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# A $-8$ mV/ $+15$ mV Double Polarity Piezoelectric Transformer-Based Step-up Oscillator for Energy Harvesting Applications

Antonio Camarda, Marco Tartagni, *Member, IEEE*, and Aldo Romani

**Abstract**—This work presents two circuit topologies of battery-less integrated boost oscillators suitable for kick-starting electronic systems in fully discharged states with ultra-low input voltages, in the context of energy harvesting applications based on thermoelectric generators, by coupling a piezoelectric transformer in a feedback loop. With respect to prior work, the first presented solution is a double polarity integrated circuit designed in a  $0.18$   $\mu\text{m}$  CMOS technology able to boost ultra-low positive and negative voltages without the need of switching matrixes. The circuit exploits a CMOS inverter made up of low threshold transistors, and also includes a hysteretic voltage monitor consuming only  $\sim 15$  nW to enable an external circuit. The minimum achieved positive and negative oscillation voltages are  $+15$  mV and  $-8$  mV, which to the best of the authors' knowledge, are among the lowest start-up voltages achieved in literature up to now without using magnetic components. Moreover, the input impedance in the range of several k $\Omega$  makes the presented solution suitable also for high impedances sources such as rectennas. The second presented circuit, designed in a  $0.32$   $\mu\text{m}$  CMOS technology, exploits an input stage based on depletion-mode MOSFETs in a common source stage configuration and achieves high-step ratios up to  $\sim 60$ .

**Index Terms**— boost circuit, bootstrap circuit, double polarity, energy harvesting, oscillator, piezoelectric transformer, rectenna, thermoelectric generator, voltage monitor

## I. INTRODUCTION

WIRELESS Sensors Networks (WSN) are groups of widespread sensors communicating between each other through wireless channels. They have the duty of collecting, transmitting and storing spatially distributed data about several environmental physical quantities such as temperature, humidity or air-pressure. Battery-powered WSN generally use a certain amount of energy for collecting and sending data to the central unit (few nJ/bit) [1]. The consumed power can be reduced by implementing architectures exploiting specific blocks such as wake-up radios or by reducing the overall

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duty-cycle of operation (i.e., the duration of the on-state compared to sleep-mode) of the network of sensors [2].

In any case, lifetime is one of the most important and critical aspects of energy autonomous systems. Technology scaling has played its role by reducing circuits dimensions and thus making systems less power-demanding. Nevertheless, this is not sufficient, given that relying on batteries is not the optimal solution because their replacement can require very high maintenance costs. In order to achieve fully energy autonomous systems, it is possible to conceive battery-less systems that convert in an electric form the harvested energy available in the environment in several forms such as sunlight, wind, heat, vibrations or RF [3] [4] [5] [6]. However, environmental energy sources often provide extremely low voltages, whereas switched-mode boost converters or charge pumps might require a higher minimum input voltage typically in the order of the threshold voltage of the used transistors in order to produce a usable supply voltage for a WSN. To cite an example, photovoltaic cells (PVCs) provide different output power levels depending on the illumination conditions which may vary in the range of two orders of magnitude [7], and the provided output voltage can be as low as  $\sim 200$  mV in low illumination conditions. Radio Frequency (RF) signals can also be exploited to provide power to passive devices, as is the case in radio frequency identification (RFID) tags. However, the distance of power transmission typically is limited to several meters [8], and at such high distances the amplitude of the received voltage of rectennas (i.e. rectifying antennas) can be as low as  $\sim 100$ - $200$  mV. Moreover, in order to extract power, temperature gradients can be exploited as well by means of thermoelectric generators (TEGs) based on the Seebeck effect (Fig. 1): when two different metals are electrically connected and subject to a temperature gradient, a voltage difference arises at the terminals of the connection. The electrical series-connection and the thermal parallel connection of such thermocouples increases the output voltage at the expenses of a higher output resistance. The size of these devices is in the order of few  $\text{cm}^2$  with output voltages around  $10$ - $50$  mV/ $^{\circ}\text{C}$  and output resistances as low as  $\sim 0.4$   $\Omega$  [9].

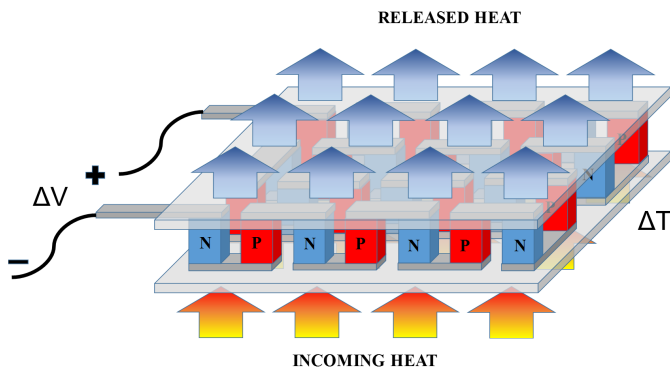


Fig. 1: Representation of a TEG. An array of electrically series-connected and thermally parallel-connected thermocouples outputs a voltage proportional to the temperature difference.

In most cases, the output voltage of the aforementioned transducers is not sufficient to overcome the threshold voltage  $V_{TH}$  of the power devices in a conventional boost power converter or charge pump. Then, in energy harvesting (EH) systems (Fig. 2), since the main power converter requires a minimum supply voltage, it can be activated only once a storage element (e.g. a capacitor) has been charged to this level by a separate kick-start voltage booster circuit. The kick-start voltage booster is a circuit able to function at very low voltages levels down to few tens of mV by exploiting normally-ON transistors such as JFETs or depletion-mode MOSFETs (DeplMOS). The main purpose of such circuit is the charge of the storage capacitor to a voltage level sufficient to provide the minimum supply voltage for a conventional power converter to perform the power conversion.

The problem of kick-starting from ultra-low voltages down to few tens of mV has been extensively addressed in literature. The first approach was presented by Damaschke in 1997 [10] and is depicted in Fig. 3. The Ultra-Low Voltage Source (ULVS) is modeled through a voltage  $V_{IN}$  with an output resistance  $R_S$ . A normally-ON transistor ( $M_1$ , DeplMOS or JFET), coupled with a Magnetic Transformer (MT), is required to provide an overall negative resistance in parallel to the LC tank, made by the secondary inductance and the equivalent capacitance  $C_G$  seen at the gate of the transistor  $M_1$ , when  $V_{IN}$  does not exceed few tens/hundreds mV. The positive feedback due to the MT thus produces an oscillation which is then rectified through the Voltage Doubler (VD) made by  $C_{PUMP}$ , the diodes  $D_1$  and  $D_2$  and the storage capacitor  $C_{STORE}$ . The minimum input voltage  $V_{IN,MIN}$  to achieve oscillation is dependent from the transformer turn-ratio  $N$  and the losses in the MT [11]. This particular topology has been adopted for RF harvesting purposes [12] and in commercial products operating from 20 mV with a 1:100 MT [13]. The cascade-connection of several MTs allows to reduce the minimum voltage to trigger an oscillation down to  $\sim 7$  mV, at the expenses of system compactness [14].

Other ultra-low voltage systems were presented in literature, but an initial voltage (e.g. provided by an external energy source or by a battery) of several hundreds of mV is necessary [15] [16].

A minimum input voltage of 35 mV was achieved through

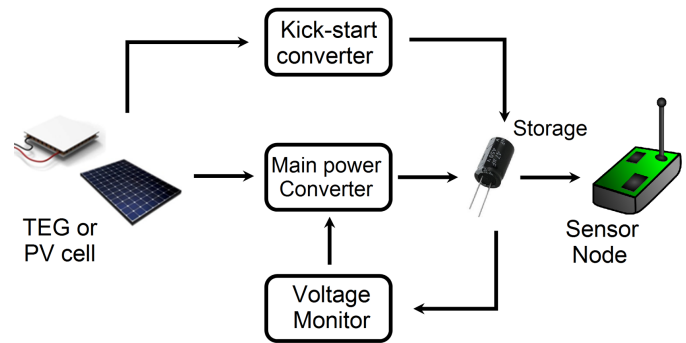


Fig. 2: Diagram of an energy autonomous system based on energy harvesting. A kick-start boost converter provides the initial voltage and energy required to turn on a more efficient power converter.

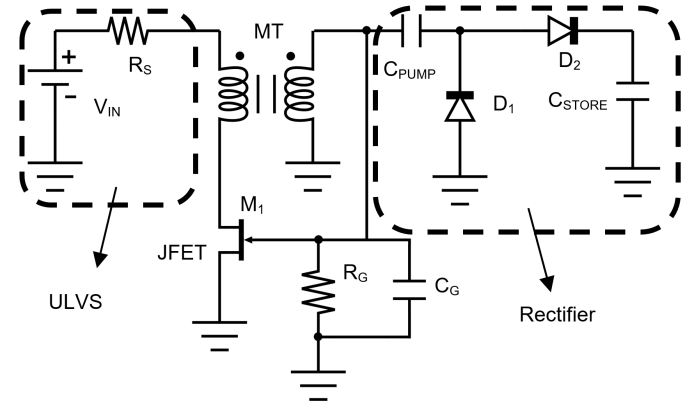


Fig. 3: Schematic of the conventional Armstrong oscillator with output rectification.

exploitation of a motion-activated switch [17]. Furthermore, input voltages down to 10 mV were also recently achieved [18] [19], but the kick-start is aided by means of a 32-stages Dickson charge pump boosting the output of a secondary harvester providing a much higher voltage (rectenna).

Other solutions are related to the design of specific integrated circuits (ICs) exploiting Forward Body Biasing (FBB) techniques [20], which aim at reducing the threshold voltage of transistors in a charge pump. Such technique allows to lower the start-up voltage down to 0.18 V. However, the main drawback is the higher leakage currents that degrade the whole efficiency of the converter. Dynamic Body Biasing (DBB) techniques [21] represent an improvement with respect to FBB, however both FBB and DBB require deep wells in the process in order to have a floating body connection.

Another technique with the purpose of starting from voltages lower than the normal threshold voltage of transistors consists in tuning the  $V_{TH}$  of the transistors at the end of the fabrication step. This technique [22] was applied to the p-channel MOS transistors (pMOS) placed in a ring oscillator providing the clock signal for a conventional Dickson charge-pump (DCP) and achieved 95 mV operation. This technique essentially provides additional fixed charge in the transistors' gate dielectrics fixing the  $V_{TH}$ . Among other types of low-voltage circuits, a battery-less integrated harvesting system with minimum voltage of 50 mV was recently presented [23], in which the oscillation is achieved through a cross-coupled nMOS LC oscillator exploiting the high  $Q$ -factor of a commercial inductor. Moreover, an ultra-low



quiescent power boost converter for sources down to 70 mV is presented in [24]: the quiescent power is lowered down to 544 pW with a peak power conversion efficiency of about 56%. However, it still relies on magnetic components and requires the injection of extra energy for kick-starting. A reconfigurable buck/boost converter suitable for solar-cells is presented in [25]. The quiescent power consumption is as low as 3 nW but the minimum start-up voltage is  $\sim 140$  mV. It is worth remarking that the efficiency as well as the minimum activation voltage of such systems generally depend on the “quality” of magnetic components, whose scaling poses several constraints during the design phase [26]. In [27], an ultra-low voltage of 21 mV, with minimum input power of 5.8  $\mu$ W and efficiency higher than 70% is achieved, exploiting 1:1 MTs. However, such values are achieved with 10 mH coils.

Piezoelectric transformers (PTs) are highly efficient power conversion devices that are generally used for high voltage applications such as Cold Cathode Fluorescent Lamps (CCFL) [28] [29] or highly efficient resonant converters [30] [31]. Nevertheless, their use in ultra-low voltage applications has been demonstrated: a novel start-up approach based on the use of PTs instead of MTs has been recently presented [32]-[34]. Such bootstrap circuit was successfully adopted for kick-starting a fixed-frequency boost converter, from a fully discharged state, performing a power conversion from a TEG with efficiencies around more than 40% in the range 30-50 mV [35]. However, such systems work only with positive voltages. It is worth remarking that PTs in a feedback loop for oscillation boosting purposes were also previously reported [36], however the PT in [36] works close the resonance frequency and the circuit is a Pierce-type oscillator that requires an inductor to achieve oscillation. Differently, in [32] [33] the PT works close to its parallel resonance (or anti-resonance) frequency and magnetic components are not necessary to trigger an oscillation.

A first contribution of this work is a new double-polarity circuit topology of integrated boost oscillator designed and fabricated in a UMC 0.18  $\mu$ m CMOS technology, able to boost positive and negative voltages without the necessity of switching the polarity of the input voltage through switching matrixes [37]. The circuit achieves a minimum positive start-up voltage of 15 mV by exploiting PTs in a feedback loop. When polarized by a negative voltage, the circuit is able to oscillate at  $-8$  mV. To the best of the authors’ knowledge, these are among the lowest values reported in literature without magnetic components and without any battery or external energy contributions. The output voltage with a  $\pm 15$  mV input is around 150 mV, achieved with a rectifying VD; however, higher voltages can be obtained through use of voltage multipliers or conventional multi-stage charge pumps. With respect to prior work [33], the input stage is based on a CMOS inverter gate made up of low-threshold native MOSFETs (low- $V_{TH}$  or NatMOS) rather than a common source (CS) stage. One of the main advantages is that NatMOS are generally more common in ICs with respect to JFETs or DeplMOS. This solution, based on CMOS inverter input stage, differently from [33], integrates also a nano-power

hysteretic voltage monitor (VM) sinking only  $\sim 9$  nA at 1.6 V. The aim of such VM is to connect a load  $R_{LOAD}$  to  $C_{STORE1}$  when  $V_{STORE1} > V_H$ , and to disconnect such load when  $V_{STORE1} < V_L$ , with  $V_H$  and  $V_L$  the switching thresholds of the VM.

Moreover, the circuit presented in prior work [33] is suitable only for low input impedances sources such as TEGs, whereas another relevant feature of the CMOS inverter-based circuit is the capability to work with high impedances sources in the range of several  $k\Omega$  such as rectennas.

An additional contribution of this work is an improved integrated version of the circuit shown in [33], based on DeplMOS rather than JFETs and implemented in a STMicroelectronics 0.32  $\mu$ m CMOS process. This circuit is able to self-start from voltages as low as 31 mV and to provide an output voltage of 5 V, with a maximum step-up ratio of  $\sim 60$ . The minimum voltage is comparable to that reported in [33], but in this specific case, no magnetic component was used, given that the main purpose of the research is to avoid the insertion of any magnetic components. Moreover, through the exploitation of a higher  $Q$ -factor PT, with respect to [33], it was possible to modulate the voltage gain of the system in order to improve further the performances.

The presented boost oscillators are not intended as stand-alone converters because of their poor efficiency. As shown in Fig. 2, their duty is to kick-start an energy autonomous system from a fully discharged state and to enable a more efficient power converter, once a sufficient voltage is achieved on the storage element. In fact, the presented circuits do not aim at extracting the maximum power from the TEG, which is the target of the main boost converter, but rather at working in proximity of the open circuit output voltage to reduce voltage drops that would compromise the gain of the amplifier stage, so as to minimize the operating voltage. Fig. 4 depicts the numerically simulated I-V (current-voltage), I-P (current-power), R-P (resistance-power) and R-V (resistance-voltage) characteristics of a TEG providing 100 mV in open circuit with an internal resistance  $R_S = 10 \Omega$ . In order to operate in the maximum power point (MPP), the boost oscillator should draw a current causing the output voltage of the TEG to be equal to half of its open-circuit voltage. However, in this condition, if the boost oscillator has a minimum operating voltage  $V_{IN,MIN}$ , the TEG should provide an open circuit voltage equal to  $2 \cdot V_{IN,MIN}$ . Hence, for these converters, the MPP operation is in contrast with the purpose of lowering the minimum activation voltage. On the other hand, biasing the TEG close to its open-circuit voltage will reduce the extracted power but will also ensure a higher voltage for supplying the kick-start converter.

The paper is organized as follows: Section I introduces the work, Section II presents the circuits, Section III deals with experimental validation, and Section IV concludes the paper.

## II. DESCRIPTION OF STEP-UP OSCILLATORS

### A. The PT

Table I reports the electromechanical parameters of the Rosen-type PT used in our work, provided by Noliac (Fig. 5).

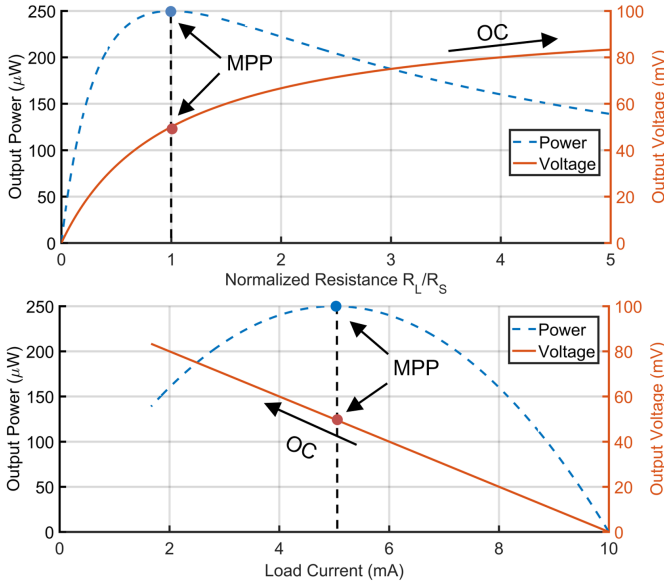


Fig. 4: Characteristics R-V, R-P, I-P and I-V of a TEG providing 100 mV in open circuit with a series resistance of 10Ω.

TABLE I  
ELECTRO-MECHANICAL PARAMETERS OF THE TWO-PORTS  
BUTTERWORTH-VAN DYKE NETWORK OF THE ADOPTED SAMPLE.

Mode	$C_{IN}$ (nF)	$L_M$ (μH)	$R_M$ (Ω)	$C_M$ (nF)	$N$	$C_{OUT}$ (pF)	$f_s$ (kHz)	$f_p$ (kHz)
1	136.65	635	0.21	14.7	55.6	17.35	58.8	61.23
2	136.65	220.8	1.14	11.2	58.3	17.35	110.4	114.14

$C_{IN}$  and  $C_{OUT}$  are the PT input and output capacitance respectively. The electromechanical transduction is modeled through the branch composed by the inductance  $L_M$ , the capacitor  $C_M$ , and the resistance  $R_M$ . The factor  $N$  represents the stress-ratio from input to output (the equivalent of the turns-ratio of a MT). Additional details concerning the modeling and behavior of PTs, can be found in [38]-[40].

The voltage transfer function of a PT is that of a two-poles system with extremely high-quality factor ( $Q$ -factor) [33]:

$$A_{V,PT}(s) = \frac{C_{M2}\omega_s^2}{NC_{OUT} \left( s^2 + \frac{\omega_s s}{Q} + \omega_s^2 \right)}, \quad (1)$$

where  $C_{M2} = C_M \cdot N^2 C_O / (C_M \cdot N^2 C_O)$  is the open-circuit mechanical capacitance,  $Q = \omega_s L_M / R_M$  is the quality factor and  $\omega_s = (L_M C_M)^{-1/2}$  is the resonance pulsation. The resonance frequency can be defined as:  $f_s = \omega_s / 2\pi$ .  $C_O$  is the equivalent capacitance seen at the output capacitance, that is  $C_O = C_{OUT} + C_{VD} + C_G$ ;  $C_{VD}$  accounts for the capacitive load effect of the VD around 0 V (due to the ultra-low voltage in input, the DC level of several nodes can be safely approximated to ground), whereas  $C_G$  is the capacitive load effect of the CS or the CMOS inverter.

As explained later, for CS-based step-up oscillators, it is possible to put an additional capacitance  $C_{OPT}$  in parallel to  $C_{IN}$ . In this case, we define  $C_{IN2} = C_{IN} + C_{OPT}$ . The anti-resonance frequency  $f_p$  of the PT can be expressed as:

$$f_p = \frac{\omega_p}{2\pi} = 2\pi \left( L_M \frac{C_{M2} \cdot C_{IN2}}{C_{M2} + C_{IN2}} \right)^{-1} \quad (2)$$

The values of both  $f_s$  and  $f_p$  for the PT used in our work are reported in Table I. At both  $f_s$  and  $f_p$  the input impedance of a

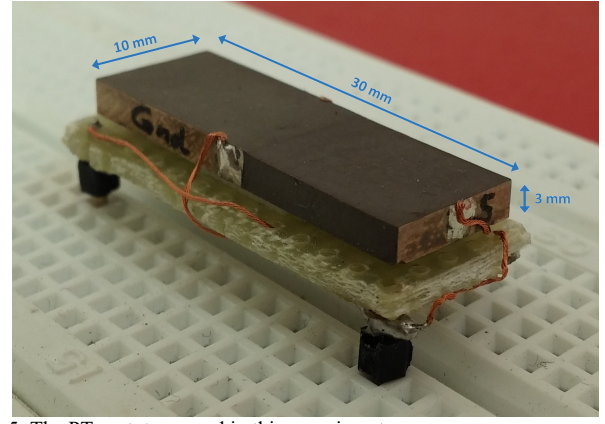


Fig. 5: The PT prototype used in this experiment.

PT is approximately real, with the only difference that at  $f_s$  the impedance is at its minimum, whereas at  $f_p$  is at its maximum. Between  $f_s$  and  $f_p$ , the input impedance is inductive, whereas for frequencies much lower than  $f_s$  and much higher than  $f_p$ , such impedance is capacitive and dominated by  $C_{IN}$  [33].

### B. Step-up oscillator with CMOS inverter as input stage and output voltage monitor

Fig. 6 depicts the schematic of the proposed step-up oscillator for ULV sources. Due to its capacitive behavior in DC, the PT has to be connected in parallel with the input inverting stage. Differently, if the PT is replaced with a MT in the circuit in Fig. 6, the connection node between  $M_1$  and  $M_2$  ( $M\_INV$  output) would be shorted to ground and the circuit does not oscillate. The ULVS (TEG/rectenna) is represented by an equivalent voltage source  $V_{IN}$  and a series resistance  $R_S$ .

The mode of operation of the feedback loop is similar as the circuit presented in [33]. In this specific case, the CS is replaced by a CMOS inverter ( $M\_INV$ ) made up by complementary NatMOS transistors. A voltage monitor (VM) is connected to the storage capacitor  $C_{STORE1}$  of 4.7 μF. The purpose of the hysteretic VM is to connect a load  $R_{LOAD}$  to  $C_{STORE1}$  when a certain voltage  $V_H$  is sensed across  $C_{STORE1}$ , and to disconnect such load when  $V_{STORE1}$  becomes lower than a threshold  $V_L < V_H$ .  $C_{PUMP}$  is an external capacitor of 200 pF. The second load capacitor  $C_{STORE2}$  was set to 100 nF. The resistor  $R_P$  (~10MΩ, placed off-chip) placed between the input and output of  $M\_INV$  is necessary to correctly polarize the stage in its high-gain region, as it is done in Pierce-type oscillators. The three series pn diodes  $D_A$ ,  $D_B$  and  $D_C$  connected at the gate of  $M_1$  have the function of clamping voltages higher than the maximum gate-source voltage  $V_{GS}$  (1.8 V) allowed by the input transistors. The  $D_{ESD}$  diode clamps the negative peak of the oscillation at ~-0.5 V, hence the oscillation will have a positive DC offset. If  $D_{ESD}$  is replaced by three series diodes with same polarity, oscillating voltage would have ~0 V DC offset. It is possible to produce an oscillation with a negative DC offset as well, by replacing  $D_A$ ,  $D_B$  and  $D_C$  with a single diode and by replacing  $D_{ESD}$  with three series-connected diodes with the same polarity of  $D_{ESD}$ . The size of both  $M_1$  and  $M_2$  is  $W/L=4166$ ,  $M_1$  and  $M_2$  have an interdigitated layout, with  $W$  and  $L$  being the width and channel length respectively.

At the beginning of the oscillation, a small-signal analysis can be applied by linearizing the circuit around the bias point

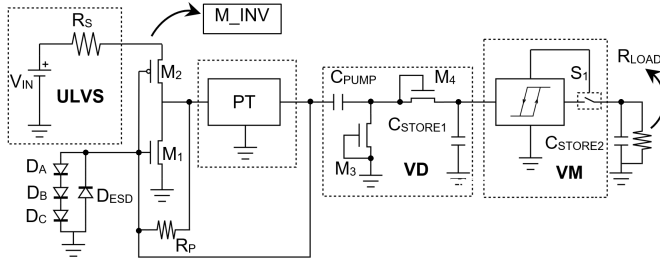


Fig. 6: Schematic of the PT-based step-up oscillator with a CMOS inverter as input stage and an integrated voltage monitor (VM).

$V_{IN}/2$  is the stage is designed to be symmetric with respect to  $V_{IN}$ . In this case, the voltage gain of the CMOS inverter is given by:

$$A_{V,INV}(V_{IN}) \cong -(g_{m,n} + g_{m,p}) \cdot (r_{out,n} // r_{out,p}), \quad (3)$$

where  $g_{m,n}$  and  $g_{m,p}$  are the transconductances of the nMOS and pMOS respectively, whereas  $r_{out,n}$  and  $r_{out,p}$  are the output resistances of the transistors. The operator “//” corresponds to the parallel combination of two resistances (e.g.  $R_A // R_B = (R_A^{-1} + R_B^{-1})^{-1}$ ). It is worth remarking that while the DeplMOS or JFETs in the CS stage work in the triode region above threshold [33], the low threshold devices in the CMOS inverter work in deep subthreshold region, given that the gate-to-source voltage does not exceed few tens of mV, whereas the threshold voltage is considerably higher. Furthermore, in the subthreshold region the transconductance has an exponential behavior with respect to the gate-source voltage, rather than the conventional linear dependence. In the CMOS inverter, the devices work in subthreshold region, where the current is due to diffusion effects rather than drift effects (as in above-threshold conditions). This implies that the  $g_m$  depends almost linearly on  $V_{IN}$  for  $V_{IN} < n \cdot V_t$ , where  $V_t = kT/q$  is the thermal voltage ( $\sim 26$  mV) and  $n$  is a number depending on the channel length; differently the dependence of  $g_m$  on  $V_{IN}$  is much lower when  $V_{IN} > n \cdot V_t$ , because all the electrons diffusing from source are collected at the drain. By means of Spice simulations, we found that  $n \cdot kT/q \approx 50$  mV, hence we expect a maximum of the voltage step-up ratio around 50 mV.

One of the most important advantages of a CMOS inverter stage compared to the conventional CS stage reported in [33], is that both transistors contribute to the voltage gain through their  $g_m$ . Fig. 7 shows a picture of the fabricated die. The size of the presented boost oscillator, is  $260 \mu\text{m} \times 170 \mu\text{m}$ .

Assuming that all blocks in Fig. 6 are unidirectional, then the return-ratio is a safe approximation of the loop gain.

In order to find the loop gain  $T(s)$ , it is necessary to find the voltage gains of the stages involved in the loop as well as the load effects.

As explained in [33], due to its high  $Q$ -factor, at a frequency extremely close to  $f_p$  the PT provides approximately  $\pi$  radians of phase shift, which combined with the additional  $\pi$  radians provided by the inverting input CMOS inverter stage, determine an overall phase shift of roughly  $2\pi$ . This necessary (but not sufficient) condition is required by the Barkhausen phase criterion to achieve oscillation. Thus, it can be safely stated that the presented system oscillates at a frequency

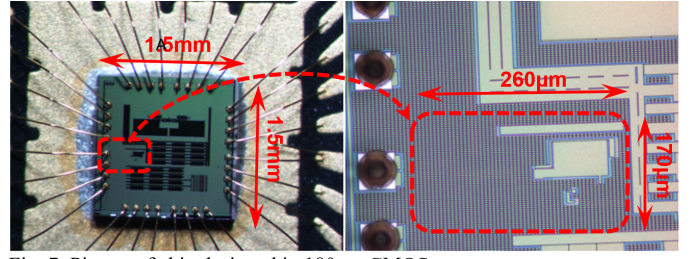


Fig. 7: Picture of chip designed in 180nm CMOS.

$f_{osc} = \omega_{osc}/2\pi$  extremely close to  $f_p$  if a sufficient loop gain is provided (Barkhausen gain criterion).

If a PT is able to vibrate at different modes, the prediction of the vibration mode from an AC linear analysis is not possible if the Barkhausen criteria are satisfied in more than a single frequency. As explained in [33], if the modes are separated by a factor two, as it happens in this case, the oscillation might start at the first mode, but the nonlinearities of the circuit will shift the oscillation to the second mode. Generally, when the Barkhausen criteria are satisfied in more than a frequency, the oscillation mode is conditioned by the shape of the non-linearity of the circuit. To cite an example, an abrupt saturation excited by a sinusoid of frequency  $f_0$  produces a square-wave with only odd higher-order harmonics ( $3 \cdot f_0, 5 \cdot f_0, 7 \cdot f_0$  etc.). Such kind of non-linearity cannot excite the mode of the PT placed at  $2 \cdot f_0$ . To cite another example, the absolute value function instead, produces a DC component plus even harmonics of the fundamental mode ( $2 \cdot f_0, 4 \cdot f_0, 6 \cdot f_0$  etc.). In order to predict the oscillation mode, when the Barkhausen criteria are satisfied in more than one frequency, a large-signal analysis is required. However, this goes beyond the scopes of the paper.

Equation (1) with the substitution  $s = j2\pi f$ , evaluated at  $f = f_p$ , assumes the form [32]:

$$A_{V,PT}(f_p) \cong -C_{IN2} / N \cdot C_O. \quad (4)$$

Equation (4) suggests that the PT voltage gain at the anti-resonance frequency is independent from the  $Q$ -factor, in a first approximation.

The input impedance of the PT at  $f = f_p$  is equal to [33]:

$$Z_{IN,PT}(f_p) \cong \frac{L_M C_{M2}}{R_M (C_{M2} + C_{IN2}) C_{IN2}} \cong R_M \left( \frac{Q C_{M2}}{C_{IN2}} \right)^2. \quad (5)$$

From (5) we see that  $Z_{IN,PT}$  is inversely proportional to  $R_M$  or directly proportional to  $Q$ .

In order to determine the loop gain, it is necessary to assess the load effect of the PT on the M\_INV stage. This is equal to:

$$A_{V,L}(f_p) \triangleq \frac{Z_{IN,PT}(f_p)}{Z_{IN,PT}(f_p) + r_{out,INV}}. \quad (6)$$

The loop gain  $T$  is then found by multiplying (3), (4) and (6), combined with (5).  $T$  is a monotonically increasing function of  $V_{IN}$ , hence there exists a minimum value of  $V_{IN}$ ,  $V_{IN,min}$ , which satisfies the Barkhausen gain criterion. The value of  $V_{IN,min}$  is found by solving the relation  $|T(V_{IN,min})|=1$ .

In this specific case, we modeled the system as a phase-shift oscillator, because the PT is a two-port linear system. Nevertheless, it can be modeled also as a negative resistance





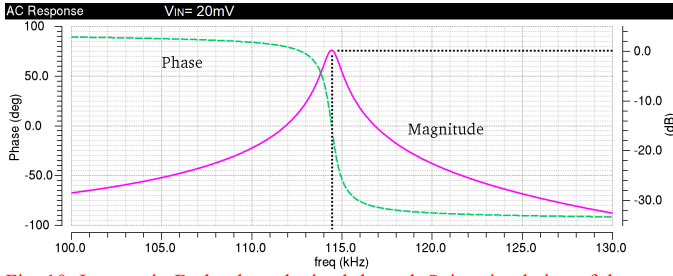


Fig. 10: Loop gain Bode-plots obtained through Spice simulation of the start-up circuit when polarized by a positive input voltage.  $V_{IN,min}$  is around 20 mV.

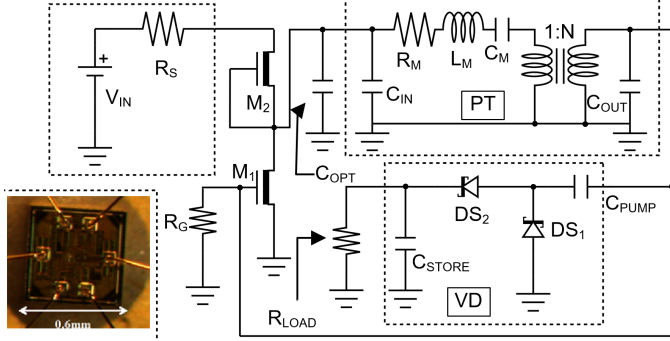


Fig. 11: Schematic of the PT-based boost oscillator with a common-source input stage. PT is the Piezoelectric Transformer, VD is the Voltage Doubler. The die also is shown.

600  $\mu\text{m} \times 600 \mu\text{m}$  including the pad frame. The circuit is an integrated version of the circuit based on discrete JFETs previously presented in [33], and is implemented in a STMicroelectronics 0.32  $\mu\text{m}$  CMOS technology. Two DeplMOS build the input inverting CS stage,  $M_1$  is the amplifier transistor and  $M_2$  is the load transistor. In our specific case,  $M_1$  and  $M_2$  have equal sizes with an aspect ratio  $(W/L) = 150$ . Higher  $W/L$  ratios lead to higher transconductances while lowering the overall output resistance, hence (3) achieves saturation. However, bigger transistors increase the load capacitance at the PT output port, lowering the loop gain, given that the gate capacitance is proportional to  $W \times L$ .  $C_{OPT}$  is a 400 nF capacitor necessary to further increase the loop gain, as explained in [33]. The VD rectifies and boosts the oscillation at the PT's output node  $V_{OUT,PT}$ . The diodes  $DS_1$  and  $DS_2$  in the VD are integrated Schottky diodes, whereas the resistance  $R_G$  ( $\sim 140\text{M}\Omega$ , placed off-chip) is used to provide a  $\sim 0$  V polarization at the gate of  $M_1$ . The storage capacitor  $C_{STORE}$  is a 4.7  $\mu\text{F}$  low-leakage polypropylene capacitor, whereas  $C_{PUMP}$  is a  $\sim 100$  pF integrated capacitor. The resistor  $R_G$  must be high enough to not to provide noticeable load effects at the PT output node, given that PTs are very load-dependent devices [43], but at the same time it must be capable to polarize the gate around 0 V. **Since PTs are very load dependent devices,  $R_G$  can be lowered down to 100  $\text{M}\Omega$ , without producing noticeable load effects at the PT's output port. Lower values can be used, but the value of  $V_{IN,min}$  increases accordingly.**

By looking at (5) and (6), it is possible to note that a higher  $Q$  does not necessarily bring to a higher magnitude of  $T$ , because if  $Q \rightarrow \infty$  then (6) asymptotically converges to 1, moreover the PT gain is independent on  $Q$  (see (4)). However, a higher  $Q$  PT gives the possibility to modulate  $Z_{IN,PT}$ , if

$Z_{IN,PT} \gg r_{OUT,CS}$  by inserting a capacitance  $C_{OPT}$  in parallel to the input port of the PT, increasing the value of (4), without affecting the value of (6). This is a main difference with respect to [33]: in our experiment, we used a PT with a  $Q$  almost double with respect to that used in [33]. As explained in [33] the value of  $|T(s=j\omega_p)|$  can be expressed as:

$$|T(s=j\omega_p)| = k \cdot V_{IN} \cdot \frac{\alpha Q / C_{IN2}^2}{\alpha Q / C_{IN2}^2 + Z_1} \cdot C_{IN2}, \quad (10)$$

where  $k$ ,  $\alpha$  are constants and  $Z_1$  is the output impedance of the previous stage (the CS in this specific case). If  $\alpha Q / C_{IN2}^2 \gg Z_1$ , then  $\partial T / \partial C_{IN2} > 0$ , hence  $C_{IN2}$  can be increased up to the value  $C_{IN2(MAX)} = \alpha Q / Z_1$ . If  $C_{IN2} > C_{IN2(MAX)}$ , then  $\partial T / \partial C_{IN2} < 0$ : high  $Q$ -factors PTs, with low values of  $C_{IN}$  are necessary in order to modulate its input impedance and voltage gain, hence the maximum value of  $C_{OPT}$  is:

$$C_{OPT(MAX)} = C_{IN2(MAX)} - C_{IN} = \frac{\alpha Q}{Z_1} - C_{IN} \quad (11)$$

Equations (3), (4), (5) and (6) can be used to model the behavior of the circuit, in which the differential parameters of  $M_{INV}$  have to be replaced with those of the CS in order to find the minimum input voltage to trigger an oscillation.

Considering the electromechanical parameters of the second mode reported in Table I,  $f_{OSC}$  is found to be 108.1 kHz, with an estimated  $C_G + C_{VD}$  around 10 pF.

The inverting voltage gain of the CS stage can be written as:

$$A_{V,CS}(s) \cong -g_{m1} \cdot r_{out1} // r_{out2} = -k_1 \cdot V_{IN} \cdot r_{OUT,CS}, \quad (12)$$

where  $g_{m1} = -k_1 \cdot V_{IN}$  is the transconductance of  $M_1$ , the parameter  $k_1$  ( $\text{A}/\text{V}^2$ ) is the current gain factor of the transistor,  $r_{OUT,CS} = r_{out1} // r_{out2}$  is the output resistance of the CS stage with  $r_{out1}$  and  $r_{out2}$  being the differential output resistances of  $M_1$  and  $M_2$  respectively. It is worth remarking that due to the very low voltages involved, the transistors are polarized in deep triode (linear) region. Since the devices are polarized in their linear region, the differential output resistance of each transistor equals the DC resistance seen from the source. The current drained from  $V_{IN}$  is equal to  $V_{IN} / (r_{out1} + r_{out2})$ , when the oscillation is about to start. If  $M_1$  and  $M_2$  have the same size, then we may safely state that  $r_{out1} = r_{out2}$ , hence the amplifying transistor  $M_1$  is polarized with a voltage  $V_{IN}/2$ .

The main difference between the CS stage made by DeplMOS (or JFETs) and the CMOS inverter stage is that the gain of the CS stage is proportional to  $V_{IN}$ , for  $V_{IN}$  up to several hundreds of mV, because the transistors operate above  $V_{TH}$  whereas in the CMOS inverter stage, such linear dependence is hold for  $V_{IN}$  up to few tens of mV, because the transistors operate much below  $V_{TH}$  as explained before.

### III. EXPERIMENTAL RESULTS

Measurements were performed in order to test the performances of the proposed circuits and to validate the effectiveness of the proposed approaches.

#### A. Step-up oscillator with CMOS inverter as input stage and integrated voltage monitor

A TTI EL302T Power Supply was used to emulate the TEG.

A FLUKE45 digital Multimeter was used to measure the average current drawn by the circuit. In the ammeter configuration, at very low voltage levels down to few tens of mV, such ammeter presents about  $12\ \Omega$  of series resistance  $R_S$ , more than one order of magnitude of some TEGs [44]. A Tektronix MSO 2024 Oscilloscope was used to visualize and sample the waveforms. Fig. 12 shows the schematic of the measurement set-up. A capacitor  $C_P = 1.2\ \mu\text{F}$  supplies the AC current while keeping the input voltage of the Circuit Under Test (CUT) at a constant level. The average drawn current can be measured through the ammeter.

Fig. 13 depicts the measured unloaded DC transfer function of the main CMOS inverter of the boost oscillator made up by  $M_1$  and  $M_2$  (see Fig. 6). The inverter was designed to have the highest voltage gain around  $0\ \text{V}$ . The DC transfer function was obtained by applying in input a symmetric duty-cycle triangular waveform at  $100\ \text{Hz}$  of frequency with  $15\ \text{mV}$  of DC offset and  $60\ \text{mV}$  of peak-to-peak amplitude generated by a Keysight 33220A Function Generator. The supply voltage was set to  $30\ \text{mV}$ . The unloaded voltage gain of the stage is about  $-0.7\ (\text{V}/\text{V})$  around  $0\ \text{V}$ , being at the least one order of magnitude higher than that of a typical common source made with JFETs or DeplMOS [33].

Fig. 14 depicts the measured characteristic of the hysteretic VM. The actual values of the threshold voltages are  $V_H = 1.65\ \text{V}$  and  $V_L = 1.1\ \text{V}$ . The capacitance  $C_{STORE2}$  was set to  $100\ \text{nF}$  in order to minimize the charge-sharing effect between  $C_{STORE1}$  and  $C_{STORE2}$ . The load connected to  $C_{STORE2}$  was  $R_{LOAD} = 1\ \text{M}\Omega$ , roughly corresponding to the intrinsic consumption of recent energy harvesting circuits such as the one in [45]. The  $\Delta V_{STORE} = 150\ \text{mV}$  in Fig. 14, depends on the fact that the boost oscillator is switched from an unloaded mode to a loaded mode ( $\sim 1\ \text{M}\Omega$ ). In addition, the value of the series resistance  $R_S$  does not affect  $V_{IN,MIN}$  given that at  $15\ \text{mV}$  the average current consumed by the circuit is only  $I_{IN} \cong 5\ \mu\text{A}$ , corresponding to an input impedance seen from the source equal to  $\sim 3\ \text{k}\Omega$ . The drained power from the source is  $P_S = V_{IN,MIN} \times I_{IN} = 75\ \text{nW}$ . It is worth noting that as the TEG scales,  $R_S$  increases. Since the presence of  $R_S$  produces a voltage divider together with the input impedance of the oscillator, as rule of thumb, we might safely neglect the voltage drop on  $R_S$  if this resistance lower than  $1/20$  of the input impedance of the oscillator, otherwise the effects of such resistance need to be taken into account considering the effective voltage drop on it.

Fig. 15 depicts a view of the start-up of the system achieved with a Peltier Cooler Multicomp MCPE1-03108NC-S [44] used as a TEG. One side of the TEG was put in contact with a cup containing hot water at  $\sim 60^\circ\text{C}$ , whereas the other side was exposed to an environmental temperature of around  $25^\circ\text{C}$ . No load was connected to  $C_{STORE}$  (buffered node). The system starts oscillating at  $V_{IN} = 15\ \text{mV}$ .  $V_{STORE1}$  reaches about  $350\ \text{mV}$  when  $V_{IN}$  is about  $25\ \text{mV}$ . This start-up voltage is higher than the declared minimum start-up voltage of recent buck-boost converters like that in [45] or the TI bq25504 [46]. Moreover, we remark that at  $V_{IN} = 15\ \text{mV}$  with multistage charge-pumps or Villard/Greinacher rectifiers the output voltage  $V_{STORE1}$  can be potentially increased even at such low

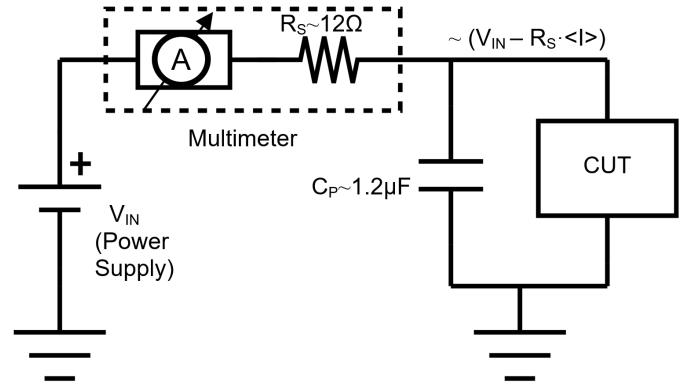


Fig. 12: Schematic of the measurement set-up of the Circuits Under Test, used for both the designed circuits.

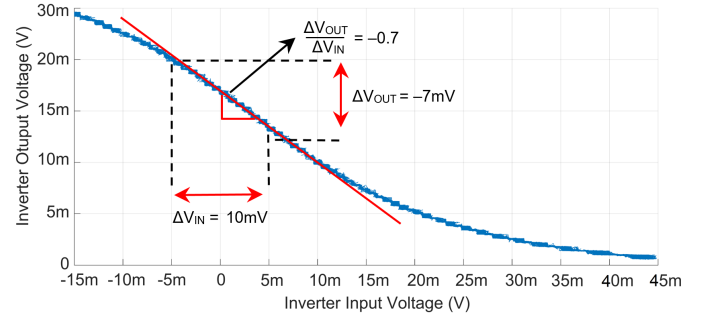


Fig. 13: Measured unloaded DC Voltage transfer characteristic of the CMOS inverter input stage of PT-based step-up oscillator.

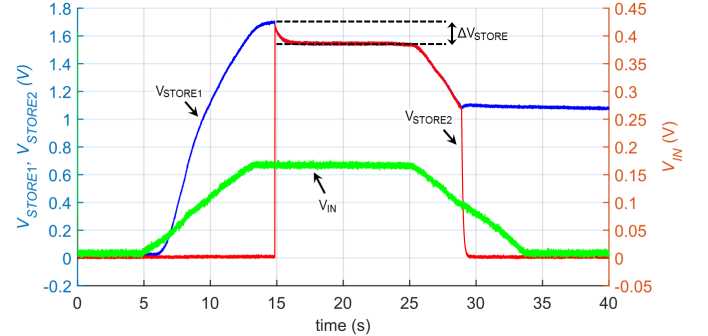


Fig. 14: Start-up of the step-up oscillator with the CMOS inverter as input stage: input and output voltage, and characteristic of the VM. The two threshold voltages are respectively  $V_H = 1.65\ \text{V}$  and  $V_L = 1.1\ \text{V}$ .  $R_{LOAD} = 1\ \text{M}\Omega$ .

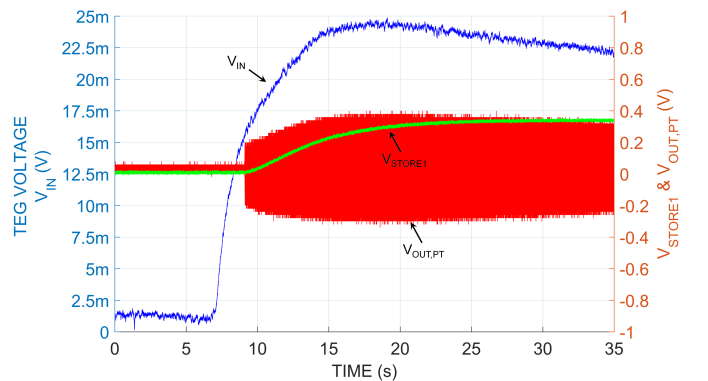


Fig. 15: Start-up of the step-up oscillator based on CMOS inverter obtained with a TEG Multicomp MCPE1-03108NC-S. The system starts oscillating at  $15\ \text{mV}$ . No-load connected on  $C_{STORE}$ .

input voltage levels [18] [19]. Concerning [46], it is worth remarking that the converter is able to harvest energy from sources down to  $80\ \text{mV}$  only once started, but the minimum input voltage from an off-state is  $330\ \text{mV}$ . Moreover, through



Spice Simulations we noted that, if the NatMOS are replaced with standard- $V_{TH}$  CMOS devices with same size and channel length, the minimum voltage to start an oscillation is around 250 mV, hence using low-threshold devices is one of critical points of the presented solution. However, if the low-voltage source is able to provide output voltages higher in the range of 250-300 mV, it might have more sense to implement a low-voltage ring oscillator driving a multi-stage charge pump.

The presented system is able to work also if a negative TEG voltage is provided without any need to switch its polarity at the input of the voltage booster (see Fig. 16). This happens because the input inverter is made of low-threshold MOSFETs whose channel is still conducting when supplied by a small negative voltage.

In this specific case,  $M\_INV$  (see Fig. 8) has a different behavior: the signal from the input of  $M\_INV$  to its output is experiences a  $\sim 0$  degrees of phase shift. This happens because the transconductances are negative: since the polarization of  $V_{IN}$  is reversed, the current inside the transistors is reversed as well. The increase in the  $V_{GS}$  injects more electrons in the channel that can be collected at the source, hence a positive  $\Delta V_{GS}$  produces a negative  $\Delta I_{DS}$ ; hence according to (3) the gain of the stage is positive. The oscillation is triggered in the PT capacitive window at a frequency  $f_{OSC} < f_s$ , because in this window the phase shift of the PT voltage gain is  $\sim 0$ .

By looking at Fig. 16, we see the circuit start-up achieved with a negative TEG voltage: 800 mV at the  $V_{STORE1}$  node are obtained when  $V_{IN} = -50$  mV. Moreover, the system oscillates at the first mode of the PT as depicted in Fig. 17: we are able to note also that the oscillation waveforms at the PT input and output port respectively are mostly in phase: an oscillation is maintained when  $V_{IN} = -8$  mV. However, in this condition, the oscillation has only 130 mV of peak-to-peak amplitude.

Fig. 18a depicts the measured average current  $I_{IN}$  drained by the source as well as the measured voltage  $V_{STORE1}$  for different values of  $V_{IN}$ , whereas Fig. 18b depicts the asymptotic voltage step-up ratio and the oscillator input impedance  $Z_{IN,OSC}$  seen from the source. A maximum voltage step-up ratio of  $\sim 17.5$  is obtained around 40 mV.

$Z_{IN,OSC}$  is in the range of few  $k\Omega$ , meaning that the oscillator based on the CMOS inverter, differently from the version with the CS stage, is also suitable for more resistive power sources such as UHF rectennas [47]. As a demonstration of the statement, we performed an additional measurement with the power supply with a series resistor  $R_S = 3.3$   $k\Omega$ , emulating a rectenna. In this case the circuit self-starts around  $V_{IN} = 50$  mV, given that most of the voltage drop is located on  $R_S$ . The circuit has an input impedance  $Z_{IN} \cong 1$   $k\Omega$ , that is a value congruent with Fig. 18b. When  $V_{IN} \cong 220$  mV, the circuit is polarized with  $V_{IN2} \cong 50$  mV, providing  $\sim 1$  V at the node  $V_{STORE1}$  (Fig. 19), with a drained current  $I_{IN} \cong 48$   $\mu A$ .

Concerning Fig. 18, it is worth noting that when  $V_{IN} \cong 150$  mV, the amplitude of the oscillation saturates due to the presence of the three protecting diodes  $D_A$ ,  $D_B$  and  $D_C$  (see Fig. 6) at the gate of the main inverter, hence the increase of  $V_{IN}$  has almost no effect on the output voltage, given that the oscillation amplitude cannot increase. Moreover, for negative

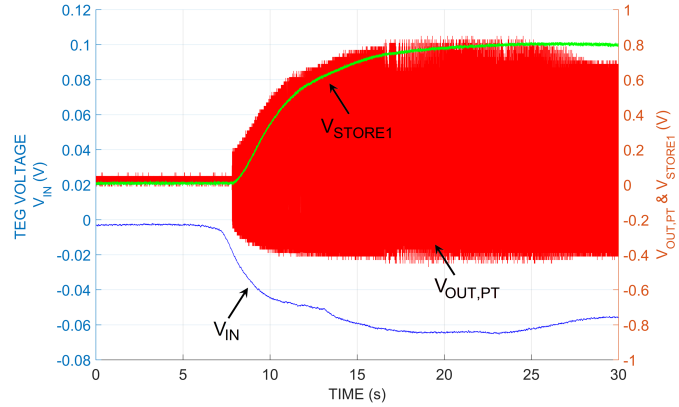


Fig. 16: Start-up of the step-up oscillator with a CMOS inverter input stage obtained with a TEG Multicomp MCPE1-03108NC-S for negative input voltages. The system starts oscillating at  $-8$  mV. A moving average was applied to  $V_{IN}$  in order to remove the oscilloscope noise.

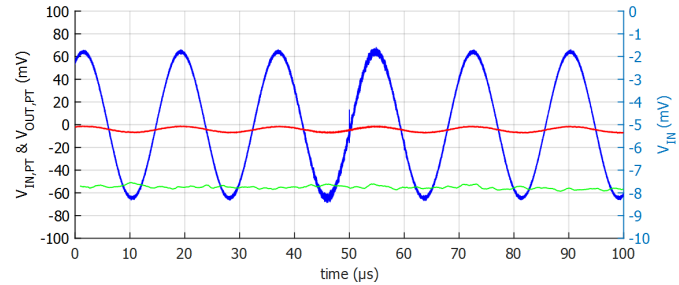


Fig. 17: Input and output oscillation at the PT ports, for  $V_{IN} \cong -8$  mV. The signals are almost in phase.

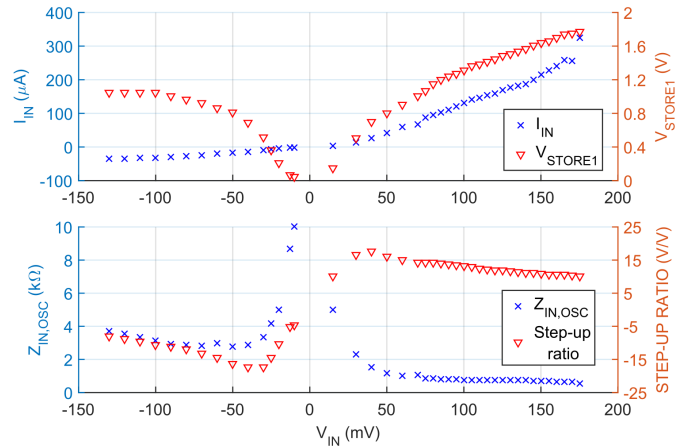


Fig. 18 (a)  $V_{STORE1}$  and average current  $I_{IN}$  sourced from the TEG as a function of the TEG voltage; (b) oscillator Input Impedance  $Z_{IN,OSC}$  and Voltage Step-Up ratio.

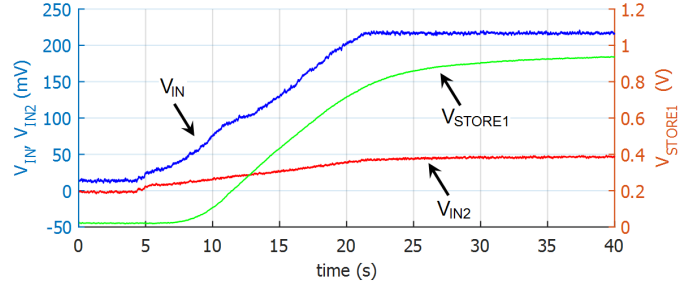


Fig. 19  $V_{STORE1}$ ,  $V_{IN}$  and  $V_{IN2}$  when a  $R_S = 3.3$   $k\Omega$ .  $V_{STORE1} \cong 1$  V when  $V_{IN} \cong 220$  mV (emulated rectenna).

values of  $V_{IN}$ , the maximum value of  $V_{STORE1}$  is around 1.05 V, when  $V_{IN} = -100$  mV. However,  $V_{STORE1} = 1$  V is obtained when  $V_{IN} = -80$  mV.

Fig. 20 depicts the behavior of the VM when a load  $R_{LOAD}$  of approximately 120 k $\Omega$  is connected in parallel to  $C_{STORE2}$  (100 nF). This is a worst-case in terms of load conditions and it accounts for more complex circuits, e.g. the in-rush current of the conventional DC/DC of Fig. 2. As in Fig. 20, the law describing the voltage across  $C_{STORE2}$  can be written as  $V_{STORE2} = 1.7 \cdot \exp(-\Delta t / R_{LOAD} \cdot C_T) = 1.1V$ , hence the time interval  $\Delta t$ , between two consecutive actions of the VM is given by  $\Delta t = \ln(1.7/1.1) \cdot R_{LOAD} C_T \cong 0.25s$  (Fig. 20 shows  $\sim 0.3s$ ). The capacitance  $C_T = C_{STORE1} + C_{STORE2}$ . In fact, the system is still able to sustain such worst-case load for about 0.3 s before the VM detaches it from the boost oscillator. If higher stored energy is needed,  $C_{STORE1}$  can be increased and  $\Delta t$  in (13) will increase accordingly. In a real application, this consumption represents the current drawn by the conventional power converter when it's being turned on and before it has started extracting power.

Fig. 21 depicts the oscillation at the PT input port (CMOS inverter output port) and PT output port (M\_INV input port) obtained at  $V_{IN} = 40mV$ . The phase difference is about  $-175^\circ$ , meaning that the PT is working close  $f_p$ , as explained in [33].

Concerning the efficiency of the circuit, it is not possible to find the efficiency as a function of  $V_{IN}$  at the node  $V_{STORE1}$  because the VM keeps the load disconnected from the boost oscillator. Moreover, the node  $V_{STORE2}$  is connected to the load when  $V_{STORE1} \cong 1.65V$  occurring at  $V_{IN} \cong 140mV$  (see Fig. 18a). The increase of  $V_{IN}$  does not affect significantly the voltage at the node  $V_{STORE1}$ , because the oscillation amplitude saturates due to the presence of the protecting diodes (see Fig. 6). However, in Table II we report some measured efficiency values for different loads connected at the  $V_{STORE2}$  node. The letters B and C indicate the phases as in Fig. 9.

### B. Step-up oscillator with common-source amplifier stage

The measurements were performed according the set-up depicted in Fig. 12. The measured steady-state oscillation frequency is 103.7 kHz, whereas the estimated frequency from an AC analysis 108.1 kHz. The discrepancy is due to the fact that in the steady-state operation the load seen from the PT is different from the load seen at the beginning of the oscillation, that is the condition in which the oscillation frequency was estimated. Typical oscillation waveforms for the PT-based step-up oscillator with normally-on transistors at the PT input port (CS output port) and PT output port (CS input port) can be found in [33].

One thing worth of attention is that, although the current-voltage equations a n-type DeplMOS are identical to that of a n-type JFET (they both have a negative  $V_{TH}$ , or "pinch-off" voltage), there is one intrinsic difference. In JFET-based oscillators such as the one in [33], the peak value of the oscillation at the gate-source voltage node cannot exceed  $\sim 0.5V$ , otherwise the p-n junction between the gate and source of the JFET can be partially turned on, thus causing the loss of the transistor effect. As a matter of fact, in [33], the oscillation in order to be sustained has a DC offset value much lower than 0V: this implies that the output oscillation falls below the pinch-off voltage of  $M_1$  and once it

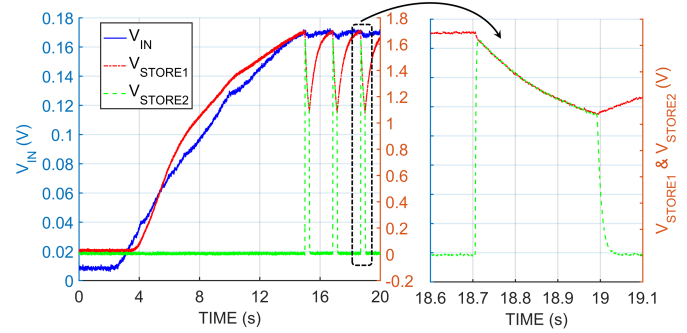


Fig. 20: Behavior of the VM in step-up oscillator based on a CMOS inverter when a 120k $\Omega$  load is connected in parallel to  $C_{STORE2}$ . When  $V_{STORE2}$  and  $V_{STORE1}$  are short-circuited, they get discharged according to the discharging law of a RC circuit.

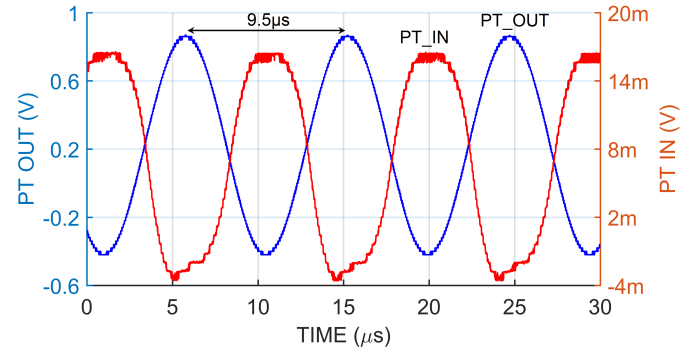


Fig. 21: Oscillation waveforms at the PT output port (CMOS inverter input) and PT input port (CMOS inverter output), when  $V_{IN} = 40mV$ . The filtering effect of the PT is noticeable.

TABLE II  
EFFICIENCY VALUES FOR DIFFERENT LOAD AT THE  $V_{STORE2}$  NODE

$V_{IN}$ (mV)	$I_{IN}$ (B) ( $\mu A$ )	$P_{IN}$ (B) ( $\mu W$ )	$I_{IN}$ (C) ( $\mu A$ )	$P_{IN}$ (C) ( $\mu W$ )	$R_{LOAD}$ ( $M\Omega$ )	$V_{OUT}$ (B/C) (V)	$P_{OUT}$ (C) ( $\mu W$ )	$Eff$ (C) (%)
140	210	29.4	300 $\mu$	42	0.9	1.7/1.5	2.5	6
140	210	29.4	240 $\mu$	33.6	5	1.7/1.65	0.55	1.7
140	210	29.4	220 $\mu$	30.8	10	1.7/1.68	0.28	0.9
140	210	29.4	215 $\mu$	30.1	20	1.7/1.7	0.145	0.5
140	210	29.4	212 $\mu$	29.7	50	1.7/1.7	0.058	0.2
140	210	29.4	$\sim 210\mu$	$\sim 29.4$	150	1.7/1.7	0.019	0.07

is in steady-state, the minimum voltage (or power) to sustain oscillation is considerably lower than the minimum voltage (or power) to start the oscillation (hysteretic behavior), because the input circuit behaves as an AB amplifier (the current is off during a portion of the oscillation period, see [33]). With DeplMOS, since there is not any upper boundary limit to the gate-source voltage, the oscillation has a  $\sim 0V$  offset (see Fig. 22) and it might slightly fall (or maybe will not fall at all) below the threshold voltage of the transistor (according to the oscillation amplitude which in turn depends on  $V_{IN}$ ): in this case the voltage (or power) required to start the oscillation is essentially the voltage (or power) to sustain the oscillation (no hysteretic behavior). In other words, a JFET-based oscillator as a higher robustness versus time variations of temperature gradients, compared to a DeplMOS-based oscillator.

Fig. 22 shows the start-up of the circuit.  $V_{IN}$  is slowly ramped up to 48 mV. Differently from the CS stage reported in [33], the oscillation has not any DC offset. The circuit starts oscillating at  $V_{IN} \cong 36mV$ . In such conditions, with no load connected in parallel to  $C_{STORE}$ ,  $V_{STORE}$  reaches about 2.7 V, corresponding to a stored energy in  $C_{STORE}$  equal to:



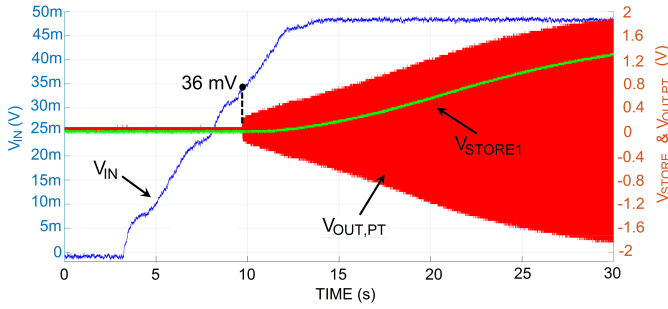


Fig. 22: Start-up of the step-up circuit with a CS input stage and  $R_S=12\Omega$ .  $V_{IN}$  is referred to the left y-axis, whereas  $V_{STORE}$  and  $V_{OUT,PT}$  are referred to the right y-axis. The start-up occurs at  $V_{IN}=36\text{mV}$ . No load on  $V_{STORE}$ .

$$E_{STORE} = \frac{1}{2} C_{STORE} (V_{STORE})^2 \cong 17\mu\text{J}. \quad (13)$$

The average current consumption of the circuit in steady state operation is approximately equal to  $500\ \mu\text{A}$  at  $V_{IN}=36\text{mV}$ . As a consequence, the voltage drop on  $R_S$  is equal to  $6\text{mV}$ . When the power supply (or TEG, see Fig. 23) is directly connected to the circuit, the start-up was found to be around  $31\text{mV}$ , which is a value compatible with the results obtained in [33] with the aid of an inductor, but we remark that in this case the insertion of any magnetic component was avoided, and that we exploited an additional capacitance  $C_{OPT}$  in parallel to the PT input port, to modulate the total loop gain, in accordance with (10).

In this case, we used the aforementioned Peltier Cooler MCPE103108NC-S  $18.8\text{W}$  [44]. This device has an internal resistance of  $0.35\ \Omega$  leading to a negligible voltage drop. Fig. 23 depicts the start-up of the circuit achieved with the aforementioned Peltier device. As it can be noted, the start-up value is still found to be about  $31\text{mV}$ , confirming that lower values of  $R_S$  yield lower activation voltages.

Fig. 24 depicts the measured  $V_{STORE}$  obtained when  $V_{IN}=V_{IN,MIN}=31\text{mV}$  as a function of the load and also the power delivered to the load defined as  $P_{LOAD}=V_{STORE}^2/R_{LOAD}$ . The maximum output power occurs at around  $38\text{M}\Omega$  of load; however, once again the purpose of such boost circuit is not to maximize the output power, but minimizing the operating input voltage. This type of circuit is intended for driving loads in the range of  $100\text{M}\Omega$ , like for example, an ultra-low power voltage monitor, or the gate of a power switch in a standard power converter, as it was done in [35]: once the main converter is turned on, the boost oscillator is switched-off.

Fig. 25a depicts the measured current drawn from the source  $I_{IN}$ : at input voltages around  $70\text{mV}$ , currents in the range of  $1\text{mA}$  are flowing from the source. Fig. 25a depicts also the asymptotic value of  $V_{STORE}$  as a function of the input voltage  $V_{IN}$ . In Fig. 16, an equivalent  $R_{LOAD} \cong 10\text{M}\Omega$  was connected, due to the oscilloscope probe. We remark that an ultra-low power VM presents a significantly higher load resistance (such as  $\sim 180\text{M}\Omega$  for the VM presented in this work). Then, the obtained results should be intended as in a worst-case condition. Fig. 25b is strictly related to Fig. 25, and depicts both the Voltage step-up ratio reaching almost 60 at  $V_{IN}=85\text{mV}$ , whereas the average oscillator input impedance seen from the source is around  $70\ \Omega$ , confirming that the devices are working in their deep triode region.

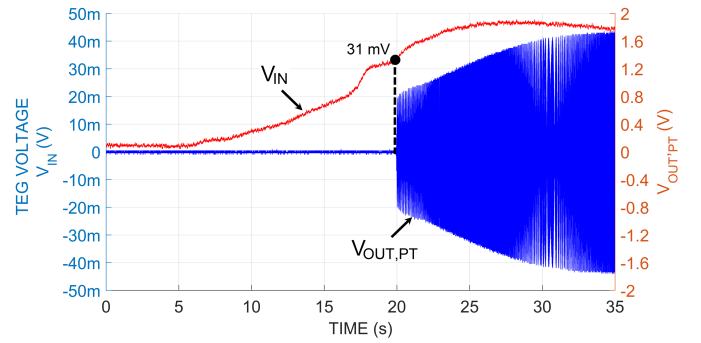


Fig. 23: Start-up of the step-up converter with a CS input stage with a Peltier Cooler used as a TEG (MCPE1-03108NC-S). The start-up voltage is  $\sim 31\text{mV}$ .

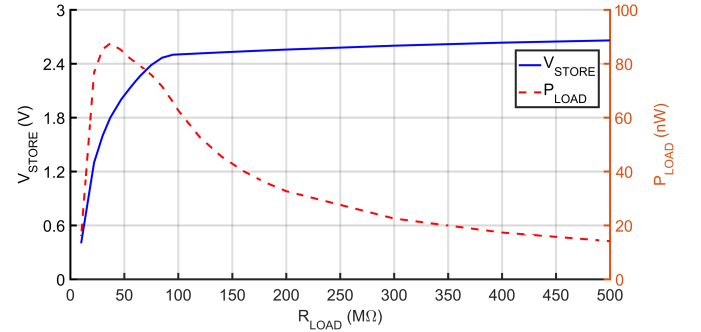


Fig. 24:  $V_{STORE}$  and  $P_{LOAD}$  as a function of  $R_{LOAD}$  for the step-up oscillator with a CS input stage.  $V_{STORE}$  is referred to the left y-axis, whereas  $P_{LOAD}$  is referred to the right y-axis. Graph obtained with  $V_{IN}=31\text{mV}$  (power supply).

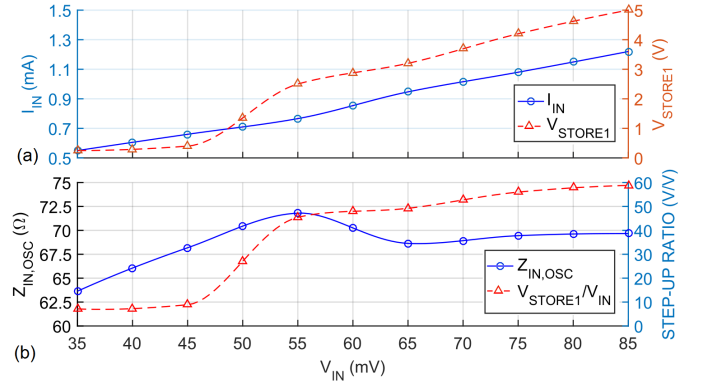


Fig. 25 a)  $I_{IN}$  and  $V_{STORE}$  for the circuit with a CS input stage with a worst case  $R_{LOAD} \cong 10\text{M}\Omega$  placed between  $V_{STORE}$  and ground nodes. The current consumption ranges from  $\sim 0.5\text{mA}$  up to  $1.2\text{mA}$ , whereas  $V_{STORE}$  ranges from  $250\text{mV}$  up to  $5\text{V}$  when  $V_{IN}$  varies from  $35$  to  $85\text{mV}$ . b) Input impedance  $Z_{IN,OSC}$  and voltage Step-up ratio of the oscillator with a CS input stage.

Fig. 26 depicts the measured efficiency for two cases: a) the efficiency is measured for a single input voltage ( $V_{IN,MIN}$ ), and the  $R_{LOAD}$  is varied from  $10\text{M}\Omega$  up to  $500\text{M}\Omega$ . The maximum efficiency is found for loads lower than  $50\text{M}\Omega$ , however such circuit is intended for higher loads; b) the load is kept fixed at  $10\text{M}\Omega$ , and then  $V_{IN}$  is varied from  $35\text{mV}$  to  $85\text{mV}$ : in this case, the efficiency increases as  $V_{IN}$  increases, because the oscillation grows in amplitude and might switch-off the transistors.

That the presented system is intended for start-up purposes, hence its efficiency is poor. In [35] the system was used to activate an inductor-based fixed-frequency boost DC/DC converter with efficiencies higher than  $40\%$  in the range of  $30\text{--}50\text{mV}$ , for loads of  $30\text{ k}\Omega$  up to  $60\text{ k}\Omega$ . By comparison in [48], efficiencies up to  $\sim 60\%$  with output power in the  $\text{mW}$

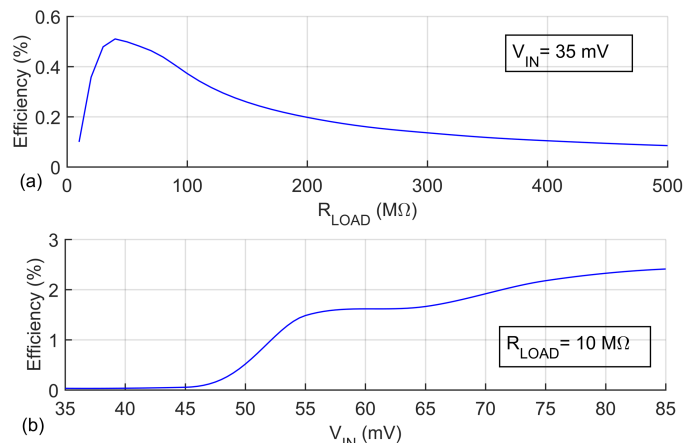


Fig. 26: a) Efficiency of the system as a function of the load for a fixed  $V_{IN}$ ; b) efficiency of the system as a function of the input voltage for a fixed  $R_{LOAD}$ .

range, are achieved exploiting a transformer-reuse technique. The presented topology in this work is fully compatible with a transformer-reuse technique, by using the same PT, after the start-up phase in a resonant power converter for microwatt applications.

#### IV. CONCLUSIONS

This work has presented two circuit topologies of ICs for EH purposes from ultra-low DC sources exploiting PTs for kick-starting without need of any battery.

The circuit based on the CMOS inverter requires common low- $V_{TH}$  MOSFETs instead of negative- $V_{TH}$  MOSFETs used in the CS inverting stage and is able to work with double polarity input signals without the need of switching matrixes. As a matter of fact, start-up values down to +15 mV and -8 mV were measured, with maximum output voltages of 1.75 V for positive polarization (achieved at 140 mV), and 1.05 V (achieved at -100 mV). Moreover, the design of the system was completed through a VM consuming only 15 nW.

Another advantage is that the boost oscillator based on the CMOS inverter has higher input impedance, in the range of several k $\Omega$ , with respect to the inverting CS stages made up by DeplMOS polarized in their deep triode region. This lowers the average power consumed from the source and makes the circuit thus less sensitive to variations of the parasitic series resistance of the harvesting source. The higher input impedance makes this type of oscillator also suitable for harvesting from rectennas as well, given that typical equivalent output resistance of such harvesters is in the order of several k $\Omega$  as well. This assertion has been demonstrated with measurements performed with a 3.3 k $\Omega$  series resistance. The CS-based step-up oscillator is not suitable for rectennas because of its low input impedance (tens of  $\Omega$ ), that would lead the Rectenna working close to its short circuit condition.

The circuit with the input stage made by a CS stage formed by two DeplMOS achieves a minimum start-up voltage of about 31 mV without using any magnetic components, and can be successfully supplied by a standard miniature Peltier Cooler used as a TEG. One of the advantages of this circuit is the use of integrated Schottky diodes in the rectifier, with lower voltage drops compared to conventional p-n diodes or diode-connected MOSFETs, and with higher voltages in the

storage capacitor. Step-up ratios as high as 58 were obtained, when the input voltage is around 85 mV. On the other side, the main drawback is the average current drained by the source (around 1 mA at 70 mV of input voltage), which is reflected in a very low input impedance, between 60  $\Omega$  and 70  $\Omega$ , given that the transistors in CS stage are polarized in their deep triode region. The low oscillator input impedance seen from the source makes this circuit more sensitive to the parasitic series resistance of the energy harvesting transducer. This makes this particular implementation more suitable for TEG harvesters with output series resistances in the range of few Ohms.

In the experiment, a discrete PT was used. However, the adopted PT was conceived to handle power levels in the order of several W. A tailored PT design for the voltage step-up application, possibly implemented with MEMS technologies for wafer-level or package-level integration, focused in shrinking its dimensions, might bring interesting results. Moreover, according to the electrodes layout and choice of operating mode (thickness or “33” mode, extensional or “31”, or both as it happens in Rosen-type PTs) depending on the aspect ratio of the PT, it is possible to design the frequency at which the system should oscillate. Such design should aim at implementing a PT capable of high voltage gains while handling powers in the range from few  $\mu$ W to few mW, which is the typical power managed in EH systems.

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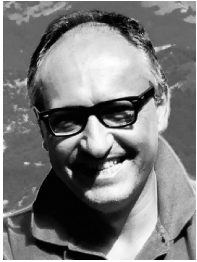
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