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4 × 4 Integrated Switches Based on On-Chip Wireless Connection through Optical Phased Arrays

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Abstract: Optical Wireless Networks on-Chip are an emerging technology recently proposed to improve the interconnection between different processing units in densely integrated computing architectures. In this work, we propose a 4 × 4 optical wireless switch (OWS) based on optical phased arrays (OPAs) for broadband reconfigurable on-chip communication. The OPA and OWS design criteria are reported. Moreover, the performances of the OWS are analyzed and optimized considering the electromagnetic propagation in on-chip multilayer structures, with different thicknesses of the cladding layer. The effect on the OWS behavior of a non-ideal distribution of the power in input to the OPA is also investigated by designing a 1 × 7 beam splitter, based on a single-stage multi-mode interference (MMI) device to be used as a single element of the OWS. Then, the MMI output signals are considered in input to the transmitting OPAs and the OWS performances are evaluated.

Keywords: wireless optical interconnects; optical antennas; optical phased arrays; on-chip interconnection



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1. Introduction

The efficient interconnection between different cores or different chiplets is a key enabler for increasing the performances of densely integrated computing architectures [1,2]. Thanks to the parallel execution, the distribution of data-intensive computational tasks to different processing units can increase the overall system performances, provided that the communication bottleneck is prevented by guaranteeing an efficient communication between the processing units. Networks on Chip (NoCs) are the state-of-the-art technology for interconnection through electrical links and switches. Although NoCs allow for efficient communication, the increase in their performances can be limited by latency, multi-hop communication, and scalability.

The need for shifting the on-chip communication paradigm has motivated intensive research toward alternative technologies, ranging from the wireless interconnection at millimeter wave frequencies to the integration of an integrated optical communication layer [3–6].

Wireless Networks on-chip (WiNoCs) exploit, at the chip scale, the advantages of wireless communication, which offers the possibility of reducing the network topological constraints, of avoiding inter-router hops, and of reducing the latency [7–9]. Although the use of very high frequencies (e.g., in the THz range) allows one to reduce the size of the radiators [10,11], the on-chip integrability of miniaturized antennas still requires intensive technological research.

The currently available and well-assessed nanophotonic technologies have enabled the development of Optical Networks on-Chip (ONoCs). In ONoCs, an on-chip integrated

optical layer is exploited for high-bitrate and low-latency communications. Moreover, the implementation of wavelength division multiplexing (WDM) schemes in ONoCs is particularly promising for increasing the overall data rate. In fact, WDM allows for the parallel transmission of multiple channels, each associated with a different wavelength, on the same waveguide. In this way, the overall data rate depends on both the individual optical channel rate and the number of channels, which can be contemporarily routed through the photonic NoC. Among the other components necessary for the WDM transmission, such as sources, detectors, multiplexers/demultiplexers, switches, multiwavelength filters, and modulators [12–19], one of the key components for the ONoC is the switching element that represents the basic building block to be arranged in switching matrixes.

The switching element can typically be based on resonant or non-resonant devices, with narrow- or broad-band behavior [20–29]. Nonetheless, in ONoCs, a broadband operation is advisable to have the contemporary routing of the WDM signals between the multiple processing units [24,30]. In fact, the use of multiple narrow-band devices, resonating at different WDM channel wavelengths, can require a considerable power overhead to finely tune the resonances, especially when the number of channels is increased to reach highly aggregated data rates.

A further and alternative on-chip communication technology that has been recently proposed exploits optical wireless communication at the chip scale [31–36]. The Optical Wireless Networks on-Chip (OWiNoCs) are based on the standard and mature photonic integrated circuit technologies, and they can conjugate the advantages of both optical and wireless communication. The combination of optical wired and wireless transmission can reduce the network complexity and the number of switching elements necessary for linking the transmitters and the receivers, also alleviating the network topological constraints.

In this paper, we propose a broadband 4×4 switch, based on on-chip wireless connection through Optical Phased Arrays (OPA), for WDM signal routing in OWiNoCs. The concept of the Optical Wireless Switch (OWS) has been proposed by the authors in [37,38], where 1×5 and 3×3 switches were reported. The beam radiated by the input OPA can be steered toward the addressed receivers by applying a phase shift α to the signals feeding the seven antennas in the OPA. This behavior ensures the reconfigurability of the OWS. The main contributions of this paper can be summarized as follows:

- (1) The proposed 4×4 OWS configuration, based on the use of OPAs with seven antennas, increases the number of transmitters and receivers that can be connected by the same wireless switch. This adds a further building block for on-chip wireless interconnection networks, thus opening new possibilities for the network design space exploration.
- (2) The design criteria for the OPA are identified and the OWS performances are analyzed and optimized by the Finite Difference Time Domain (FDTD) numerical simulations.
- (3) The effect of multipath propagation in the multi-layer on-chip structure is evaluated, showing that the device performances can be optimized, in terms of insertion loss and crosstalk, by varying the cladding layer thickness.
- (4) The effect on the OWS behavior of a non-ideal distribution of the power in input to the OPA is investigated. For this purpose, we first report the results of the design of a 1×7 beam splitter, based on a Multi-Mode Interference (MMI) device. Then, the MMI output signals are considered in input to the transmitting OPA to evaluate the effect of the non-uniform distribution of the OPA input on the performances of the OWS. This analysis shows that the OWS performances are not significantly affected by a non-ideal distribution of power in input to the OPA.

In Section 2, the OPA radiation characteristics are investigated and the performances of the OWS are analyzed by the Finite Difference Time Domain (FDTD) numerical simulations.

In Section 3, the effect on the OWS behavior of a non-ideal distribution of the power in input to the OPA is investigated. Finally, the discussion of the results and the conclusions are reported in Sections 4 and 5.

2. 4×4 Optical Wireless Switch

The proposed reconfigurable optical wireless switch is schematized in Figure 1. It exploits optical phased arrays made of seven antennas to connect four inputs (i.e., I_i with $i = 1, 2, 3, 4$) with four outputs (i.e., O_i with $i = 1, 2, 3, 4$). A wavelength division multiplexing signal, with M channels associated with M different wavelengths, is launched in input to one of the OPAs (e.g., I_1), and it is distributed to the seven antennas through a 1×7 beam splitter. Each antenna receives, in input, all the WDM channels, radiating them in the surrounding space.

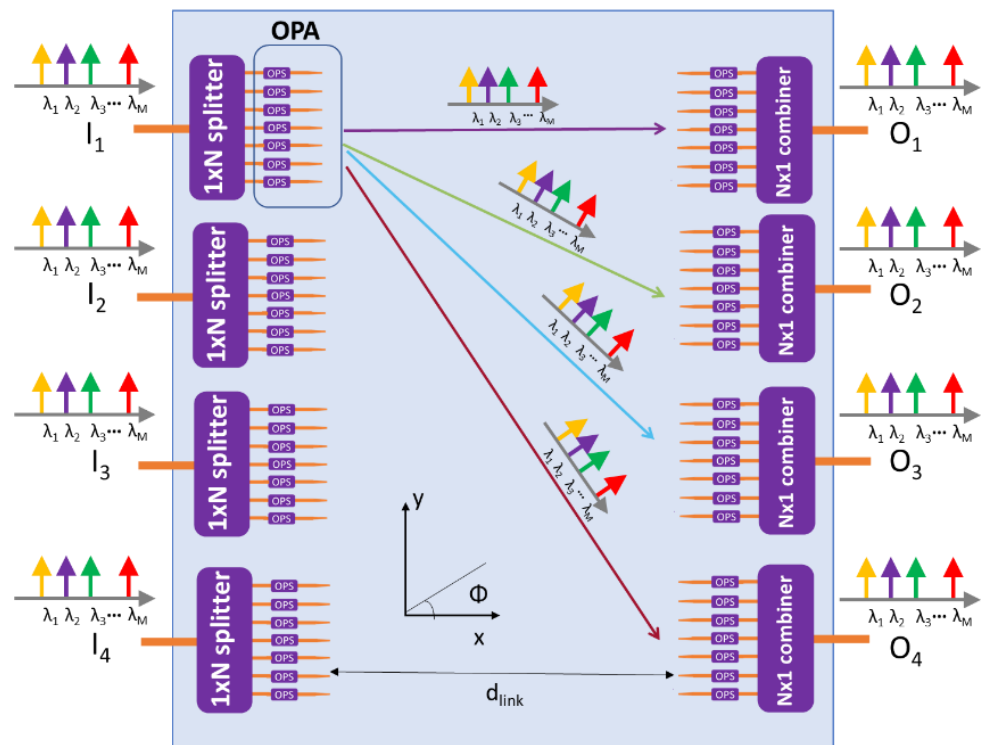


Figure 1. Scheme of the 4×4 Optical Wireless Switch (OWS), which connects four-input with four-output Optical Phased Arrays (OPAs). The signal in each of the seven antennas in the OPAs is suitably phase-shifted by Optical Phase Shifters (OPSs) to guarantee beam steering.

The beam radiated by the OPA, given by the superposition of the electromagnetic field radiated by each antenna, can be steered toward the addressed receiver, as schematized by the colored arrows in Figure 1, thanks to the phase shift α induced by the Optical Phase Shifters (OPSs). To guarantee efficient communication, it is also necessary to virtually steer the radiation diagram of the receiving OPA toward the transmitter. This can be accomplished by applying a corresponding phase shift $-\alpha$ also at the addressed receiver. The received signal is then recombined in the output waveguide by an $N \times 1$ combiner.

The OWS operation principle is similar to that of the 3×3 configurations proposed by the authors in [37], but here, seven antennas are used in the OPA to increase the number of inputs and outputs. In fact, by suitably choosing the phase shift α , a 4×4 connectivity can be achieved with the seven-antenna OPAs.

In the following subsections, the radiation characteristics of the seven-antenna OPA are first analyzed, and then the full device performances are evaluated.

2.1. Optical Phased Array Radiation Diagram

The OPAs at the transmitting and receiving nodes are made by aligning seven taper antennas along the y axis. Each antenna is obtained by inversely tapering a standard SOI waveguide (cross-section height $h = 220$ nm and width $w = 450$ nm), terminated on a small tip (length $l = 1$ μm , and width $w_T = 130$ nm) [34]. The distance between the antennas in

the OPA is chosen to be equal to $d = \lambda_m$, where λ_m is the wavelength in the surrounding medium, where the radiation occurs. According to antenna theory, the distance $d = \lambda_m$ guarantees that the OPA exhibits only a main radiation lobe, which can be steered by suitably phase-shifting the antenna input signals [39]. In particular, considering a uniform linear array configuration, the phase shift α between two adjacent antennas is constant.

The characteristics of the taper antenna have already been investigated by the authors in [37,38], through three-dimensional FDTD simulations with standard near-to-far field transformation [40]. This antenna radiates mainly in the direction of the mode propagation (i.e., the x axis in the reference system in Figure 1). Moreover, its gain increases with the taper length L_T , with a corresponding decrease in the half-power beam width, i.e., the angular aperture of the radiated beam. Further details on the radiation characteristics of the single-taper antenna can be found in [37,38].

The overall radiation diagram of a uniform linear array made of N_a identical antennas can be obtained through the multiplication of the electromagnetic field radiated by the single antenna by the array factor (AF), expressed as [38,39]:

$$AF = \sum_{q=1}^{N_a} a_q e^{-j(q-1)[kd \sin(\phi) + \alpha]} \tag{1}$$

where a_q is the excitation amplitude of each element, and $k = 2\pi n_m / \lambda$ is the propagation constant in a homogeneous medium with the refractive index n_m at the wavelength λ . The angle ϕ shown in the scheme of Figure 1 is defined considering the spherical coordinate reference system, where r is the modulus of the position vector \mathbf{r} that identifies the calculation point, θ is the inclination angle (between \mathbf{r} and the z axis), and ϕ is the azimuthal angle (angle of rotation from the x axis). The behavior of a one-dimensional array is well described considering the plane identified by the direction x of maximum radiation for the single antenna and the axis y of the alignment of the array. Therefore, in the following, the radiation characteristics of the OPA are reported considering the xy plane, and they are represented as a function of the angle ϕ .

According to Equation (1), Figure 2 shows the gain of the OPA as a function of the angle ϕ calculated for an array of $N_a = 7$ taper antennas and for different values of the phase shift α . The distance between two adjacent antennas in the OPA is equal to $d = \lambda_m$, whereas the taper length is $L_T = 2 \mu\text{m}$. The gain $G(\theta, \phi)$ is calculated as [39]:

$$G(\theta, \phi) = 4\pi \frac{I(\theta, \phi)}{P_{in}}$$

where $I(\theta, \phi)$ is the radiation intensity in spherical coordinates and in the far-field region, and P_{in} is the total power in the input to the array, given by the sum of the power values in the input to each antenna waveguide.

For each value of the phase shift α , the radiation diagram of the OPA exhibits a main lobe and several side lobes and radiation nulls in the xy plane, i.e., the plane identified by the direction x of maximum radiation for the single antenna and the axis y of alignment of the array. The phase shift values $\alpha = \pm p 360^\circ / N_a$, with $p = 1, 2, 3$, are chosen to steer the main beam on the same positions of the nulls of the broadside ($\alpha = 0^\circ$) array, thus minimizing the crosstalk. In Figure 2, each of the seven main lobes, highlighted by a different number, corresponds to an addressable receiver. As analyzed in [38], the magnitude of the main radiation lobes, for the different values of the phase shift α , follows the envelope of the single-antenna radiation diagram. Therefore, according to the single-antenna radiation characteristics, the taper length $L_T = 2 \mu\text{m}$ (i.e., short-taper condition) has been chosen to minimize the variation in the maximum gain of the steered beams and, therefore, the difference in the power received by the different output OPAs.

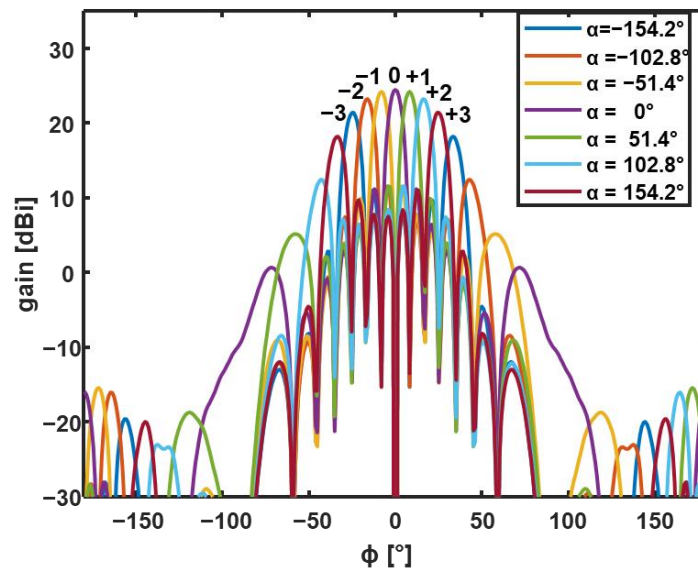


Figure 2. Gain as a function of the angle ϕ calculated for an array of $N_a =$ seven taper antennas and for different values of the phase shift α . The distance between two adjacent antennas in the OPA is equal to $d = \lambda_m$, whereas the taper length is $L_T = 2 \mu\text{m}$.

2.2. 4×4 OWS Operation Principle and Performances

Considering the radiation diagrams shown in Figure 2, the different connections between the transmitting and the receiving nodes of the 4×4 OWS can be explained.

When $\alpha = 0^\circ$, the main radiation occurs at $\phi = 0^\circ$, i.e., in front of the OPA. Therefore, considering the scheme in Figure 1, if I_1 is transmitting, the receiver O_1 can be illuminated when $\alpha = 0^\circ$. Similarly, when $\alpha = +p360^\circ/N_a$, with $p = 1, 2, 3$, the receivers O_i , with $i = 2, 3, 4$, can be addressed.

The transmitter I_2 can be connected to O_1 when $\alpha = -360^\circ/N_a = 51.4^\circ$, whereas it can be connected to O_2 when $\alpha = 0^\circ$ and to the other receivers O_3 and O_4 when α is, respectively, $\alpha = +360^\circ/N_a = 51.4^\circ$ and $\alpha = +2 \cdot 360^\circ/N_a = 102.8^\circ$.

Extending the same concepts to all the transmitters and receivers, the connectivity table (see Table 1) for the 4×4 OWS can be derived. Table 1 reports the phase shifts α required for connecting the four I_i transmitters with the four receivers O_i .

Table 1. Phase shifts α required for connecting the four I_i transmitters with the four receivers O_i .

Input \ Output	O_1	O_2	O_3	O_4
I_1	$\alpha = 0^\circ$	$\alpha = 51.4^\circ$	$\alpha = 102.8^\circ$	$\alpha = 154.2^\circ$
I_2	$\alpha = -51.4^\circ$	$\alpha = 0^\circ$	$\alpha = 51.4^\circ$	$\alpha = 102.8^\circ$
I_3	$\alpha = -102.8^\circ$	$\alpha = -51.4^\circ$	$\alpha = 0^\circ$	$\alpha = 51.4^\circ$
I_4	$\alpha = -154.2^\circ$	$\alpha = -102.8^\circ$	$\alpha = -51.4^\circ$	$\alpha = 0^\circ$

A numerical model of the 4×4 OWS was implemented by the FDTD method [40] considering the electromagnetic propagation in a multilayer medium, which is representative of the on-chip structure of photonic integrated circuits. In particular, we first considered a multilayer structure corresponding to the sample fabricated and characterized by the authors in [41] for the evaluation of point-to-point wireless links. It consists of a standard Silicon-on-Insulator (SOI) sample, with a bottom layer of bulk silicon Si (with a refractive index $n_{Si} = 3.457$), an overlying layer of silicon dioxide SiO_2 (with a height equal to $3 \mu\text{m}$ and a refractive index $n_{\text{SiO}_2} = 1.445$), and another layer of Si (220 nm -thick), where the standard waveguides and the antennas are patterned. A further thin layer of borophosphorous tetraethyl orthosilicate (BPTEOS) (of a height equal to 300 nm and a refractive index $n_{\text{BPTEOS}} = 1.453$) was considered on top of the antennas to improve the homogeneity of the

refractive index around the antennas. The BPTEOS is then covered with a UV26 polymer cladding layer (with a thickness $t = 3.78 \mu\text{m}$ and a refractive index $n_{\text{UV26}} = 1.526$). The last layer of the multilayer stack, considered in the simulations, is air. Both the bottom bulk Si layer and the top air layer are considered as semi-infinite by using Perfectly Matched Layer (PML) boundary conditions.

The behavior of the 4×4 OWS can be fully described by considering the connection between the transmitter I_1 and the four receivers O_i , with $i = 1, 2, 3, 4$, thanks to the symmetry of the radiation diagrams. Figure 3 shows the transmittances T_{O_i} in dB, calculated as a function of the wavelength, at the receiving OPAs, i.e., O_1, O_2, O_3 , and O_4 , when the transmitting OPA I_1 is excited with phase shifts: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$. The link distance $d_{\text{link}} = 70 \mu\text{m}$ was arbitrarily chosen. Figure 3 agree with the connectivity table (Table 1), as the addressed receivers exhibit high transmittance values. Nonetheless, a small part of the signal is also received by the non-addressed outputs, thus originating crosstalk.

To better quantify the OWS performances, we define the insertion loss as:

$$IL_j = -10 \cdot \text{Log}_{10}(T_{O_j}) \quad (2)$$

where T_{O_j} is the transmittance at the addressed output O_j , and we define the crosstalk as:

$$XT_{i,j} = 10 \text{Log}_{10} \left(\frac{T_{O_i}}{T_{O_j}} \right) \quad (3)$$

where T_{O_i} is the transmittance of a non-addressed port. The arrows in Figure 3 highlight, for each phase-shift, the curves from which the maximum $XT_{i,j}$ is calculated according to Equation (3).

Considering Figure 3, the worst-case insertion loss and crosstalk occur when the link between the transmitter and the further receiver O_4 is considered (Figure 3d). It is worth pointing out that the simulations take into account the propagation in the multilayer structure, which is affected by multiple reflections at the media interfaces. This phenomenon causes constructive or destructive interference that depends on the multilayer characteristics, such as layer thicknesses, refractive indices, and the OPA radiation diagram [41,42].

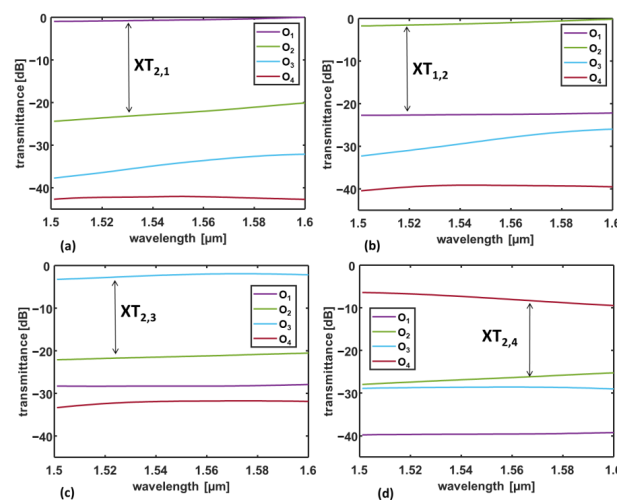


Figure 3. Transmittance in dB, calculated as a function of the wavelength, at the receiving OPAs, i.e., O_1, O_2, O_3 , and O_4 , when the transmitting OPA I_1 is excited with phase shifts: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$. The simulated device exploits reconfigurable OPAs made of $N =$ seven taper antennas with a taper length $L_T = 2 \mu\text{m}$. The link distance is $d_{\text{link}} = 70 \mu\text{m}$, and the UV26 cladding thickness is $t = 3.78 \mu\text{m}$.

Therefore, the thickness of the cladding layer is one of the parameters that can be engineered to improve the OWS performances, especially in terms of insertion loss.

In Figures 4 and 5, the results of the parametric analysis of the insertion loss and crosstalk, as a function of the UV26 layer thickness, are reported. In particular, Figure 4 shows the insertion loss IL_j in dB, calculated as a function of the wavelength and of the cladding thickness t , at the receiving OPAs, i.e., O_1 , O_2 , O_3 , and O_4 , when the transmitting OPA I_1 is excited: (a) IL_1 , (b) IL_2 , (c) IL_3 , and (d) IL_4 .

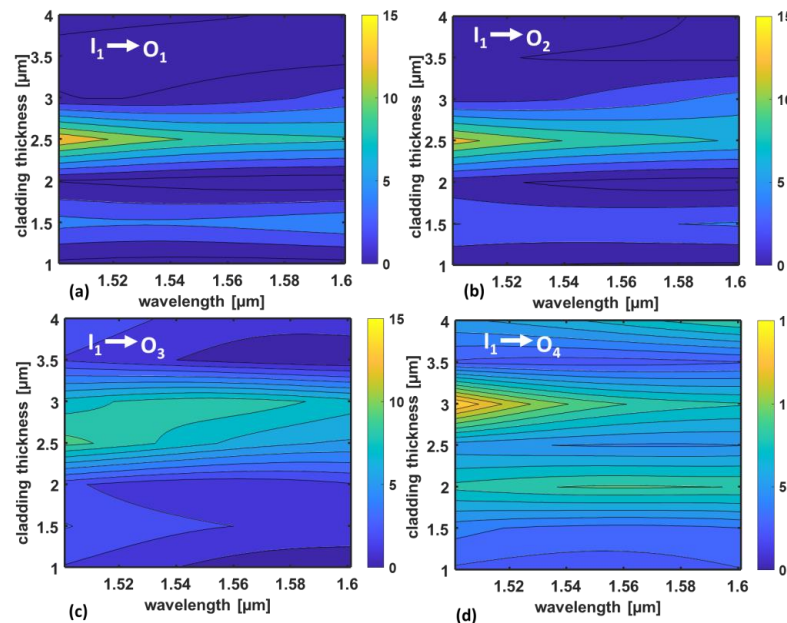


Figure 4. Insertion loss in dB, calculated as a function of the wavelength and of the cladding thickness, at the receiving OPAs, i.e., O_1 , O_2 , O_3 , and O_4 , when the transmitting OPA I_1 is excited: (a) IL_1 , (b) IL_2 , (c) IL_3 , and (d) IL_4 . The simulated device exploits reconfigurable OPAs made of $N =$ seven taper antennas with a taper length $L_T = 2 \mu\text{m}$. The link distance is $d_{\text{link}} = 70 \mu\text{m}$.

In Figure 4, for fixed values of the wavelength, the variation in the insertion loss with the cladding thickness is mainly due to the different multipath contribution that causes constructive or destructive interference. The worst-case insertion loss occurs around the thickness $t = 2.5 \mu\text{m}$ when the receivers O_1 , O_2 , and O_3 are considered, whereas it occurs around $t = 3 \mu\text{m}$ for the receiver O_4 , which, in general, exhibits higher insertion loss values. Considering Figure 4, for many of the values of the thickness t , the variation in IL_j with the wavelength is limited to less than 3 dB, thus confirming that a wideband operation of the OWS can be achieved. This is particularly relevant when WDM communication schemes are implemented. Therefore, the layer thickness can be optimized to reduce the insertion loss while ensuring broadband communication. In particular, the configuration that minimizes the insertion loss at the further receiver O_4 (i.e., insertion $IL_4 < 2.8 \text{ dB}$ in the whole wavelength range) corresponds to $t = 1 \mu\text{m}$.

Moreover, Figure 5 reports the maximum crosstalk in dB, calculated as a function of the wavelength and of the cladding thickness, at the receiving OPAs, i.e., O_1 , O_2 , O_3 , and O_4 , when the transmitting OPA I_1 is excited: (a) XT_{21} , (b) XT_{12} , (c) XT_{23} , and (d) XT_{24} . To ease the readability of Figure 5, the color scales have been adapted to the range of variation of the XT values. Considering the XT values shown in Figure 5, the crosstalk remains, in general, below -15 dB . For the thickness $t = 1 \mu\text{m}$, the worst-case crosstalk XT_{23} is below -18 dB in the whole wavelength range.

The results reported above consider, in the input to OPA, an ideal 1×7 beam splitter that equally divides the input signal into the OPA antennas. In the following, we aim to investigate the effect on the OWS performances of a non-ideal behavior of the 1×7 splitter.

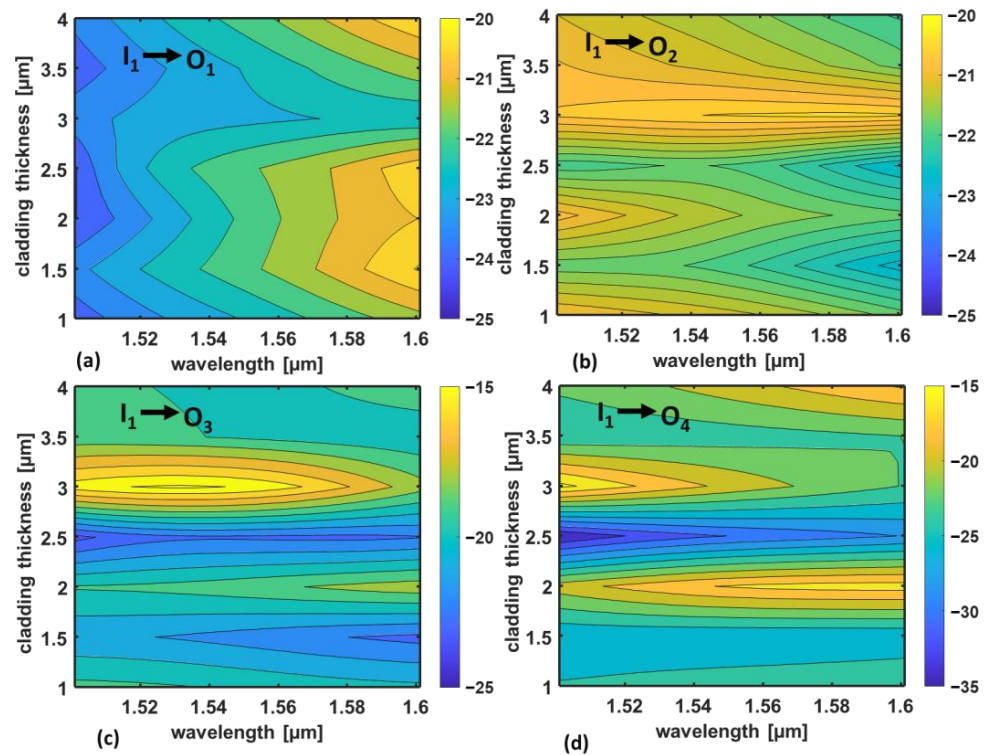


Figure 5. Crosstalk in dB, calculated as a function of the wavelength and of the cladding thickness, at the receiving OPAs, i.e., O_1 , O_2 , O_3 , and O_4 , when the transmitting OPA I_1 is excited: (a) XT_{21} , (b) XT_{12} , (c) XT_{23} , and (d) XT_{24} . The simulated device exploits reconfigurable OPAs made of $N =$ seven taper antennas with a taper length $L_T = 2 \mu\text{m}$. The link distance is $d_{\text{link}} = 70 \mu\text{m}$.

The phase shift necessary for the optical wireless switch operation can be achieved in silicon waveguides by using optical phase shifters based either on thermo-optic or plasma-optic effects [43]. Thermally controlled waveguide phase shifters could be preferred because they are based on relatively simple and robust structures, and their fabrication is less prone to errors. Conversely, plasma-optic phase shifters are expected to outperform the thermal effect in terms of switching speed.

For example, reference [44] reports the design and the experimental results of thermo-optic phase-shifters based on TiN metal and N^{++} -doped silicon. The phase shifters proposed in [44] show negligible losses and have very stable power consumption. Considering the results reported in [44], the power P_π necessary to achieve a π phase shift with thermo-optic phase-shifters based on TiN metal is equal to $P_\pi = 21.4 \text{ mW}$.

If P_q is the power required for phase shifting each antenna in the array (identified by the index q), the total power for steering the optical phased array with $N_a =$ seven antennas is $P_{\text{OPA}} = \sum_{q=1}^{N_a} P_q$.

Indeed, the power P_q depends on the chosen phase shifter technology and on the array requirements. The power required for achieving a generic phase shift γ can be roughly estimated as $P_i = P_\pi \gamma / \pi$, given the almost linear dependence of the phase shift on the thermo-optic heater power.

Considering Equation (1), the phase shift applied to the q th antenna is equal to $\gamma_q = (q - 1)\alpha$. If γ_q is greater than 360° , the angle γ_q can be converted into the corresponding angle in the $[0^\circ, 360^\circ]$ range by the following equation:

$$\gamma_q^{\text{eq}} = \gamma_q - \left\lfloor \frac{\gamma_q}{2\pi} \right\rfloor 2\pi \tag{4}$$

where q is the index denoting the different antennas in the array and $\lfloor \gamma_q/2\pi \rfloor$ is the greatest integer less than or equal to $\gamma_q/2\pi$ (i.e., floor operation).

Considering a thermal phase shifter, as in [44], the power necessary to phase shift each antenna in the array can be calculated as:

$$P_q = \frac{P_\pi}{\pi} \gamma_q^{\text{eq}} \quad (5)$$

In the analyzed optical wireless switch, the phase shift values $\alpha = \pm p360^\circ/N_a$, with $p = 1, 2, 3$, are chosen to steer the main beam on the same positions of the nulls of the broadside ($\alpha = 0^\circ$) array, thus minimizing the crosstalk. The power necessary to steer the optical phased array of the transmitter I_1 toward the receiver O_4 is $\sum_{q=1}^{N_a} P_q = 128$ mW. To guarantee efficient communication between I_1 and O_4 , it is also necessary to virtually steer the radiation diagram of the receiving array toward the transmitter. This can be accomplished by applying a corresponding phase shift $-\alpha$ also at the addressed receiver. Therefore, the connection of one transmitter and one receiver would require approximately $P_{\text{tot}} = 2 \cdot P_{\text{OPA}} = 256$ mW. Similar values of the total power are obtained for the connection of the other transmitters and receivers.

The power consumption of the switch should be evaluated considering the aggregated data rate that can be achieved. In fact, the common approach in the network on-chip analysis is to consider the energy-per-bit as a figure of merit, rather than the power. The proposed optical wireless switch is a broadband device, and the power required for signal routing is independent of the number of WDM channels allocated in the bandwidth. For example, considering a WDM signal with a channel spacing $\Delta\lambda = 0.8$ nm, virtually 120 channels can be allocated in the simulated 100 nm bandwidth of the optical wireless switch. If the individual optical channel data rate is 50 Gb/s, the estimated energy per bit is equal to about 42 fJ/b. Optimized phase shifter configurations (e.g., based on plasma-optic effect) or higher data rates could result in lower energy-per-bit consumption.

3. Effect of Non-Uniform Power Distribution at the Transmitting OPA Antennas

An ideal 1×7 beam splitter would equally distribute the input power to the seven antennas. However, in actual devices, the power distribution at the $1 \times N$ output of the beam splitter can be non-uniform. A possible implementation of the $1 \times N$ beam splitter that brings the signal to the OPSs, and then to the OPA antennas, can be made by cascading multiple 1×2 Y junctions to increase the number of outputs [45,46]. A 1×2 Y junction keeps the two outputs in phase while equally dividing the input power into the two waveguides. Another possible implementation can exploit the cascade of multiple Multi-Mode Interference (MMI) devices [47,48]. In the case of 1×2 MMIs, the power is equally divided at the two outputs, and the signals are in phase.

Possible errors introduced by the $1 \times N$ splitter by a non-uniform power distribution to the OPA antennas can change the performances of the OWS. Additionally, errors in the phase shift between the antennas can have an effect on the switch performances, but they can be compensated by the optical phase shifters. In Si waveguides, OPSs can be implemented by exploiting either thermo-optic or plasma-optic effects [43,44].

Here, we aim at evaluating the effect on the OWS performances of a non-uniform power distribution to the input OPA antennas. In particular, we define a case study considering a beam splitter based on a 1×7 MMI. This device can be a more compact solution for the beam splitting functionality in a single stage. The MMI beam splitter is based on general interference, which can be explained through self-imaging in multi-mode waveguides [49,50]. Thanks to the self-imaging property, the input field profile is reproduced in multiple images in the multimode section of the MMI, at periodic intervals [50]. Therefore, by suitably choosing the length of the MMI, the input signal can be split into multiple output waveguides.

In the following subsections, we report the results of the design of a 1×7 MMI, focusing on the analysis of the uniformity of the output power at the seven ports. Moreover,

we evaluate the effect of non-uniform distribution of the power in the input to the antennas by simulating the whole OWS.

3.1. 1 × 7 Multi-Mode Interference Beam Splitter

The proposed 1 × 7 MMI configuration is schematized in Figure 6. The waveguides are made of silicon. Moreover, the same multilayer of the OWS is considered in the simulations, with the UV26 cladding thickness $t = 1 \mu\text{m}$. The MMI device is, indeed, patterned in the antenna layer, and the input and output waveguides are standard SOI waveguides (cross-section height $h = 220 \text{ nm}$ and width $w = 450 \text{ nm}$). Input and output tapers are considered to improve the output power uniformity, as discussed in [49], with the taper width $w_T = 1 \mu\text{m}$. The distance between the output waveguides was arbitrarily chosen to be equal to $d_A = 1.50 \mu\text{m}$, but in the design of an overall integrated circuit, it can be customized to meet the topological constraints, e.g., for fabrication. The width of the MMI region is chosen to be equal to $w_{\text{MMI}} = 7 \cdot d_A$. The optimal length of the MMI can be estimated by the following equation [50]:

$$L_{\text{MMI}} = L_{\pi} \frac{3}{4N_A} = \frac{\lambda_0}{2(n_1 - n_2)} \frac{3}{4N_A} \tag{6}$$

where L_{π} is defined as the beat length of the two lowest-order modes of the multimode region, N_A is the number of outputs, and λ_0 is the design wavelength (i.e., $1.55 \mu\text{m}$), whereas n_1 and n_2 are the effective refractive indices of the two lowest-order modes of the MMI.

The length of the MMI, $L_{\text{MMI}} = 28.6 \mu\text{m}$, estimated by Equation (6), must be optimized by the simulation of the full device, taking into account the non-uniformity of the power at the seven outputs. The output power non-uniformity NU, which is a function of the wavelength, can be quantified as:

$$\text{NU}(\lambda) = \max_i(T_{A_i}(\lambda)) - \min_i(T_{A_i}(\lambda)) \tag{7}$$

where T_{A_i} with $i = 1, 2, \dots, 7$ are the transmittances in dB calculated at the seven output ports of the MMI. Indeed, the non-uniformity is zero when the output transmittances T_{A_i} are all equal, and it is greater than zero when the power distribution is non-uniform.

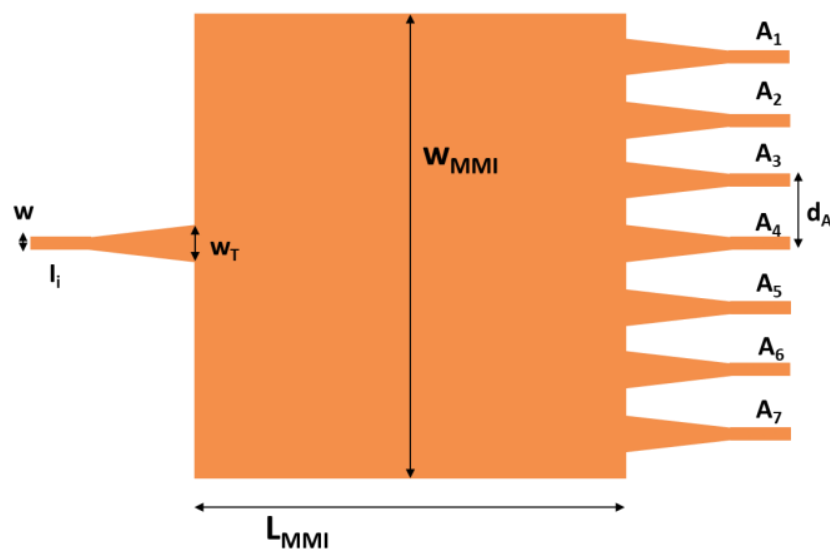


Figure 6. Scheme of the 1 × 7 MMI beam splitter. The waveguides are made of silicon, and the same multilayer of the OWS is considered in the simulations with the UV26 cladding thickness $t = 1 \mu\text{m}$.

The parametric analysis, as a function of the MMI length L_{MMI} , was performed by three-dimensional FDTD simulations. Figure 7 shows the transmittances T_{Ai} calculated at the output ports (in dB) as a function of the MMI length and at the wavelength $\lambda_0 = 1.55 \mu\text{m}$. Given the symmetry of the device with respect to the central output A_4 , only four curves are visible in Figure 7. Considering Figure 7, the MMI length that gives the lowest variation of the transmittances (i.e., calculated non-uniformity $NU(\lambda_0) = 0.44 \text{ dB}$ at the design wavelength $\lambda_0 = 1.55 \mu\text{m}$) is equal to $L_{MMI} = 26.5 \mu\text{m}$.

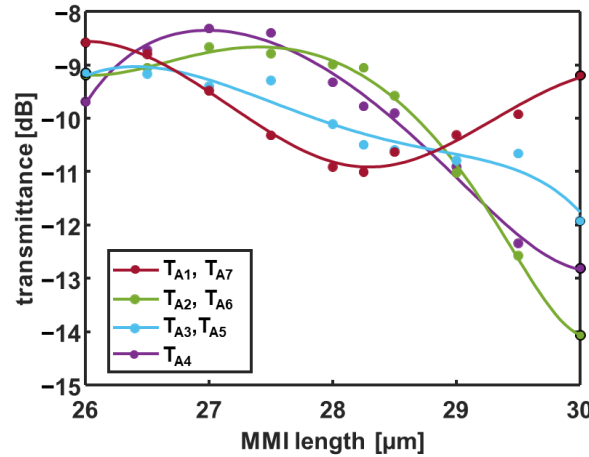


Figure 7. Transmittances T_{Ai} at the output ports in dB as a function of the MMI length (calculated values (dots) and polynomial fitting (solid curves)) at the wavelength $\lambda = 1.55 \mu\text{m}$. The same multilayer of the OWS is considered in the FDTD simulations with the UV26 cladding thickness $t = 1 \mu\text{m}$.

It is worth pointing out that the wavelength behavior of the $1 \times N$ beam splitter can affect the bandwidth of the overall OWS; therefore, it is necessary to analyze the non-uniformity also as a function of the wavelength.

Figure 8a shows the non-uniformity NU calculated in dB as a function of the MMI length and of the wavelength. Moreover, Figure 8b shows the non-uniformity $NU(\lambda_0)$ calculated at the design wavelength $\lambda_0 = 1.55 \mu\text{m}$ (dashed curve) and the maximum non-uniformity in the bandwidth NU_{max} (solid curve) calculated, for each value of L_{MMI} , as:

$$NU_{max} = \max_{\lambda} (NU(\lambda))$$

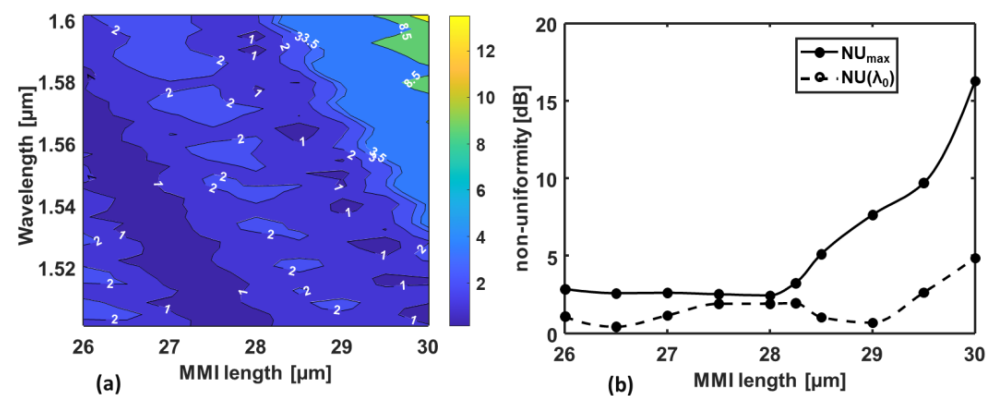


Figure 8. (a) Non-uniformity NU_{MMI} at the output ports calculated in dB as a function of the MMI length and of the wavelength; (b) non-uniformity $NU(\lambda_0)$ at the design wavelength $\lambda_0 = 1.55 \mu\text{m}$ (calculated values (dots) and polynomial fitting (dashed curve)), and maximum non-uniformity in the bandwidth NU_{max} (calculated values (dots) and polynomial fitting (solid curve)) as a function of the MMI length. The same multilayer of the OWS is considered in the FDTD simulations with the UV26 cladding thickness $t = 1 \mu\text{m}$.

As Figure 8 shows, longer values of the MMI length L_{MMI} give higher non-uniformity. Considering Figure 8b, the length value that gives the lowest non-uniformity in the whole analyzed wavelength range, i.e., $\text{NU}(\lambda) < \text{NU}_{\text{max}}$ with $\text{NU}_{\text{max}} = 2.46$ dB, is $L_{\text{MMI}} = 28$ μm . The corresponding non-uniformity at the design wavelength is $\text{NU}(\lambda_0) = 1.93$ dB.

The 1×7 MMI divider also contributes to the overall insertion loss of the OWS. For the MMI length $L_{\text{MMI}} = 28$ μm , the overall insertion loss of the 1×7 MMI, calculated by summing up the power coupled at the seven output waveguides, is about $\text{IL}_{\text{MMI}} = 1.27$ dB. Moreover, the calculated back reflection at the input port of the 1×7 MMI is below -14 dB in the whole frequency range.

Considering the scheme in Figure 1, each transmitter I_i is equipped with an OPA that radiates the whole signal spectrum (i.e., all the WDM channels) toward the same addressed receiver. At the addressed receiver, all the seven antennas of the OPA contribute to the received power. If the proper phase relation is ensured, as assumed in this paper, the signals received by each antenna in the receiving OPA constructively sum up in the output waveguide. Therefore, the $N \times 1$ combiner is intended to constructively add the signals received by the seven antennas in the receiving OPA and to deliver the resulting signal at the output waveguide.

The worst-case insertion loss from the transmitting to the receiving OPA is about 2.8 dB, which accounts for the wireless connection between the transmitting OPA and the receiving one. Considering the scheme in Figure 1, other elements can contribute to the OWS insertion loss, such as the 1×7 beam splitter and the 7×1 combiner. Additional loss contributions can be introduced by the propagation loss in silicon waveguides (approximately 0.1 dB/cm [21]) and by the optical phase shifters. For example, the thermo-optic phase shifters proposed in [48] exhibit negligible losses ($\text{IL}_{\text{OPS}} < 0.01$ dB).

As discussed in [49], further improvement of the 1×7 MMI performance can be achieved by optimizing the taper width. However, the full optimization of the MMI goes beyond the scope of this paper, which mainly intends to focus on the analysis of the effect on the OWS performance of a non-uniform power distribution at the input OPA. Therefore, the studied MMI configurations, whose output signals are assumed in the input to the transmitting OPA, give a reference to quantify the possible non-uniformity of the OPA inputs. It is also worth pointing out that the output signals of the MMI are not in the phase, but their phase difference can be compensated by the OPSs.

3.2. OWS Performance with Non-Uniform Input Power

In order to study the effect on the OWS behavior of a non-uniform power distribution at the input OPA, we decided to consider the case study of the 1×7 MMI beam splitter. As discussed above, this device can be preferred to the cascade of multiple 1×2 splitters because it is more compact. As a drawback, the output power distribution can be non-uniform, in the sense that the power is not equally divided in the seven outputs. Based on the 1×7 MMI analysis, we first consider the modification of the OPA radiation diagram, calculated through antenna theory [39]. Then, the results of the FDTD simulations of the whole OWS are discussed.

Based on the MMI analysis reported above, we define a case study considering the amplitudes at the MMI outputs when $L_{\text{MMI}} = 28$ μm . As mentioned above, this configuration exhibits a non-uniformity $\text{NU}(\lambda_0) = 1.93$ dB at the design wavelength and $\text{NU}(\lambda) < 2.46$ dB in the whole wavelength range.

Considering the calculation of the OPA radiation diagrams, the variation in the input amplitudes can be taken into account by changing the a_q coefficients in the theoretical formula of Equation (1). The array factor AF is then multiplied by the far-field of the single-taper antenna, calculated by FDTD simulation.

In the analyzed case, the seven a_q coefficients have been chosen to be equal to the amplitudes of the MMI output signals, calculated at the wavelength $\lambda_0 = 1.55$ μm . Figure 9 shows the gain as a function of the angle ϕ for an array of $N_a =$ seven taper antennas calculated for the uniform (solid curve) and the non-uniform (dashed curves) input am-

plitude distributions. The four values of the phase shift α necessary for the description of 4×4 OWS operation are considered: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$. In Figure 9 the represented ϕ range was restricted between -90° and 90° to ease the readability of the curves.

As can be seen in Figure 9, the non-uniform amplitude at the OPA antennas slightly changes the level of the side lobes, but it does not significantly change the main lobe level. The main lobe is mainly responsible for the communication between the transmitter and the addressed receiver. Therefore, since it is unaltered, no significant change in the OWS insertion loss is expected for all the phase shift values. Conversely, a change in the crosstalk is expected since it is mainly associated with the sidelobes.

In order to analyze the OWS behavior, the full device has been simulated by the FDTD method. In the FDTD model of the OWS, the transmitting OPA was fed by seven mode sources, each injecting a TE-guided mode into the OPA antennas [40]. The amplitudes of the modes were set equal to the MMI output amplitudes, calculated at the wavelength $\lambda_0 = 1.55 \mu\text{m}$, thus taking into account the non-uniform power distribution at the 1×7 MMI outputs. We have chosen to apply a straightforward simulation approach in which the amplitudes of the seven mode sources do not change with the wavelength.

Figure 10 shows the transmittances in dB, calculated as a function of the wavelength, at the receiving OPAs, i.e., O_1, O_2, O_3 , and O_4 , when the transmitting OPA I_1 is excited with uniform (solid curves) and non-uniform (dashed curves) input signals. The phase shift values considered in Figure 10 are: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$.

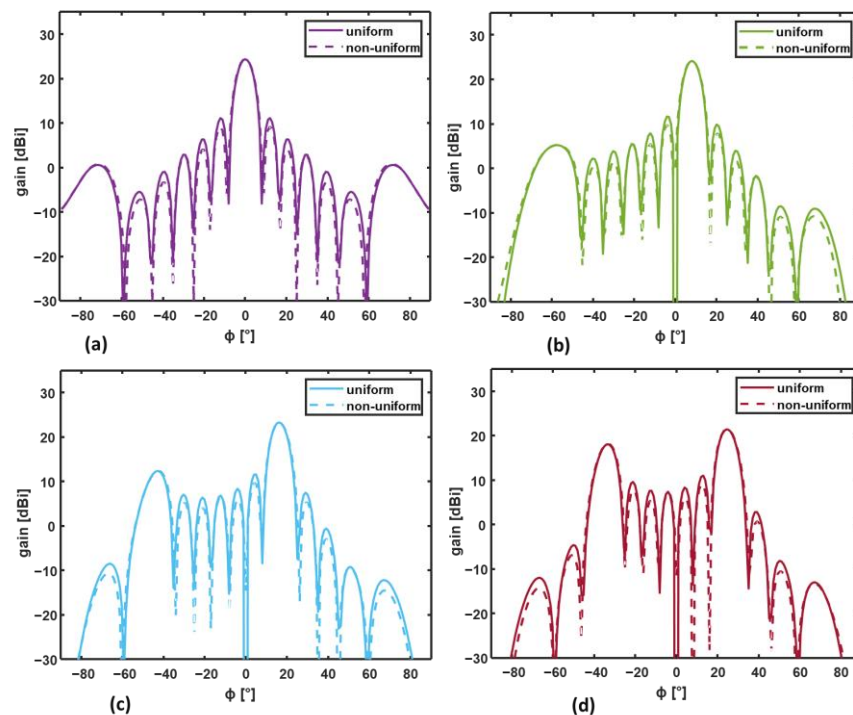


Figure 9. Gain as a function of the angle ϕ for an array of $N_a =$ seven taper antennas calculated for uniform (solid curve) and non-uniform (dashed curves) input power distributions. Different values of the phase shift α are considered: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$. The distance between two adjacent antennas in the OPA is equal to $d = \lambda_m$, whereas the taper length is $L_T = 2 \mu\text{m}$.

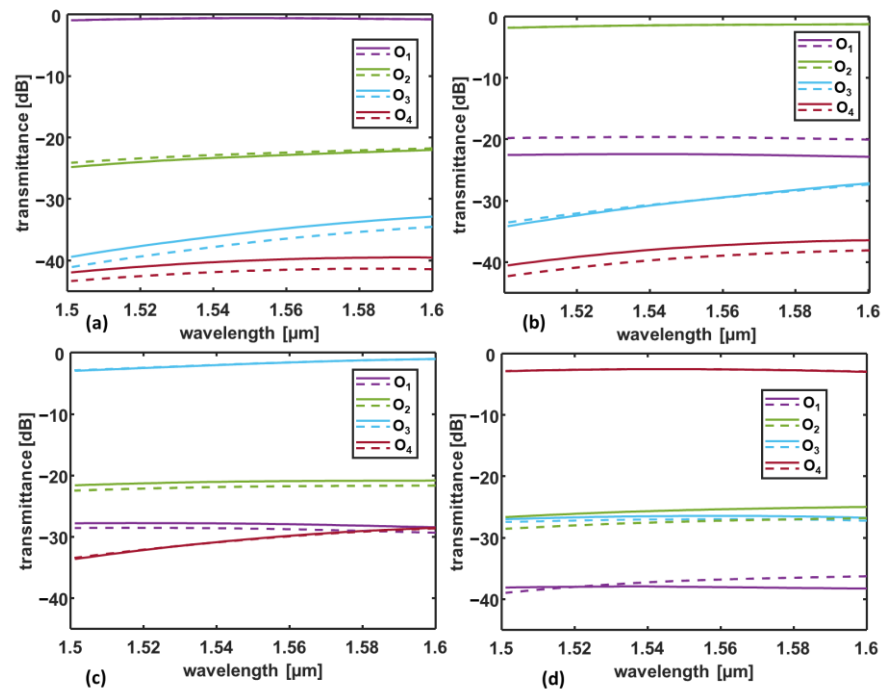


Figure 10. Transmittance in dB, calculated as a function of the wavelength, at the receiving OPAs, i.e., O_1 , O_2 , O_3 , and O_4 , when the transmitting OPA I_1 is excited with uniform (solid curves) and non-uniform (dashed curves) input signals. The phase shift values are: (a) $\alpha = 0^\circ$, (b) $\alpha = 51.4^\circ$, (c) $\alpha = 102.8^\circ$, and (d) $\alpha = 154.2^\circ$. The simulated device exploits reconfigurable OPAs made of $N =$ seven taper antennas with a taper length $L_T = 2 \mu\text{m}$. The link distance is $d_{\text{link}} = 70 \mu\text{m}$, and the UV26 cladding thickness is $t = 1 \mu\text{m}$.

Figure 10 shows that the transmittances at the addressed ports do not significantly change when a non-uniform excitation is considered, as expected by the analysis of the radiation diagrams in Figure 9. Conversely, the change in the transmittances at the non-addressed outputs is more pronounced. In particular, a 2.8 dB worsening of the crosstalk XT_{12} is found for the connection between the transmitter I_1 and the receiver O_2 (Figure 10b), which now corresponds to the worst-case crosstalk. Nonetheless, the crosstalk XT_{12} is still below -18 dB in the whole wavelength range.

4. Discussion

The proposed 4×4 OWS increases the number of transmitters and receivers that can be connected by the same wireless switch with respect to the configurations previously proposed in [37,38]. This configuration widens the available building blocks for the implementation of optical wireless networks on-chip.

The numerical analysis of the OWS has shown that the worst-case insertion loss is $IL = 2.8$ dB for the wireless interconnection. The evaluated worst-case crosstalk is -18 dB, and the bandwidth is about 100 nm. The reported results have shown that, thanks to the multipath propagation, the insertion losses can be minimized by suitably choosing the cladding layer thickness for the same switch configuration.

Other state-of-the-art broadband solutions for photonic switching consist of arranging 2×2 Mach–Zehnder interferometers into higher-order switching topologies. Owing to their operating principles, Mach–Zehnder interferometers are able to switch multiple wavelengths simultaneously [24,51]. 2×2 Mach–Zehnder switches are typically organized into larger $N \times N$ switching topologies through carefully optimized connectivity patterns such as Benes or dilated-Benes networks [30,52]. Even though a direct comparison of performances is not straightforward, the crosstalk and insertion losses of the proposed 4×4 OWS are comparable with those reported in [30,52]. As a drawback, the power

consumption is higher, but it is partly compensated by the broadband behavior, which leads to an acceptable energy-per-bit requirement of about 42 fJ/b.

As future directions of the research, the performances of the 4×4 OWS can be further improved by considering optimized phase shifter configurations, e.g., based on the plasma-optic effect, which can reduce the power requirement. Moreover, the OWS insertion losses can be further improved by optimizing the design of the beam splitters and combiners and by considering different single antennas in the OPAs.

In the proposed configuration, the OPAs exploit seven dielectric taper antennas. Nonetheless, the proposed design criteria can be applied to virtually any kind of integrated antennas. Other possible implementations of integrated dielectric antennas have been proposed in [34,53], reaching high gain values for the single-antenna radiation (e.g., greater than 20 dBi). The use of high-gain single radiators can increase the overall gain of the OPA, thus contributing to the insertion loss reduction. As a drawback, these antennas have sizes in the order of 15 μm or 25 μm , thus increasing the device footprint. Other, more compact antenna configurations can be implemented through integrated plasmonic antennas. Different configurations of plasmonic antennas have been proposed in the literature for on-chip wireless communications [54–57]. The plasmonic antennas have very small dimensions (a few microns), but they suffer non-negligible losses due to plasmonic mode propagation in the lossy metal. These losses limit the maximum achievable antenna gain, which is, in general, in the order of 10 dBi or 15 dBi. A wide range of possibilities are therefore available for studying new OWS configurations based on the identified design criteria.

The proposed OWS technology, capable of implementing broadband, low-latency communication among multiple transmitters and receivers, can find application in both long-haul optical communications and short-reach optical interconnects [58,59]. Integrated optical multi-port switches, ensuring reconfigurability and high bandwidth, can find application, for example, in photonic internet network nodes or in the interconnection of multiple chiplets in 2.5D manycore systems. Further future innovative applications can be investigated in quantum communication networks to reduce latency and increase the overall bitrate [60–63].

5. Conclusions

A 4×4 OWS based on reconfigurable OPAs has been proposed for on-chip wireless interconnection. Each OPA exploits seven taper antennas to radiate an optical beam toward the addressed receiving OPA, thanks to the beam steering capability.

The OWS performances have been analyzed and optimized considering the electromagnetic propagation in on-chip multilayer structures, with different thicknesses of the cladding UV26 layer. The optimized configuration exhibits worst-case insertion loss and crosstalk below 2.8 dB and -18 dB, respectively, in the whole wavelength range. The broadband operation of the OWS, which covers the entire C-band, can allow for the implementation of WDM communication schemes. Moreover, all the allocated WDM channels can be switched at the same time, thus making the power requirement for signal routing independent of the number of WDM channels.

As a further step of the analysis, the effect on the OWS behavior of a non-ideal distribution of the power in the input to the OPA has been investigated. For this purpose, a 1×7 beam splitter, based on a single-stage MMI device, has been designed. Then, the MMI output signals were considered in the input to the transmitting OPA. The designed 1×7 MMI is characterized by a non-uniformity of the output signals below 2.46 in the whole wavelength range. Considering this non-uniform signal distribution at the antennas of the input OPA, the performances of the OWS do not change significantly, showing a good robustness to input amplitude errors.

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Writing—original draft preparation: G.C., G.B., F.F., M.B., D.B., V.T. and V.P. Writing—review and editing: G.C., G.B., F.F., M.B., D.B., V.T. and V.P. Supervision: G.C. All authors have read and agreed to the published version of the manuscript.

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