR2680176
 - [LN12] ______, Regularity of flat free boundaries in two-phase problems for the p-Laplace operator, Ann. Inst. H. Poincaré C Anal. Non Linéaire 29 (2012), no. 1, 83–108, DOI 10.1016/j.anihpc.2011.09.002. MR2876248
 - [MZ97] Jan Malý and William P. Ziemer, Fine regularity of solutions of elliptic partial differential equations, Mathematical Surveys and Monographs, vol. 51, American Mathematical Society, Providence, RI, 1997. MR1461542
 - [Man86] Juan Jose Manfredi, Regularity of the gradient for a class of nonlinear possibly degenerate elliptic equations, ProQuest LLC, Ann Arbor, MI, 1986. Thesis (Ph.D.)—Washington University in St. Louis. MR2635642

- [MW08] Sandra Martínez and Noemi Wolanski, A minimum problem with free boundary in Orlicz spaces, Adv. Math. 218 (2008), no. 6, 1914–1971, DOI 10.1016/j.aim.2008.03.028. MR2431665
- [MW14] Diego Moreira and Lihe Wang, Singular perturbation method for inhomogeneous nonlinear free boundary problems, Calc. Var. Partial Differential Equations 49 (2014), no. 3-4, 1237–1261, DOI 10.1007/s00526-013-0620-x. MR3168631
- [PV05a] Arshak Petrosyan and Enrico Valdinoci, Geometric properties of Bernoulli-type minimizers, Interfaces Free Bound. 7 (2005), no. 1, 55–77, DOI 10.4171/IFB/113. MR2126143
- [PV05b] _____, Density estimates for a degenerate/singular phase-transition model, SIAM J. Math. Anal. **36** (2005), no. 4, 1057–1079, DOI 10.1137/S0036141003437678. MR2139200
- [Tol84] Peter Tolksdorf, Regularity for a more general class of quasilinear elliptic equations, J. Differential Equations 51 (1984), no. 1, 126–150, DOI 10.1016/0022-0396(84)90105-0. MR727034
- [Ura68] N. N. Ural'ceva, Degenerate quasilinear elliptic systems, Zap. Naučn. Sem. Leningrad. Otdel. Mat. Inst. Steklov. (LOMI) 7 (1968), 184–222 (Russian). MR0244628
- [Val06] Enrico Valdinoci, Flatness of Bernoulli jets, Math. Z. **254** (2006), no. 2, 257–298, DOI 10.1007/s00209-006-0947-5. MR2262703

SERENA DIPIERRO: DEPARTMENT OF MATHEMATICS AND STATISTICS, UNIVERSITY OF WESTERN AUSTRALIA, 35 STIRLING HWY, CRAWLEY WA 6009 AUSTRALIA

Email address: serena.dipierro@uwa.edu.au

Fausto Ferrari: Dipartimento di Matematica, Università di Bologna, Piazza di Porta S.Donato 5, 40126, Bologna-Italy

Email address: fausto.ferrari@unibo.it

Nicolò Forcillo: Dipartimento di Matematica, Università di Bologna, Piazza di Porta S.Donato 5, 40126, Bologna-Italy

Email address: nicolo.forcillo2@unibo.it

ENRICO VALDINOCI: DEPARTMENT OF MATHEMATICS AND STATISTICS, UNIVERSITY OF WESTERN AUSTRALIA, 35 STIRLING HWY, CRAWLEY WA 6009 AUSTRALIA

Email address: enrico.valdinoci@uwa.edu.au