

# XXL-HSC: Host properties of X-ray detected AGNs in XXL clusters

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## ABSTRACT

**Context.** There is compelling evidence that AGNs are strongly influenced by their environment, from their host galaxies to immense structures such as galaxy clusters. Therefore, studying the AGN population of clusters is essential, as both large-scale structures and AGN play key roles in galaxy evolution, though the interactions between these elements are still not well understood.

**Aims.** The primary objective of this study is to unravel the different factors that may significantly affect the triggering of AGN activity in cluster galaxies, including galaxy merging and interactions with other galaxies, and ram pressure from the hot intracluster medium.

**Methods.** For our purposes, we used 82 X-ray detected AGN found within a  $4r_{500}$  radius of 164 X-ray detected and spectroscopically confirmed galaxy clusters in the northern  $25 \text{ deg}^2$  field of the XXL survey, up to a redshift of  $z \sim 1$ . This field is also covered by deep optical observations of the Hyper Suprime-Cam, mounted on the 8 m *Subaru* Telescope, which allows for a reliable morphological classification of galaxies. We thoroughly investigated the morphology of X-ray AGN host galaxies, using both Statmorph software and visual inspection, in an attempt to discover disturbances as indications of interactions that could lead to AGN triggering. Furthermore, using the X-ray hardness ratio, the optical spectra and the spectral energy distributions of the X-ray sources, we have studied the obscuration and other AGN properties, as well as the star formation rate of the hosts as further indicators of interactions.

**Results.** We found a moderately significant, at the  $2\sigma$  confidence level, higher fraction of X-ray AGN in galaxy clusters hosted by merging or disturbed galaxies, compared to non-active cluster galaxies or X-ray AGN in the field. This excess is primarily localised in the cluster outskirts (between 1 and  $2r_{500}$ ). Also, we discovered a higher number of X-ray-hard (hence, possibly obscured) AGN in clusters than in the field, at the  $2\sigma$  confidence level, particularly in the outskirts. These findings further support the idea that galaxy mergers and interactions may serve as mechanisms for the triggering and obscuration of AGN activity.

**Conclusions.** The relatively high number of disturbed, merging, and possibly obscured AGN hosts in cluster outskirts suggests that galaxy merging and interactions are key drivers in triggering AGN activity in these outer regions of clusters.

**Key words.** galaxies: active – galaxies: clusters: general – galaxies: interactions – galaxies: nuclei

## 1. Introduction

Supermassive black holes (SMBHs) are at the forefront of modern astrophysical research today not only because they are hosted by every massive galaxy in the local Universe, but also because the evolution of the SMBH and its host galaxy appears tightly linked (e.g. Gültekin et al. 2009; Zubovas & King 2012). All SMBHs are thought to undergo active phases, the so-called active galactic nucleus (AGN) phases, during which they accrete the surrounding gas and emit an immense amount of energy. Theoretical models and simulations have proposed that dur-

ing this active phase, the active nucleus produces a feedback wind that can explain the co-evolution of the SMBH and its host galaxy (e.g. Schawinski et al. 2009; Cen & Chisari 2011). Therefore, the study of AGNs is essential for understanding the cosmic history of accretion into SMBH and their relation to the host galaxy.

One specific research direction for investigating the cosmic history of AGN evolution is studying AGN as a function of their environment. Several studies have provided evidence that AGN are affected by both their immediate surroundings (e.g. Maiolino et al. 1997; Dultzin-Hacyan et al. 1999; Sorrentino et al. 2006; González et al. 2008; Silverman et al. 2008; Dultzin et al. 2008; Koulouridis et al. 2006a,b, 2013;

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Manzer & De Robertis 2014; Silva et al. 2021; Duplancic et al. 2021; Pierce et al. 2023; Li et al. 2023) and by their large-scale environment (e.g. Constantin et al. 2008; Stroe et al. 2020; Ceccarelli et al. 2022; Hashiguchi et al. 2023; Muñoz Rodríguez et al. 2023; Koulouridis et al. 2016a, 2024; Toba et al. 2024; de Vos et al. 2024). In addition, early studies reported overdensities of X-ray point-sources in clusters with respect to the field (Cappi et al. 2001; Molnar et al. 2002; D’Elia et al. 2004; Branchesi et al. 2007; Gilmour et al. 2009), and others have spectroscopically verified the existence of a large population of AGN in clusters and argued on their probable evolution with redshift (e.g. Martini et al. 2002, 2007, 2009; Johnson et al. 2003). Therefore, it is crucial to thoroughly investigate the AGN population of galaxy clusters, as both the immense structure and the powerful nucleus seem to play an important role in galaxy evolution.

However, the interplay among the immense cluster and the powerful nucleus is still not well understood. This uncertainty arises from the various physical mechanisms that may influence galaxies and SMBHs within clusters. Numerous studies have demonstrated that the AGN fraction in member galaxies of massive clusters ( $M > 10^{14} M_{\odot}$ ) is lower than the respective fraction in field galaxies (e.g. Kauffmann et al. 2004; Gavazzi et al. 2011; Ehlert et al. 2013, 2014; Mishra & Dai 2020; Beyoro-Amado et al. 2021). This is likely caused by ram pressure stripping (RPS). In more detail, galaxies within the dense intracluster medium (ICM) of clusters are subjected to intense pressure, leading to the efficient ram pressure stripping of their gas. This reduces the availability of cold gas necessary for fueling AGN, resulting in a diminished cold gas reservoir required to trigger nuclear activity (e.g. Gunn & Gott 1972; Cowie & Songaila 1977; Giovanelli & Haynes 1985; Popesso & Biviano 2006; Chung et al. 2009; Haines et al. 2012; Sabater et al. 2013; Jaffé et al. 2015; Poggianti et al. 2017a). The impact of RPS is expected to be proportional to the cluster’s mass and inversely proportional to the galaxy’s mass (e.g. Boselli et al. 2022). In support of these expectations, studies on poor clusters and groups reported that AGN activity in group galaxies is at least as frequent as in the field (Sabater et al. 2012; Koulouridis et al. 2014, 2018, hereafter **XXL Paper XXXV**). However, Bufanda et al. (2017) did not detect differences in the fraction of X-ray luminous AGN ( $L_X > 10^{43} \text{ erg sec}^{-1}$ ) between groups and clusters for 432 clusters from the Dark Energy Survey (DES) up to  $z = 0.95$ . We note that Poggianti et al. (2017a) proposed that RPS might also act as a triggering mechanism for AGN activity in cluster members. Furthermore, the so-called jellyfish galaxies (Chung et al. 2009; Bekki 2009; Poggianti et al. 2017a), which are conspicuously affected by RPS, were found to host a significantly higher number of AGN than similar field galaxies (Peluso et al. 2022). Finally, many studies have revealed a positive evolution of the AGN fraction in cluster galaxies with redshift (e.g. Kocevski et al. 2009; Fassbender et al. 2012; Martini et al. 2013; Bufanda et al. 2017; Hashiguchi et al. 2023), while low-mass protoclusters at higher redshifts may potentially contain a higher number of AGN (Lehmer et al. 2013; Krishnan et al. 2017; Gatica et al. 2024; Vito et al. 2024).

Contrary to the suppression of AGN observed in the centre of massive clusters, an increase has been found in the outskirts of the cluster (e.g. Johnson et al. 2003; Branchesi et al. 2007; Koulouridis et al. 2014). However, the results vary depending on the different selection of clusters and AGN samples. In particular, Ruderman & Ebeling (2005) discovered a mild excess of X-ray sources between 1.5 and 3 Mpc in massive clusters that span the redshift range of  $z = 0.3\text{--}0.7$ . However, the excess was

found only in dynamically relaxed clusters, while no excess was found in the outskirts of disturbed clusters. These findings were confirmed more recently in the optical band by Stroe & Sobral (2021), in a sample of 14 clusters ( $z \sim 0.15\text{--}0.31$ ) that span a wide range of masses and dynamical states. They found that the H $\alpha$ -detected AGN fraction peaks in the outskirts of relaxed clusters ( $\sim 1.5\text{--}3$  Mpc). In addition, **XXL Paper XXXV** revealed a significant overdensity of spectroscopically confirmed X-ray AGN in the outskirts of low-mass clusters ( $M_{500}^1 < 10^{14} M_{\odot}$  and  $0.1 < z < 0.5$ ) from the XXL Survey (Adami et al. 2018, hereafter **XXL Paper XX**), while no excess was confirmed for higher cluster masses. At higher redshifts, a similar excess of X-ray AGN was also reported by Fassbender et al. (2012) between 4 and 6 arcmin from the centres of 22 massive clusters ( $0.9 < z < 1.6$ ). Koulouridis & Bartalucci (2019) confirmed a highly significant excess, at the 99.9% confidence level, of X-ray point-like sources in the outskirts ( $2\text{--}2.5r_{500}$ ) of the five most massive,  $M_{500}^{SZ} > 10^{14} M_{\odot}$ , and distant,  $z \sim 1$ , galaxy clusters in the *Planck* and South Pole Telescope (SPT) surveys. Very recently, Koulouridis et al. (2024) reported a significant excess of X-ray AGN in the outskirts of relaxed clusters, compared both to non-relaxed clusters and to the field. Furthermore, a similar excess in cluster outskirts ( $\sim 3r_{500}$ ) was also recently found in the *Magneticum* simulations (Rihtaršič et al. 2024). Nevertheless, some studies have found that AGN have no special position inside galaxy clusters (e.g. Gilmour et al. 2009; Ehlert et al. 2015), unless only the most powerful optical AGNs are considered. In addition, Muñoz Rodríguez et al. (2023) demonstrated that the observed excess of X-ray sources in the outskirts of massive clusters ( $M > 5 \times 10^{14} M_{\odot}$ ) at  $z \sim 1$  might be caused by projection effects.

The increase reported by various works may be attributed to a corresponding increase in the galaxy merging rate, which is favoured by the lower galaxy velocities in the outskirts compared to the centre, as well as low-mass groups ( $M_{500} < 10^{14} M_{\odot}$ ) when compared to massive clusters (e.g. Ehlert et al. 2015; Lopes et al. 2017; Gordon et al. 2018). During a merger, the involved galaxies experience tidal forces that can lead to the redistribution of their stellar and gaseous components, triggering star formation and potentially fueling AGN. Generally, merging represents a significant process in the hierarchical formation and evolution of galaxies. As galaxies traverse the dense environments of clusters, gravitational interactions and dynamical friction facilitate their coalescence. In addition, notably in the outskirts of clusters, galaxies experience less intense ICM pressure, preserving their cold gas supplies and thus maintaining the conditions favourable for AGN activity. Furthermore, nuclear activity in the outskirts may be triggered by interactions between the host galaxy and the cluster itself during passage through virial shocks (e.g. Keshet & Reiss 2018).

Another possibility, which explains the excess of AGN in the outskirts of clusters, is that AGN enter the cluster environment along with infalling small groups. Galaxies typically do not enter clusters as isolated entities. Instead, they often enter as part of smaller groups that offer a more favourable environment for AGN triggering, namely, they are pre-processed. Pre-processed indicates that these galaxies have already experienced significant interactions and evolution before entering the cluster (e.g. Fujita 2004; Haines et al. 2015; Sengupta et al. 2022;

<sup>1</sup> The  $M_{500}$  cluster mass refers to the total mass of a galaxy cluster within a spherical region where the average density is 500 times the critical density of the universe at the cluster’s redshift. The corresponding radius,  $r_{500}$ , is the radius at which this density contrast is reached.

Lokas 2023). However, Koulouridis et al. (2024) found no evidence that the prevalence of X-ray AGN in clusters is affected by the presence of X-ray-detected infalling groups, or that their location is correlated with the positions of these infalling groups.

In summary, the environmental dichotomy between the cluster centre and the outskirts results in a spatial variation of AGN activity. However, the physical mechanisms that trigger nuclear activity in cluster galaxies are still debated, since the AGN frequency appears to be affected by multiple factors. In this context, we have studied the morphology of cluster galaxies hosting X-ray-detected AGN up to a redshift of  $z \sim 1$ . They are located out to a distance of  $4r_{500}$  radii from the centres of X-ray detected clusters in the northern  $25 \text{ deg}^2$  XXL survey field, which is also covered by deep Hyper Suprime-Cam (HSC) optical observations. Our aim is to investigate the role of galaxy merging and interactions in the triggering of AGN activity. In addition, we have studied the obscuration and the accretion power of the AGN, as well as the star formation rate (SFR) of the hosts, as further indicators of the physical mechanisms that drive their triggering and evolution.

The outline of this paper is as follows. In Sect. 2 we discuss the data preparation and sample selection. The methodology is described in 3, and the results are presented in Sect. 4. Our discussion and conclusions are presented in Sect. 5, and a summary is provided in Sect. 6. Throughout this paper we assume a Planck cosmology Planck Collaboration XIII (2016) of  $H_0 = 67.8 \text{ h km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_m = 0.308$ , and  $\Omega_\Lambda = 0.692$ .

## 2. Sample selection

In this section, we describe the surveys and the data used to study the properties of X-ray AGN in galaxy clusters of the XMM-XXL northern field. We used the latest X-ray catalogue from the XXL survey along with deep HSC imaging. Additionally, we exploited all available spectroscopic and photometric data from a multitude of surveys, to study the spectral properties and the spectral energy distributions (SEDs) of our sources.

### 2.1. XMM-XXL survey

The XXL survey (Pierre et al. 2016) is the largest XMM-Newton programme to date totaling  $\sim 6.9$  Msec. It covers two extragalactic fields of about  $25 \text{ deg}^2$  each, XXL-North (XXL-N) and XXL-South (XXL-S), at a point-source sensitivity of  $6 \times 10^{-15} \text{ ergs}^{-1} \text{ cm}^{-2}$  in the  $[0.5\text{--}2] \text{ keV}$  band (completeness limit). The main goals of the survey are to provide constraints on the dark energy equation of state from the space-time distribution of clusters of galaxies and to serve as a pathfinder for future, wide-area X-ray missions. In the current work, we have used the latest XXL v4.3 catalogue. The cluster selection criteria are described in (Pacaud et al. 2006) and the 365 cluster catalogue is presented in XXL Paper XX. The creation of the XMM-XXL X-ray point-source catalogue, along with the optical counterpart matching and associated multi-wavelength data, is described in Chiappetti et al. (2018), also known as XXL Paper XXVII. However, in this work we used the more recent internal release of the catalogue obtained with the V4.2 XXL pipeline that contains in total 15547 X-ray sources.

Spectroscopic redshifts in XXL-N were obtained with large spectroscopic surveys, such as SDSS, VIPERS (Guzzo et al. 2014), and GAMA (Liske et al. 2015), and from a large campaign with the AAOmega spectrograph mounted on the Anglo-Australian Telescope described in Lidman et al. (2016), also known as XXL Paper XIV. Other smaller-scale spectro-

scopic observations (e.g. Koulouridis et al. 2016a) complement the sample. The photometric redshifts used are described in Fotopoulou et al. (2016), also known as XXL Paper VI.

### 2.2. Hyper Suprime-Cam (HSC)

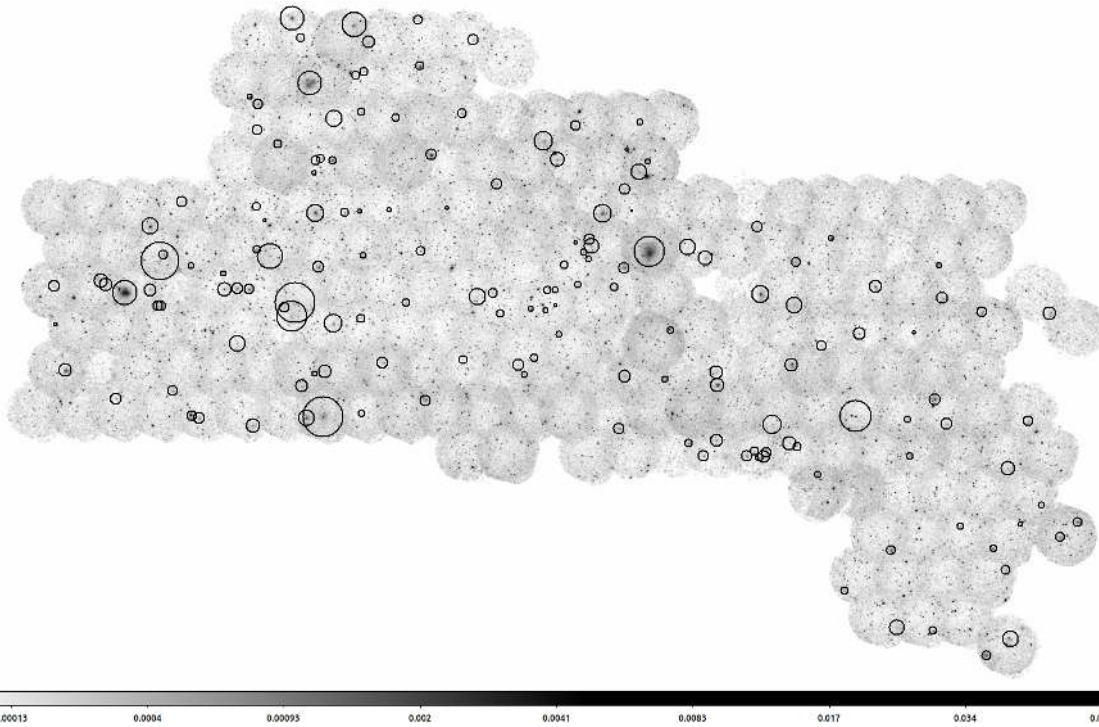
The HSC (Miyazaki et al. 2018) is a  $1.77 \text{ deg}^2$  imaging camera with a pixel scale of 0.168 arcsec mounted at the prime focus of the 8.2 m Subaru Telescope. It is composed of 116 charge-coupled devices (CCDs; 104 for science, four for the auto guider, and eight for focus monitoring). This facility is operated by the National Astronomical Observatory of Japan on the summit of Maunakea (Hawaii, USA). The HSC Subaru Strategic Program (HSC-SSP; Aihara et al. 2018) is a three-layered survey (wide, deep, and ultradeep) of  $1400 \text{ deg}^2$  in five different bands (grizy) and four narrow filters. For our purposes, we used imaging and photometric data from the deep HSC survey, which covers most of the northern XMM-XXL survey, reaching a depth of  $r \sim 27$ .

### 2.3. X-ray detected AGN in XXL clusters

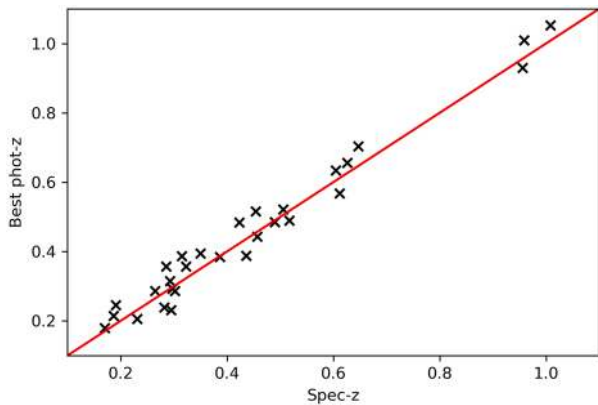
Our aim was to combine X-ray data from the XXL survey with HSC imaging to perform a reliable visual inspection of the X-ray AGN host galaxies up to a high redshift ( $z \sim 1$ ). The HSC footprint covers only the northern XXL field, therefore our sample comprises 164 clusters. The frequency of AGN activity in most of these clusters, up to redshift  $z = 0.5$ , was studied thoroughly in XXL Paper XXXV. We initially selected all point-like X-ray sources within a projected cluster-centric distance of  $4r_{500}$ . We used the  $r_{500,MT}$  values when available, from Table 5 in XXL Paper XX, derived after a spectral fit of the X-ray observations. In the case where a spectral fit was not possible, we used the  $r_{500,scal}$  values derived from scaling relations (XXL Paper XX, Table F.1). The XXL-N field and the location of the 164 X-ray detected clusters are presented in Fig. 1.

Subsequently, only sources above a luminosity threshold of  $L_{X[0.5\text{--}10] \text{ keV}} > 10^{42} \text{ erg s}^{-1}$  at the redshift of each cluster were kept in the sample. This luminosity threshold indicates that X-ray emission from point-like sources is most likely due to an AGN rather than other sources such as X-ray binaries or star formation. Finally, only galaxies with concordant spectroscopic or photometric redshift with the redshift of the clusters were kept in the sample. In more detail, in case of available spectroscopic redshifts ( $\sim 50\%$  of our sample), we require the line-of-sight (los) relative velocity to the average cluster redshift,  $\Delta v = v_{\text{los}} - \langle v \rangle$ , to be less than  $2000 \text{ km/sec}$ . When only a photometric redshift is available, the relative los distance should not exceed  $0.1(1 + z_{\text{cl}})$ , where  $z_{\text{cl}}$  is the average cluster redshift. This threshold allows the selection of the sources with the most reliable photometric redshifts (XXL Paper VI, see Figure 2), while it excludes catastrophic outliers. Furthermore, if the source was outside the 68% confidence interval of the photometric redshift probability distribution (PDZ) around the median value, or if this interval was larger than 0.5, the source was excluded from our analysis. The latter criterion ensures that sources with flat PDZs, and therefore unreliable redshift estimation, are excluded from the analysis. In Fig. 2 we plot the spectroscopic vs. photometric redshifts for the sources that would have been included in our sample even if the spectroscopic redshift was not available. While our stringent criteria verify the purity of our sample, some true X-ray AGN that belong to the clusters inevitably will be missed.

The final selected main sample comprises 82 X-ray AGN up to  $4r_{500}$  radii from the centres of the clusters. Moreover, we assembled four control samples to enable a thorough



**Fig. 1.** XXL-N count-rate map. We overlay the positions of 164 X-ray detected clusters used in this paper. Circles mark the position and they represent the  $r_{500}$  radius of each cluster.



**Fig. 2.** Spectroscopic versus photometric redshift scatter plot of the spectroscopically confirmed X-ray detected AGN in clusters. The sources were selected based on specific photometric redshift criteria, as described in Sect. 2.3. The red line denotes the equality.

comparative analysis: (i) 1987 field X-ray-detected AGN in XXL-N and (ii) 1914 cluster galaxies selected based on their spectroscopic redshift within  $4r_{500}$  radii of our 164 clusters. From these two control samples, we extracted two smaller subsamples of (iii) 166 field X-ray AGN and (iv) 208 cluster galaxies respectively, in order to visually inspect their morphology and optical spectra. Control samples (i) and (iii) for field X-ray AGN are formulated to mirror the main sample in terms of both redshift and X-ray luminosity distribution. Control samples (ii) and (vi) for cluster galaxies are designed to emulate the main sample in terms of redshift and stellar mass distribution. A summary of the samples is presented in Table 1, while the redshift distributions of the main sample and the two small subsamples (iii) and (iv) are presented in Fig. 3.

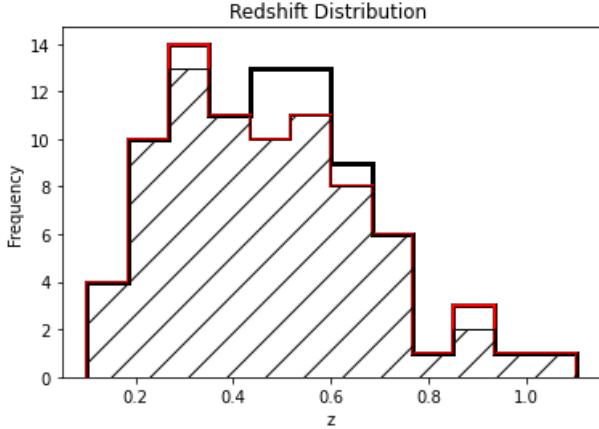
**Table 1.** Summary of the samples.

Sample (1)	Description (2)	Size (3)
Main	cluster X-ray AGN	82
Control (i)	field X-ray AGN	1987
Control (ii)	cluster galaxies	1914
Control (iii)	field X-ray AGN	166
Control (iv)	cluster galaxies	208

**Notes.** Control samples (iii) and (iv) were directly derived from the larger respective samples (i) and (ii) in a manner that preserves their properties and enables visual inspection.

### 3. Methodology

In this section we describe the methodology we followed to treat our data. In order to examine in detail the properties of X-ray AGN in clusters, we divided the circum-cluster area into four concentric annuli centred on the X-ray peak of the diffuse emission. We selected each annulus to have a width of  $r_{500}$  radius. Employing any  $r_{\Delta}$  radius is crucial when studying the impact of the cluster environment on AGN and their host galaxies, as it provides a direct link to the physical conditions at each galaxy's location (Koulouridis & Bartalucci 2019). We can assume that similar conditions prevail within the specific annulus of other clusters, regardless of the actual physical or projected distance. The  $r_{500}$  radius is also useful for direct comparison with previous results as it is used extensively in the literature. Then, we stacked the number counts of X-ray AGN found in each respective annuli of all clusters. We consider the first  $r_{500}$  annulus to be the centre of the cluster, while the second annulus the cluster outskirts. The motivation to use these specific boundaries comes from the fact that the radius of  $2r_{500}$  roughly coincides with the virial



**Fig. 3.** Normalised redshift distribution. The hatched area corresponds to the main sample of 82 X-ray detected AGN in clusters while the black line and the red line correspond to the 166 field X-ray AGN (sub-sample ii) and the 208 clusters galaxies (sub-sample iv), respectively.

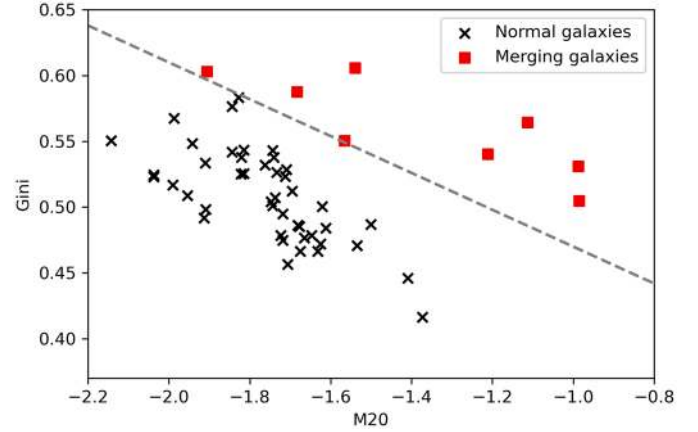
radius and the splash-back radius<sup>2</sup>. Therefore, we can assume that galaxies within this radius are bound to the cluster potential, while outside they are not yet influenced by any effect of the dense cluster environment. Nevertheless, for any comparison with the field in the current work we have used independent control samples of X-ray AGN (see Sect. 2.3). We note that within an area roughly encompassing the inner half of the  $r_{500}$  radius, therefore  $\sim 25\%$  of the total area of the central annulus, the high X-ray background caused by the diffuse ICM emission may hinder the detection of low-luminosity point-like sources (e.g., Bhargava et al. 2023).

### 3.1. Morphological analysis of X-ray AGN host galaxies

Our goal is to investigate whether merging is more frequent in X-ray AGN hosts within clusters than in the field or when compared to inactive cluster galaxies. To this end, we used HSC imaging and photometry in order to examine and determine the morphology and the immediate environment of all galaxies in our samples. Specifically, we were investigating for signatures of mergers, or interactions with neighbouring galaxies. To find merging systems in all samples, we used Statmorph (Rodríguez-Gomez et al. 2019). This code is designed for measuring the morphological properties of galaxies, especially for the non-parametric morphological indicators. In this work, we employed Gini coefficient and  $M_{20}$  diagnostics to select mergers (Lotz et al. 2004). The Gini coefficient evaluates the bias of the light distribution in a galaxy, where larger values indicate a more concentrated light distribution. On the other hand,  $M_{20}$  assesses the distribution of the brightest twenty percent of the galaxy’s light, where higher values suggest more extended bright regions. To find the merging systems, we applied the Gini- $M_{20}$  classification to the HSC  $i$ -band image (see Yanagawa et al. in prep. for a full description of merger classification for HSC sources with Statmorph). In particular, we classified as mergers all galaxies that fall above the line that divides merging and non-merging systems (Lotz et al. 2008; Rodríguez-Gomez et al. 2019):

$$\text{Gini} = -0.14 M_{20} + 0.33 \quad (1)$$

<sup>2</sup> The splash-back radius of a galaxy cluster is the boundary that marks the outer edge of its gravitational influence, where the accreted matter reaches its farthest point after falling into the cluster for the first time.



**Fig. 4.** Gini- $M_{20}$  classification diagram of galaxy morphology. Black crosses and red squares mark the position of the X-ray AGN in clusters; we excluded those with bad measurements. The hatched line (Equation (1)) separates “normal” and merging galaxies.

Sources that were flagged as having a problem with basic measurements for various reasons (Rodríguez-Gomez et al. 2019, see Section 4) were discarded from the above analysis as undefined (U). In Fig. 4 we plot the positions of the X-ray AGN host galaxies on the Gini- $M_{20}$  diagram and the dividing line.

Additionally, for the main sample of 82 X-ray AGN and for the two small control subsamples (iii) and (iv) we performed a visual examination of all HSC images and classified the AGN host galaxies as merging or disturbed (M), Ellipticals (E), Spirals (S) and Undefined (U). In more detail, we classified galaxies as disturbed (D) if they showed clear signs of disturbed morphology, the presence of tidal tails, or noticeable substructures. In addition, mergers should exhibit two visibly separate cores. We note that a few galaxies classified as “U” by Statmorph due to poorly masked secondary sources were recovered as merging systems when the redshifts of the two cores were also taken into account. Conversely, cases of projection that were incorrectly identified as mergers were excluded. Thus, while the results of Statmorph and visual classification are not identical, they give consistent statistical results. Visually small galaxies or bright QSOs are classified as undefined. However, only two cases among the 82 X-ray AGN in the main sample exhibit such powerful QSO activity that it obstructs the morphological classification of their host galaxies. Three examples of X-ray AGN merging or disturbed hosts can be found in Fig. 5. A mosaic of the full sample of the 82 X-ray-detected AGN in clusters can be found in Table A.2.

### 3.2. Multi-wavelength properties

#### 3.2.1. Optical spectra analysis

Active Galactic Nuclei can be classified into two main types, broad-line (BL) and narrow-line (NL) AGN, according to the width of Balmer emission lines in their optical spectrum. In more detail, in AGN models a torus of gas and dust surrounds the central SMBH, obscuring the nucleus from certain viewing angles. Narrow-line AGN are believed to be obscured by the torus, which blocks the direct radiation from the central AGN region so that the broad components of the Balmer lines cannot be detected. However, many studies suggest that below a specific accretion rate of material into the black hole the broad-line region (BLR) might be absent (Nicastro 2000; Nicastro et al.



**Fig. 5.** RGB colour images based on the *gri* filters of HSC. Top panel: X-ray AGN merging host of 4XLSU J023147.2-045702 with a disturbed morphology ( $z = 0.190$ ). Middle panel: X-ray AGN merging host of XLSSU J023324.9-053106 with two distinct cores ( $z = 0.436$ ). Bottom panel: X-ray AGN disturbed host of XLSSU J021649.4-032831 ( $z = 0.21$ ). The image scale is approximately  $30'' \times 30''$ .

2003; Bian & Gu 2007; Elitzur & Ho 2009; Bianchi et al. 2012, 2017; Elitzur et al. 2014; Elitzur & Netzer 2016; Koulouridis 2014; Koulouridis et al. 2016b). Therefore, some of the low-luminosity sources of our sample might be “true” narrow-line AGN, intrinsically lacking the BLR. In particular, almost 25% of our sources satisfy the criteria of accretion rate and luminosity described in Marinucci et al. (2012), Koulouridis et al. (2016b) that would classify them as potentially lacking the BLR. Furthermore, additional obscuration or dilution, caused by the host galaxy, might affect the AGN classification, regardless of the inclination of the torus (e.g. Lagos et al. 2011; Gkini et al. 2021)

We classified the majority of sources with available optical spectra as either broad-line or narrow-line AGN based on the width of Balmer emission lines, using optical spectra primarily obtained from SDSS. All spectra in the redshift range of this study include at least the  $H\beta$  region, while most of them include also the  $H\alpha$ . In particular, we classify all sources with a FWHM of the Balmer lines less than 500 km/s and no evidence for any broadening with respect to forbidden lines (such as [OIII], [NII]) as narrow-line AGN. Nevertheless, a non-negligible fraction presents the typical spectrum of an absorption-line galaxy (ALG), while we were not able to reliably classify a few of them due to poor spectrum quality. In Fig. 6 we present an example of a broad- and a narrow-line spectrum. A summary of the classification is presented in Table 2, and individual classifications are given in Table A.1.

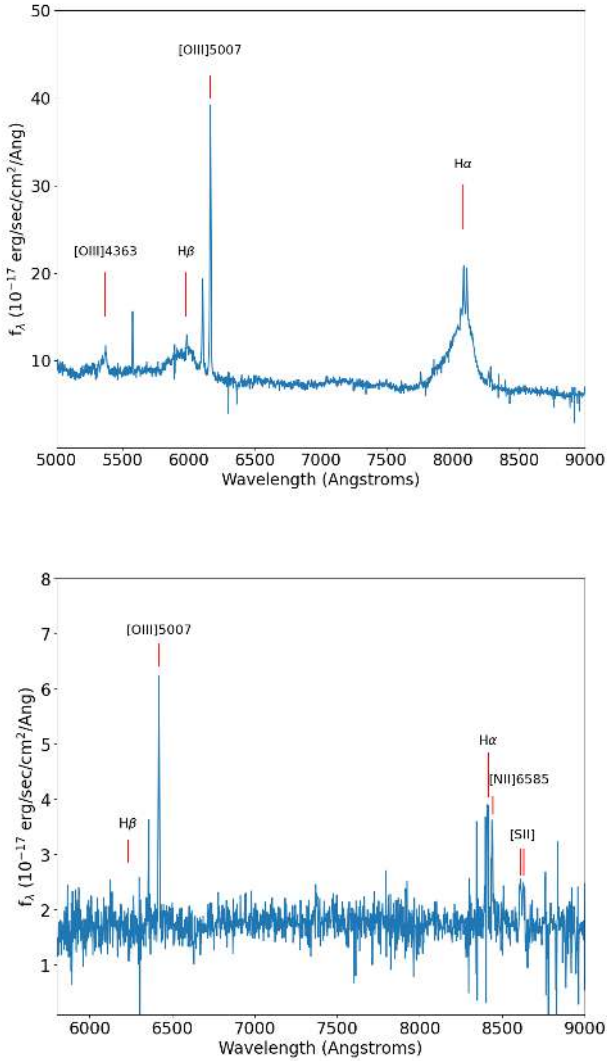
### 3.2.2. X-ray hardness ratio

The X-ray hardness ratio (HR) refers to a measure of the relative intensity of X-rays at different energies. It is typically calculated as the ratio of the counts detected in two different energy bands with the equation:

$$HR = \frac{H - S}{H + S}, \quad (2)$$

where  $H$  is the hard- and  $S$  is the soft-band count-rate. In our case, we use the [0.5–2] keV band as the soft and the [2–10] keV as the hard. The X-ray hardness ratio can serve as a proxy for absorption from the obscuring torus, assuming a spectral modelling for the continuum emission. For our purposes, we adopt a simplified model for the AGN emission, assuming a single power law fit to the data. The torus absorbs softer X-rays preferentially over harder ones due to its composition and density. Consequently, the X-ray hardness ratio, which compares the counts of X-rays at different energy bands, can indicate the degree of obscuration. High hardness ratio values indicate significant obscuration, as softer X-rays are absorbed more effectively, resulting in a relatively greater detection of harder X-rays. To compute the HR values of our sources, we used the Bayesian method described by Park et al. (2006).

In the current work, this method allows an estimation of obscuration for the full sample, contrary to the similar analysis using optical spectra where we only have data for approximately half of the sample. However, we note that although the X-ray and the optical obscuration are well correlated, they exhibit significant scatter (Jaffarian & Gaskell 2020) and overall, there is no a one-to-one correlation between them (e.g. Masoura et al. 2020, also known as **XXL Paper XL**). Furthermore, there are also indications that the X-ray spectral index depends on Eddington rate (e.g. Bianchi et al. 2009), and that the soft X-ray band, [0.5–2] keV, at low redshift may include contributions from the soft excess, leading to the underestimation of obscuration. Nevertheless, for the purpose of statistical comparisons in the current



**Fig. 6.** Top panel: Broad-line AGN spectrum (type-1 AGN) at redshift  $z = 0.231$  (XLSSU J020139.1-050118). Bottom panel: Narrow-line AGN spectrum (type-2 AGN) at redshift  $z = 0.282$  (XLSSU J021007.1-060459).

work, the use of HR as an indicator of obscuration is instructive. In Table 2 we quote the average HR values of BL and NL AGN, and ALGs, along with their optical classification. There is a good agreement between the two obscuration proxies with the BL AGN exhibiting significantly lower HR values than the NL population. The average HR of ALG sources does not indicate any heavy absorption in the X-ray band. In these sources, we assume that the AGN optical emission is diluted by the strong stellar continuum of their host galaxies.

### 3.2.3. Spectral Energy Distribution (SED)

The AGN (e.g. accretion power) and the host-galaxy (stellar mass and SFR) properties of the sources were derived through SED fitting techniques using the Code Investigating GALaxy Emission *CIGALE* algorithm (Boquien et al. 2019; Yang et al. 2020, 2022). The SED fitting analysis is described in detail in Pouliaxis et al. (2020, 2022a). In brief, we used the stellar population synthesis model defined in Bruzual & Charlot (2003) assuming the initial mass function by Salpeter (1955) and a constant solar metallicity ( $Z = 0.02$ ) for the stellar emission in addi-

**Table 2.** AGN classification.

Annulus (1)	BL (2)	NL (3)	ALG (4)
1	3 (−0.50)	6 (−0.07)	4 (−0.55)
2	1 (−0.44)	6 (0.10)	1 (−0.76)
3–4	4 (−0.66)	12 (−0.32)	4 (−0.32)

**Notes.** (1)  $r_{500}$  annulus in which the X-ray AGN are located, (2)–(4) Number of broad-line (BL) AGN, narrow-line (NL) AGN, and absorption line galaxies (ALG), as classified by their optical spectra. In the parentheses we report the average HR, computed as described on Sect. 3.2.2.

tion to a delayed star-formation history (with a functional form  $\text{SFR} \propto t \times \exp(-t/\tau)$ ) that includes a star formation burst, no longer than  $\tau = 20$  Myr (Małek et al. 2018; Buat et al. 2019). Furthermore, we used the attenuation law of the stellar emission by Charlot & Fall (2000) and we modelled the dust emission of the galaxy with the templates of Dale et al. (2014) without including the AGN emission. Finally, we used the SKIRTOR model (Stalevski et al. 2012, 2016) for the AGN emission at the different wavelengths without including the X-ray or radio modules. The parameter space for the several modules in the SED fitting process was adopted from Mountrichas et al. (2024).

It is important in our analysis, to have reliable measurements of the AGN properties along with both the global  $M_*$  and the SFR of the their host galaxies. For that purpose, we required our sources to have low reduced  $\chi^2$  ( $\chi_r^2$ ) that is indicative of the goodness of the SED fitting process. To this end, we excluded sources that have  $\chi_r^2 > 5$  (e.g., Mountrichas et al. 2019; Buat et al. 2021; Pouliaxis et al. 2022b).

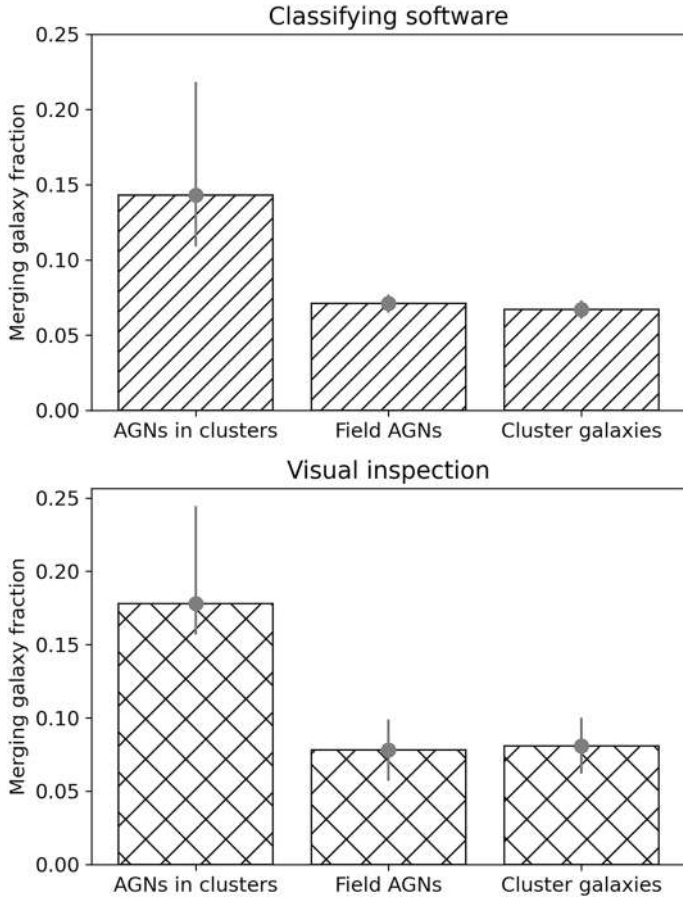
## 4. Results

### 4.1. X-ray-detected AGN in merging and disturbed galaxies

In the current study, the merging fraction in each sample is defined as the number of merging and disturbed galaxies divided by the total number of sample galaxies. The results are presented in Fig. 7. Using either Statmorph classifying software (see Sect. 3.1) or visual classification, the merging fraction of X-ray AGN hosts within clusters is significantly higher, at the  $2\sigma$  confidence level, than in field X-ray AGN or in non-active cluster galaxies.

To identify possible trends in the spatial distribution of merging hosts, we followed the methodology described in Sect. 3 to segregate the area around clusters. To improve the statistics, we merge the third and the fourth annuli, which we consider to be the field. This decision was made following our initial inspection, which confirmed that the results from the two outer annuli consistently matched those of the field AGN control samples (i and iii). Our results are plotted in Fig. 8. For comparison, we plot only the results of the large control sample of cluster galaxies (ii), as they coincide with sample (iv) but offer better statistics.

Our analysis revealed that the majority of merging and disturbed X-ray AGN hosts reside in cluster outskirts. In particular, their fraction in the second annulus is significantly higher, at the  $2\sigma$  confidence level, than the respective fraction in non-active cluster galaxies and field X-ray AGN. The merging fraction in the two outer annuli ( $0.12^{+0.08}_{-0.05}$ ) is consistent with the corresponding fraction ( $0.071 \pm 0.006$ ) in the control samples.



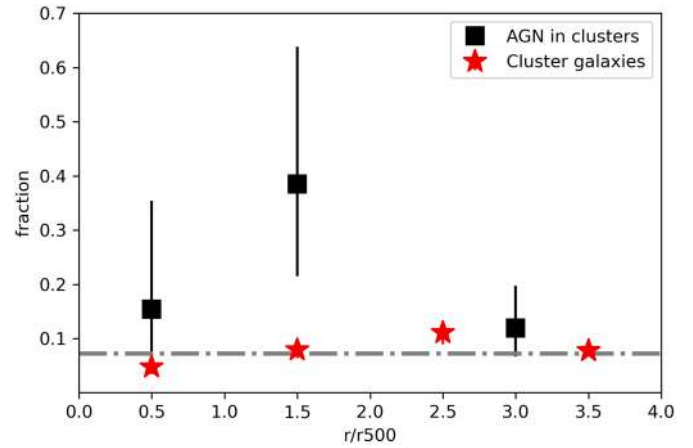
**Fig. 7.** Fraction of merging and disturbed X-ray AGN host galaxies (main sample) in comparison with the corresponding fraction in control samples (i) and (ii) (top panel), and (iii) and (iv) (bottom panel). Error bars indicate the  $1\sigma$  confidence limits for small numbers of events (Gehrels 1986). A significant excess of merging galaxies, at the  $2\sigma$  confidence level, is found in the sample of X-ray AGN in clusters.

This confirms our assumption that beyond the  $2r_{500}$  radius, AGN behave similarly to the field population, as previously reported in *XXL Paper XXXV* for the XXL sample below  $z < 0.5$ . However, the difference between the outskirts and the cluster centre is not significant at any confidence level, due to the small number of sources. The statistical results of the visual classification, including spiral and elliptical hosts, are presented in Fig. 9, while individual classification and images in Table A.1 and Table A.2, respectively.

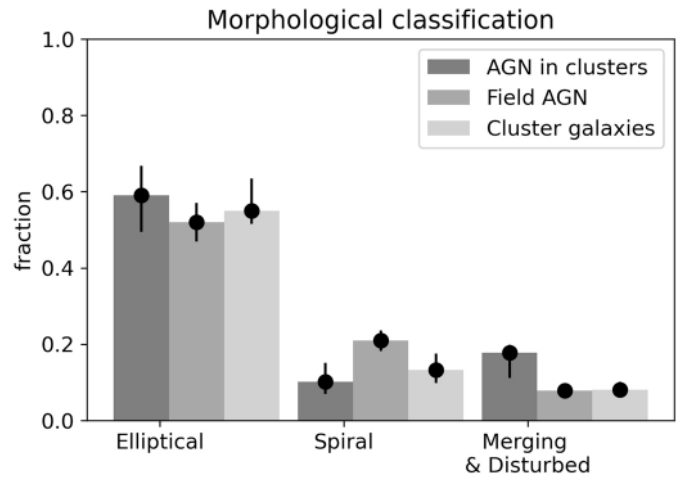
Finally, in Fig. 10 we plot the average X-ray luminosity ( $L_{X[0.5-10]\text{keV}}$ ) of AGN within clusters, which is a direct output of the XXL pipeline as described in Faccioli et al. (2018), also known as *XXL Paper XXIV*. It highlights how their radiative output varies with distance from the cluster centre. Similarly to the merging fraction, the average X-ray luminosity is three times higher than in the field, at the  $\sim 1.5\sigma$  confidence level, while the outer annuli are in agreement with the field value. This may further support the idea that the relatively high number of galaxy mergers in the cluster outskirts contribute to the triggering of AGN.

#### 4.2. AGN Obscuration

In this section, we examine the obscuration of X-ray AGN in our samples. Our aim was to investigate potential trends that may



**Fig. 8.** Fraction of merging X-ray AGN in clusters (main sample) as a function of distance from the cluster centre. For comparison, we overlay the respective fraction for the 1914 non-AGN cluster galaxies (control sample (ii)). The hatched line represents the average value of the merging fraction derived from the sample of 1987 field X-ray AGN (control sample (i)). Error bars indicate the  $1\sigma$  confidence limits for small numbers of events (Gehrels 1986). The results indicate a significant excess, at the  $2\sigma$  confidence level, of merging AGN hosts in cluster outskirts.

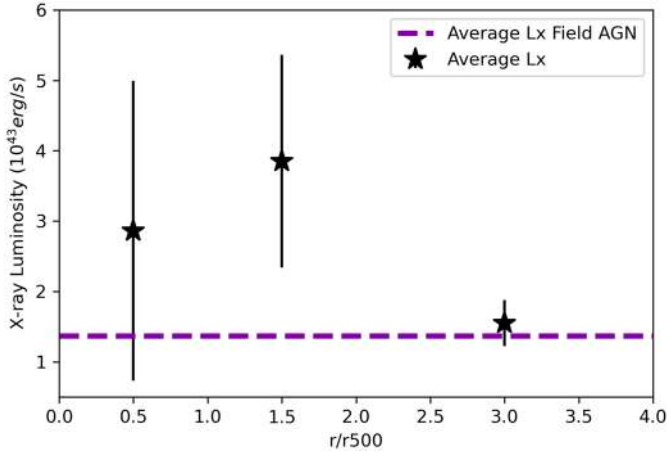


**Fig. 9.** Visual morphological classification of X-ray AGN hosts in clusters in comparison with field X-ray AGN (iii) and non-AGN cluster galaxies (iv). The classification was based on visual inspection of HSC images. Dominance of elliptical shapes was expected, particularly among cluster galaxies. The careful selection of control-sample galaxies based on redshift and stellar mass ensured the inclusion of similar galaxy types across all samples. Error bars indicate the  $1\sigma$  confidence limits for small numbers of events (Gehrels 1986).

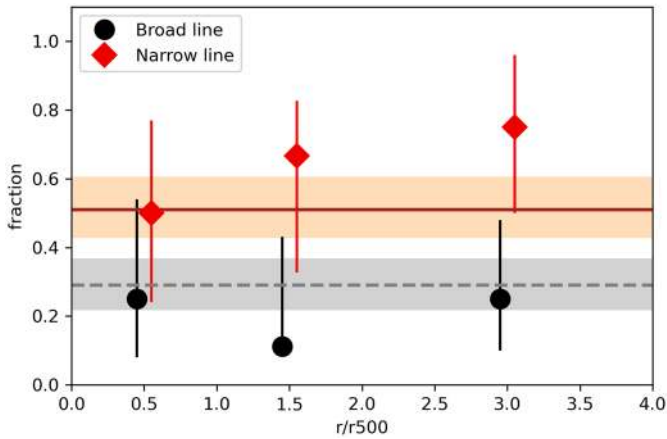
reveal the physical conditions responsible for triggering AGN in the dense environments of galaxy clusters. We use two indicators of obscuration, as described in Sect. 3; the AGN optical type (type-1/type-2) and the X-ray hardness ratio.

##### 4.2.1. Optical AGN classification

First, we investigate potential trends related to AGN classification based on the optical spectra of our X-ray-detected AGN. In total, we find eight broad-line and 22 narrow-line AGN. Another six sources have an absorption-line galaxy (ALG) spectrum, typical of elliptical galaxies. However, most of the ALG present some weak emission lines, especially in the  $H\alpha$  region. Optical



**Fig. 10.** Average X-ray luminosity ( $L_{X[0.5-10] \text{ keV}}$ ) of X-ray AGN in clusters. The hatched line represents the average value of  $L_x$  derived from the sample of 1987 field X-ray AGN (control sample (i)). Error bars denote the  $1\sigma$  confidence limits.



**Fig. 11.** AGN optical classification of X-ray AGN in clusters (black circles for broad-line and red squares for narrow-line AGN) compared to the control sample of field X-ray AGN (dashed black line for broad-line and continuous red line for narrow-line AGN). The classification and detailed results are described in Sect. 3.2.1 and Sect. 4.2.1. Error bars and shaded areas indicate the  $1\sigma$  confidence limits for small numbers of events (Gehrels 1986).

emission from these AGN, which are hosted mostly by massive elliptical galaxies, could be diluted by the stellar continuum. In addition, these AGN may be intrinsically weak, since the X-ray luminosity for four out of six is below  $3 \times 10^{42}$  erg/sec, placing them in the first quartile. The results are presented in Fig. 11. We were unable to identify any significant trends in the cluster-centric radial distribution of optical AGN types, likely due to the small sample size and resulting large uncertainties.

The ratio of type-2 to type-1 AGN identified in the current study is consistent with the field ratio observed in the Local Universe (e.g. Maia et al. 2003) and in high-redshift clusters ( $z \sim 1$ ) (Mo et al. 2018). This is in good agreement with a similar study of spectroscopically confirmed X-ray AGN in 19 galaxy clusters (Koulouridis et al. 2024), within a narrow redshift range  $0.16 \leq z \leq 0.28$ . However, our findings diverge from recent results obtained from local clusters in the WINGS and Omega-WINGS surveys. In those studies, a notably higher optical type-2 to type-1 fraction was reported (approximately 10 to 1) compared to the field (Marziani et al. 2023). However, the sample

selections in these studies differ significantly, so any comparisons should be approached with caution.

Interestingly, Koulouridis et al. (2024) reported that three out of the four broad-line AGN within  $2r_{500}$  were found in the central  $r_{500}$  annulus. They argued that close to the cluster core, AGN activity may be triggered either by the influence of strong RPS, as suggested by studies of “jellyfish” galaxies (Poggianti et al. 2017b; Peluso et al. 2022), or by tidal shocks, as galaxies pass through cluster pericentre. Our results support these findings, since three out of the four broad-line sources up to  $2r_{500}$  were found again in the same region. We note, that the cluster mass distribution of the current sample, mostly comprising groups and poor clusters, is markedly different from the massive clusters used in Koulouridis et al. (2024). In addition, the SED analysis showed that the accretion power of these three broad-line sources is higher than that of the rest of the X-ray AGN sample, and also of the other five broad-line AGN found farther from the cluster centre. We note that the fraction of broad-line AGN in the first two annuli is comparable to that in the last two annuli, which represent the field. However, the environment is drastically different.

We note that only half of our X-ray AGN have available spectroscopic data. Consequently, in the next section, we present the results of a statistical analysis based on the X-ray hardness ratio of our entire source sample.

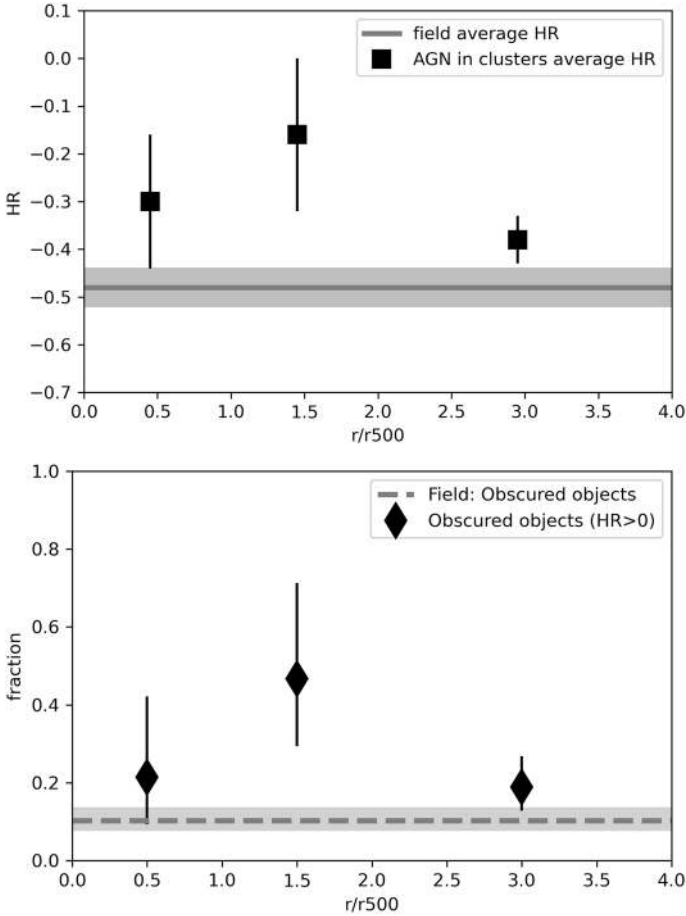
#### 4.2.2. X-ray hardness ratio

The X-ray hardness ratios of cluster and field X-ray AGN were computed following the methodology described in Sect. 3.2.2. For our purposes, we then derived the mean HR in each cluster annulus. The results are presented in Fig. 12. We cannot claim any difference between the outskirts and the central annulus because of the large uncertainties. Nevertheless, The mean HR values indicate that X-ray AGN in cluster outskirts are more obscured compared to those in the outer annuli ( $2-4r_{500}$ ), at the  $1\sigma$  confidence level, or in the field, at the 90% confidence level (Gehrels 1986).

The higher average obscuration observed in the outskirts compared to the field may result from a higher number of obscured sources or, alternatively, from the presence of a few heavily obscured sources. To determine the source of obscuration, we needed to estimate the fraction of sources that are obscured in each annulus. To this end, we used PIMMS online tool<sup>3</sup> to compute the HR value that would best define a threshold between obscured and unobscured sources based on the column density of the obscuring material. A reasonable value of the column density above which we can assume that a source is obscured is  $10^{22} \text{ cm}^{-2}$ . The HR that corresponds to this value depends on redshift and the slope ( $\Gamma$ ) of the X-ray spectrum. Our computations within the redshift range of our sources and for various values of  $\Gamma$  result in an average HR = 0 as a reasonable threshold. The results are plotted in the bottom panel of Fig. 12. The fraction of obscured sources in the outskirts of clusters are significantly higher than the corresponding fraction in the field at the  $2\sigma$  confidence level. However, we cannot confirm any statistically significant difference between the outskirts and the centre of clusters due to small number statistics.

The results from the HR analysis are seemingly not in agreement with the results from the optical spectra analysis presented in the previous section. However, the two analyses are not

<sup>3</sup> <https://heasarc.gsfc.nasa.gov/cgi-bin/Tools/w3pimms/w3pimms.pl>



**Fig. 12.** Top panel: Average hardness ratio (black squares) of X-ray detected AGN in clusters divided in three cluster-centric annuli, 0–1R<sub>500</sub>, 1–2R<sub>500</sub> and 2–4R<sub>500</sub> with their standard 1σ errors. The grey line corresponds to the average hardness ratio value of field AGN (the shaded area denotes the 1σ error range). Evidently, the average HR in the 1–2R<sub>500</sub> area is significantly higher than in the field, at the 90% confidence level (Gehrels 1986). Bottom panel: Obscured X-ray AGN fraction in clusters. We consider as obscured AGN the objects with HR > 0 (see Sect. 4.2.2). The dashed line represents the fraction of the obscured X-ray AGN in the field (the shaded area denotes the 1σ error range). There is a significant excess of X-ray AGN in cluster outskirts, at the 2σ confidence level.

directly comparable, since the samples are different. The current sample includes many faint sources ( $m_i > 21$ ) without spectroscopic data, possibly more obscured than the spectroscopic part of the sample. More importantly, as described in Sect. 3.2.2, the correlation between these two obscuration proxies presents large scatter.

In summary, the average obscuration is higher in cluster outskirts than in the field, at the 2σ confidence level, due to a larger number of obscured AGN. The obscuration excess coincides with the merging excess shown in Fig. 8, further supporting galaxy merging as the main AGN triggering mechanism in the outskirts.

#### 4.3. AGN and host galaxy properties derived from SED fitting

The SED analysis of the X-ray AGN in our cluster sample may reveal some interesting environmental trends. Specifically, we statistically examine the black hole mass ( $M_{\text{BH}}$ ), the black hole

accretion rate (BHAR), and the Eddington rate ( $\lambda_{\text{edd}}$ ) of the AGN, as well as the SFR of their host galaxies.

In order to estimate the  $M_{\text{BH}}$ , we used the following equation from Kormendy & Ho (2013) that takes into account the stellar-mass ( $M_*$ ) dependence:

$$M_{\text{BH}} = 4.9 \times 10^{-3} \left( \frac{M_*}{10^{11} M_{\odot}} \right)^{0.14} M_*, \quad (3)$$

where  $M_*$  is the quantity “bayes.stellar.mstar” from CIGALE. The average BHAR and  $\lambda_{\text{edd}}$  may provide insight into the mass growth rates of black holes across different cluster regions. For the calculation, we used the following equation from Yang et al. (2023), Pouliaxis et al. (2024):

$$\text{BHAR} = \frac{L_{\text{disk}}(1 - \varepsilon)}{\varepsilon c^2} = 1.59 \times \frac{L_{\text{disk}}}{10^{46} \text{ erg s}^{-1}} [M_{\odot} \text{ yr}^{-1}], \quad (4)$$

where  $L_{\text{disk}}$  is the quantity “agn.accretion.power” from our SED fits, which is the viewing-angle-averaged intrinsic accretion-disk luminosity. Also,  $\varepsilon$  is the radiative efficiency, the fraction of the accreted mass converted into radiation. For comparison, we derive the  $\lambda_{\text{edd}}$  of sources both from the SEDs and from the X-rays, assuming  $L_{\text{disk}} \sim L_{\text{bol}}$ . In particular, the  $\lambda_{\text{edd}}$  calculated from the SEDs is:

$$\lambda_{\text{edd}} = \frac{L_{\text{disk}}}{L_{\text{edd}}}, \quad (5)$$

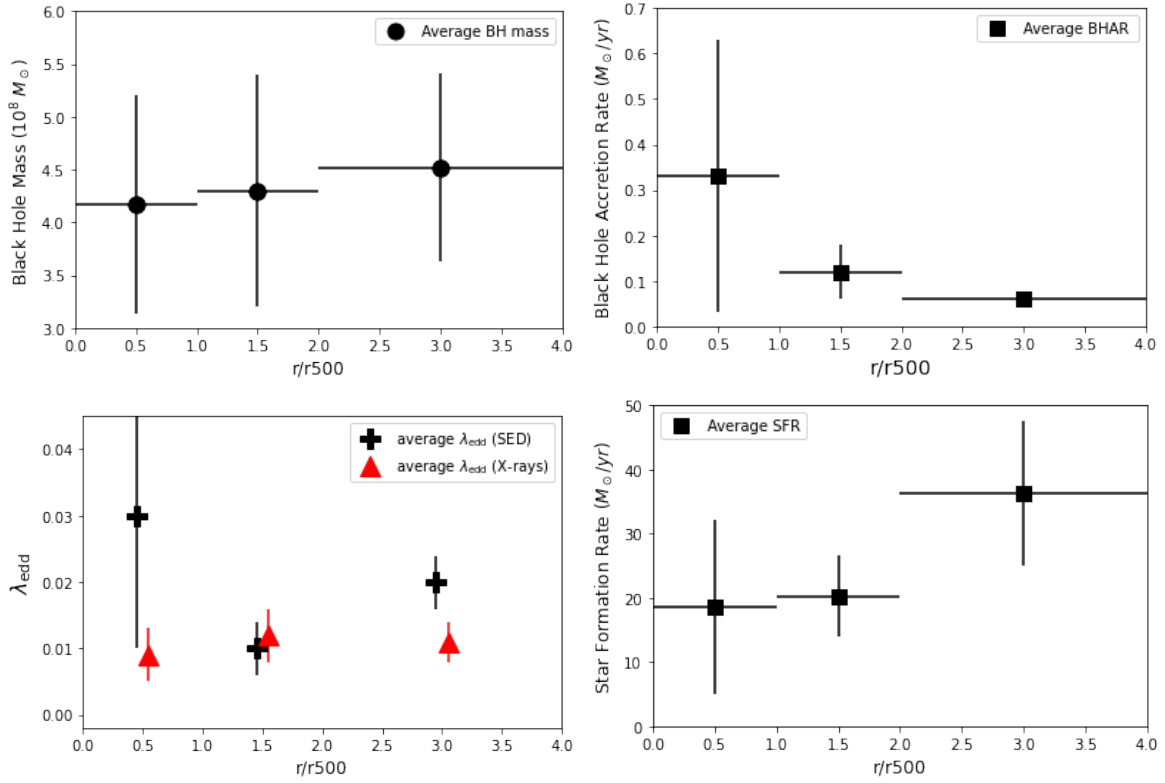
while from the X-rays it is

$$\lambda_{\text{edd}} = \frac{Kx(Lx) \times Lx}{L_{\text{edd}}}, \quad (6)$$

where  $L_{\text{edd}} = 1.3 \times 10^{38} (M_{\text{BH}}/M_{\odot})$  and  $Kx(Lx)$  is the bolometric correction as described in Duras et al. (2020).

The top panels of Fig. 13 illustrate the average values of the  $M_{\text{BH}}$  and the BHAR as a function of cluster-centric distance. We found an average  $M_{\text{BH}} \sim 4.5 \times 10^8 M_{\odot}$  and BHAR  $\sim 10^{-1} M_{\odot}/\text{yr}$  for all X-ray AGN in all annuli. Similarly, we find no statistically significant differences in the Eddington rate of AGN in clusters (bottom left panel of Fig. 13), either calculated from the SEDs or directly from the X-rays. However, we found evidence that the average SFR (bottom right panel of Fig. 13) is decreasing toward the cluster centres, although the uncertainties are large. This is in agreement with results from both the local and high-redshift Universe, where the mean SFR was always found to be higher in field galaxies than in the cluster centres (Treu et al. 2003; Poggianti et al. 2006; Raichoor & Andreon 2014; Haines et al. 2015). The decreasing SFR is probably an indication of the ram pressure induced by the ICM and frequent galaxy-galaxy interactions, which are very effective in stripping the gas of infalling galaxies (e.g. Larson et al. 1980; Cole et al. 2000; Balogh et al. 1999; Kawata & Mulchaey 2008; Boselli et al. 2022), transforming them into inactive ellipticals.

In addition, sources closer to cluster centres have slightly larger stellar masses ( $\text{med}(M_*) \approx 6.5 \times 10^{10} M_{\odot}$ ) compared to AGN in the 2–4 annuli ( $\text{med}(M_*) \approx 5.1 \times 10^{10} M_{\odot}$ ). Further analysis of the SFR-stellar mass relation of our AGN relative to the main sequence (MS) of the star-forming galaxies is presented in Appendix C. The SEDs plotted in Fig. 14 are examples (with a good fit of  $\chi^2 < 5$ ) of different categories of host galaxies of our sample, namely a spiral, an elliptical, and a merging host galaxy, respectively.



**Fig. 13.** Summary plot of X-ray AGN and host galaxy properties as a function of cluster-centric distance. From top left to bottom right panel: Average  $M_{\text{BH}}$ , BHAR,  $\lambda_{\text{Edd}}$  of AGN, and average SFR of their host galaxies. Error bars denote the  $1\sigma$  confidence limits.

## 5. Discussion and conclusions

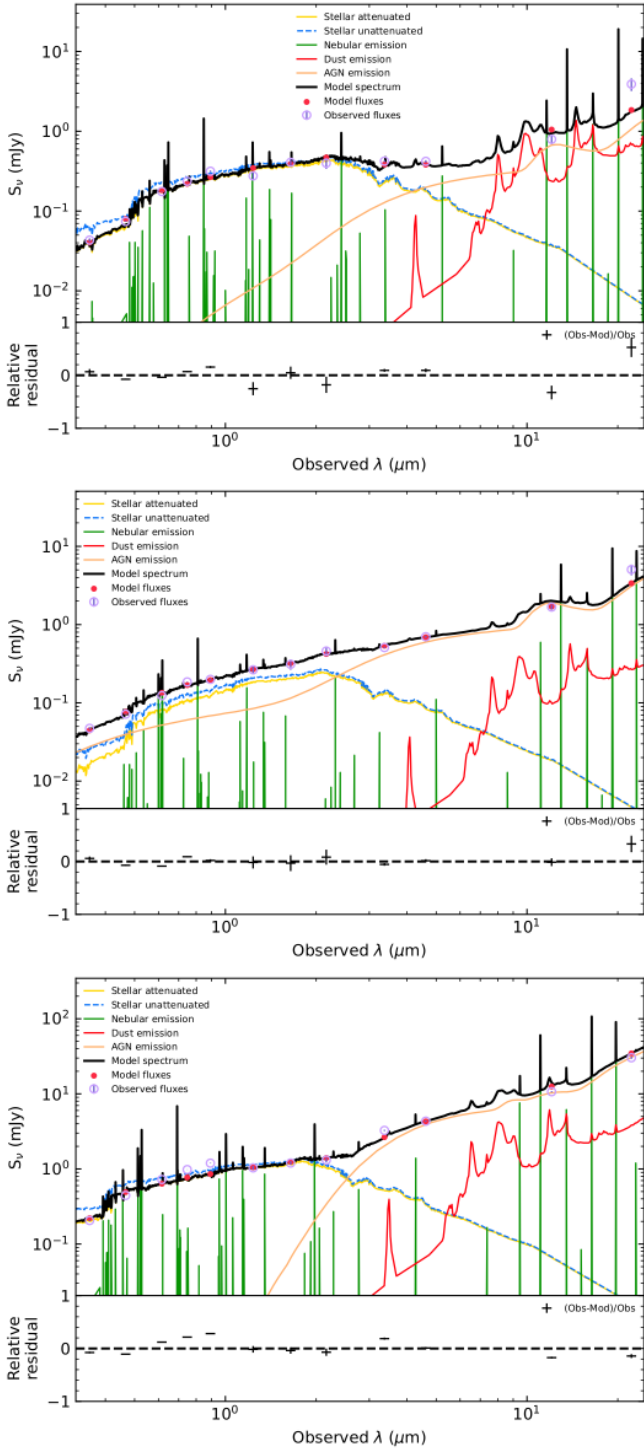
In the current study, our aim was at investigating the role of specific properties of X-ray-detected AGN in galaxy clusters, such as the merging fraction and the morphology of their host galaxies, obscuration, star formation and accretion rate. To this end, we have used a main sample of 82 X-ray AGN detected within  $4r_{500}$  radii of XXL clusters. Moreover, we assembled four control samples (consisting of field X-ray AGN and non-AGN cluster galaxies) to enable a thorough comparative analysis.

We discovered that cluster galaxies that host X-ray AGN are significantly more often merging or disturbed, at the  $2\sigma$  confidence level, than similar X-ray AGN hosts in the field and cluster galaxies without X-ray detected AGN. These findings align with earlier research, which indicates that the interactions occurring during mergers may play a crucial role in triggering the formation of AGN in dense environments (e.g. Ehlert et al. 2015; Koulouridis et al. 2024). Mergers can trigger the formation of new stars, as well as the recirculation of gas and dust. Furthermore, they are thought to be a very efficient way to trigger AGN activity, as they can provide a large amount of gas and dust to fuel the SMBH (e.g. Silk & Rees 1998; Granato et al. 2004; Häring & Rix 2004; Springel et al. 2005; Koulouridis et al. 2006b; Hopkins et al. 2008; Zubovas & King 2012; Steffen et al. 2023; Li et al. 2023; Comerford et al. 2024; Duan et al. 2024; La Marca et al. 2024; Bonaventura et al. 2025).

Likewise, we found that the excess of merging and disturbed AGN hosts can be largely attributed to galaxies in cluster outskirts (between  $r_{500}$  and  $2r_{500}$ ). This result supports the scenario in which the excess of AGN activity in the outskirts of clusters (Johnson et al. 2003; Ruderman & Ebeling 2005;

Branchesi et al. 2007; Fassbender et al. 2012; Koulouridis et al. 2014, 2024; Koulouridis & Bartalucci 2019; Hashiguchi et al. 2023) can be attributed to a high frequency of galaxy interactions and merging. While high-velocity dispersions in massive clusters may diminish the probability of galaxy interactions (e.g. Arnold et al. 2009; Manzer & De Robertis 2014), Haines et al. (2012) proposed the existence of two separate groups of X-ray AGN in the outskirts: newly infalling galaxies and those that have already passed the pericentre and are nearing the apocentre of their orbit, termed the back-splash population. These galaxies likely have low velocities, potentially facilitating interactions and mergers with other galaxies in the outskirts. Supporting our findings, a recent study of mergers in six low-redshift clusters suggested that post-merging galaxies likely merged in low-velocity environments, such as clusters outskirts and dynamically relaxed clusters (Kim et al. 2024).

A high merging frequency may result in increased absorption of the AGN emission, as obscuring material is driven towards the galaxy nucleus (e.g. Hopkins et al. 2008; Koulouridis et al. 2014, 2016b; La Marca et al. 2024). Therefore, investigating the obscuration of the sources may provide some further evidence of the mechanisms that trigger AGN in clusters. To this end, we have followed two different approaches to estimate the intrinsic absorption in the samples under study: investigating (a) the optical spectra of the sources that have spectroscopic data ( $\sim 50\%$  of the sample), and (ii) the X-ray hardness ratio (full sample). The classification of the optical spectra showed a predominance of narrow emission-line spectra among AGN situated in galaxy clusters. However, the broad to narrow-line ratio in clusters (up to  $2R_{500}$ ) is not dissimilar to the respective ratio in the field. Nevertheless, three out of four broad-line sources are located within the central  $R_{500}$  annulus, including the only powerful



**Fig. 14.** SED examples of AGN host galaxies. Top panel: Broad-line passive/disturbed spiral galaxy at  $z = 0.445$  (4XLSSU J021118.8-042516). Middle panel: Narrow-line elliptical galaxy at  $z = 0.054$  (4XLSSU J022700.7-042022). Bottom panel: Narrow-line merging galaxy at  $z = 0.959$  (4XLSSU J021616.8-045033).

QSO (with  $AP = 5 \times 10^{44}$  erg/sec,  $L_X = 3 \times 10^{44}$  erg/sec and no visible host galaxy) found in this sample. Despite the small number statistics, this is an interesting trend corroborated by recent similar results from a different sample of 19 clusters below  $z = 0.5$  (Koulouridis et al. 2024). As suggested in that publication, this trend may indicate a different triggering mechanism or accretion rate, between the cluster’s core and its periph-

ery. In more detail, AGN activity near the cluster centre might arise from two primary triggers: the influence of strong RPS, as indicated by research on “jellyfish” galaxies (Poggianti et al. 2017b; Peluso et al. 2022), and tidal shocks occurring as galaxies traverse the pericentre of a cluster. It is important to note that the SED analysis revealed a significantly higher value of accretion power in these broad-line AGN. This indicates that the processes driving the growth of these SMBH are more intense, providing further support for the AGN triggering by RPS scenario. The accretion power rate is a crucial factor in understanding the dynamics and evolution of these AGN and influences the emission and the overall energetic output in cluster cores. This heightened activity within the core region suggests that a number of centrally located AGN may have more efficient energy feedback, which could have significant implications for the surrounding cluster environment.

Regarding the X-ray hardness ratio distribution of the full sample, our analysis indicates that the emission coming from AGN located in cluster outskirts is significantly more absorbed when compared either to the control sample of field X-ray AGN (at the 90% confidence level) or to the  $2-4R_{500}$  region (at the  $2\sigma$  level). Furthermore, we confirmed that this higher obscuration is due to a high number of obscured sources in the outskirts and not just a few highly obscured ones. These findings align with the noted increase in galaxy merging and interactions within the outskirts. The excess of obscured sources in this region suggests that the dense, dynamic environment, characterised by frequent merging events, plays a crucial role in the process of AGN triggering.

## 6. Summary

The conclusions that can be drawn from the above analysis are the following:

- We identified a significant excess, at the  $2\sigma$  confidence level, of X-ray AGN in galaxy clusters hosted by merging or disturbed galaxies, compared to non-active cluster galaxies or X-ray AGN in the field. This excess can be localised in the cluster outskirts (between 1 and  $2r_{500}$  radius). Previous studies suggest a connection between galaxy interactions and the triggering of AGN in cluster environments. We argue that galaxy merging and interactions are more likely to happen in the outskirts of clusters, causing a higher frequency of AGN triggering than in field galaxies and thus, leading to the observed excess.
- The hardness ratio distribution indicates higher level of obscuration in clusters than in the field, specifically in the outskirts, further supporting that galaxy merging and interactions may act as an AGN triggering mechanism.
- The ratio of narrow-line to broad-line AGN in clusters is comparable to that in the field. Interestingly, three out of four broad-line AGN, including the only QSO, are found close to the cluster centre (within  $r_{500}$ ). SED analysis of these broad-line AGN revealed that their accretion power is among the highest in our sample. These results are possibly indicating AGN triggering in infalling galaxies by RPS.

Future research should focus on maximizing AGN sample sizes within meticulously characterised cluster datasets. Our future plans include thorough studies of AGN in large X-ray-selected cluster samples, such as X-CLASS (Koulouridis et al. 2021), XCS (Mehrtens et al. 2012; Giles et al. 2022) and eROSITA (Bulbul et al. 2024).

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## Appendix A: Inventory of host galaxies of X-ray detected AGN in clusters

Table A.1 provides a list of the 82 X-ray detected AGN in our cluster sample, while Table A.2 display their colour images.

**Table A.1.** 82 X-ray detected AGN in clusters within  $4r_{500}$  radius.

Cluster name XLSSC	$r_{500}$	source name 4XLSSU	RA (J2000) degrees	Dec (J2000) degrees	spec/phot	$z$	host morphology	Spectral type
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)
025	1	J022521.1-043950	36.337835	-4.663836	spec	0.265	E	NL
030	1	J022310.5-041249	35.795024	-4.214393	spec	0.626	M	ALG
040	1	J022206.2-043251	35.523711	-4.546156	spec	0.323	E	ALG
080	1	J021819.1-052343	34.579399	-5.396029	spec	0.647	E	ALG
082	1	J021046.1-060854	32.692832	-6.148447	spec	0.428	E	NL
091	1	J023147.2-045702	37.947556	-4.950985	spec	0.190	M	NL
111	1	J021234.2-053545	33.142287	-5.595251	spec	0.302	S	NL
114	1	J020139.1-050118	30.413549	-5.021844	spec	0.231	E	BL
117	1	J021235.9-053210	33.14998	-5.536541	spec	0.299	E	BL
142	1	J021856.0-052611	34.733459	-5.435817	spec	0.448	E	U
168	1	J022935.0-055210	37.396398	-5.869275	spec	0.293	U	QSO BL
183	1	J022016.9-045645	35.070248	-4.946167	spec	0.517	E	NL
187	1	J021631.8-042958	34.132286	-4.499856	spec	0.454	E	ALG
194	1	J021648.5-043318	34.202471	-4.555327	spec	0.41	S/tidal	NL
040	2	J022206.5-042909	35.52714	-4.485877	spec	0.315	E	ALG
049	2	J022347.1-043346	35.946002	-4.563383	spec	0.49	D	NL
071	2	J022238.1-050100	35.659184	-5.016476	phot	0.87	E	-
071	2	J022236.1-050143	35.650138	-5.027919	spec	0.844	D	NL
083	2	J021035.8-061027	32.649452	-6.17388	spec	0.433	E	NL
089	2	J022832.6-044608	37.134301	-4.768676	phot	0.60	E	-
097	2	J021325.6-060329	33.357639	-6.058701	spec	0.694	U	U
101	2	J020838.3-042511	32.15988	-4.419889	spec	0.753	E	NL
105	2	J023324.9-053106	38.35469	-5.517847	spec	0.436	M	NL
107	2	J020534.3-073708	31.392291	-7.618196	phot	0.43	no HSC	-
110	2	J021413.9-053405	33.557274	-5.568639	phot	0.440	no HSC	NL
116	2	J021047.8-060354	32.699967	-6.064865	phot	0.53	E	-
130	2	J022053.8-052538	35.224761	-5.42698	phot	0.54	E	-
156	2	J020305.0-070948	30.771721	-7.163623	phot	0.33	D/tidal	-
158	2	J021118.8-042516	32.828806	-4.422185	spec	0.445	S/D	BL
048	3	J022236.1-032639	35.651316	-3.444693	spec	1.008	E	NL
083	3	J021117.5-061916	32.825396	-6.320064	phot	0.43	E	-
083	3	J021052.9-061809	32.721253	-6.302944	spec	0.423	E	NL
085	3	J021125.7-061936	32.858153	-6.327543	phot	0.42	U	-
085	3	J021153.7-061033	32.974242	-6.177046	spec	0.420	E	U
101	3	J020853.8-042937	32.225226	-4.493658	phot	0.75	E	-
139	3	J021649.4-032831	34.206554	-3.474726	phot	0.21	E/tidal	-
157	3	J020336.4-070010	30.901106	-7.003745	phot	0.59	E	-
159	3	J020919.5-051152	32.329839	-5.198636	spec	0.612	M	NL
163	3	J021007.1-060459	32.529514	-6.083835	spec	0.282	E	NL
172	3	J020613.0-054957	31.554222	-5.832453	phot	0.43	E	-
183	3	J022029.8-044657	35.123919	-4.782901	phot	0.51	E/D	-
200	3	J020115.4-064331	30.314141	-6.725744	phot	0.32	E	-
001	4	J022445.4-035509	36.188901	-3.919289	spec	0.605	U	QSO BL
003	4	J022750.3-032106	36.958818	-3.352099	phot	0.84	U	-
008	4	J022519.4-035444	36.331233	-3.912253	spec	0.299	ring/D	BL
011	4	J022700.7-042022	36.753235	-4.339044	spec	0.053	E	NL
018	4	J022430.5-050842	36.127355	-5.144935	spec	0.322	S/tidal	NL
029	4	J022418.9-041316	36.078919	-4.221791	spec	1.057	E	U
030	4	J022254.3-041629	35.725574	-4.274752	spec	0.63	E	NL
056	4	J021537.1-045005	33.904892	-4.834451	spec	0.350	E	ALG
064	4	J021818.4-045843	34.576028	-4.978517	phot	0.87	E	-
067	4	J021835.9-053758	34.649637	-5.632743	spec	0.387	E	BL
071	4	J022255.1-045328	35.730888	-4.891314	phot	0.83	E/U	-
077	4	J021731.0-032444	34.380138	-3.41251	phot	0.20	S	-

**Table A.1.** Continued.

078	4	J021616.8-045033	34.069856	-4.842637	spec	0.959	E/M	NL
078	4	J021610.6-045232	34.044313	-4.874921	spec	0.956	E/D	NL
091	4	J023138.0-051420	37.908463	-5.238549	spec	0.187	E	NL
093	4	J020629.7-064905	31.624312	-6.81841	phot	0.42	E	-
093	4	J020614.4-065635	31.561428	-6.94357	phot	0.43	U	-
097	4	J021304.4-060037	33.268811	-6.009716	phot	0.69	E	-
104	4	J022841.1-055724	37.171036	-5.956224	spec	0.297	E	ALG
106	4	J020516.2-055230	31.317255	-5.875662	spec	0.295	E	NL
107	4	J020553.3-073743	31.473609	-7.627636	phot	0.43	no HSC	-
109	4	J020909.6-062617	32.289923	-6.437786	phot	0.49	E	-
109	4	J020930.7-062542	32.378151	-6.428463	phot	0.49	E	-
110	4	J021353.0-053917	33.470966	-5.655553	phot	0.45	E	-
116	4	J020954.7-060712	32.479197	-6.120809	phot	0.53	U	-
124	4	J021759.0-045551	34.496463	-4.931106	spec	0.506	E/tidal	ALG
128	4	J022342.4-030411	35.927244	-3.069659	phot	0.48	E	-
135	4	J021453.6-035929	33.724019	-3.991403	phot	0.37	U	-
137	4	J021815.7-034141	34.565714	-3.694404	phot	0.29	U	-
149	4	J022959.9-045716	37.498909	-4.954494	spec	0.286	U	NL
158	4	J021144.9-041843	32.937205	-4.311633	phot	0.44	E/D	-
158	4	J021043.2-042509	32.679161	-4.420212	phot	0.45	E	-
169	4	J023044.0-053605	37.68341	-5.601003	phot	0.49	S	-
169	4	J023007.4-054902	37.530904	-5.81804	phot	0.49	E	-
182	4	J022542.5-032024	36.426675	-3.340322	spec	0.170	E	BL
183	4	J021922.3-045943	34.844415	-4.995945	phot	0.51	E	-
183	4	J021921.5-045838	34.840431	-4.977854	spec	0.512	E	U
184	4	J022115.7-040901	35.315354	-4.150097	phot	0.81	U	-
187	4	J021543.7-042456	33.932907	-4.414839	spec	0.457	E	ALG
199	4	J020019.5-064750	30.082133	-6.797771	phot	0.33	E	-

**Notes.** (1) XXL Cluster name, (2)  $r_{500}$  annulus where the X-ray detected AGN is located, (3) X-ray point-source ID (4) right ascension of the optical counterpart, (5) declination of the optical counterpart, (6) flag for spectroscopic (spec) or photometric (phot) redshift, (7) redshift, (8) morphological classification of the host galaxy, (9) optical spectral classification: narrow-line (NL), broad-line(BL), absorption-line galaxy (ALG), undefined (U)

**Table A.2.** RGB colour images of X-ray-detected AGN host galaxies, based on the gri filters of HSC.

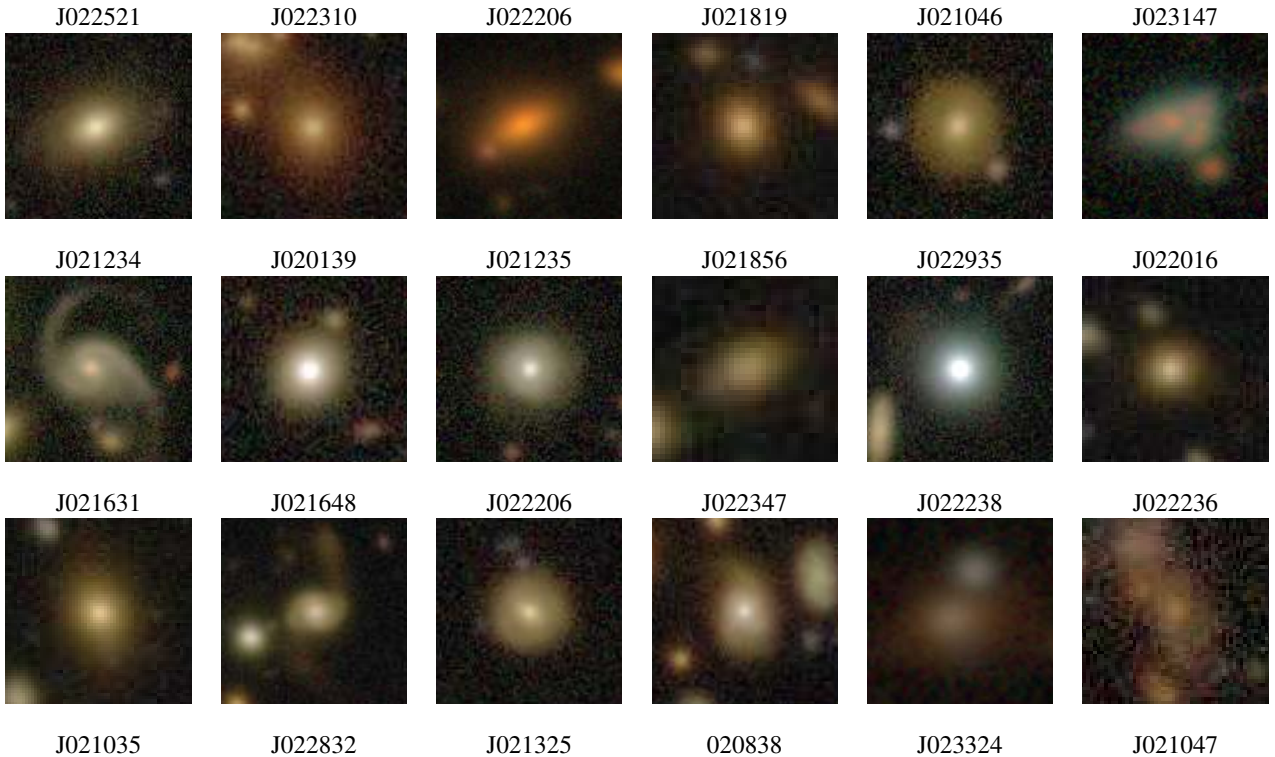


Table A.2. Continued.

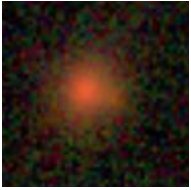
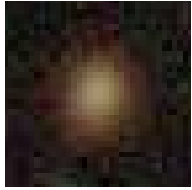
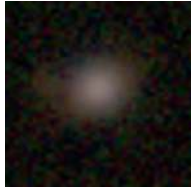

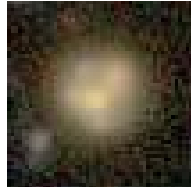


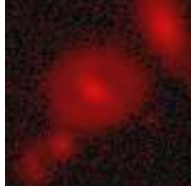


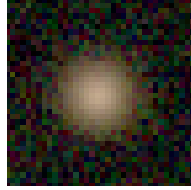
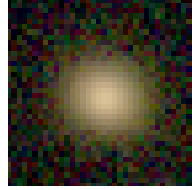
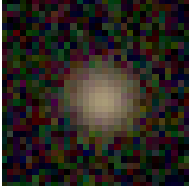

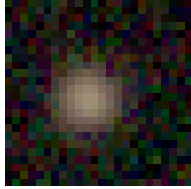

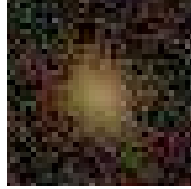
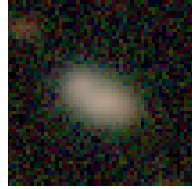


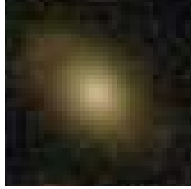
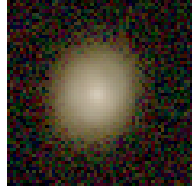
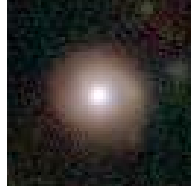
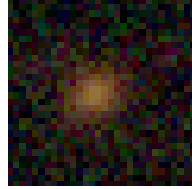




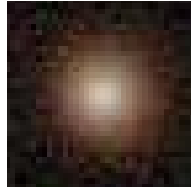

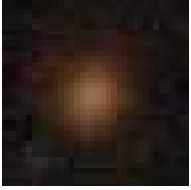
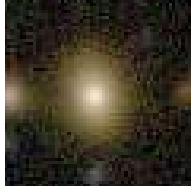


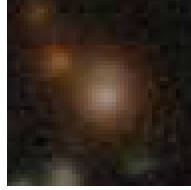
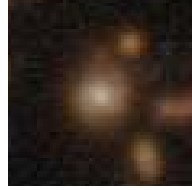
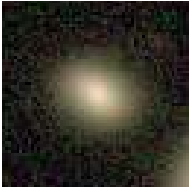
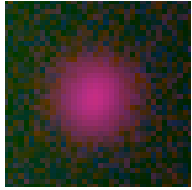
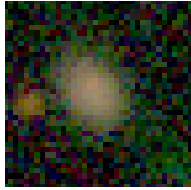
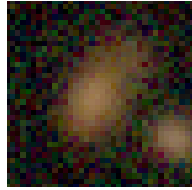

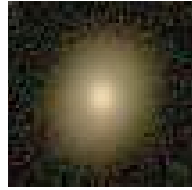
					
J022053	J020305	J021118	J022236	J021117	J021052
					
J021125	J021153	J020853	J021649	J020336	J020919
					
J021007	J020613	J022029	J020115	J022445	J022750
					
J022519	J0222700	J022430	J022418	J022254	J021537
					
J021818	J021835	J022255	J021731	J021616	J021610
					
J023138	J020629	J020614	J021304	J022841	J020516
					
J020909	J020930	J021353	J020954	J021759	J022342

Table A.2. Continued.

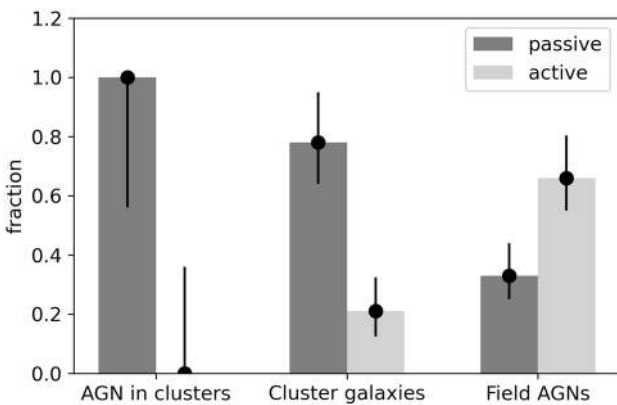


## Appendix B: The case of spiral galaxies

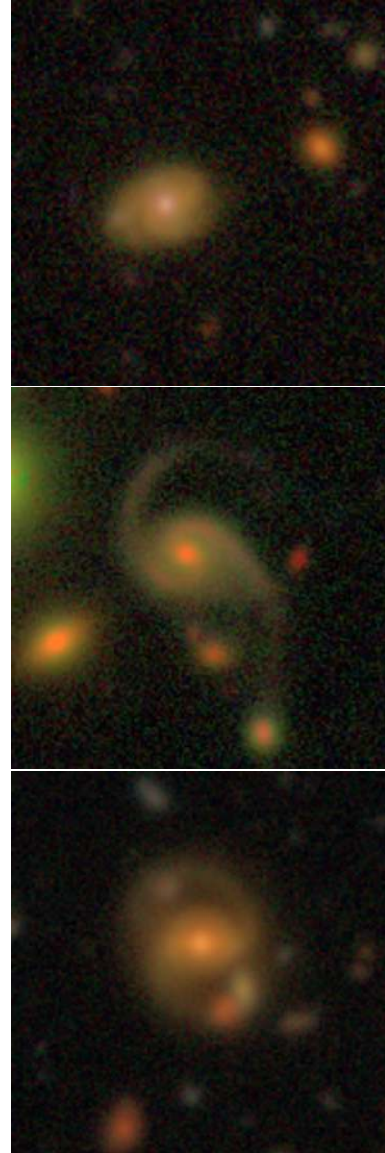
In this section, we focus specifically on the spiral host galaxies of X-ray AGN in clusters, comparing them with the spiral galaxy population in the control samples. We categorised them into two types: passive (red) and active (blue). The colours were defined using HSC *gri* photometric magnitudes. We define a galaxy as red and thus passive when  $g - r > 0.5$  (Daoutis et al. 2023; Shimakawa et al. 2022). Passive spiral galaxies are characterised by a lack of star formation, in sharp contrast to active spiral galaxies (Shimakawa et al. 2022), which are also the vast majority. Our results are presented in Fig. B.1. Despite the limited number of spiral host cases, we observe a clear trend in the results. Specifically, in the main sample of X-ray AGN in clusters, all spiral hosts are passive, whereas in the field, active hosts dominate. We also examined the spiral host galaxies in the control sample of cluster galaxies, finding mostly passive cases with only a few active ones. Additionally, we reviewed the SFR of our sample galaxies using their SED, as described in Sect. 3.2.3. Because of the low number of spiral hosts, we use all derived SEDs even if  $\chi_r^2 > 5$ . The results confirm the low-level of star forming activity (average SFR  $\approx 0.8 M_\odot/\text{yr}$ ) in the spiral galaxies of our sample.

This is likely a consequence of environmental processes such as ram pressure stripping. Quenching mechanisms probably suppress star formation more rapidly than the time needed for morphological transformation through bulge growth and disc fading (e.g. Kelkar et al. 2019; Martínez et al. 2022; Oxland et al. 2024). These results suggest that the conditions within clusters, including interactions with other galaxies and the intracluster medium, play a significant role in quenching star formation and rendering these spiral hosts inactive.

We conclude that all spiral X-ray AGN hosts within clusters are passive in terms of star formation, as defined by colour or SED analysis. This is also true for the majority of non-AGN spiral cluster galaxies, but opposite from what we observe in the field. Although the number of spiral hosts is small, we argue that this is likely a result of ram pressure stripping.



**Fig. B.1.** Fraction of passive vs. active spiral host galaxies in the main sample of x-ray AGN in clusters comparing with both the control samples, with cluster galaxies and with x-ray AGN from the field. The classification has been made based on visual inspection from the HSC images. The sample is statistically small but there is a clear trend, specifically in the case of the x-ray AGN in clusters we find only passive spiral AGN host galaxies and found up to  $2r_{500}$  cluster region. Error bars indicate the  $1\sigma$  confidence limits for small numbers of events (Gehrels 1986).

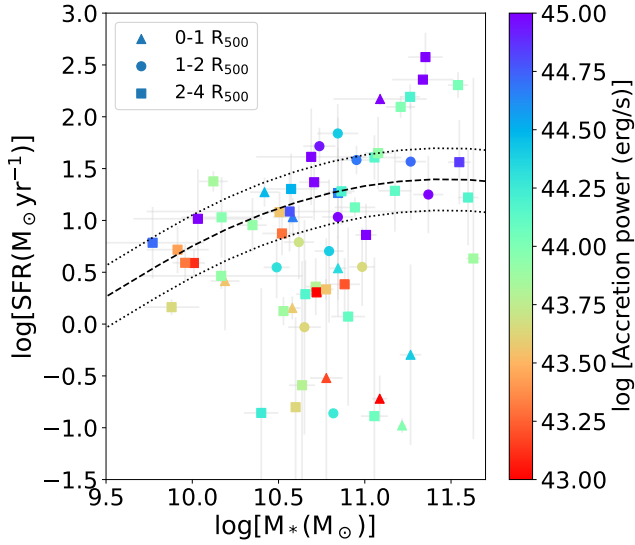


**Fig. B.2.** Passive spiral host galaxies of X-ray AGN. RGB colour images are produced by *gri* broad-band imaging from the HSC-SSP PDR3. From top to bottom: spectroscopically confirmed cluster members at redshift  $z = 0.445$  (4XLSSU J021118.8-042516),  $z = 0.302$  (4XLSSU J021234.2-053545) and  $z = 0.322$  (4XLSSU J022430.5-050842). The image scale is approximately  $30'' \times 30''$ .

## Appendix C: SFR-stellar mass relation of AGN in clusters

In this section we explore the SFR-stellar mass relation of our sources relative to the main sequence (MS) of the star-forming galaxies. Star-forming galaxies show a tight correlation between their stellar mass ( $M_*$ ) and SFRs, known as the MS of star-forming galaxies (e.g., Whitaker et al. 2012; Speagle et al. 2014; Schreiber et al. 2015). This relation is valid through a wide range of redshifts (Schreiber et al. 2016). Hence, we used the SED derived properties of our samples to examine the location of the X-ray AGN with respect to the MS.

In Fig. C.1, we plot our sources in the SFR- $M_*$  plane using the estimated SFR and  $M_*$  values for our three samples, as indicated in the legend. The sources are colour-coded with the accretion power (AP) of the AGN ("agn.accretion.power" derived



**Fig. C.1.** SFR as a function of  $M_*$ , colour-coded based on the AGN accretion power. The different shapes represent the AGN samples used in our analysis in different distances from the cluster centres as indicated in the legend. The dashed line represents the main sequence of star-forming galaxies obtained by [Schreiber et al. \(2015\)](#) with median redshift value  $z = 0.5$ . The dotted lines correspond to the uncertainties defined as  $\pm 0.3$  dex.

Furthermore, our results revealed that 50% of the sources in the centre of clusters lie inside or above the MS of the star-forming galaxies, while this percentage increases to about 65% in the case of sources in the field.

from CIGALE). It provides the intrinsic (unextinct) luminosity of the AGN disc averaged in all directions (e.g. [Yang et al. 2023](#)). We compared the source positions relative to the MS, using for the latter the analytical expression of equation 9 of [Schreiber et al. \(2015\)](#) (dashed line). For this calculation, we used the median redshift of our sample ( $z_{\text{med}} = 0.5$ ). The sources that have SFRs within 0.3 dex from the [Schreiber et al. \(2015\)](#) SFR (dotted lines) are considered to lie within the MS. Using the best SFR and  $M_*$  values, we find that the majority of the sources inside or above the MS (that is about 56% of the full sample) have in general higher AGN accretion power ( $2.8 \times 10^{44}$  erg/s) compared to those that lie below the MS ( $6.8 \times 10^{43}$  erg/s). To evaluate whether the sources below and those above or inside the MS come from the same parent distribution, we performed a two-side Kolmogorov-Smirnov (KS) test. The p-value is 0.00053, indicating that the two distributions are different. These results agree with previous studies (e.g. [Aird et al. 2019](#); [Pouliasis et al. 2022b](#)) that may indicate that higher accreting SMBHs may reside in host galaxies with enhanced SFR.

In clusters this might be linked to the high frequency of merging and disturbed galaxies. Indeed, we found evidence that host galaxies that are visually identified as disturbed exhibit relatively high values of SFR. This elevated SFR appears to be associated with substantial AP. In more detail, the same hosts that show signs of disturbance, possibly due to interactions or mergers, are also the ones where both vigorous star formation and intense accretion activity are present. This is in agreement with previous results (e.g. [Storchi-Bergmann & Schnorr-Müller 2019](#); [Silva et al. 2021](#); [Cezar et al. 2024](#)) and suggests a connection between the dynamic state of these galaxies and their energetic processes, both in forming new stars and in feeding their central black holes.