



# Seismic Design and Performances of Frame Structures Connected to a Strongback System and Equipped with Different Configurations of Supplemental Viscous Dampers

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The paper investigates the dynamic behavior of structural systems obtained by connecting a moment-resisting frame structure with a vertical rigid truss pinned at the base, known in literature as “strongback,” and equipped with added fluid-viscous dampers. The strongback, designed in order to remain in the elastic field under strong seismic ground motion, acts as a mast by imposing to the structure a linear lateral deformed shape. By regularizing the lateral drift profile of the structure, the strongback limits undesired effects such as weak-storey mechanisms, damage concentration and residual drifts. In addition, when supplemental dampers are inserted in the structure, a considerable amount of energy can be dissipated, thus reducing the peak seismic response. The aim of the work is twofold: i) to provide analytical formulations for the preliminary design of added dampers based on the Generalized Single Degree Of Freedom (GSDOF) concept, and ii) to evaluate the increase in energy dissipation capabilities for selected dampers configurations thanks to the presence of the strongback. The formulas are developed for different configurations of added viscous dampers: dampers inserted within the frame between all or selected consecutive storeys (inter-storey placement) and dampers located at the base of the strongback to realize a rigid “dissipative tower.” The effectiveness of the dampers configurations is evaluated through dynamic time-history analyses.

**Keywords:** strongback, frame structure, generalized single degree of freedom, supplemental viscous dampers, equivalent damping ratio

## INTRODUCTION

Traditional steel frame structures have typically limited resistance towards seismic events, with the tendency to form soft- and weak-storey mechanisms, damage concentration, P- $\Delta$  effects and residual drifts. These issues can be particularly relevant for conventional concentrically-braced frames (Khatib et al., 1988; Rai and Goel 1997; Sabelli 2001; Tremblay 2003; Hines and Appel 2007; Uriz and Mahin 2008; Hines et al., 2009; Lai et al., 2010; Chen and Mahin 2012).

In order to mitigate such undesired effects, some specific considerations during the design phase of frame structures should be made following the Performance-Based Seismic Design (PBSD) approach (Bertero and Bertero 2002) together with the adoption of new technological solutions.

Among others, possible solutions are represented by the use of enhanced dissipative braces such as: steel bracing systems (Foutch et al., 1987; Whittaker et al., 1990), buckling-restrained braces and zippers (Khatib et al., 1988; Tirca and Tremblay 2004; Kiggins and Uang 2006, Yang et al., 2008; Yang et al., 2009; Yang et al., 2010; Stavridis and Shing 2010), rocking or uplifting systems (Clough and Huckelbridge 1977; Kelly and Tsztsoo 1977; Uriz and Mahin 2008; Deierlein et al., 2011), viscous dampers (Constantinou et al., 1998; Christopoulos and Filiatrault 2006; Silvestri et al., 2010; Castaldo and De Iuliis 2014; Palermo et al., 2018), hysteretic devices (Trombetti et al., 2004; Palermo et al., 2014; Palermo et al., 2015) or shear links (Nuzzo et al., 2018).

An alternative solution is to couple the frame system with strong rocking walls, able to reduce the damage concentration in the frame, as well as to uniform the deformed shape under seismic action, thus reducing the concentration of drifts in the lower storeys (Alavi and Krawinkler 2004). Analytical and numerical studies have been proposed by several researchers to investigate the seismic behavior of strong rocking wall-frame systems, consisting of an external rigid wall pinned at the base and rigidly connected at each floor to the frame (Qu et al., 2011; Wu et al., 2016). Qu et al. (2012) proposed a dissipative solution combined to strong rocking walls through horizontal yielding dampers able to dissipate the seismic-induced energy through the relative vertical displacement between the frame and the wall at each floor level. Analytical formulations of the wall-frame-damper system considered as a continuous model were proposed in Sun et al. (2018). Pan et al. (2018) proposed the use of infill rocking wall-frame systems to reduce damage concentration of frame structures subjected to seismic action. Their effectiveness has been verified through quasi-static cyclic tests on a prototype structure. The same system has also been studied in Wu et al. (2017) with specific focus on its reparability after damage. Wang et al. (2018) studied the seismic response of infill rocking walls equipped with yielding dampers at the base, referred to as plastic-hinge-supported wall system.

The insertion of viscous dampers at the base of an external pinned rigid truss structure was proposed (Roia et al., 2013) and patented as the “dissipative tower” concept. A further work by Gioiella et al. (2018) investigated the dynamic behavior of such structural solution by means of complex damping theory and provided a case-study application. The results proved that the solution is effective in terms of energy dissipation.

A concept similar to strong rocking walls was proposed by Lai and Mahin (2014) and referred to as “strongback system.” In detail, the first idea has been developed as a rigid steel bracing within the frame structure, to avoid weak- and soft-storey mechanisms and thus reduce the concentration of damage induced by seismic action in the first storeys. Numerical and experimental studies on the same system have been developed by Simpson and Mahin (2018) and further implemented in the work done by Toorani (2020). Palermo et al. (2018a) adapted the same idea on an external pinned-based rigid steel braced truss structure linked to the frame at each floor and investigated the coupled static behaviour of a frame structure connected to this external strongback system. The aim of the work was to obtain expressions

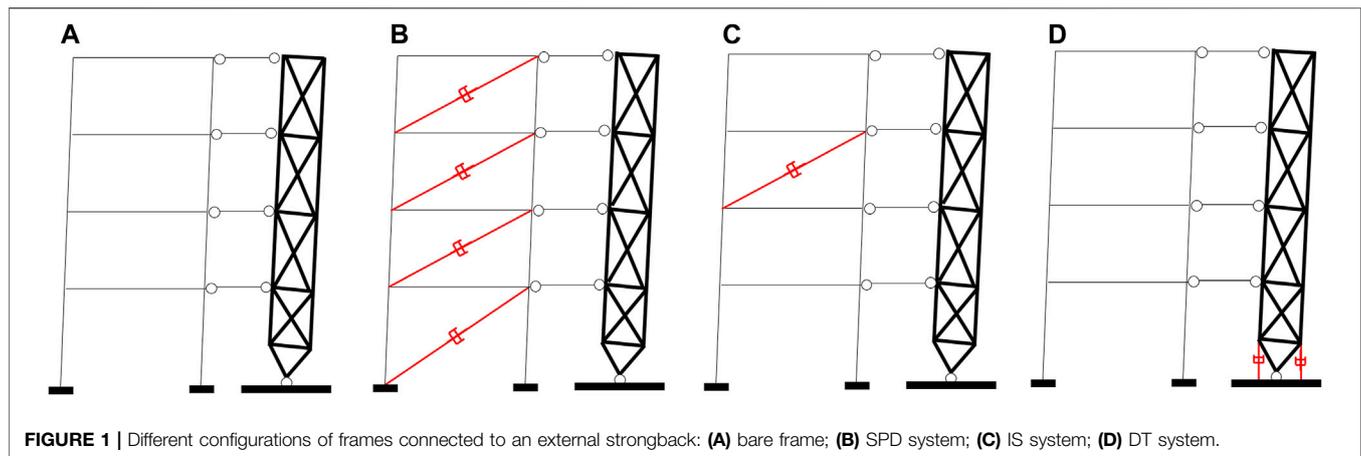
of the mutual actions exchanged between the frame and the external strongback system. The results indicated that both the distribution and the amplitude of such mutual actions are strongly influenced by the beam-to-column stiffness ratio. In particular, in case of flexible beams, large mutual actions and internal frame forces concentrate in the lower storeys, while the upper storeys remain practically unloaded.

In this paper, the dynamic behavior of frames connected to an external strongback (SB) system and equipped with different configurations of supplemental viscous dampers is investigated. In addition to the dissipative tower configuration, another configuration of added viscous dampers is considered, namely viscous dampers inserted within two consecutive storeys of the frame (inter-storey placement). In detail, different configurations of inter-storey viscous dampers are studied, either at all storeys or only at selected ones. The aims are i) to provide analytical formulations for the preliminary design of the added dampers, based on the Generalized Single Degree Of Freedom (GSDOF) approach recently discussed in the work by Palermo et al. (2020), and ii) to evaluate the increase of energy dissipation capabilities for the considered dampers configurations due to the presence of the strongback. Indeed, the presence of the rigid strongback, which imposes a linear lateral deformed shape, is expected to induce a more uniform along-the-height energy dissipation, especially for non-uniform dampers configurations.

## PROBLEM FORMULATION AND OBJECTIVE

Since the aim of the study is to investigate the fundamental dynamic properties of frame structures connected to a strongback and equipped with supplemental viscous dampers, simplified idealized systems are here considered, similarly to the ones considered in the work by Palermo et al. (2018) and Palermo et al. (2020). Henceforth, the term “system” will be used to indicate the frame connected to the external strongback and equipped with supplemental dampers. For this specific purpose, it is convenient to idealize the strongback as an external rigid truss pinned at the base and connected to the adjacent frame through horizontal rigid trusses (**Figure 1A**). It is assumed that the strongback is composed by rigid elements so that it rotates around the pinned base as a rigid body.

Various configurations of supplemental viscous dampers are here investigated: i) diagonal viscous dampers inserted within the frame between all consecutive storeys with damping coefficients proportional to the storey’s lateral stiffnesses (**Figure 1B**), thus resulting in a Stiffness Proportional Damping (SPD) system (Trombetti and Silvestri 2004), ii) diagonal viscous dampers inserted within the frame in a unique floor, according to the so-called Inter-Storey damper placement (**Figure 1C**, such system will be referred to as IS system); iii) viscous dampers placed at the base of the strongback to realize the so-called rigid Dissipative Tower (**Figure 1D**, such system will be referred to as DT system). As previously mentioned, for the IS system the presence of the rigid strongback should be capable of enhancing the dissipative properties of the damping system with respect to



**FIGURE 1** | Different configurations of frames connected to an external strongback: **(A)** bare frame; **(B)** SPD system; **(C)** IS system; **(D)** DT system.

the corresponding systems without the strongback, that are characterized by quite reduced effectiveness (Palermo et al., 2020). In a DT system, the two viscous dampers located at the base of the strongback dissipate energy through a couple of damping forces acting along the vertical direction. The vertical velocity of the dampers can be determined from the rigid rotation of the strongback, as it will be clarified in the next section.

In this work, all added viscous dampers are assumed to have a linear force-velocity constitutive behavior and to be supported by infinitely stiff braces. Regarding these two assumptions, for design purposes, the following observations are made. It is common practice to first define the size of the dampers through a linear model, and then identify the non-linear mechanical properties of commercial manufactured dampers by means of specific energetic criteria (Christopoulos and Filiatrault 2006; Silvestri et al., 2010; Palermo et al., 2018). Similarly, the stiffness of the elastic brace connecting the viscous damper to the structure should be properly designed (Silvestri et al., 2010) to avoid detrimental effects (Castaldo and De Iuliis 2014).

As illustrated in **Figure 2**, the geometrical configuration of frames connected to the strongback can be described by assuming a system of coordinates  $(x, z)$  with origin at the base of the strongback (namely point O in **Figure 2**). The  $i$ -th storey is located at a height  $z_i$ . The lateral displacement of the  $i$ -th floor is indicated with  $u_i$ . It is assumed that the frame elements have a linear elastic behavior. The lateral stiffness of the  $i$ -th storey is denoted as  $k_i$ . The axial deformability of the frame members is neglected. Shear-type frames are considered. The kinematics of the strongback is completely described by the angle of rotation  $\theta$ . According to these assumptions, when subjected to a generic set of external lateral forces, the system will globally develop a linear lateral floor displacement profile. The linear deformed shape of the frame is guaranteed by mutual actions  $H_i$  which are exerted between the frame and the strongback. In the work by Palermo et al. (2018b), analytical expressions of the mutual actions were derived for different frame behaviors (namely moment resisting shear-type frames and hinged frames), different distributions of static external forces and storey lateral stiffnesses  $k_i$ .

Under the above assumptions, from a global dynamic equilibrium perspective, the system can be rigorously treated, without introducing approximations, as a GSDOF system

(Chopra 2001; Palermo et al., 2020), since its motion is governed and fully described by the rigid rotation angle of the strongback.

In the next sections, the dynamic behavior of the different damped systems with strongback is investigated with the purpose of determining their fundamental dynamic properties (natural frequency and damping ratio), their dissipative properties and their performances under earthquake excitation in comparison to systems without strongback.

## DYNAMIC EQUILIBRIUM OF UNDAMPED AND DAMPED FRAMES CONNECTED TO A STRONGBACK

### The Undamped System

In the present section, the attention is focused on the undamped system, namely the system without added dampers. In free vibration conditions, according to D'Alembert's principle (Chopra 2001), the dynamic system is in equilibrium, at each time instant, under the following dynamic actions: the inertia forces  $F_{I,i}$ , the elastic resisting forces  $F_{S,i}$  and the mutual actions  $H_i$  ( $i$  indicates the generic  $i$ -th storey). **Figure 2** displays the system cut vertically along the connections between the frame and the strongback to evidence these mutual forces.

As demonstrated in the work by Palermo et al. (2020), the dynamic equilibrium of a generic (either undamped or damped) system can be studied by introducing two different GSDOF system, namely a translational GSDOF (derived by imposing the global translational equilibrium, referred to as Generalized Translational Oscillator–GTO) and a rotational GSDOF (obtained by imposing the global rotational equilibrium at the base, referred to as Generalized Rotational Oscillator–GRO). The approach requires an assumed deformed shape vector  $\mathbf{d}$  representative of a given mode shape. However, when the system is connected to a rigid strongback, the shape of the lateral deformed shape is known and defined by a linear vector  $\mathbf{d}$ . In this case, as demonstrated in the work by Palermo et al. (2020), the two GSDOF systems are characterized by the same fundamental dynamic properties, e.g. period of vibration and damping ratio (see **Supplementary Appendix S1**).

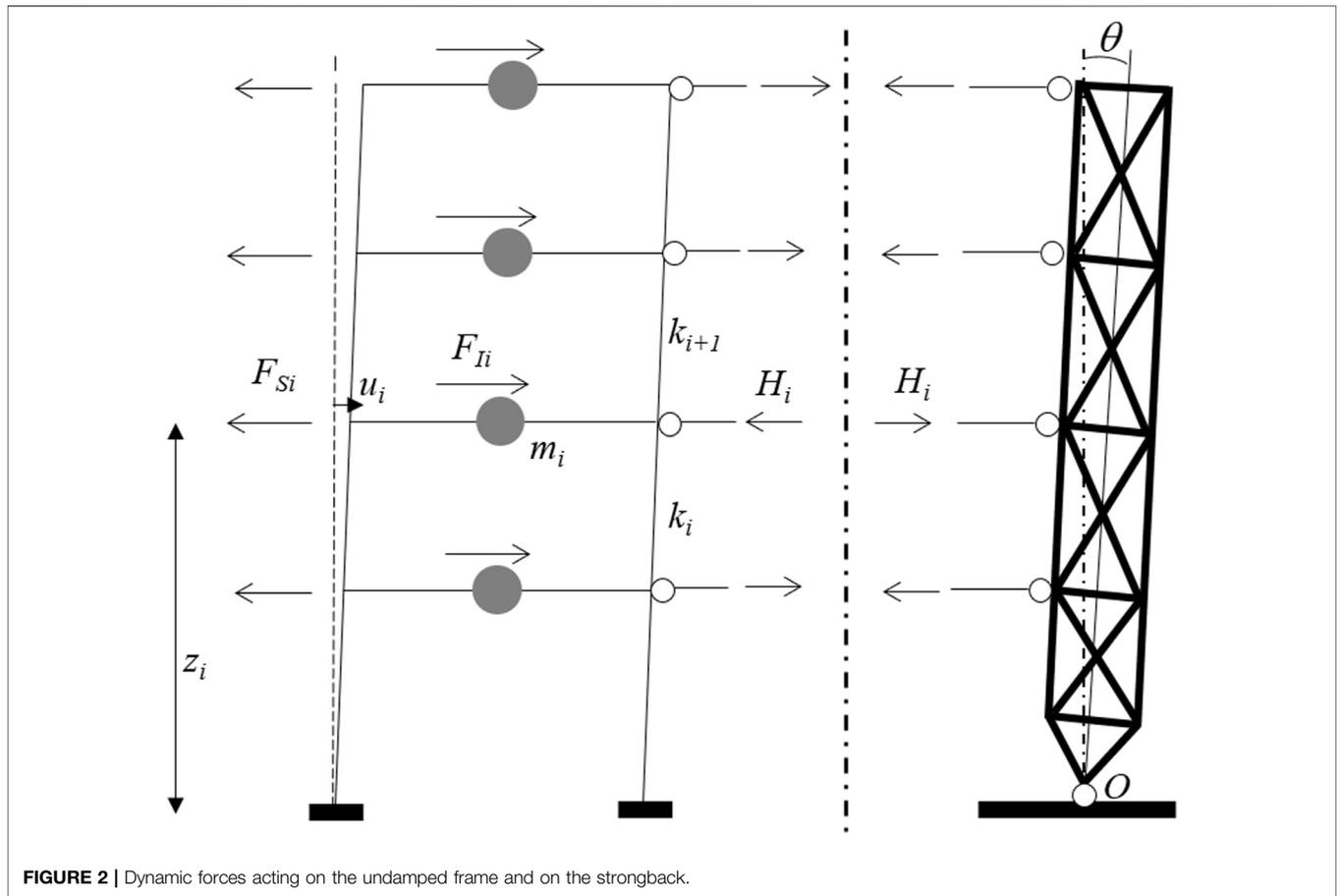


FIGURE 2 | Dynamic forces acting on the undamped frame and on the strongback.

Henceforth, for the specific case of an undamped frame structure connected to a strongback, only the global rotational equilibrium equations of the frame and of the strongback (see Figure 2) will be considered:

$$\begin{cases} \sum_{i=1}^N F_{L,i} \cdot z_i + \sum_{i=1}^N F_{S,i} \cdot z_i + \sum_{i=1}^N H_i \cdot z_i = 0 & \text{(a)} \\ \sum_{i=1}^N H_i \cdot z_i = 0 & \text{(b)} \end{cases} \quad (1)$$

where:

$$F_{L,i} = m_i \cdot \ddot{u}_i \quad (2)$$

$$F_{S,i} = \begin{cases} k_N \cdot \delta_N & \text{for } i = N \\ k_i \cdot \delta_i - k_{i+1} \cdot \delta_{i+1} & \text{for } i = 1, 2, \dots, N - 1 \end{cases} \quad (3)$$

and  $\delta_i = u_i - u_{i-1}$  is the  $i$ -th inter-storey drift.

Since the third term of Eq. 1a is null due to Eq. 1b, 1a can be rearranged as follows:

$$\sum_{i=1}^N (m_i \cdot \ddot{u}_i \cdot z_i) + \sum_{i=1}^{N-1} (k_i \cdot \delta_i - k_{i+1} \cdot \delta_{i+1}) \cdot z_i + k_N \cdot \delta_N \cdot z_N = 0 \quad (4)$$

Considering a uniform inter-storey height equal to  $h$ , then the inter-storey drift is also uniform and equal to  $\delta$ . In this case, the  $i$ -th floor height  $z_i$  and the lateral displacement  $u_i$  can be expressed as follows:

$$z_i = h \cdot i \quad (5)$$

$$u_i = \delta \cdot i \quad (6)$$

Substitution of Eq. 6 in Eq. 4 leads to:

$$\left( \sum_{i=1}^N (m_i \cdot z_i \cdot i) \right) \ddot{\delta} + \left( \sum_{i=1}^{N-1} (k_i - k_{i+1}) \cdot z_i + k_N \cdot z_N \right) \delta = 0 \quad (7)$$

Eq. 7 represents the equation of motion of the GSDOF system in the generalized coordinate  $\delta$ . The same equation can be also expressed in compact notation or, alternatively, in matrix notation ( $m$  representing the mass matrix and  $k$  the stiffness matrix) using the angle  $\theta$  ( $\theta = \delta/h$ ) as the degree of freedom (Palermo et al., 2020):

$$M^* \ddot{\delta} + K^* \delta = 0 \quad (8)$$

$$(z^T m d) \ddot{\theta} + (z^T k d) \theta = 0 \quad (9)$$

where  $z$  is the vector that groups the coordinates  $z_i$  of the different floors.

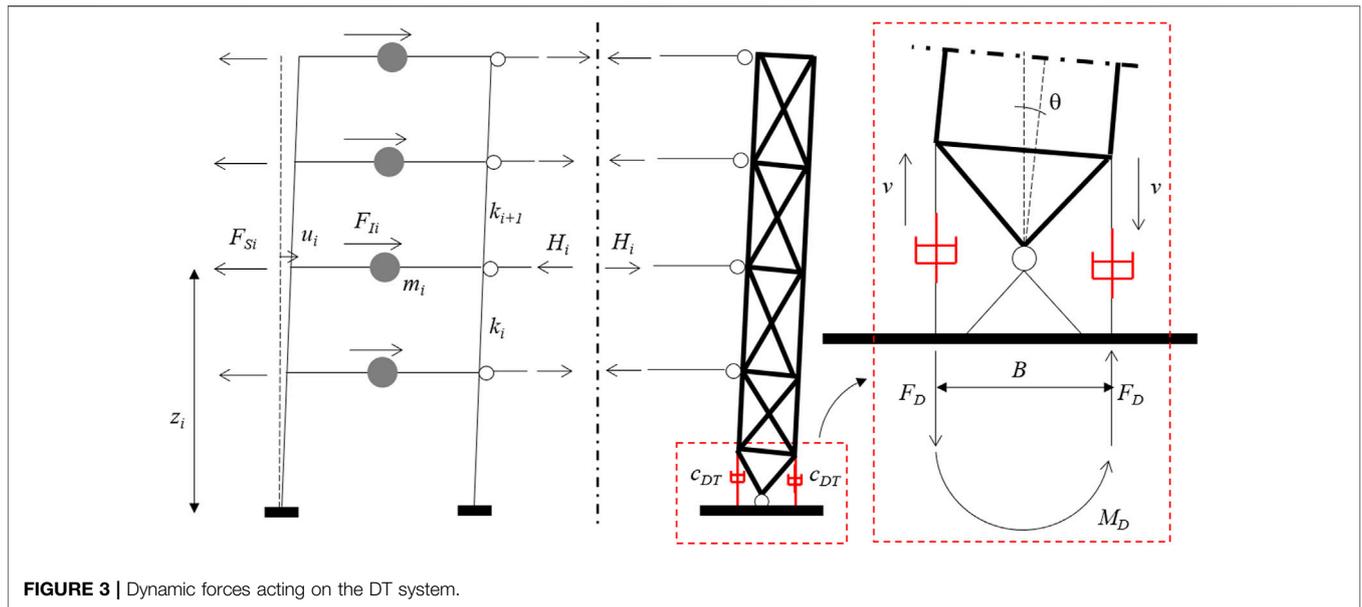


FIGURE 3 | Dynamic forces acting on the DT system.

The expressions of the generalized mass  $M^*$ , generalized stiffness  $K^*$  and circular frequency  $\omega = \sqrt{\frac{K^*}{M^*}}$  of the GSDOF system are reported in the work by Palermo et al. (2020).

### The Damped Systems

In the present section, the attention is focused on a frame connected to the strongback and equipped with a generic damping system, for instance with either a SPD, an IS or a DT system. In the general case, the dynamic equilibrium of the whole system can be formulated using the same approach as the one adopted for the undamped case, but including also the contribution of the damping forces. The damping system can be characterized by the damping matrix  $\mathbf{c}$ . By doing so, the following global dynamic rotational equilibrium equation governs the behavior of the frame structure:

$$(\mathbf{z}^T \mathbf{m} \mathbf{d}) \ddot{\theta} + (\mathbf{z}^T \mathbf{c} \mathbf{d}) \dot{\theta} + (\mathbf{z}^T \mathbf{k} \mathbf{d}) \theta = 0 \tag{10}$$

Alternatively, in the case of uniform inter-storey height  $h$ :

$$M^* \ddot{\delta} + C^* \dot{\delta} + K^* \delta = 0 \tag{11}$$

The expression of the generalized damping coefficient  $C^*$  depends on the specific configuration of the damping system. The expression of  $C^*$  for the case of the SPD and IS damping systems (without the strongback) were derived in the work by Palermo et al. (2020), for both cases of a generic deformed shape and a linear deformed shape. The latter corresponds to the case of a frame connected with a strongback.

The dynamic behavior of the DT damping system is here investigated in detail (Figure 3) with the purpose of deriving the corresponding dynamic equation of motion. In this case, the energy is dissipated through the forces exerted by the two viscous dampers located at the base of the strongback. Each damper has a damping coefficient equal to  $c_{DT}$  and acts along the vertical direction. From

simple geometrical considerations (see Figure 3), the vertical displacement  $v$  of each damper can be related to the strongback rigid rotation  $\theta$  (or, similarly, to the first inter-storey drift  $\delta_1$ ) through the following relationship:

$$v = \left(\frac{B}{2}\right) \cdot \dot{\theta} = \left(\frac{B}{2 \cdot h_1}\right) \cdot \dot{\delta}_1 \tag{12}$$

By indicating with  $\dot{\delta}_1 = \omega \delta_1$  the first inter-storey velocity, each of the two dampers exerts an equal and opposite damping force equal to:

$$F_D = c_{DT} \cdot \dot{v} = c_{DT} \cdot \left(\frac{B}{2 \cdot h_1}\right) \cdot \dot{\delta}_1 \tag{13}$$

The couple of damping forces  $F_D$  produces a dissipative bending moment  $M_D$  equal to:

$$M_D = F_D \cdot B = c_{DT} \cdot \left(\frac{B^2}{2 \cdot h_1}\right) \cdot \dot{\delta}_1 \tag{14}$$

The whole system is globally in equilibrium if the following system of rotational equilibrium equations is satisfied:

$$\begin{cases} \sum_{i=1}^N F_{I,i} \cdot z_i + \sum_{i=1}^N F_{S,i} \cdot z_i + \sum_{i=1}^N H_i \cdot z_i = 0 & \text{(a)} \\ \sum_{i=1}^N H_i \cdot z_i = M_D & \text{(b)} \end{cases} \tag{15}$$

Substituting Eq. 15b and Eq. 14 in Eq. 15a leads to the following expression of the global equilibrium equation of the frame:

$$\sum_{i=1}^N (m_i \cdot \ddot{u}_i \cdot z_i) + \sum_{i=1}^{N-1} (k_i \cdot \delta_i - k_{i+1} \cdot \delta_{i+1}) \cdot z_i + k_N \cdot \delta_N \cdot z_N + c_{DT} \cdot \left(\frac{B^2}{2 \cdot h_1}\right) \cdot \dot{\delta}_1 = 0 \tag{16}$$

**TABLE 1** | Damping coefficients and damping ratios in the general case and for the uniform systems.

System	Generalized damping coefficient	Damping ratio
SPD: general case	$C^* = \sum_{i=1}^{N-1} (C_{h,i+1} - C_{h,i}) \cdot z_i + C_{h,N} \cdot z_N$	$\xi = \frac{\sum_{i=1}^{N-1} (C_{h,i+1} - C_{h,i}) \cdot z_i + C_{h,N} \cdot z_N}{2 \cdot \left( \sum_{i=1}^N (m_i \cdot z_i \cdot i) \right) \cdot \omega}$
SPD: uniform case	$C^* = c_h \cdot N \cdot h$	$\xi = \frac{c_h}{2 \cdot m \cdot \omega} \cdot \left( \frac{6}{(N+1) \cdot (2N+1)} \right)$
IS: general case	$C^* = c_{h,i} \cdot h_i$	$\xi = \frac{c_{h,i} \cdot h_i}{2 \cdot \left( \sum_{i=1}^N (m_i \cdot z_i \cdot i) \right) \cdot \omega}$
IS: uniform case	$C^* = c_h \cdot h$	$\xi = \frac{c_h}{2m\omega} \left( \frac{6}{N(N+1)(2N+1)} \right)$
DT: general case	$C^* = c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right)$	$\xi = \frac{c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right)}{2 \cdot \left( \sum_{i=1}^N (m_i \cdot z_i \cdot i) \right) \cdot \omega}$
DT: uniform case	$C^* = c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right)$	$\xi = \frac{c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right)}{4 \cdot m \cdot \omega \cdot N} \cdot \left( \frac{6}{(N+1)(2N+1)} \right)$

Considering a uniform inter-storey height  $h$ , then the first inter-storey drift  $\delta_1$  is equal to  $\delta$ . Substituting Eq. 5 and Eq. 6 in Eq. 16 leads to the following equation of motion of the DT damped system:

$$\left( \sum_{i=1}^N (m_i \cdot z_i \cdot i) \right) \ddot{\delta} + c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right) \cdot \dot{\delta} + \left( \sum_{i=1}^{N-1} (k_i - k_{i+1}) \cdot z_i + k_N \cdot z_N \right) \delta = 0 \tag{17}$$

By comparing Eq. 17 with Eq. 11 it is clear that the expression of the generalized damping coefficient for the DT system results equal to  $C^* = c_{DT} \cdot \left( \frac{B^2}{2 \cdot h_1} \right)$ .

### EQUIVALENT DAMPING RATIO FOR THE DIFFERENT DAMPED SYSTEMS

Since Eq. 11 and Eq. 17 represent the equations of motion of damped GSDOF systems in the degree of freedom  $\delta$ , from fundamentals of structural dynamics it is possible to obtain the relationships relating the equivalent damping ratio with the damping coefficients. Table 1 summarizes these expressions for the different damped systems considered in the present study, considering both a generic frame and a uniform frame, namely a frame with uniform storey height  $h$ , uniform distribution of floor mass  $m$  and storey stiffness  $k$  and with dampers having the same damping coefficient  $c$ . The analytical expressions are based on an assumed linear lateral deformed shape.

The analytical expressions of the equivalent damping ratios collected in Table 1 allow to compare the effectiveness of the different damped systems in terms of energy dissipation capabilities. For this purpose, it is useful to evaluate the total damping coefficient  $c_{tot} = \sum_{k=1}^{n_d} c_k$  (namely the sum of all damping coefficients of each individual damper  $c_k$ )

necessary to obtain a given target damping ratio  $\bar{\xi}$ . For instance, for a uniform SPD system, the total damping coefficient is equal to  $c_{tot,SPD} = N \cdot c_h = N \cdot c \cdot \cos^2 \alpha$ , (where  $c_h$  refers to the horizontal component of the damping coefficient and  $\alpha$  the angle of inclination of the damper with respect to the horizontal line) while for a uniform DT system it is equal to  $c_{tot,DT} = 2 \cdot c_{DT}$ . Clearly for the uniform IS system, the total damping coefficient  $c_{tot}$  is coincident with the damping coefficient of the single damper  $c$ . In detail, their analytical expressions are here reported:

$$c_{tot,SPD} = \frac{1}{3} \cdot \bar{\xi} \cdot m \cdot \omega \cdot N \cdot (N + 1) \cdot (2N + 1) \tag{18}$$

$$c_{tot,IS} = \frac{1}{3} \cdot \bar{\xi} \cdot m \cdot \omega \cdot N \cdot (N + 1) \cdot (2N + 1) \tag{19}$$

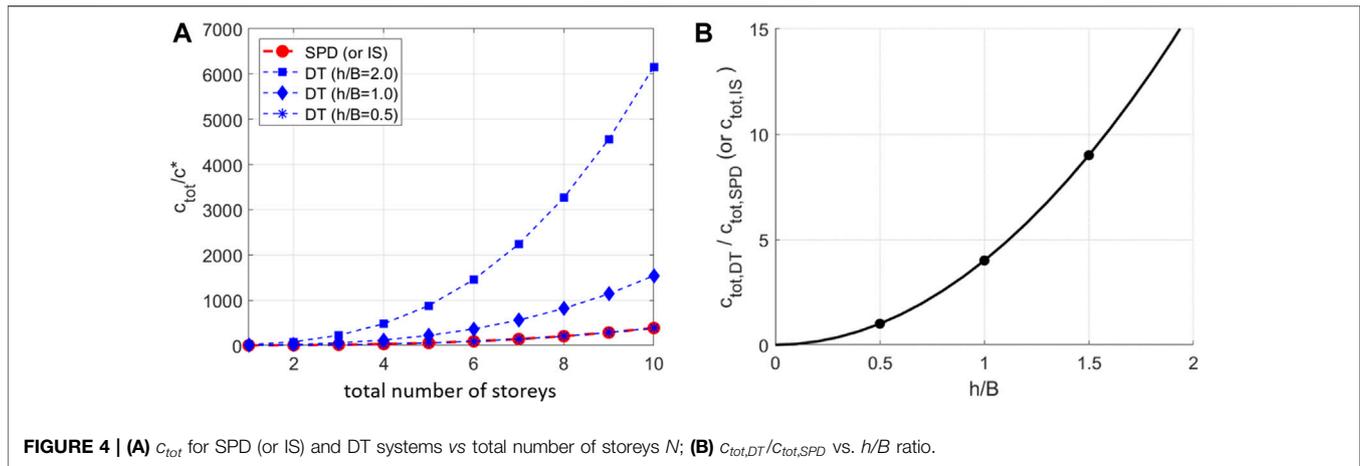
$$c_{tot,DT} = \frac{4}{3} \cdot \left( \frac{h}{B} \right)^2 \cdot \bar{\xi} \cdot m \cdot \omega \cdot N \cdot (N + 1) \cdot (2N + 1) \tag{20}$$

Firstly, it can be noted that the expressions for  $c_{tot,SPD}$  and  $c_{tot,IS}$  are identical. Moreover,  $c_{tot,SPD}$  (or equivalently  $c_{tot,IS}$ ) and  $c_{tot,DT}$  have a similar dependency on  $m$ ,  $\omega$  and  $N$ . It should be noted that  $c_{tot,DT}$  depends on the geometrical aspect ratio  $h/B$ . Therefore, the ratio between  $c_{tot,DT}$  and  $c_{tot,SPD}$  (or equivalently  $c_{tot,IS}$ ) depends uniquely on  $h/B$ , i.e.  $\frac{c_{tot,DT}}{c_{tot,SPD}} = 4 \cdot \left( \frac{h}{B} \right)^2$ .

Figure 4A displays the trends of  $c_{tot,SPD}$  (or equivalently  $c_{tot,IS}$ ) and  $c_{tot,DT}$  as a function of the total number of storeys  $N$ , normalized by  $c^* = 2 \cdot \bar{\xi} \cdot m \cdot \omega$  (corresponding to the expression of  $c_{tot,SPD}$ , or equivalently  $c_{tot,IS}$ , for  $N = 1$ ) thus allowing for a comparison of two systems leading to the same damping ratio. Figure 4B displays the trend of  $c_{tot,DT}/c_{tot,SPD}$  (or equivalently  $c_{tot,IS}$ ) as a function of  $h/B$ . The black circles of Figure 4B indicate selected values at  $h/B = 0.5, 1.0$  and  $1.5$ . For the case of  $h/B = 0.5$ ,  $c_{tot,DT}$  is equal to  $c_{tot,SPD}$  (or equivalently  $c_{tot,IS}$ ) meaning that the two systems lead to the same damping ratio (i.e., they dissipate the same amount of energy) with the same amount of total damping coefficient. For  $h/B$  values larger than  $0.5$ ,  $c_{tot,DT}$  becomes larger than  $c_{tot,IS}$ , meaning that the DT system needs a larger value of total damping coefficient to dissipate the same amount of energy dissipated by an IS system.

The GSDOF approach here used to determine the damping ratio of the damped frames connected to a strongback can be easily applied to any added dampers configuration, since, thanks to the presence of the strongback, the along-the-height lateral deformed shape remains linear and is not affected by the presence of the added dampers, as instead occurs for the case of frames without strongback. In fact, for particular dampers configurations, as those with diagonal dampers inserted in selected storeys where highly non-classical damping mechanism governs the energy dissipation process, the effectiveness of the damping system cannot be captured by the dynamics of the GSDOF system, as shown in the work of Palermo et al. (2020).

In the next section, specific dampers configurations are investigated through numerical simulations with the main purposes of i) verifying the effectiveness of the formulas derived for the different damped frames connected with a strongback and ii) comparing the seismic performances of the same damping system with and without the strongback.



## NUMERICAL SIMULATIONS

### Analyzed Systems and Input

A 6-storey one-bay uniform frame is considered as a case study. The frame columns are made with European HE320B profiles ( $J_{max} = 3.08 \cdot 10^8 \text{ mm}^4$ ,  $J_{min} = 9.27 \cdot 10^7 \text{ mm}^4$ ), while the beams are made with European IPE400 profiles ( $J_{max} = 2.31 \cdot 10^8 \text{ mm}^4$ ,  $J_{min} = 1.32 \cdot 10^7 \text{ mm}^4$ ). The floor mass  $m$  is set equal to 1,440 kN/g ( $g$  is the gravity acceleration). The systems without the strongback are indicated with the initials BF (acronym for bare frame), while the systems connected with the strongback are indicated with the initials SB (acronym for strongback).

The following BF systems are considered:

- BF-NAKED: bare frame structure;
- BF-SPD: bare frame structure equipped with the SPD system;
- BF-IS1: bare frame structure equipped with an IS system with the damper placed at the first storey;
- BF-IS3: bare frame structure equipped with an IS system with the damper placed at the third storey;
- BF-IS6: bare frame structure equipped with an IS system with the damper placed at the top (sixth) storey.

The following SB systems are considered:

- SB-NAKED: bare frame structure connected to the external rigid strongback;
- SB-SPD: bare frame structure connected to the external rigid strongback and equipped with the SPD system;
- SB-IS1: bare frame structure connected to the external rigid strongback and equipped with an IS system with the damper placed at the first storey;
- SB-IS3: bare frame structure connected to the external rigid strongback and equipped with an IS system with the damper placed at the third storey;
- SB-IS6: bare frame structure connected to the external rigid strongback and equipped with an IS system with the damper placed at the top (sixth) storey;

- SB-DT: bare frame structure connected to the external rigid strongback and equipped with a uniform DT damping system.

For the SB systems, the structural members of the strongback are modelled as rigid elements in order to guarantee a linear along-the-height lateral deformed shape. The FE model of all analyzed systems (displayed in **Figure 5**) are developed with the computer software (SAP, 2000) v23.1.0.

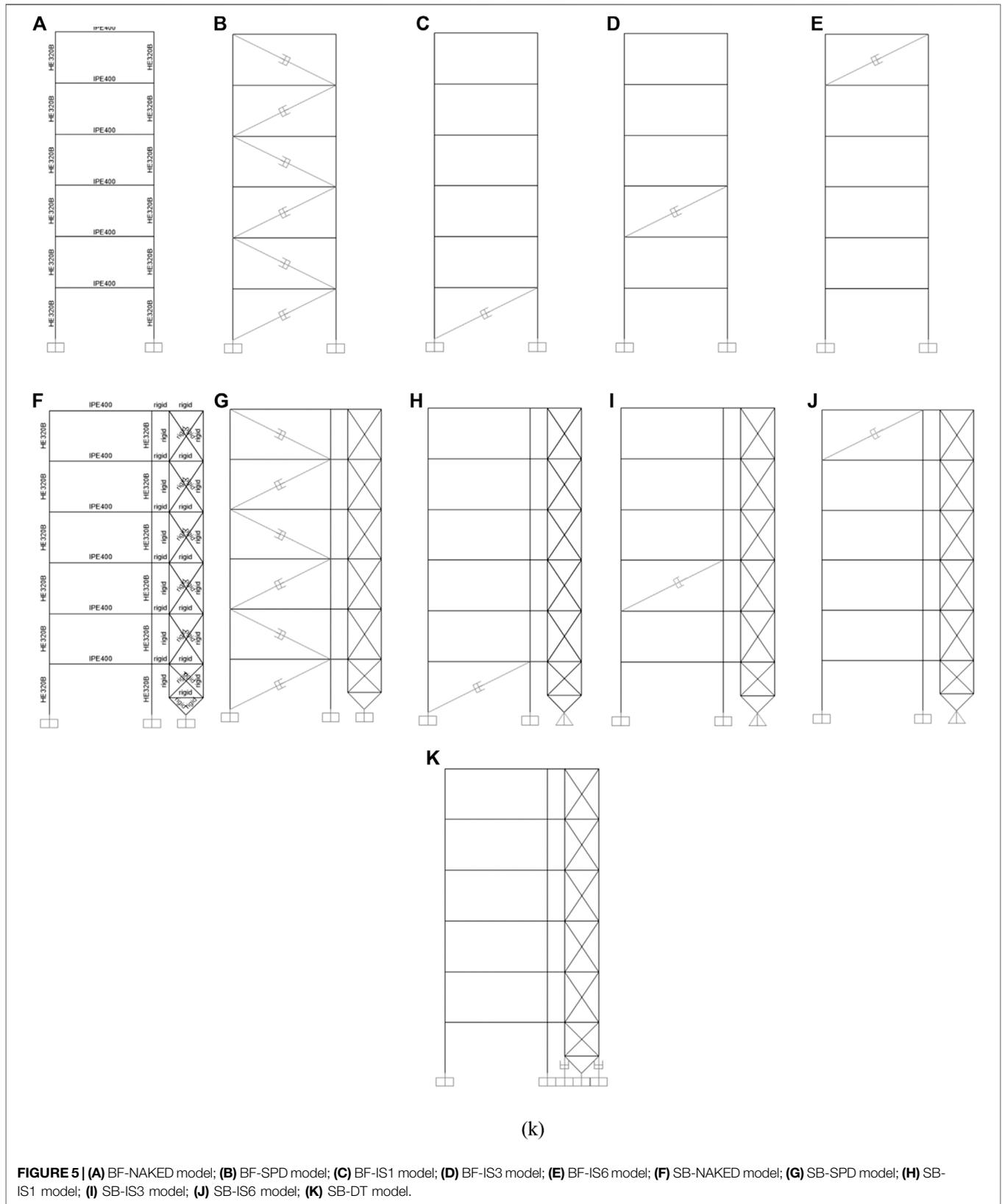
The first-mode deformed shape of the uniform BF-NAKED and SB-NAKED models are reported in **Figure 6**. The BF-NAKED model has a fundamental period equal to 1.28 s and a first-mode deformed shape that is typical of moment-resisting frames with stiff beams (**Figure 6A**). As expected, the first deformed shape of the SB-NAKED model is instead linear (**Figure 6B**) and the fundamental period is slightly lower than that of the bare frame (1.20 vs. 1.28 s).

The viscous dampers are sized to obtain a specific target viscous damping ratio ( $\bar{\xi}_v$ ) equal to 30% (the total target damping ratio  $\bar{\xi}$  is evaluated as the sum of the viscous damping ratio  $\bar{\xi}_v$  and the inherent damping ratio  $\xi_h = 5\%$ ). The values of the damping coefficients corresponding to each target damping ratio  $\bar{\xi}_v$  are reported in **Table 2**, as computed according to the analytical formulas of **Table 1**.

Dynamic time-history analyses are performed by: i) reproducing free damped vibrations starting from a given initial deformed configuration with linear along-the-height floor initial displacement and zero initial velocity (snap-back tests) and ii) applying at the base of the analyzed systems the El Centro S00E record (Imperial Valley 1940 earthquake).

### Evaluation of the Damping Properties Through Snapback Tests

Numerically simulated snap-back tests are performed with the purpose of estimating the equivalent damping ratio from the response in free vibration. For each system, the damping ratio is calculated according to the logarithmic decrement method (Chopra 2001) considering the time-history response of the top storey. As expected, the damping ratios obtained from the time-



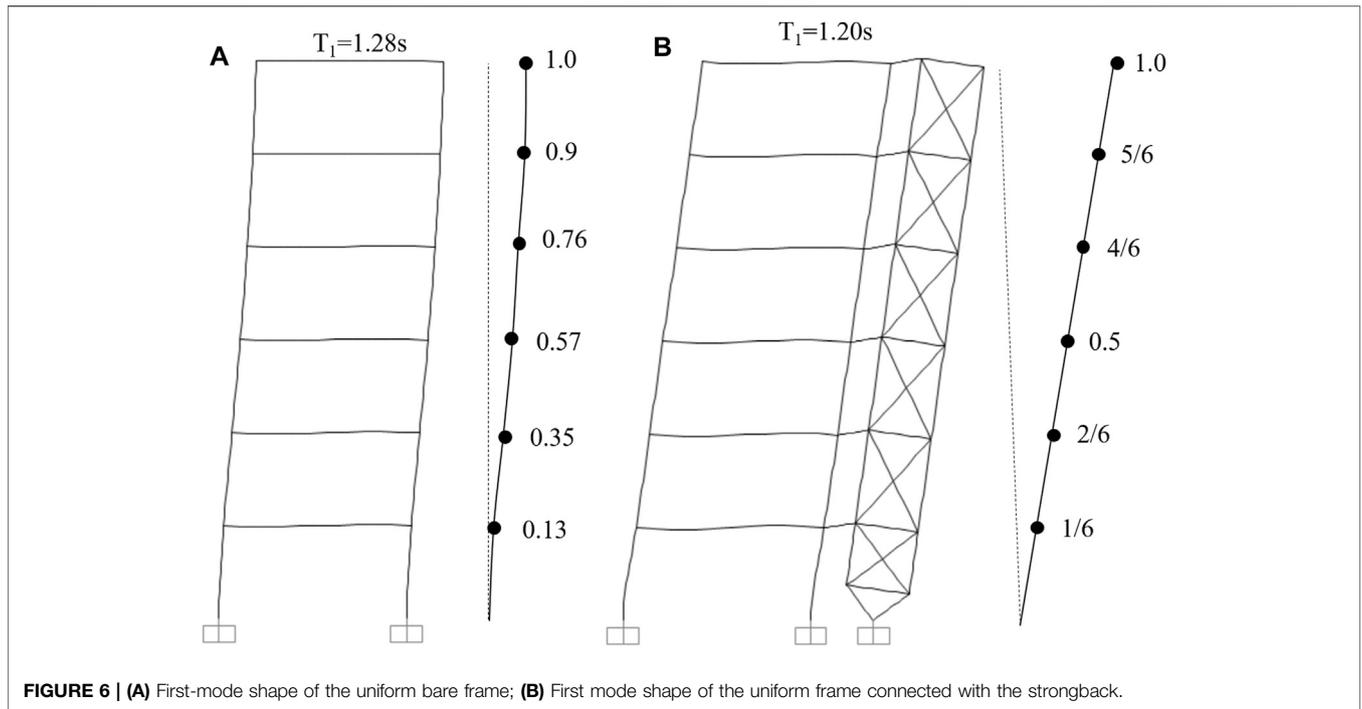


FIGURE 6 | (A) First-mode shape of the uniform bare frame; (B) First mode shape of the uniform frame connected with the strongback.

TABLE 2 | Damping coefficients for the case study.

System	$\bar{\xi}_v$	$c$ [ $\frac{MN \cdot s}{m}$ ]
BF-SPD	30%	1,340
BF-IS1, -IS3, -IS6	30%	8,040
SB-SPD	30%	1,428
SB-IS1, -IS3, -IS6	30%	8,566
SB-DT	30%	30,876

history responses of all the damped SB systems and of the BF-SPD system (between 33% and 35%, Figure 7A) are in quite good agreement with the target values (35%). On the contrary, the response in free vibration of the three BF-IS systems (Figure 7B) are quite different from the expected ones and, in some cases, it is not possible to determine the damping ratio according to the logarithmic decrement method (for instance: the response of system BF-IS3). Anyway, even for the cases where the free vibration responses are similar to that of a damped SDOF system (for instance: both the BF-IS1 and the BF-IS6 systems), the values of the damping ratios are far lower than the target ones estimated from the corresponding GSDOF systems. In particular, the response of the BF-IS6 system is quite close to the response of the BF-NAKED system, thus indicating an almost ineffective damping system (similar findings were also obtained in the work by Palermo et al., 2020).

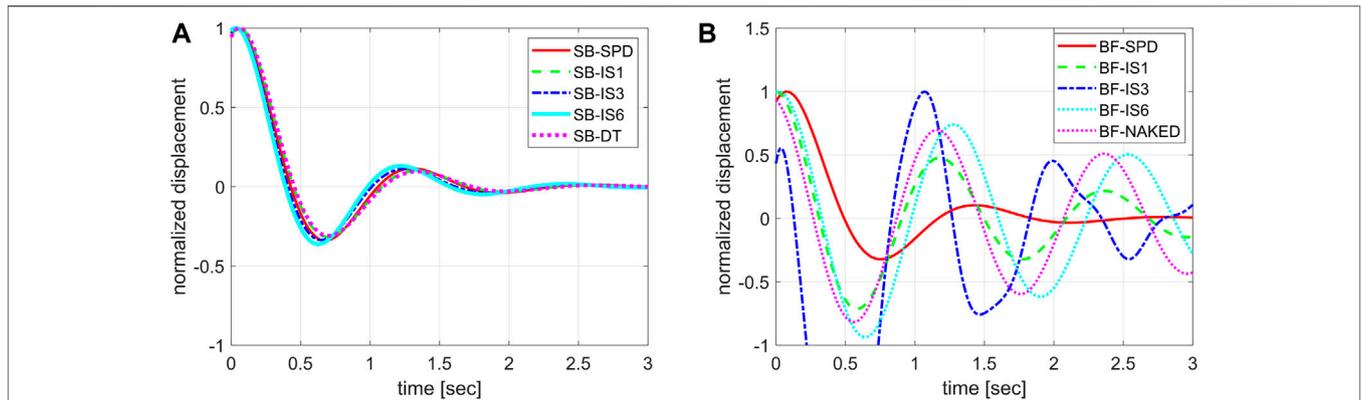
### Comparing the Seismic Performances of the Systems With and Without Strongback

The seismic performances of selected damped systems are compared considering the peak inter-storey drift ratio (*ID*)

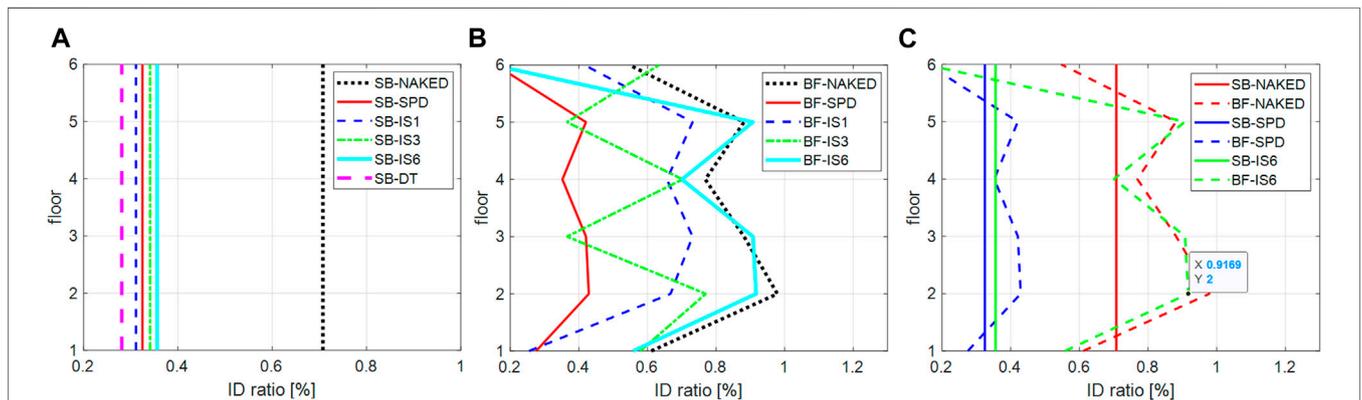
profiles under the El Centro ground motion. The purpose is to evaluate the influence of the presence of the strongback on the effectiveness of the different damping configurations.

Figures 8A,B display the *ID* profiles for all SB and BF systems. Then, a comparison between the *ID* profiles for selected models (namely SB-NAKED, BF-NAKED, SB-SPD, BF-SPD, SB-IS6, and BF-IS6) is displayed in Figure 8C. As expected, for the three SB models, the *ID* profile is uniform along the height, with values in the range of 0.25–0.35 for the damped systems and around 0.7 for the SB-NAKED model. As expected, for the naked model the *ID* values are not constant along the height and exhibit a very large variability with values ranging from 0.2 (at the peak storey for the BF-SPD model) to almost 1.0 (for both BF-NAKED and BF-IS6 models). The average *ID* value of the BF-NAKED model is around 0.7, thus close to that of the SB-NAKED. In general, the three BF-IS models exhibit larger *ID* values with respect to those observed for the BF-SPD model, thus indicating a reduced effectiveness in terms of energy dissipation.

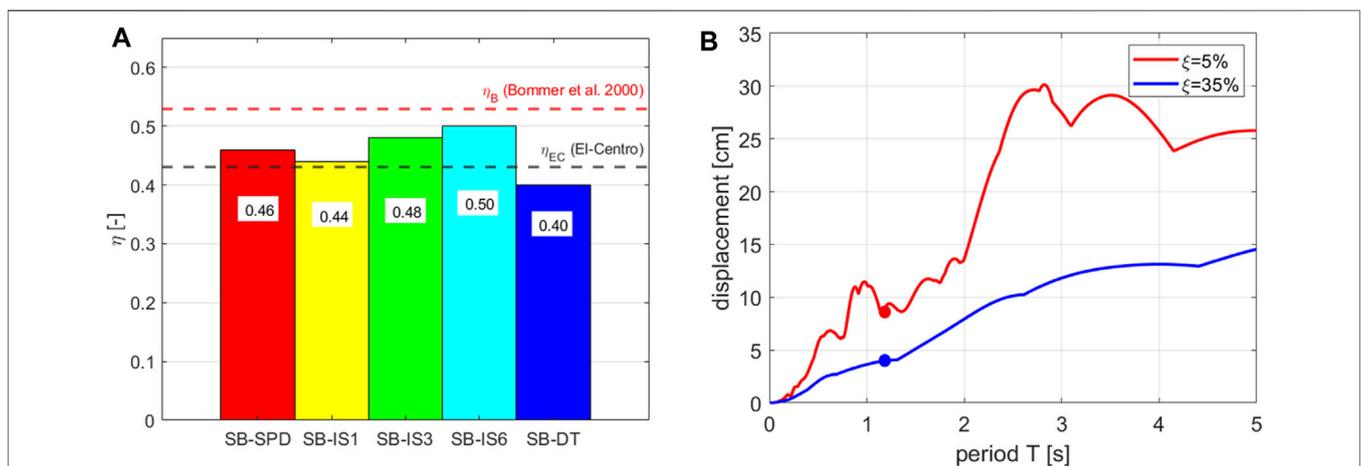
For the SB models, also the values of the *ID* damping reduction factors are computed as the ratio between the peak inter-storey drifts of one damped SB model and the SB-NAKED model. The values are reported in the histogram of Figure 9A. For sake of comparison, the histogram also includes the values of the damping reduction factor obtained i) according to the well-known formulation introduced by Bommer et al., 2000 (leading to a value of  $\eta_B = 0.53$ ) and ii) as the ratio of the 5% and 35% ordinates of the El-Centro displacement spectra (Figure 9B), as computed for a period equal to the fundamental period of the structure (leading to a value of  $\eta_{EC} = \frac{S_{d,6\%}(T_1)}{S_{d,5\%}(T_1)} = 0.43$ ). It can be noted that all values reported in the histogram are close to each other, indicating that the different SB damped systems dissipate a similar amount of energy.



**FIGURE 7 | (A)** Time-history response from snap-back test for the damped SB systems; **(B)** Time-history response from snap-back test for the BF systems.



**FIGURE 8 | (A)** ID for the SB systems. **(B)** ID for the BF systems. **(C)** Comparison of ID for the SB and BF systems.



**FIGURE 9 | (A)** Histogram of the damping reduction factors of the ID for the SB systems. **(B)** Displacement spectra for El-Centro ground motion.

It should be noted that the obtained numerical results are always affected by model uncertainties that lead to unavoidable discrepancies between the theoretical predictions, either analytical or numerical, and the actual response of the structure (especially in the case of non-linear models, Castaldo et al., 2020) which hence should be considered in design applications.

## CONCLUSION

The paper illustrates the main results of a study aimed at investigating the seismic design and performances of frame structures connected with an external strongback system and equipped with different configurations of added viscous dampers. For this study, the strongback system is modelled as an external pinned-base rigid vertical truss system connected with the frame and capable of imposing a linear along-the-height lateral deformed shape, thus limiting potentially dangerous effects associated to soft-storey mechanisms. The supplemental dampers provide extra energy dissipation capabilities to reduce the whole peak displacement response under earthquake excitation. The dynamic behavior of this coupled system is studied in detail for selected configurations of added viscous dampers, including i) dampers inserted within the frame at all storeys leading to a Stiffness Proportional Damping (SPD) system, ii) dampers inserted between two consecutive storeys only (“Inter-Storey” placement, IS system), and iii) dampers inserted at the base of the pinned-base strongback to realize a so-called rigid “Dissipative Tower” (DT system). For the special case of uniform systems (namely systems with uniform inter-storey height and uniform distributions of floor masses, lateral stiffness and viscous dampers), analytical design formulations are derived on the basis of the Generalized Single Degree of Freedom (GSDOF) approach. It provides insights into the dynamic behavior of the coupled system and allows to obtain useful

results for the preliminary seismic design of the added viscous dampers in terms of total amount of damping coefficient necessary to achieve a given target damping ratio. The results indicate that the DT system becomes more effective than both SPD and IS systems when the geometrical aspect ratio (inter-storey height over dampers arm) is smaller than 0.5. Finally, the effectiveness of the proposed design equations to size the added dampers are evaluated by means of numerical snap-back and earthquake simulations carried out on uniform SPD, IS and DT systems. It is shown that, especially for the IS systems, the presence of the strongback allows to increase the effectiveness of the damping system, thanks to the uniform dissipation of energy along the whole height of the structure.

## DATA AVAILABILITY STATEMENT

The datasets presented in this article will be available upon requests directed to the corresponding author.

## AUTHOR CONTRIBUTIONS

Conceptualization: MP, VL, and SS; Methodology: MP, Visualization: MP and VL; Writing–Original draft: MP, VL, and SS; Writing–Reviewing and Editing: GG; Resources: SS; Supervision: TT.

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## APPENDIX 1

A  $N$ -storey uniform shear-type frame connected to a rigid strongback is here considered. The system is referred to as SYSTEM A. The floor mass is indicated with  $m$ , while the storey lateral stiffness is indicated with  $k$ . According to the GSDOF analogy (Palermo et al., 2020), the system can be reduced to two equivalent SDOF systems, by imposing either the global translational equilibrium (Generalized Translational Oscillator - GTO) or the global rotational equilibrium (Generalized Rotational Oscillator - GRO) of the systems of dynamic forces acting on the frame:

$$\sum_{i=1}^N f_{L,i} + \sum_{i=1}^N f_{S,i} + \sum_{i=1}^N H_i = 0, \quad (\text{A1})$$

$$\sum_{i=1}^N f_{L,i} \cdot z_i + \sum_{i=1}^N f_{S,i} \cdot z_i + \sum_{i=1}^N H_i \cdot z_i = 0. \quad (\text{A2})$$

Since the strongback is capable of imposing a unique and known linear deformed shape  $\mathbf{d}$ , the two resulting GSDOF systems are hereafter demonstrated to be characterized by the same circular frequency of vibration.

The attention is focused on the global translational equilibrium equation (Eq. A1) leading to the GTO. For the sake of deriving the circular frequency of vibration of the GTO, it is convenient to introduce SYSTEM B, namely a shear-type frame system with the same uniform floor mass  $m$ , connected with the rigid strongback but with a fictitious non-uniform storey lateral stiffness  $k_i^*$ . The distribution of the storey lateral stiffness  $k_i^*$  has two main properties: (i) the sum of all  $k_i^*$  is equal to  $\sum_{i=1}^N k_i^* = N \cdot k$  that guarantees the same global lateral stiffness and period of vibration for both SYSTEM A and SYSTEM B; (ii) the free vibration response is characterized by

null values of mutual actions:  $H_i = 0 \forall i$ . By definition, the global translational equilibrium equation of SYSTEM B reduces to:

$$\sum_{i=1}^N f_{L,i} + \sum_{i=1}^N f_{S,i}^* = 0, \quad (\text{A3})$$

where:

$$f_{S,i}^* = \begin{cases} (k_i^* - k_{i+1}^*) \cdot \delta & \text{for } i = 1, 2, \dots, N-1 \\ k_N^* \cdot \delta & \text{for } i = N \end{cases}, \quad (\text{A4})$$

is the elastic force acting at the  $i$ -th storey of the SYSTEM B. The general analytical expression of  $k_i^*$  leading to null mutual actions in free vibration has been derived in the work by Palermo et al. (2018b) and is here reported for the specific condition of  $\sum_{i=1}^N k_i^* = N \cdot k$ :

$$k_i^* = \frac{6k}{2N^2 + 3N + 1} \cdot \frac{(N-i+1) \cdot (N+i)}{2}. \quad (\text{A5})$$

Substitution of Eqs. A4, A5 and Eq. 1 into Eq. A1 leads to:

$$\begin{aligned} \left( m \cdot \sum_{i=1}^N i \right) \cdot \ddot{\delta} + k_1^* \cdot \delta &= \left[ m \cdot \left( \frac{N(N+1)}{2} \right) \right] \ddot{\delta} + \left[ \frac{6Nk}{2N^2 + 3N + 1} \right. \\ &\quad \left. \cdot \left( \frac{N+1}{2} \right) \right] \delta \\ &= 0. \end{aligned} \quad (\text{A6})$$

It is easy to recognize that the two terms in brackets are the generalized mass  $M_T$  and the generalized stiffness  $K_T$  of the GTO. By definition, the circular frequency of the GTO results equal to:  $\omega_T = k/m[6/(N+1)(2N+1)]$ , that is coincident with the expression of circular frequency of the GRO  $\omega_R$  as reported in Table 1 of the work by Palermo et al. (2020).