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MIXED LOCAL AND NONLOCAL ELLIPTIC OPERATORS: REGULARITY AND MAXIMUM PRINCIPLES

STEFANO BIAGI, SERENA DIPIERRO, ENRICO VALDINOCI, AND EUGENIO VECCHI

ABSTRACT. We develop a systematic study of the superpositions of elliptic operators with different orders, mixing classical and fractional scenarios. For concreteness, we focus on the sum of the Laplacian and the fractional Laplacian, and we provide structural results, including existence, maximum principles (both for weak and classical solutions), interior Sobolev regularity and boundary regularity of Lipschitz type.

1. INTRODUCTION

The goal of this paper is to develop a systematic study of mixed operators. The word “mixed” refers here to the differential (or pseudo-differential) order of the operator, and to the type of the operator, which combines classical and fractional features.

Though many of the techniques that we present here are rather “general”, for the sake of concreteness, and not to hide the main flow of ideas by technical complications, we focus here on an operator which deals with the coexistence of a Laplacian and a fractional Laplacian, given by

$$(1.1) \quad \mathcal{L} := -\Delta + (-\Delta)^s, \quad \text{for some } s \in (0, 1).$$

Here, $(-\Delta)^s$ is the nonlocal operator defined as

$$(1.2) \quad (-\Delta)^s u(x) = c_{N,s} \text{ P.V. } \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{N+2s}} dy,$$

where $c_{N,s}$ is a suitable normalizing constant, whose explicit expression is given by

$$c_{N,s} = \left(\int_{\mathbb{R}^N} \frac{1 - \cos(\zeta_1)}{|\zeta|^{N+2s}} d\zeta \right)^{-1},$$

and, as customary, “P.V.” stands for the Cauchy’s principal value, see e.g. formulas (3.1) and (3.2) in [25].

We recall that the literature presents several variants of the fractional Laplacian, including one that is defined in terms of the eigenbasis and spectrum of the Dirichlet Laplacian on a bounded domain and one in which the singular integral only extends to the given domain. These are quite different operators from the one in (1.2), see e.g. [2, Sections 2.1, 2.2, 2.3, 4.2 and 4.3] for a list of similarities and differences between these fractional operators.

Operators as in (1.2) arise naturally from the superposition of two stochastic processes with different scales (namely, a classical random walk and a Lévy flight): roughly speaking, when a particle can follow either of these two processes according to a certain probability, the associated limit diffusion equation is described by an operator of the form described in (1.1): see in particular

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the appendix in [28] for a thorough discussion of this phenomenon and [30] for the description of a biological population in an ecological niche modeled by a mixed operator.

In view of this motivation, we think that operators as in (1.1) will enjoy a constantly rising popularity in applied sciences, also to study the different impact of “local” and “nonlocal” diffusions in concrete situations (e.g. how different types of “regional” or “global” restrictions may reduce the spreading of a pandemic disease, see e.g. [32]). Other classical applications include heat transport in magnetized plasmas (see [13]).

The mathematical study of operators with different order is not new in itself, and indeed the literature already presents results concerning, among the others, the theory of viscosity solutions (see [37, 38, 6, 11, 21, 4, 5]), the Aubry-Mather theory for sums of different fractional Laplacians (see Remark 5.6 in [41]), regularization effects of Cahn-Hilliard equations (see [18]), numerics ([12]), probability and stochastics (see [19, 20, 42]), symmetry results for mixed range phase transitions (see [16]), porous medium equations (see [24]), decay estimates for parabolic equations (see [31]), specific Liouville theorems for systems of equations driven by sums of fractional Laplacians (see [40, 3]), fractional damping effects (see [23]), and Bernstein-type regularity results (see [15]).

Though conceptually different, the problems related to mixed order equations are closely related in spirit to the ones of variable order equations (see [39, 8, 9, 47]), which have themselves concrete significance in applied sciences (see e.g. [48, 36]).

The main focus here is on the operator in (1.1) with the aim of obtaining a number of *structural results* (based on techniques which we plan to extend to more general situations in future works). These results deal with distributional as well as classical solutions, and they can be grouped into four categories: existence, maximum principles, interior regularity, and boundary regularity.

Let us now describe the main results in further detail. First of all, we will introduce a suitable notion of weak solution associated to the operator in (1.1). In light of the mixed nature of the problem, this step already presents some caveats, due to the possible choices of including or excluding the external data within the classical Sobolev spaces framework. Our setting for weak solutions will be described in Definition 2.1, and then exploited to obtain, via the Lax-Milgram’s theorem, the following existence result, of very classical flavor. In all the forthcoming statements, we tacitly understand that $\Omega \subseteq \mathbb{R}^N$ is a *bounded open set with C^1 boundary*.

Theorem 1.1. *Given $f \in L^2(\Omega)$, there exists a unique weak solution $u_f \in H^1(\mathbb{R}^N)$ of*

$$\begin{cases} \mathcal{L}u = f & \text{in } \Omega, \\ u|_{\mathbb{R}^N \setminus \Omega} = 0. \end{cases}$$

Furthermore, there exists a $\mathbf{c} > 0$ such that

$$\|u_f\|_{H^1(\mathbb{R}^N)} \leq \mathbf{c} \|f\|_{L^2(\Omega)}.$$

We then focus on the maximum principles associated to the operator in (1.1). Their formulation is slightly different for weak and classical solutions. More precisely, we present a weak maximum principle for weak solutions, as follows:

Theorem 1.2. *Let $u \in H^1(\mathbb{R}^N)$ weakly satisfy $\mathcal{L}u \geq 0$ in Ω . If $u \geq 0$ a.e. in $\mathbb{R}^N \setminus \Omega$, then $u \geq 0$ a.e. on Ω .*

For classical solutions, this statement is strengthened in the following result:

Theorem 1.3. *Let $u \in C(\mathbb{R}^N, \mathbb{R}) \cap C^2(\Omega, \mathbb{R})$, with*

$$\int_{\mathbb{R}^N} \frac{|u(x)|}{1 + |x|^{N+2s}} dx < \infty.$$

Suppose that

$$\begin{cases} \mathcal{L}u \geq 0 & \text{pointwise in } \Omega, \\ u \geq 0 & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$

Then

$$(1.3) \quad u \geq 0 \quad \text{in} \quad \Omega.$$

Furthermore,

$$(1.4) \quad \text{if there exists } x_0 \in \Omega \text{ such that } u(x_0) = 0, \text{ then } u \equiv 0 \text{ throughout } \mathbb{R}^N.$$

As customary, the statement in (1.3) can be considered as a weak maximum principle and the one in (1.4) as a strong maximum principle. The difference between the weak maximum principle in Theorem 1.2 and that in (1.3) is in the assumptions required on u (the Sobolev setting being considered in Theorem 1.2 and the classical one in Theorem 1.3).

Though these maximum principles are of classical flavor, we remark that their validity depends sensibly on the type of the operator and on the setting of the data. In particular, we will show in Appendix A that the maximum principle does not hold if the external condition “ $u \geq 0$ in $\mathbb{R}^N \setminus \Omega$ ” is replaced by “ $u \geq 0$ on $\partial\Omega$ ”: namely, classical boundary conditions are not enough to guarantee the validity of maximum principles for the operator in (1.1), notwithstanding the presence of the classical Laplacian in this operator.

Furthermore, we will show that these maximum principles do not hold when one of the elliptic operators in (1.1) has the “wrong sign”, e.g. for operators of the type $\Delta + (-\Delta)^s$.

The fact that the weak maximum principles in Theorem 1.2 and in (1.3) of Theorem 1.3 hold for \mathcal{L} but not for similar operators with the “wrong sign” is very reasonable, especially in view of some potential-theoretic results of the early '60s (see, e.g., [14, 22]). In fact, in the aforementioned papers is proved that: if

$$A : C_0^2(\mathbb{R}^N) \rightarrow C(\mathbb{R}^N)$$

is a linear operator, then the next two conditions are *equivalent*:

(i) A is continuous (with respect to suitable topologies) and

$$\begin{cases} u \in C_0^2(\mathbb{R}^N), \\ u(x) = \inf_{\mathbb{R}^N} u \leq 0 \end{cases} \implies Au(x) \leq 0;$$

(ii) for every $u \in C_0^2(\mathbb{R}^N)$, we have

$$(1.5) \quad \begin{aligned} Au = & - \sum_{i,j=1}^N a_{i,j}(x) \frac{\partial^2 u}{\partial x_i \partial x_j}(x) + \sum_{i=1}^N b_i(x) \frac{\partial u}{\partial x_i}(x) + c(x)u(x) \\ & + \int_{\mathbb{R}^N} s(x, dy) [u(x) - u(y) - \mathbf{1}_{\{|x-y| \leq 1\}} \cdot \langle \nabla u(x), x - y \rangle], \end{aligned}$$

where $A(x) = (a_{i,j}(x))$ is *positive definite* for every $x \in \mathbb{R}^N$, $c \leq 0$ on the whole of \mathbb{R}^N and s is a measurable kernel on \mathbb{R}^N satisfying the properties

- $s(x, \{x\}) = 0$;
- for every positive function $f \in C_0(\mathbb{R}^N)$, the map

$$x \mapsto \int_{\mathbb{R}^N} s(x, dy) |y - x|^2 f(y) dy \quad \text{is measurable.}$$

In particular, property (i) (which shall be used in the proof of (1.3)) holds for \mathcal{L} but not for similar operators with “wrong sign”. It is worth mentioning that in the papers [14, 22] it is not studied the validity of a weak maximum principle as in (1.4) (i.e., the possibility of “propagating” the sign of u from $\mathbb{R}^N \setminus \Omega$ into Ω). On the other hand, the equivalence of (i) and (ii) is exploited in the recent paper [3] to characterize all the operators of the form (1.5) for which a Liouville-type theorem holds.

We devote part of this paper to the interior and boundary regularity properties of solutions. Though a variety of different directions can be taken for this, we focus here on the interior regularity theory in Sobolev spaces. The main result in this sense goes as follows:

Theorem 1.4. *Let $m \in \mathbb{N} \cup \{0\}$ and $f \in H^m(\Omega)$. Assume that $u \in H^1(\mathbb{R}^N)$ is a weak solution of*

$$\mathcal{L}u = f \quad \text{in } \Omega.$$

Then $u \in H_{\text{loc}}^{m+2}(\Omega)$.

The proof of Theorem 1.4 requires some technical improvements with respect to the classical energy methods and bootstrap arguments, since the fractional operator prevents the possibility of taking derivatives of the equation. To overcome this difficulty, we will exploit truncation arguments and difference quotients in a suitable way.

As for the boundary regularity, for concreteness we focus on the linear growth and Lipschitz regularity for solutions in convex domain. Our result can be summarized as follows:

Theorem 1.5. *Assume that Ω is strictly convex and let $\bar{C} > 0$. Let $u \in H^1(\mathbb{R}^N)$ be such that*

$$(1.6) \quad \begin{cases} \mathcal{L}u \leq \bar{C} & \text{in } \Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \Omega, \\ u \leq \bar{C} & \text{in } \mathbb{R}^N. \end{cases}$$

Then, there exist $C, \ell > 0$ such that, for every $p \in \partial\Omega$, we have that

$$(1.7) \quad u(x) \leq C \bar{C} |x - p| \quad \text{for a.e. } x \in B(p, \ell).$$

As customary, the notation $B(p, \ell)$ here denotes the Euclidean ball centered at p with radius ℓ .

Interestingly, the boundary regularity in Theorem 1.5 is stronger than in the case of the fractional Laplacian, in which the solution is in general not better than C^s , see [43].

As a byproduct of Theorem 1.5, one obtains also existence and regularity results, as given by the following result:

Theorem 1.6. *Assume that Ω is strictly convex. Let $m \geq \frac{N}{2} + 3$ and $f \in C^m(\Omega, \mathbb{R}) \cap L^\infty(\Omega)$. Then, there exists a unique classical solution $\mathbf{u}_f \in C(\mathbb{R}^N, \mathbb{R}) \cap C^2(\Omega, \mathbb{R})$ of*

$$(1.8) \quad \begin{cases} \mathcal{L}\mathbf{u}_f = f & \text{in } \Omega, \\ \mathbf{u}_f|_{\mathbb{R}^N \setminus \Omega} = 0, \end{cases}$$

with

$$\int_{\mathbb{R}^N} \frac{|\mathbf{u}_f(x)|}{1 + |x|^{N+2s}} dx < \infty.$$

Moreover, this solution \mathbf{u}_f satisfies the following additional properties:

- (i) $\mathbf{u}_f \in H^1(\mathbb{R}^N)$;
- (ii) $\|\mathbf{u}_f\|_{L^\infty(\mathbb{R}^N)} \leq \mathbf{c} \|f\|_{L^\infty(\Omega)}$;
- (iii) *for every $p \in \partial\Omega$ there exists $\ell > 0$ such that*

$$|\mathbf{u}_f(x)| \leq \mathbf{c} \|f\|_{L^\infty(\Omega)} \cdot |x - p| \quad \text{for all } x \in \Omega \cap B(p, \ell).$$

Here, $\mathbf{c} > 0$ is a constant independent of \mathbf{u}_f .

Let us comment a bit about the “philosophy” of the regularity theory developed in this paper. At first, in view of the classical regularity theory for the Laplacian, one could believe that adding an extra fractional Laplacian to an already elliptic problem would just produce “expected” outcomes, as a lower order effect. For instance, one could argue that interior regularity results (such as Theorem 1.4, as well as the forthcoming Theorem 4.1) could be a consequence of the classical theory. Specifically: on the one hand, it is known that, if $g \in L^2(\Omega)$ and $u \in H^1(\Omega)$ are such that

$$(1.9) \quad \begin{cases} -\Delta u = g & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

then $u \in H_{\text{loc}}^2(\Omega)$ and the corresponding norm is controlled, up to constants, by $\|u\|_{L^2(\Omega)} + \|g\|_{L^2(\Omega)}$.

On the other hand, if we have a weak solution u to the mixed problem, one can set $g := f - (-\Delta)^s u$ and we reduce the situation to the classical problem as formulated in (1.9).

These observations yield the regularity in class $H_{\text{loc}}^2(\Omega)$ with norm controlled by $\|u\|_{L^2(\Omega)} + \|(-\Delta)^s u\|_{L^2(\Omega)} + \|f\|_{L^2(\Omega)}$, leading to even better statements than the ones provided here (e.g. in the forthcoming estimate (4.1)), *but only when s is below $1/2$* (and above this threshold the norm in $H^s(\Omega)$ would be too strong compared with the initial regularity of the solution in $H^1(\Omega)$).

That is: for the low-range of s , the nonlocal part of the operator can be treated as a perturbation of the classical elliptic problem but for the high-range of fractional exponents these tricks seem to be unavailable and the operator must be treated on its own terms (roughly speaking, one can efficiently consider the fractional Laplacian as a lower order perturbation only when the fractional exponent is “sufficiently small”).

This is perhaps not a merely technical aspect of the problem: for instance, a distinction of this sort will pop up also in the construction of the barrier constructed to prove Theorem 1.5, since for s below $1/2$ one would not need the iteration exploited to reabsorb the nonlocal contributions and the proof would significantly simplify.

Similar occurrence of different phenomena according to the fractional threshold $1/2$ appear in the literature in several descriptions of nonlocal problems, see e.g. [17, 44, 27, 29] and the references therein.

For all these reasons, though in principle other approaches could be possible (relying e.g. on intermediate estimates to bootstrap, scaled norms and interpolation theory), we think it is useful and instructive to develop a series of analytic tools which account for mixed operators in their whole complexity rather than limiting our vision to perturbation methods from the classical cases.

After this article was completed and posted online, we have received the very interesting preprint [1], which considered a nonlinear problem of mixed type. The motivations, methodologies and results obtained are different from the ones in this paper, but Theorem 1.1 in [1] is related to Theorem 1.6 here. Complementary to this result, we stress that Theorem 1.4 in [1] shows that there exists no classical solutions of (1.8) when f is merely in $L_{\text{loc}}^\infty(\Omega)$.

Finally, we mention that similar results could be obtained via the probabilistic methods and the Green representation formulas dealt with in [19, 20]. For regularity results involving mixed order diffusive operators in different directions see also [4, 34, 35].

The proof of the boundary regularity result in Theorem 1.5 relies on the introduction of an explicit barrier. To the best of our knowledge, this barrier is completely new in the literature, and its construction is based on an iterative method of introducing “correctors” to recursively compensate the terms produced by the action of the nonlocal operator. We think that this technique of iteratively canceling “the nonlocal tails” is interesting in itself and can produce other results in greater generality.

Some of the methodologies and motivations presented in this paper will be also exploited in [10] to analyze the qualitative properties of solutions in specific problems.

The rest of this article is organized as follows. In Section 2 we introduce the functional framework in which we work, discussing in particular the notion of weak solutions and giving the existence result in Theorem 1.1. Section 3 is devoted to the maximum principles, and to the proofs of Theorems 1.2 and 1.3, and Section 4 to the regularity theory, and to the proofs of Theorems 1.4, 1.5 and 1.6.

Then, in Appendix A we collect some counterexamples to the maximum principle.

2. EXISTENCE OF WEAK SOLUTIONS

In this section we establish some basic facts on the existence of weak solutions for the Dirichlet problem associated with \mathcal{L} , as defined in (1.1), that is,

$$(2.1) \quad \begin{cases} \mathcal{L}u = f & \text{in } \Omega, \\ u|_{\mathbb{R}^N \setminus \Omega} = 0. \end{cases}$$

Throughout the sequel, we tacitly understand that $\Omega \subseteq \mathbb{R}^N$ is a bounded open set with C^1 boundary. Moreover, $s \in (0, 1)$ is a fixed parameter, and $(-\Delta)^s$ is as in (1.2). In studying the solvability of (2.1), a ‘natural’ space to consider is the following

$$(2.2) \quad \mathbb{X}(\Omega) := \{u \in H^1(\mathbb{R}^N) : u \equiv 0 \text{ in } \mathbb{R}^N \setminus \Omega\}.$$

We observe that, in view of the regularity assumption on Ω , the space $\mathbb{X}(\Omega)$ is contained in $H^1(\mathbb{R}^N)$ and is isomorphic to $H_0^1(\Omega)$ via the ‘zero-extension’ map defined as

$$\mathcal{E}_0 : H_0^1(\mathbb{R}^N) \rightarrow \mathbb{X}(\Omega), \quad \text{such that} \quad \mathcal{E}_0(u) := u \cdot \chi_\Omega.$$

As a consequence, $\mathbb{X}(\Omega)$ is endowed with a structure of (real) Hilbert space by the scalar product defined as follows

$$\langle u, v \rangle_{\mathbb{X}} := \int_{\Omega} \langle \nabla u, \nabla v \rangle \, dx \quad \text{for all } u, v \in \mathbb{X}(\Omega).$$

The norm associated with $\langle \cdot, \cdot \rangle_{\mathbb{X}}$ is

$$\|u\|_{\mathbb{X}} := \|\nabla u\|_{L^2(\Omega)} \quad \text{for all } u \in \mathbb{X}(\Omega),$$

and $C_0^\infty(\Omega, \mathbb{R})$ is dense in $\mathbb{X}(\Omega)$.

Furthermore, the classical Poincaré inequality holds in \mathbb{X} : more precisely, there exists a constant $\mathbf{c} > 0$ such that

$$(2.3) \quad \|u\|_{H^1(\mathbb{R}^N)} \leq \mathbf{c} \|u\|_{\mathbb{X}} \quad \text{for all } u \in \mathbb{X}(\Omega).$$

After all these preliminaries, we can give the following definition.

Definition 2.1. Let $f \in L^2(\Omega)$. We say that a function $u \in H^1(\mathbb{R}^N)$ is a *weak solution of the equation*

$$(2.4) \quad \mathcal{L}u = f \quad \text{in } \Omega$$

if, for every test function $\varphi \in C_0^\infty(\Omega, \mathbb{R})$, one has

$$(2.5) \quad \int_{\Omega} \langle \nabla u, \nabla \varphi \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy = \int_{\Omega} f \varphi \, dx.$$

Furthermore, one can more generally say that a function $u \in H^1(\mathbb{R}^N)$ weakly satisfies $\mathcal{L}u \geq f$ in Ω if, for every nonnegative $v \in \mathbb{X}(\Omega)$, one has

$$(2.6) \quad \int_{\Omega} \langle \nabla u, \nabla v \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+2s}} \, dx \, dy \geq \int_{\Omega} f v \, dx.$$

One can say that u weakly satisfies $\mathcal{L}u \leq f$ in Ω if $v := -u$ weakly satisfies $\mathcal{L}v \geq f$ in Ω .

Finally, if u is a weak solution of (2.4) and $u \in \mathbb{X}(\Omega)$, we say that u is a *weak solution of problem (2.1)*.

Remark 2.2. We notice that, if $u \in H^1(\mathbb{R}^N)$ is any weak solution of (2.4) (for some $f \in L^2(\Omega)$), from the density of $C_0^\infty(\Omega, \mathbb{R})$ in $\mathbb{X}(\Omega)$ we get

$$(2.7) \quad \begin{aligned} \int_{\Omega} \langle \nabla u, \nabla v \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+2s}} \, dx \, dy \\ = \int_{\Omega} f v \, dx \quad \text{for all } v \in \mathbb{X}(\Omega). \end{aligned}$$

In particular, if u_f is a solution of (2.1) (so that $u_f \in \mathbb{X}(\Omega)$), we are entitled to use u_f as a test function in (2.7), obtaining

$$\int_{\Omega} |\nabla u_f|^2 \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u_f(x) - u_f(y))^2}{|x - y|^{N+2s}} \, dx \, dy = \int_{\Omega} f u_f \, dx.$$

Having introduced the functional framework in which we work, we now prove the basic existence result in Theorem 1.1.

Proof of Theorem 1.1. We consider on the space $\mathbb{X}(\Omega)$ the bilinear form B and the linear map F defined, respectively, as follows:

$$(\star) \quad B(u, v) := \int_{\Omega} \langle \nabla u, \nabla v \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(v(x) - v(y))}{|x - y|^{N+2s}} \, dx \, dy,$$

$$(*) \quad F(u) := \int_{\Omega} f u \, dx.$$

Using (2.3) and the fact that $H^1(\mathbb{R}^N)$ is continuously embedded into $H^s(\mathbb{R}^N)$, we see that both B and F are (well-posed and) continuous with respect to the Hilbert structure of $\mathbb{X}(\Omega)$. Moreover, we have that

$$B(u, u) = \|\nabla u\|_{L^2(\Omega)}^2 + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|u(x) - u(y)|^2}{|x - y|^{N+2s}} \, dx \, dy \geq \|\nabla u\|_{L^2(\Omega)}^2 = \|u\|_{\mathbb{X}(\Omega)}^2,$$

for all $u \in \mathbb{X}(\Omega)$. Hence, we are in the position to apply Lax-Milgram's theorem, ensuring the existence of a unique function $u_f \in \mathbb{X}(\Omega)$ such that

$$(2.8) \quad B(u_f, v) = F(v) \text{ for every } v \in \mathbb{X}(\Omega),$$

further satisfying the estimate

$$\begin{aligned} \|u_f\|_{\mathbb{X}(\Omega)} &\leq \mathbf{c} \|F\|_{(\mathbb{X}(\Omega))^*} = \mathbf{c} \sup_{\|v\|_{\mathbb{X}(\Omega)}=1} |F(v)| \leq \mathbf{c} \sup_{\|v\|_{\mathbb{X}(\Omega)}=1} \int_{\Omega} |f| |v| \, dx \\ &\leq \mathbf{c} \sup_{\|v\|_{\mathbb{X}(\Omega)}=1} \|f\|_{L^2(\Omega)} \|v\|_{L^2(\Omega)} \leq \mathbf{c} \|f\|_{L^2(\Omega)}, \end{aligned}$$

with the positive constant \mathbf{c} possibly varying from line to line (actually, this argument shows that $\|u_f\|_{\mathbb{X}(\Omega)} \leq \mathbf{c} \|f\|_{H^{-1}(\Omega)}$). From (2.8) we immediately conclude that u_f is a weak solution of (2.4) (according to Definition 2.1), and the proof is complete. \square

Since one of the aims of this paper is to prove regularity results for weak solution of (2.1), it is convenient to fix the following definition. We set

$$(2.9) \quad \mathcal{C}_s(\mathbb{R}^N) := \left\{ u \in C(\mathbb{R}^N, \mathbb{R}) : \int_{\mathbb{R}^N} \frac{|u(x)|}{1 + |x|^{N+2s}} \, dx < \infty \right\}.$$

Definition 2.3. Let $f : \Omega \rightarrow \mathbb{R}$. We say that a function $u : \mathbb{R}^N \rightarrow \mathbb{R}$ is a *classical solution* of (2.1) if it satisfies the following properties:

- (1) $u \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$;
- (2) $\mathcal{L}u = f$ pointwise in Ω ;
- (3) $u(x) = 0$ for every $x \in \mathbb{R}^N \setminus \Omega$.

If u fulfills only (1) and (2), we say that u is a *classical solution* of (2.4).

We observe that, if $u \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$ is a classical solution of (2.1), it follows from (1) and (3) in Definition 2.3 that

$$(2.10) \quad u \in L^\infty(\mathbb{R}^N), \quad \text{with } \|u\|_{L^\infty(\mathbb{R}^N)} = \max_{x \in \bar{\Omega}} |u(x)|.$$

Remark 2.4. As it is very well-known, the (linear) space $\mathcal{C}_s(\mathbb{R}^N)$ is 'good' for dealing with the fractional Laplacian. In fact, if $u \in C^2(\Omega, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$, it is possible to compute $(-\Delta)^s u(x)$ pointwise for every $x \in \Omega$, and

$$\begin{aligned} (-\Delta)^s u(x) &= c_{N,s} \text{P.V.} \int_{\mathbb{R}^N} \frac{u(x) - u(y)}{|x - y|^{N+2s}} \, dy \\ &= -\frac{c_{N,s}}{2} \int_{\mathbb{R}^N} \frac{u(x+z) + u(x-z) - 2u(x)}{|z|^{N+2s}} \, dz. \end{aligned}$$

Moreover, it is easy to check that $(-\Delta)^s u \in L^\infty_{\text{loc}}(\Omega)$.

The next remark describes the relation between weak and classical solutions.

Remark 2.5. Let $f \in L^2(\Omega)$ and let $u \in H^1(\mathbb{R}^N)$ be a weak solution of (2.4). If we further assume that $u \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$, we can compute

$$\mathcal{L}u(x) = -\Delta u(x) + (-\Delta)^s u(x) \quad \text{for every } x \in \Omega.$$

Then, it is not difficult to check that $\mathcal{L}u = f$ pointwise in Ω , so that u is a classical solution of (2.4). Conversely, if $u \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$ is a classical solution of (2.4) further satisfying $u \in H^1(\mathbb{R}^N)$, then u is also a weak solution of (2.4).

A simple consequence of Theorem 1.1 is the solvability of the non-homogeneous Dirichlet problem, as follows:

Corollary 2.6. *Let $f \in L^2(\Omega)$ and $g \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\overline{\Omega}, \mathbb{R}) \cap H^1(\mathbb{R}^N)$. Then, there exists a unique weak solution $u \in H^1(\mathbb{R}^N)$ of the non-homogeneous Dirichlet problem*

$$(2.11) \quad \begin{cases} \mathcal{L}u = f & \text{in } \Omega, \\ u|_{\mathbb{R}^N \setminus \Omega} = g. \end{cases}$$

Proof. We observe that, since $g \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\overline{\Omega}, \mathbb{R})$, the function $\mathcal{L}g = -\Delta g + (-\Delta)^s g$ can be computed pointwise in Ω , and

$$\mathcal{L}g \in L^\infty(\Omega) \subseteq L^2(\Omega).$$

Hence, we can apply Theorem 1.1 to get the existence of a unique weak solution $v \in H^1(\mathbb{R}^N)$ of

$$\begin{cases} \mathcal{L}v = f - \mathcal{L}g & \text{in } \Omega, \\ v|_{\mathbb{R}^N \setminus \Omega} = 0. \end{cases}$$

Setting $u := v + g$, it is then immediate to see that $u \in H^1(\mathbb{R}^N)$ solves (2.11), thus completing the proof of Corollary 2.6. \square

Remark 2.7. Another functional framework naturally arising with the mixed operator \mathcal{L} is given by the spaces $\mathcal{H} := H^1(\Omega) \cap H^s(\mathbb{R}^N)$ and

$$\mathcal{H}_0 := \{u \in \mathcal{H} : u \equiv 0 \text{ a.e. in } \mathbb{R}^N \setminus \Omega\}.$$

In principle, one is tempted to define a \mathcal{H} -weak solution of (2.4) as a function $u \in \mathcal{H}$ such that (2.5) holds for every $\varphi \in C_0^\infty(\Omega, \mathbb{R})$, and a \mathcal{H}_0 -weak solution of (2.1) as a \mathcal{H} -weak solution of (2.4) which belongs to \mathcal{H}_0 .

The use of Lax-Milgram's theorem would provide the existence of a unique \mathcal{H}_0 -solution v_f of (2.1) (for some $f \in L^2(\Omega)$). Nevertheless, we prefer to use the functional setting in (2.2), since it possesses better density properties for smooth compactly supported functions, allowing us to use the positive and negative parts of the solution as a test function in (2.5).

This technical aspect is crucial for us in proving a weak maximum principle as in Theorem 1.2.

Another possible functional setting consists in requiring that, for weak solutions, identity (2.5) is fulfilled for all $v \in \mathcal{H}_0$. Nevertheless, this approach causes a technical difficulty in the integration by parts formula, which is needed to prove that a classical solution of (2.4)-(2.1) is also a weak solution, thus confirming that our choice in (2.2) is likely to be structurally more robust to deal with the PDE properties of the solutions.

3. SOME MAXIMUM PRINCIPLES FOR \mathcal{L}

In this section we establish some weak/strong maximum principles for \mathcal{L} , as defined in (1.1). To begin with, we prove the weak maximum principle for *weak solutions* in Theorem 1.2.

Proof of Theorem 1.2. Arguing by contradiction, we assume that there exists a set $E \subseteq \Omega$, with positive Lebesgue measure, such that $u < 0$ a.e. on E . We then define

$$w := u_- = \max\{-u, 0\}$$

and we observe that, since $u \in H^1(\mathbb{R}^N)$ and $u \geq 0$ a.e. in $\mathbb{R}^N \setminus \Omega$, one has

$$(3.1) \quad w \in H^1(\mathbb{R}^N) \text{ and } w \equiv 0 \text{ a.e. in } \mathbb{R}^N \setminus \Omega.$$

Moreover,

$$(3.2) \quad w = -u > 0 \quad \text{on } E.$$

In particular, recalling formula (2.2), from (3.1) we have that $w \in \mathbb{X}(\Omega)$. Hence, according to formula (2.7) in Remark 2.2, we can use w as a test function in (2.6), obtaining that

$$(3.3) \quad \begin{aligned} 0 &\leq \int_{\Omega} \langle \nabla u, \nabla w \rangle dx + \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy \\ &= - \int_{\Omega} |\nabla u_-|^2 dx + \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy \\ &\leq \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy. \end{aligned}$$

On the other hand, denoting by v the positive part of u , that is $v := u_+ = \max\{u, 0\}$, we have that $u = v - w$. Therefore, utilizing (3.1) and (3.2), we have that

$$(3.4) \quad \begin{aligned} &\int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy \\ &= \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy - \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|w(x) - w(y)|^2}{|x - y|^{N+2s}} dx dy \\ &\leq \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy - \int_{E \times (\mathbb{R}^N \setminus \Omega)} \frac{|w(x) - w(y)|^2}{|x - y|^{N+2s}} dx dy \\ &= \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy - \int_{E \times (\mathbb{R}^N \setminus \Omega)} \frac{|w(x)|^2}{|x - y|^{N+2s}} dx dy \\ &< \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy. \end{aligned}$$

Also, a ‘case-by-case’ computation shows that

$$(v(x) - v(y))(w(x) - w(y)) = (u_+(x) - u_+(y))(u_-(x) - u_-(y)) \leq 0$$

for almost every $x, y \in \mathbb{R}^N$. Plugging this information into (3.4), we obtain that

$$\int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(w(x) - w(y))}{|x - y|^{N+2s}} dx dy < 0.$$

This is in contradiction with (3.3), and thus $u \geq 0$ a.e. on Ω , as desired. \square

As regards *classical solutions*, we now prove Theorem 1.3:

Proof of Theorem 1.3. First of all, we establish the weak maximum principle in (1.3). For this, we suppose by contradiction that there exists a point $\xi \in \Omega$ such that $u(\xi) < 0$. Since $u \in C(\mathbb{R}^N, \mathbb{R})$ and $\overline{\Omega}$ is compact, we can thus find $x_0 \in \overline{\Omega}$ such that

$$(3.5) \quad u(x_0) = \min_{\overline{\Omega}} u < 0.$$

Moreover, since $u \geq 0$ in $\mathbb{R}^N \setminus \Omega$, we necessarily have that $x_0 \in \Omega$, that is x_0 in an *interior* minimum point for u in Ω . Hence, $\Delta u(x_0) \geq 0$ and

$$(3.6) \quad (-\Delta)^s u(x_0) = \mathcal{L}u(x_0) + \Delta u(x_0) \geq \Delta u(x_0) \geq 0.$$

Moreover, we have that $u(x_0) \leq u(x)$ for all $x \in \mathbb{R}^N$. Thus, we get

$$(3.7) \quad (-\Delta)^s u(x_0) = c_{N,s} \text{ P.V. } \int_{\mathbb{R}^N} \frac{u(x_0) - u(y)}{|x_0 - y|^{N+2s}} dy = c_{N,s} \int_{\mathbb{R}^N} \frac{u(x_0) - u(y)}{|x_0 - y|^{N+2s}} dy \leq 0.$$

Owing to (3.6), and taking into account the fact the integrand function in (3.7) is non-positive, we then conclude that

$$u \equiv u(x_0) \text{ on } \mathbb{R}^N.$$

In particular, since $u \geq 0$ in $\mathbb{R}^N \setminus \Omega$, we get $u(x_0) \geq 0$, but this is clearly in contradiction with (3.5). Hence, $u \geq 0$ on Ω , which establishes (1.3), as desired.

Now we prove the strong maximum principle for classical solutions in (1.4). To this end, we recall (3.6) and we note that

$$(3.8) \quad 0 \leq (-\Delta)^s u(x_0) = c_{N,s} \text{P.V.} \int_{\mathbb{R}^N} \frac{u(x_0) - u(y)}{|x_0 - y|^{N+2s}} dy = -c_{N,s} \int_{\mathbb{R}^N} \frac{u(y)}{|x_0 - y|^{N+2s}} dy.$$

Additionally, by (1.3), we know that $u \geq 0$ in \mathbb{R}^N . Comparing this with (3.8), we obtain that $u \equiv 0$ throughout \mathbb{R}^N . This establishes (1.4) as desired. \square

We shall see in Appendix A that, due to the presence of the non-local term $(-\Delta)^s$, a maximum principle analogous to that in (1.4) *does not hold* if the condition $u \geq 0$ is satisfied only on $\partial\Omega$ (as in the classical case). Furthermore, we shall also show that maximum principles analogous to Theorem 1.2 and in (1.3) of Theorem 1.3 *do not hold* for

$$\mathcal{L}' := \Delta + (-\Delta)^s.$$

As a simple consequence of the weak maximum principle in (1.3) we have a uniqueness result, as follows:

Corollary 3.1. *There exists at most one classical solution of (2.1).*

Proof. Let $u_1, u_2 \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$ be two classical solutions of problem (2.1). Introducing the function $v := u_1 - u_2$, it is immediate to recognize that

- (a) $v \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$;
- (b) v is a solution of the Dirichlet problem

$$\begin{cases} \mathcal{L}u = -\Delta u + (-\Delta)^s u = 0 & \text{pointwise in } \Omega, \\ u = 0 & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$

Thus, by the weak maximum principle in (1.3) we readily conclude that $v \equiv 0$ on \mathbb{R}^N , so that $u_1 \equiv u_2$ on \mathbb{R}^N . This ends the proof. \square

4. INTERIOR AND BOUNDARY REGULARITY FOR \mathcal{L}

The main aim of this section is to prove both interior and boundary regularity for \mathcal{L} . To be more precise, we first establish interior H^m -regularity for the weak solutions of (2.4), that is Theorem 1.4; then, we prove boundary regularity for the solutions of (2.1), that is Theorem 1.5.

4.1. Interior H^m -regularity and proof of Theorem 1.4. To begin with, we prove the following H^2 -regularity theorem, which in turn will serve as the basic step to prove interior H^m -regularity in Theorem 1.4.

Theorem 4.1. *Let $f \in L^2(\Omega)$ and let $u \in H^1(\mathbb{R}^N)$ be a weak solution of equation (2.4). Then, $u \in H_{\text{loc}}^2(\Omega)$. Furthermore, given any open set V with $\overline{V} \subseteq \Omega$, there exists a constant $\Lambda > 0$, independent of u , such that*

$$(4.1) \quad \|u\|_{H^2(V)} \leq \Lambda \left(\|f\|_{L^2(\Omega)} + \|u\|_{H^1(\mathbb{R}^N)} \right).$$

Proof. Let V be a fixed open set with $\overline{V} \subseteq \Omega$, and let $\rho_0 > 0$ be such that

$$V_\rho := \{x \in \mathbb{R}^N : \text{dist}(x, V) < \rho\} \subseteq \Omega \quad \text{for every } \rho \in [0, 2\rho_0].$$

Moreover, let $\zeta \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ be a cut-off function satisfying

- (a) $\zeta \equiv 1$ on V and $\text{supp}(\zeta) \subseteq V_{\rho_0}$;
- (b) $0 \leq \zeta \leq 1$ on \mathbb{R}^N .

Finally, for every fixed $k \in \{1, \dots, N\}$ and every $0 < |h| < \rho_0$, we set

$$(4.2) \quad \varphi := D_k^{-h} (\zeta^2 D_k^h u), \quad \text{where } D_k^h w(x) := \frac{w(x + h e_k) - w(x)}{h}.$$

We notice that, since $u \in H^1(\mathbb{R}^N)$, then $D_k^h u \in H^1(\mathbb{R}^N)$. Also, in light of (a), we have that $\zeta \in C_0^\infty(V_{\rho_0}, \mathbb{R})$. As a result,

$$\hat{\varphi} := \zeta^2 D_k^h u \in H^1(\mathbb{R}^N) \quad \text{and} \quad \text{supp}(\hat{\varphi}) \subseteq V_{\rho_0} \subseteq \Omega.$$

As a consequence, by the definition of φ in (4.2), we see that $\varphi \in H^1(\mathbb{R}^N)$ and $\text{supp}(\varphi) \subseteq V_{2\rho_0} \subseteq \Omega$, which implies that $\varphi \in H_0^1(\Omega)$. Therefore, we are in the position of using φ as a test function in (2.7), obtaining (after a standard ‘integration by parts’ for difference quotients)

$$\begin{aligned} & \sum_{i=1}^N \int_{\Omega} \zeta^2 |D_k^h(\partial_{x_i} u)|^2 dx + \sum_{i=1}^N 2 \int_{\Omega} \zeta D_k^h u \partial_{x_i} \zeta D_k^h(\partial_{x_i} u) dx \\ (4.3) \quad & + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(D_k^h u(x) - D_k^h u(y))(\zeta^2(x) D_k^h u(x) - \zeta^2(y) D_k^h u(y))}{|x - y|^{N+2s}} dx dy \\ & = \int_{\Omega} f \varphi dx. \end{aligned}$$

Now, by exploiting Cauchy-Swcharz’s inequality and the classical Young inequality (with $\varepsilon = 1/2$), we obtain the following estimate

$$\begin{aligned} & \left| \sum_{i=1}^N 2 \int_{\Omega} \zeta D_k^h u \partial_{x_i} \zeta D_k^h(\partial_{x_i} u) dx \right| \leq 2 \int_{\Omega} \zeta |D_k^h u| |\nabla \zeta| |D_k^h(\nabla u)| dx \\ (4.4) \quad & \leq \frac{1}{2} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx + 2 \left(\sup_{\mathbb{R}^N} |\nabla \zeta| \right) \cdot \int_{V_{\rho_0}} |D_k^h u|^2 dx \\ & \leq \frac{1}{2} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx + C \int_{\Omega} |\nabla u|^2 dx, \end{aligned}$$

for some $C > 0$, where in the last inequality we have used the fact that

$$(4.5) \quad \int_{\mathbb{R}^N} |D_k^h \omega|^2 dx \leq 4N^2 \int_{\mathbb{R}^N} |\nabla \omega|^2 dx, \quad \text{for every } \omega \in H^1(\mathbb{R}^N)$$

(see, e.g., [33, Theorem. 3, Chapter 5.8.2]). Gathering together estimates (4.3) and (4.4), we get

$$\begin{aligned} & \frac{1}{2} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx - C \int_{\Omega} |\nabla u|^2 dx \\ (4.6) \quad & + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(D_k^h u(x) - D_k^h u(y))(\zeta^2(x) D_k^h u(x) - \zeta^2(y) D_k^h u(y))}{|x - y|^{N+2s}} dx dy \\ & \leq \int_{\Omega} f \varphi dx. \end{aligned}$$

Now we estimate the integral in the right hand side of (4.6). To this end we first observe that, recalling (4.2) and exploiting once again (4.5), we have

$$\begin{aligned} \int_{\Omega} \varphi^2 dx &= \int_{\mathbb{R}^N} |D_k^{-h}(\zeta^2 D_k^h u)|^2 dx \\ &\leq 4N^2 \int_{\mathbb{R}^N} |\nabla(\zeta^2 D_k^h u)|^2 dx = C \int_{V_{\rho_0}} |\nabla(\zeta^2 D_k^h u)|^2 dx \\ &\leq C \left(\int_{V_{\rho_0}} |D_k^h u|^2 dx + \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx \right) \\ &\leq C \left(\int_{\Omega} |\nabla u|^2 dx + \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx \right), \end{aligned}$$

up to renaming $C > 0$ from line to line. From this, using Young's inequality (with $\varepsilon = 1/(4C)$), we get

$$(4.7) \quad \begin{aligned} \left| \int_{\Omega} f \varphi \, dx \right| &\leq \varepsilon \int_{\Omega} \varphi^2 \, dx + \frac{1}{\varepsilon} \int_{\Omega} f^2 \, dx \\ &\leq \frac{1}{4} \int_{\Omega} |\nabla u|^2 \, dx + \frac{1}{4} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 \, dx + 4C \int_{\Omega} f^2 \, dx. \end{aligned}$$

By combining (4.7) with (4.6), we then obtain

$$(4.8) \quad \begin{aligned} &\frac{1}{4} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 \, dx \\ &+ \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(D_k^h u(x) - D_k^h u(y)) (\zeta^2(x) D_k^h u(x) - \zeta^2(y) D_k^h u(y))}{|x - y|^{N+2s}} \, dx \, dy \\ &\leq C \left(\int_{\Omega} f^2 \, dx + \int_{\Omega} |\nabla u|^2 \, dx \right), \end{aligned}$$

up to relabeling $C > 0$.

We now provide a careful estimate of the non-local term in the left hand side of (4.8), i.e.,

$$(4.9) \quad J_s := \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(D_k^h u(x) - D_k^h u(y)) (\zeta^2(x) D_k^h u(x) - \zeta^2(y) D_k^h u(y))}{|x - y|^{N+2s}} \, dx \, dy.$$

To this end we first notice that, with obvious algebraic manipulation, we can write

$$(4.10) \quad J_s = J_{0,s} + J_{1,s} + J_{2,s},$$

where

$$\begin{aligned} J_{0,s} &:= \int_{\mathbb{R}^N \times \mathbb{R}^N} \zeta^2(x) \cdot \frac{|D_k^h u(x) - D_k^h u(y)|^2}{|x - y|^{N+2s}} \, dx \, dy, \\ J_{1,s} &:= \int_{\mathbb{R}^N \times \mathbb{R}^N} D_k^h u(y) \zeta(x) \cdot \frac{(D_k^h u(x) - D_k^h u(y)) (\zeta(x) - \zeta(y))}{|x - y|^{N+2s}} \, dx \, dy \\ \text{and } J_{2,s} &:= \int_{\mathbb{R}^N \times \mathbb{R}^N} D_k^h u(y) \zeta(y) \cdot \frac{(D_k^h u(x) - D_k^h u(y)) (\zeta(x) - \zeta(y))}{|x - y|^{N+2s}} \, dx \, dy. \end{aligned}$$

We notice that, by exchanging the variables x and y , we see that

$$(4.11) \quad \begin{aligned} J_{0,s} &= \int_{\mathbb{R}^N \times \mathbb{R}^N} \zeta^2(y) \cdot \frac{|D_k^h u(y) - D_k^h u(x)|^2}{|y - x|^{N+2s}} \, dx \, dy \\ &= \int_{\mathbb{R}^N \times \mathbb{R}^N} \zeta^2(y) \cdot \frac{|D_k^h u(x) - D_k^h u(y)|^2}{|x - y|^{N+2s}} \, dx \, dy. \end{aligned}$$

Moreover, since $\zeta \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$, for every $y \in \mathbb{R}^N$ we have

$$(4.12) \quad \begin{aligned} &\int_{\mathbb{R}^N} \frac{|\zeta(x) - \zeta(y)|^2}{|x - y|^{N+2s}} \, dx \\ &\leq \left(\sup_{\mathbb{R}^N} |\nabla \zeta|^2 \right) \cdot \int_{\{|x-y| \leq 1\}} \frac{dx}{|x - y|^{N+2(s-1)}} + 4 \left(\sup_{\mathbb{R}^N} |\zeta|^2 \right) \cdot \int_{\{|x-y| > 1\}} \frac{dx}{|x - y|^{N+2s}} \\ &=: \mathbf{c}(\zeta, s). \end{aligned}$$

Using Young's inequality once again, (4.12) and (4.5), we can estimate $J_{1,s}$ as follows:

$$\begin{aligned}
|J_{1,s}| &\leq \frac{1}{4} \int_{\mathbb{R}^N \times \mathbb{R}^N} \zeta^2(x) \cdot \frac{|D_k^h u(x) - D_k^h u(y)|^2}{|x - y|^{N+2s}} dx dy \\
&\quad + 4 \int_{\mathbb{R}^N \times \mathbb{R}^N} |D_k^h u(y)|^2 \cdot \frac{|\zeta(x) - \zeta(y)|^2}{|x - y|^{N+2s}} dx dy \\
(4.13) \quad &\leq \frac{J_{0,s}}{4} + 4 \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |D_k^h u(y)|^2 dy \\
&\leq \frac{J_{0,s}}{4} + 4 \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |\nabla u|^2 dx.
\end{aligned}$$

Similarly, making again use of Young's inequality and (4.5), and recalling (4.11), we have the following estimate for $J_{2,s}$:

$$\begin{aligned}
|J_{2,s}| &\leq \frac{1}{4} \int_{\mathbb{R}^N \times \mathbb{R}^N} \zeta^2(y) \cdot \frac{|D_k^h u(x) - D_k^h u(y)|^2}{|x - y|^{N+2s}} dx dy \\
&\quad + 4 \int_{\mathbb{R}^N \times \mathbb{R}^N} |D_k^h u(y)|^2 \cdot \frac{|\zeta(x) - \zeta(y)|^2}{|x - y|^{N+2s}} dx dy \\
(4.14) \quad &\leq \frac{J_{0,s}}{4} + 4 \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |D_k^h u(y)|^2 dy \\
&\leq \frac{J_{0,s}}{4} + 4 \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |\nabla u|^2 dx.
\end{aligned}$$

Gathering together (4.10), (4.13) and (4.14), we obtain that

$$(4.15) \quad J_s \geq \frac{J_{0,s}}{2} - 8 \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |\nabla u|^2 dx.$$

Finally, by combining (4.8) and (4.15), and recalling (4.9), we derive

$$\begin{aligned}
&\frac{1}{4} \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx + \frac{c_{N,s}}{2} \cdot \frac{J_{0,s}}{2} \\
&\leq 4c_{N,s} \mathbf{c}(\zeta, s) \int_{\mathbb{R}^N} |\nabla u|^2 dx + C \left(\int_{\Omega} f^2 dx + \int_{\Omega} |\nabla u|^2 dx \right) \\
&\leq C \left(\int_{\Omega} f^2 dx + \int_{\mathbb{R}^N} |\nabla u|^2 dx \right),
\end{aligned}$$

up to renaming $C > 0$. From this, since $J_{0,s} \geq 0$, we get

$$\int_V |D_k^h(\nabla u)|^2 dx \leq \int_{\Omega} \zeta^2 |D_k^h(\nabla u)|^2 dx \leq C \left(\int_{\Omega} f^2 dx + \int_{\mathbb{R}^N} |\nabla u|^2 dx \right),$$

and thus, owing to [33, Theorem 3, Chapter 5.8.2], we conclude that $u \in H^2(V)$ and

$$\|u\|_{H^2(V)} \leq \Lambda \left(\|f\|_{L^2(\Omega)} + \|u\|_{H^1(\mathbb{R}^N)} \right),$$

for a suitable $\Lambda > 0$. Finally, since a careful inspection of the proof shows that Λ does not depend on u , we also obtain estimate (4.1). \square

Starting from Theorem 4.1, our next aim is to establish Theorem 1.4. In contrast to the classical case, the proof of Theorem 1.4 is not merely a bootstrap argument based on Theorem 4.1: in fact, the presence of the non-local term $(-\Delta)^s$ in \mathcal{L} prevents us to take derivatives of equation (2.4). To overcome this technical issue, we need to combine a suitable truncation argument with the use of difference quotient. This is done in the following preliminary lemmata.

In the sequel, we shall make use of the following notation: given an arbitrary set $A \subseteq \mathbb{R}^N$ and a number $\delta > 0$, we define

$$(4.16) \quad A_\delta := \{x \in \mathbb{R}^N : \text{dist}(x, A) < \delta\}.$$

Moreover, for any $\omega \in L^2(\mathbb{R}^N)$, we set

$$(4.17) \quad \omega_h(x) := D_k^h \omega(x) = \frac{\omega(x + he_k) - \omega(x)}{h},$$

with $h \in \mathbb{R}$ and $k \in \{1, \dots, N\}$.

Lemma 4.2. *Let $u \in H^1(\mathbb{R}^N)$ be a solution of (2.4), and let \mathcal{O} be an open set with $\overline{\mathcal{O}} \subseteq \Omega$. Let $\rho := \text{dist}(\mathcal{O}, \partial\Omega) > 0$. If $|h| < \rho$, then u_h solves*

$$(4.18) \quad \mathcal{L}u_h = f_h \quad \text{in } \mathcal{O}.$$

Proof. We observe that, for any $\varphi \in C_0^\infty(\mathcal{O}, \mathbb{R})$, we have that

$$(4.19) \quad \int_{\mathcal{O}} \langle \nabla u_h, \nabla \varphi \rangle dx = \int_{\mathbb{R}^N} \langle \nabla u_h, \nabla \varphi \rangle dx = - \int_{\mathbb{R}^N} \langle \nabla u, \nabla \psi_{-h} \rangle dx,$$

where we have used the notation

$$(4.20) \quad \psi_{-h}(x) := -\frac{\varphi(x - he_k) - \varphi(x)}{h}.$$

Moreover, since $\text{supp}(\varphi) \subseteq \mathcal{O}$, recalling the notation in (4.16), we have that

$$\text{supp}(\varphi(\cdot - he_k)) \subseteq \mathcal{O} + he_k \subseteq \{x \in \mathbb{R}^N : \text{dist}(x, \mathcal{O}) < |h|\} = \mathcal{O}_{|h|}.$$

Hence, if $|h| < \rho$ we obtain that

- (a) $\psi_{-h} \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$;
- (b) $\text{supp}(\psi_{-h}) \subseteq \mathcal{O}_\rho \subseteq \Omega$.

In particular,

$$(4.21) \quad \psi_{-h} \in C_0^\infty(\mathcal{O}_\rho, \mathbb{R}).$$

As a consequence, identity (4.19) can be written as

$$(4.22) \quad \int_{\mathcal{O}} \langle \nabla u_h, \nabla \varphi \rangle dx = - \int_{\Omega} \langle \nabla u, \nabla \psi_{-h} \rangle dx.$$

Furthermore, we observe that, for all $h \in \mathbb{R}$, we have

$$(4.23) \quad \begin{aligned} & \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u_h(x) - u_h(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dx dy \\ &= - \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(\psi_{-h}(x) - \psi_{-h}(y))}{|x - y|^{N+2s}} dx dy. \end{aligned}$$

Thus, by combining (4.22) with (4.23), and recalling (4.21) and the fact that u solves (2.4), we get

$$(4.24) \quad \begin{aligned} & \int_{\mathcal{O}} \langle \nabla u_h, \nabla \varphi \rangle dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u_h(x) - u_h(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dx dy \\ &= - \int_{\Omega} \langle \nabla u, \nabla \psi_{-h} \rangle dx - \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(\psi_{-h}(x) - \psi_{-h}(y))}{|x - y|^{N+2s}} dx dy \\ &= - \int_{\mathcal{O}_\rho} f \psi_{-h} dx. \end{aligned}$$

Now we observe that, if $|h| < \rho$,

$$\text{supp}(\varphi) \subseteq \mathcal{O} \subseteq \mathcal{O}_\rho \cap (\mathcal{O}_\rho - he_k),$$

and therefore, recalling (4.17) and (4.20), we have that

$$\begin{aligned} - \int_{\mathcal{O}_\rho} f \psi_{-h} dx &= \int_{\mathcal{O}_\rho} f(x) \left(\frac{\varphi(x - he_k) - \varphi(x)}{h} \right) dx \\ &= \frac{1}{h} \left(\int_{\mathcal{O}_{\rho - he_k}} f(x + he_k) \varphi(x) dx - \int_{\mathcal{O}_\rho} f(x) \varphi(x) dx \right) \\ &= \int_{\mathcal{O}} f_h \varphi dx. \end{aligned}$$

As a consequence of this and (4.24),

$$\int_{\mathcal{O}} \langle \nabla u_h, \nabla \varphi \rangle dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u_h(x) - u_h(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dx dy = \int_{\mathcal{O}} f_h \varphi dx,$$

for any $\varphi \in C_0^\infty(\mathcal{O}, \mathbb{R})$. Thus, recalling (1.1), this implies that u_h solves (4.18), as desired. \square

Thanks to Lemma 4.2, we can prove a ‘weaker version’ of Theorem 1.4.

Proposition 4.3. *Let $m \in \mathbb{N} \cup \{0\}$ and $f \in H^m(\Omega)$. If $u \in H^{m+1}(\mathbb{R}^N)$ is any weak solution of (2.4), then $u \in H_{\text{loc}}^{m+2}(\Omega)$.*

Furthermore, given any open set V with $\overline{V} \subseteq \Omega$, there exists a constant $\Lambda_m > 0$, independent of the function u , such that

$$(4.25) \quad \|u\|_{H^{m+2}(V)} \leq \Lambda_m \left(\|f\|_{H^m(\Omega)} + \|u\|_{H^{m+1}(\mathbb{R}^N)} \right).$$

Proof. We proceed by induction on $m \in \mathbb{N} \cup \{0\}$. First of all, the case $m = 0$ (that is, $f \in L^2(\Omega)$ and $u \in H^1(\mathbb{R}^N)$) is given by Theorem 4.1. Then, we assume that Proposition 4.3 holds for some $m \geq 0$, and we prove that it still holds for $m + 1$.

Let V be a fixed open set with $\overline{V} \subseteq \Omega$, and let \mathcal{O} be an open subset of Ω satisfying $\overline{V} \subseteq \mathcal{O}$ and $\overline{\mathcal{O}} \subseteq \Omega$. Moreover, let $f \in H^{m+1}(\Omega)$ and let $u \in H^{m+2}(\mathbb{R}^N)$ be a weak solution of (2.4).

Setting $\rho := \text{dist}(\mathcal{O}, \partial\Omega) > 0$, and recalling the notation in (4.17), if $|h| < \rho$ we know from Lemma 4.2 that u_h is a weak solution of the equation

$$\mathcal{L}u_h = f_h \text{ in } \mathcal{O}.$$

Moreover, since $u \in H^{m+2}(\mathbb{R}^N)$, we clearly have that $u_h \in H^{m+1}(\mathbb{R}^N)$. As a result, we are in the position of applying the inductive hypothesis to u_h , thus obtaining

- (a) $u_h \in H_{\text{loc}}^{m+2}(\mathcal{O})$ and, in particular, $u_h \in H^{m+2}(V)$;
- (b) there exists a constant $\Lambda_m > 0$, independent of h , such that

$$\|u_h\|_{H^{m+2}(V)} \leq \Lambda_m \left(\|f_h\|_{H^m(\mathcal{O})} + \|u_h\|_{H^{m+1}(\mathbb{R}^N)} \right) \quad \text{for } |h| < \rho.$$

Furthermore, we observe that, since $u \in H^{m+2}(\mathbb{R}^N)$, one has

$$\|u_h\|_{H^{m+1}(\mathbb{R}^N)} \leq \mathbf{c} \|u\|_{H^{m+2}(\mathbb{R}^N)},$$

where $\mathbf{c} > 0$ is a suitable constant which is independent of h (see, e.g., [33]). Analogously, since $f \in H^{m+1}(\Omega)$, we also have

$$\|f_h\|_{H^m(\mathcal{O})} \leq \mathbf{c} \|f\|_{H^{m+1}(\Omega)}.$$

Gathering together these facts, we obtain

$$\|u_h\|_{H^{m+2}(V)} \leq \mathbf{c} \Lambda_m \left(\|f\|_{H^{m+1}(\Omega)} + \|u\|_{H^{m+2}(\mathbb{R}^N)} \right),$$

and this estimate is *uniform* with respect to $h \in (-\rho, \rho)$. On account of [33, Theorem. 3, Chapter 5.8.2], we then easily conclude that $u \in H^{m+3}(V)$. Moreover,

$$\|u\|_{H^{m+3}(V)} \leq \mathbf{c} \Lambda_m \left(\|f\|_{H^{m+1}(\Omega)} + \|u\|_{H^{m+2}(\mathbb{R}^N)} \right).$$

This is precisely estimate (4.25), and the proof is complete. \square

To remove the assumption that $u \in H^{m+1}(\mathbb{R}^N)$ in Proposition 4.3, we need to perform a “truncation” argument: this is described in the next two lemmata.

Lemma 4.4. *Let $\mathcal{O} \subseteq \mathbb{R}^N$ be open, $\delta > 0$, and $\alpha > N$. Let $z \in L^2(\mathbb{R}^N)$ be such that*

$$(4.26) \quad z \equiv 0 \text{ a.e. on } \mathcal{O}_\delta,$$

with the notation introduced in (4.16). Then, the following facts hold.

(i) *for every fixed $x \in \mathcal{O}_{\delta/2}$, we have*

$$(4.27) \quad y \mapsto \frac{z(y)}{|x-y|^\alpha} \in L^1(\mathbb{R}^N).$$

(ii) *The function $\mathcal{I}_\alpha[z]$ defined as*

$$(4.28) \quad \mathcal{I}_\alpha[z](x) := \int_{\mathbb{R}^N} \frac{z(y)}{|x-y|^\alpha} dy$$

is of class C^∞ on $\mathcal{O}_{\delta/2}$.

Proof. (i) Let $x \in \mathcal{O}_{\delta/2}$. From (4.26) and Hölder’s inequality, we have the following estimate

$$(4.29) \quad \begin{aligned} \int_{\mathbb{R}^N} \frac{|z(y)|}{|x-y|^\alpha} dy &= \int_{\mathbb{R}^N \setminus \mathcal{O}_\delta} \frac{|z(y)|}{|x-y|^\alpha} dy \\ &\leq \left(\int_{\mathbb{R}^N \setminus \mathcal{O}_\delta} \frac{|z(y)|^2}{|x-y|^\alpha} dy \right)^{1/2} \cdot \left(\int_{\mathbb{R}^N \setminus \mathcal{O}_\delta} \frac{1}{|x-y|^\alpha} dy \right)^{1/2}. \end{aligned}$$

On the other hand, since $x \in \mathcal{O}_{\delta/2}$, it is immediate to check that

$$(4.30) \quad |x-y| \geq \frac{\delta}{2} \quad \text{for every } y \in \mathbb{R}^N \setminus \mathcal{O}_\delta.$$

By exploiting (4.30), we get from (4.29) that

$$\begin{aligned} \int_{\mathbb{R}^N} \frac{|z(y)|}{|x-y|^\alpha} dy &\leq \left(\frac{2}{\delta} \right)^{\alpha/2} \left(\int_{\{|x-y| \geq \delta/2\}} \frac{1}{|x-y|^\alpha} dy \right)^{1/2} \cdot \|z\|_{L^2(\mathbb{R}^N)} \\ &\leq \left(\frac{2}{\delta} \right)^{\alpha/2} \left(\int_{\{|w| \geq \delta/2\}} \frac{1}{|w|^\alpha} dw \right)^{1/2} \cdot \|z\|_{L^2(\mathbb{R}^N)}. \end{aligned}$$

From this, reminding that $\alpha > N$ and that $z \in L^2(\mathbb{R}^N)$, we obtain (4.27).

(ii) First of all, owing to (i), the function $\mathcal{I}_\alpha[z]$ is well-posed on $\mathcal{O}_{\delta/2}$. To prove its smoothness on $\mathcal{O}_{\delta/2}$ we show that, for every fixed $x \in \mathcal{O}_{\delta/2}$ and every N -tuple $\gamma = (\gamma_1, \dots, \gamma_N)$ of non-negative integers, one has

$$(4.31) \quad (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \mathcal{I}_\alpha[z](x) = \int_{\mathbb{R}^N \setminus \mathcal{O}_\delta} (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \left(\frac{z(y)}{|x-y|^\alpha} \right) dy.$$

To this end we first observe that, setting $|\gamma| := \sum_{k=1}^N \gamma_k$, one has

$$(4.32) \quad \left| (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \left(\frac{z(y)}{|x-y|^\alpha} \right) \right| \leq \kappa_\alpha \frac{|z(y)|}{|x-y|^{\alpha+|\gamma|}}$$

for every $x \neq y \in \mathbb{R}^N$ (here, κ_α is a positive constant only depending on α). Thus, by assertion (i) (applied to $\alpha + |\gamma|$ in place of α), we derive that

$$y \mapsto \mathcal{D}_x(y) := (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \left(\frac{z(y)}{|x-y|^\alpha} \right) \in L^1(\mathbb{R}^N)$$

for every fixed $x \in \mathcal{O}_{\delta/2}$. On account of this fact, and owing to classical results on the regularity of parameter-dependent integrals, to establish (4.31) it suffices to prove the following fact: *for every*

N -tuple $\gamma = (\gamma_1, \dots, \gamma_N)$ of non-negative integers and every $x_0 \in \mathcal{O}_{\delta/2}$, there exist $r > 0$ and a function

$$(4.33) \quad \Theta = \Theta_{\gamma, x_0, r} \in L^1(\mathbb{R}^N \setminus \mathcal{O}_\delta)$$

such that

- (a) $B(x_0, r) \subseteq \mathcal{O}_{\delta/2}$;
- (b) for every $x \in B(x_0, r)$ and every $y \in \mathbb{R}^N \setminus \mathcal{O}_\delta$, one has

$$(4.34) \quad \left| (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \left(\frac{z(y)}{|x-y|^\alpha} \right) \right| \leq \Theta(y).$$

To prove this statement, let $x_0 \in \mathcal{O}_{\delta/2}$, and $\gamma = (\gamma_1, \dots, \gamma_N) \in (\mathbb{N} \cup \{0\})^N$. Moreover, let $r > 0$ be such that $B(x_0, r) \subseteq \mathcal{O}_{\delta/2}$. We claim that there exists a constant $\mathbf{c} > 0$, only depending on x_0 and r , such that

$$(4.35) \quad \frac{|x-y|}{|x_0-y|} \geq \mathbf{c} \quad \text{for all } x \in B(x_0, r) \text{ and } y \in \mathbb{R}^N \setminus \mathcal{O}_\delta.$$

Indeed, recalling (4.30), if $x \in B(x_0, r) \subseteq \mathcal{O}_{\delta/2}$ and $y \in (\mathbb{R}^N \setminus \mathcal{O}_\delta) \cap B(x_0, 2r)$, one has

$$\frac{|x-y|}{|x_0-y|} \geq \frac{\delta/2}{2r} = \frac{\delta}{4r}.$$

On the other hand, if $y \in (\mathbb{R}^N \setminus \mathcal{O}_\delta) \setminus B(x_0, 2r)$, we have that

$$|x-x_0| < r < \frac{1}{2}|x_0-y|,$$

and therefore, by triangle inequality,

$$\frac{|x-y|}{|x_0-y|} \geq \frac{|x_0-y| - |x-x_0|}{|x_0-y|} = 1 - \frac{|x-x_0|}{|x_0-y|} \geq \frac{1}{2}.$$

Gathering together these facts, we obtain (4.35) with

$$\mathbf{c} := \min \{ \delta/(4r), 1/2 \}.$$

Now, by combining (4.32) with (4.35) we get

$$\left| (\partial_{x_1})^{\gamma_1} \dots (\partial_{x_N})^{\gamma_N} \left(\frac{z(y)}{|x-y|^\alpha} \right) \right| \leq \frac{\kappa_\alpha}{\mathbf{c}^{\alpha+|\gamma|}} \cdot \frac{|z(y)|}{|x_0-y|^{\alpha+|\gamma|}} =: \Theta_{\gamma, x_0, r}(y),$$

for every $x \in B(x_0, r)$ and every $y \in \mathbb{R}^N \setminus \mathcal{O}_\delta$, and this gives (4.34). Moreover, since $x_0 \in \mathcal{O}_{\delta/2}$, from (i) we infer that $\Theta_{\gamma, x_0, r} \in L^1(\mathbb{R}^N)$, thus showing (4.33). This ends the proof. \square

Lemma 4.5. *Let $u \in H^1(\mathbb{R}^N)$ be a solution of (2.4), and let \mathcal{O} be an open set with $\overline{\mathcal{O}} \subseteq \Omega$. Let $\rho := \text{dist}(\mathcal{O}, \partial\Omega)$ and $\zeta \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ satisfy*

- (i) $\zeta \equiv 1$ on $\mathcal{O}_{\rho/4}$;
- (ii) $\text{supp}(\zeta) \subseteq \mathcal{O}_{\rho/2}$;
- (iii) $0 \leq \zeta \leq 1$ on \mathbb{R}^N ;

with the notation introduced in (4.16).

Then, there exists $\psi \in C^\infty(\overline{\mathcal{O}}, \mathbb{R})$ such that $v := u\zeta$ is a weak solution of

$$(4.36) \quad \mathcal{L}v = f + \psi \quad \text{in } \mathcal{O}.$$

Proof. First of all, since $u \in H^1(\mathbb{R}^N)$ and $\zeta \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$, one has that $v \in H^1(\mathbb{R}^N)$. Moreover, we set $\omega := u(1 - \zeta)$ and we observe that $\omega = u - u\zeta = u - v$. Since u solves (2.4), from (2.5) we

deduce that, for any $\varphi \in C_0^\infty(\mathcal{O}, \mathbb{R})$,

$$\begin{aligned}
 \int_{\mathcal{O}} f \varphi \, dx &= \int_{\mathcal{O}} \langle \nabla u, \nabla \varphi \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy \\
 (4.37) \quad &= \int_{\mathcal{O}} \langle \nabla v, \nabla \varphi \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy \\
 &\quad + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(\omega(x) - \omega(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy.
 \end{aligned}$$

We now observe that, since $u \in H^1(\mathbb{R}^N)$ and $1 - \zeta$ is smooth on \mathbb{R}^N , both ω and $|\omega|$ are in $H^1(\mathbb{R}^N)$. Moreover, since $\zeta \equiv 1$ on $\mathcal{O}_{\rho/4}$, one has

$$(4.38) \quad \omega = u(1 - \zeta) \equiv 0 \quad \text{on } \mathcal{O}_{\rho/4}.$$

Therefore, we are in the position to apply Lemma 4.4 (with $\alpha := N + 2s > N$, $\delta := \rho/4$, and either $z := \omega$ or $z := |\omega|$), and so, recalling (4.28), we see that

$$\mathcal{I}_{N+2s}[\omega](x) = \int_{\mathbb{R}^N} \frac{\omega(y)}{|x - y|^{N+2s}} \, dy \quad \text{and} \quad \mathcal{I}_{N+2s}[|\omega|](x) = \int_{\mathbb{R}^N} \frac{|\omega(y)|}{|x - y|^{N+2s}} \, dy$$

are (well-posed and) smooth on $\mathcal{O}_{\rho/8}$. In particular, using this fact, and recalling (4.38) we obtain that, for any $\varphi \in C_0^\infty(\mathcal{O}, \mathbb{R})$,

$$\begin{aligned}
 &\int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|\omega(x) - \omega(y)| \cdot |\varphi(x) - \varphi(y)|}{|x - y|^{N+2s}} \, dx \, dy \\
 &\leq 2 \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{|\omega(x) - \omega(y)|}{|x - y|^{N+2s}} \cdot |\varphi(x)| \, dx \, dy \\
 &= 2 \int_{\mathbb{R}^N} \left(\int_{\mathbb{R}^N} \frac{|\omega(x) - \omega(y)|}{|x - y|^{N+2s}} \, dy \right) |\varphi(x)| \, dx \\
 &= 2 \int_{\mathcal{O}} \left(\int_{\mathbb{R}^N} \frac{|\omega(y)|}{|x - y|^{N+2s}} \, dy \right) |\varphi(x)| \, dx \\
 &= 2 \int_{\mathcal{O}} \mathcal{I}_{N+2s}[|\omega|](x) |\varphi(x)| \, dx \\
 &\leq 2|\mathcal{O}| \max_{\overline{\mathcal{O}}} (\mathcal{I}_{N+2s}[|\omega|]) \cdot \max_{\mathbb{R}^N} |\varphi| < \infty,
 \end{aligned}$$

where $|\cdot|$ denotes the standard Lebesgue measure in \mathbb{R}^N . Thanks to the above estimate we can apply Fubini's theorem, thus giving

$$\begin{aligned}
 &\int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(\omega(x) - \omega(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy \\
 &= 2 \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{\omega(x) - \omega(y)}{|x - y|^{N+2s}} \cdot \varphi(x) \, dx \, dy = 2 \int_{\mathcal{O}} \mathcal{I}_{N+2s}[\omega](x) \varphi(x) \, dx.
 \end{aligned}$$

Plugging this information into (4.37), we get

$$\begin{aligned}
 \int_{\mathcal{O}} f \varphi \, dx &= \int_{\mathcal{O}} \langle \nabla v, \nabla \varphi \rangle \, dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} \, dx \, dy \\
 &\quad + c_{N,s} \int_{\mathcal{O}} \mathcal{I}_{N+2s}[\omega](x) \varphi(x) \, dx.
 \end{aligned}$$

In its turn, this identity gives

$$\begin{aligned} & \int_{\mathcal{O}} \langle \nabla v, \nabla \varphi \rangle dx + \frac{c_{N,s}}{2} \int_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(v(x) - v(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dx dy \\ &= \int_{\mathcal{O}} (f - c_{N,s} \mathcal{I}_{N+2s}[\omega]) \varphi dx \quad \text{for all } \varphi \in C_0^\infty(\mathcal{O}, \mathbb{R}), \end{aligned}$$

which shows that v satisfies (4.36) (with $\psi := -c_{N,s} \mathcal{I}_{N+2s}[\omega]$). Finally, since we know that $\mathcal{I}_{N+2s}[\omega]$ is smooth on $\mathcal{O}_{\rho/8}$, we conclude that also ψ is smooth on $\overline{\mathcal{O}}$, and the proof is complete. \square

With Lemma 4.5, we can finally prove Theorem 1.4.

Proof of Theorem 1.4. As in the proof of Proposition 4.3, we proceed by induction on $m \in \mathbb{N} \cup \{0\}$. First of all, the case $m = 0$ (that is, $f \in L^2(\Omega)$) is nothing but Theorem 4.1; we then assume that the desired result holds for a certain integer $m \geq 0$, and we prove that it still holds for $m + 1$.

Let V be a fixed open set satisfying $\overline{V} \subseteq \Omega$. Moreover, let $f \in H^{m+1}(\Omega)$ and let $u \in H^1(\mathbb{R}^N)$ be a solution of (2.4). Since, in particular, $f \in H^m(\Omega)$, from the inductive hypothesis we derive that $u \in H_{\text{loc}}^{m+2}(\Omega)$, and thus

$$(4.39) \quad u \in H^{m+2}(V).$$

We now arbitrarily fix an open set $\mathcal{O} \subseteq \mathbb{R}^N$ such that $\overline{V} \subseteq \mathcal{O}$ and $\overline{\mathcal{O}} \subseteq \Omega$, and we set

$$\rho := \text{dist}(\mathcal{O}, \partial\Omega) > 0.$$

Moreover, we choose a cut-off function $\zeta \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ satisfying (i)-(ii)-(iii) in Lemma 4.5, and we define $v := u \zeta$. On account of (4.39), and since $\text{supp}(\zeta) \subseteq \Omega$, we have

$$v \in H^{m+2}(\mathbb{R}^N).$$

Moreover, from Lemma 4.5 we know that there exists a function $\psi \in C_0^\infty(\overline{\mathcal{O}}, \mathbb{R})$ such that v is a weak solution of the equation

$$\mathcal{L}v = f + \psi \quad \text{in } \mathcal{O}.$$

Since $f \in H^{m+1}(\Omega)$ and ψ is smooth on an open neighborhood of $\overline{\mathcal{O}}$, we derive that $g := f + \psi \in H^{m+1}(\mathcal{O})$. As a consequence, we can apply Proposition 4.3, thus obtaining that

$$v \in H_{\text{loc}}^{m+3}(\mathcal{O}) \text{ and, in particular, } v \in H^{m+3}(V).$$

From this, since $\zeta \equiv 1$ on $\mathcal{O}_{\rho/4} \supset V$, we conclude that

$$v \equiv u \in H^{m+3}(V),$$

and the proof is finally complete. \square

By combining Theorem 1.4 with the well-known Sobolev Embedding theorems, we immediately obtain the Corollary 4.6 below.

Corollary 4.6. *Let $m \in \mathbb{N}$ satisfy $m > N/2$ and let $f \in C^m(\Omega, \mathbb{R})$. Moreover, let $u \in H^1(\mathbb{R}^N)$ a weak solution of (2.4). Then, there exist a non-negative integer $k = k_{m,N}$ and a (unique) function $\hat{u} \in C^k(\Omega, \mathbb{R})$ such that*

$$\hat{u} \equiv u \quad \text{a.e. on } \Omega.$$

More precisely, the number k is explicitly given by

$$(4.40) \quad k = k_{m,N} = \begin{cases} [m - \frac{N}{2}], & \text{if } m - \frac{N}{2} \notin \mathbb{N}, \\ m - \frac{N}{2} - 1 & \text{if } m - \frac{N}{2} \in \mathbb{N}. \end{cases}$$

In particular, if $f \in C^\infty(\Omega, \mathbb{R})$, then $\hat{u} \in C^\infty(\Omega, \mathbb{R})$.

4.2. Boundary regularity and proofs of Theorems 1.5 and 1.6. Now that we have established interior regularity for the weak solutions of (2.4), we focus on the *boundary* regularity for the weak solutions of (2.1).

To begin with, we prove the following theorem.

Theorem 4.7. *Assume that¹ $N \geq 3$, and let $f \in L^p(\Omega)$, with $p > N/2$. Moreover, assume that there exists the weak solution $u_f \in \mathbb{X}(\Omega)$ of (2.1).*

Then, $u_f \in L^\infty(\mathbb{R}^N)$ and

$$(4.41) \quad \|u_f\|_{L^\infty(\mathbb{R}^N)} \leq C \|f\|_{L^p(\Omega)},$$

for some constant $C > 0$.

Remark 4.8. We point out that, on account of Theorem 1.1, a (unique) weak solution of (2.1) exists if $f \in L^p(\Omega)$ with $p \geq 2$.

Proof of Theorem 4.7. The proof employs the classical method by Stampacchia, as extended in the nonlocal setting, see e.g. the proof of Proposition 9 in [45] and of Theorem 2.3 in [26]. We give full details for the reader's convenience.

Let $\delta > 0$ to be conveniently chosen later on. Assuming that u_f is not identically zero (otherwise there is nothing to prove), we set

$$(4.42) \quad \tilde{u} := \frac{\delta u_f}{\|u_f\|_{L^{2^*}(\Omega)} + \|f\|_{L^p(\Omega)}} \quad \text{and} \quad \tilde{f} := \frac{\delta f}{\|u_f\|_{L^{2^*}(\Omega)} + \|f\|_{L^p(\Omega)}},$$

where $2^* := \frac{2N}{N-2}$. In this way, we have that

$$(4.43) \quad \begin{cases} \mathcal{L}\tilde{u} = \tilde{f} & \text{in } \Omega, \\ \tilde{u} = 0 & \text{in } \mathbb{R}^N \setminus \Omega. \end{cases}$$

Also, for every $k \in \mathbb{N}$, we define $C_k := 1 - 2^{-k}$ and

$$v_k := \tilde{u} - C_k, \quad w_k := (v_k)_+ := \max\{v_k, 0\}, \quad U_k := \|w_k\|_{L^{2^*}(\Omega)}^2.$$

We point out that, by the Dominated Convergence Theorem,

$$(4.44) \quad \lim_{k \rightarrow +\infty} U_k = \lim_{k \rightarrow +\infty} \|w_k\|_{L^{2^*}(\Omega)}^2 = \|(\tilde{u} - 1)_+\|_{L^{2^*}(\Omega)}^2.$$

Also, if we take $k := 0$, we see that $w_0 = (v_0)_+ = (\tilde{u} - C_0)_+ = \tilde{u}_+$, and thus

$$(4.45) \quad U_0 = \left(\int_{\Omega} w_0^{2^*}(x) \, dx \right)^{2/2^*} \leq \left(\int_{\Omega} \tilde{u}^{2^*}(x) \, dx \right)^{2/2^*} = \|\tilde{u}\|_{L^{2^*}(\Omega)}^2 \leq \delta^2,$$

which can be taken conveniently small in what follows. In addition, in $\mathbb{R}^N \setminus \Omega$ we have that $v_{k+1} = -C_{k+1} \leq 0$ and thus

$$w_{k+1} = 0.$$

We can then use w_{k+1} as test function and deduce from (4.43) that

$$(4.46) \quad \begin{aligned} & \int_{\Omega} \nabla w_{k+1}(x) \cdot \nabla \tilde{u}(x) \, dx + \iint_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(w_{k+1}(x) - w_{k+1}(y))(\tilde{u}(x) - \tilde{u}(y))}{|x - y|^{N+2s}} \, dx \, dy \\ &= \int_{\Omega} w_{k+1}(x) \tilde{f}(x) \, dx. \end{aligned}$$

We also remark that, for a.e. $x, y \in \mathbb{R}^N$, we have (see, e.g., [45, Lemma 10])

$$(4.47) \quad \begin{aligned} |w_{k+1}(x) - w_{k+1}(y)|^2 &= |(v_{k+1})_+(x) - (v_{k+1})_+(y)|^2 \\ &\leq ((v_{k+1})_+(x) - (v_{k+1})_+(y))(v_{k+1}(x) - v_{k+1}(y)) \\ &= (w_{k+1}(x) - w_{k+1}(y))(\tilde{u}(x) - \tilde{u}(y)). \end{aligned}$$

¹We observe that when $N = 1, 2$ the boundedness of the energy solutions follows directly by Sobolev Embedding.

Moreover,

$$\int_{\Omega} \nabla w_{k+1}(x) \cdot \nabla \tilde{u}(x) \, dx = \int_{\Omega \cap \{\tilde{u} > C_k\}} \nabla v_{k+1}(x) \cdot \nabla \tilde{u}(x) \, dx = \int_{\Omega} |\nabla w_{k+1}(x)|^2 \, dx.$$

From this, (4.46) and (4.47) we conclude that

$$\int_{\Omega} |\nabla w_{k+1}(x)|^2 \, dx \leq \int_{\Omega} w_{k+1}(x) \tilde{f}(x) \, dx.$$

Hence, by Sobolev Inequality,

$$(4.48) \quad U_{k+1} = \left(\int_{\Omega} |w_{k+1}(x)|^{2^*} \, dx \right)^{2/2^*} \leq C \int_{\Omega} |\nabla w_{k+1}(x)|^2 \, dx \leq C \int_{\Omega} w_{k+1}(x) |\tilde{f}(x)| \, dx,$$

for some $C > 0$. Also, $v_{k+1} \leq v_k$ and therefore

$$(4.49) \quad w_{k+1} \leq w_k.$$

Moreover, we observe that

$$w_k = (\tilde{u} - C_k)_+ = \left(\tilde{u} - C_{k+1} + \frac{1}{2^{k+1}} \right)_+ = \left(v_{k+1} + \frac{1}{2^{k+1}} \right)_+,$$

and, as a result,

$$(4.50) \quad \{w_{k+1} > 0\} = \{v_{k+1} > 0\} \subseteq \left\{ w_k > \frac{1}{2^{k+1}} \right\}.$$

We also observe that

$$2^* - \frac{2^*}{p} - 1 > 2^* - \frac{2^*}{N/2} - 1 = \frac{2N}{N-2} - \frac{4}{N-2} - 1 = 1.$$

Hence, we can define

$$(4.51) \quad q := 2^* \left(2^* - \frac{2^*}{p} - 1 \right)^{-1} < 2^*.$$

We observe that

$$q > \frac{2^*}{2^* - 1} > 1.$$

In addition,

$$(4.52) \quad \frac{1}{2^*} + \frac{1}{p} + \frac{1}{q} = 1.$$

From this, (4.49) and (4.50), using the Hölder Inequality with exponents 2^* , p and q , we deduce that

$$\begin{aligned} \int_{\Omega} w_{k+1}(x) |\tilde{f}(x)| \, dx &= \int_{\Omega \cap \{w_{k+1} > 0\}} w_{k+1}(x) |\tilde{f}(x)| \, dx \\ &\leq \|\tilde{f}\|_{L^p(\Omega)} \|w_{k+1}\|_{L^{2^*}(\Omega)} |\Omega \cap \{w_{k+1} > 0\}|^{1/q} \\ &\leq \|w_k\|_{L^{2^*}(\Omega)} \left| \Omega \cap \left\{ w_k > \frac{1}{2^{k+1}} \right\} \right|^{1/q} \\ &\leq U_k^{1/2} \left(2^{2^*(k+1)} \int_{\Omega \cap \{w_k > \frac{1}{2^{k+1}}\}} w_k^{2^*} \right)^{1/q} \\ &\leq \tilde{C}^k U_k^{1/2} U_k^{2^*/(2q)}, \end{aligned} \quad (4.53)$$

for some $\tilde{C} > 1$. We now define

$$\beta := \frac{1}{2} + \frac{2^*}{2q},$$

and we stress that

$$(4.54) \quad \beta > 1,$$

thanks to (4.51). Using this notation, we deduce from (4.48) and (4.53) that

$$U_{k+1} \leq \hat{C}^k U_k^\beta,$$

for some $\hat{C} > 1$. As a result, recalling (4.45) (and supposing $\delta > 0$ appropriately small), we conclude that

$$\lim_{k \rightarrow +\infty} U_k = 0.$$

This and (4.44) give that

$$\|(\tilde{u} - 1)_+\|_{L^{2^*}(\Omega)}^2 = 0,$$

and therefore $\tilde{u} \leq 1$. As a consequence, recalling (4.42), for every $x \in \Omega$,

$$(4.55) \quad u_f(x) \leq \frac{\|u_f\|_{L^{2^*}(\Omega)} + \|f\|_{L^p(\Omega)}}{\delta}.$$

On the other hand, by testing the equation against $u_f|_\Omega$ (see Remark 2.2), and recalling once again relation (4.52), we see that

$$\begin{aligned} \int_\Omega |\nabla u_f(x)|^2 dx &\leq \int_\Omega |\nabla u_f(x)|^2 dx + \iint_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u_f(x) - u_f(y))^2}{|x - y|^{N+2s}} dx dy \\ &= \int_\Omega u_f(x) f(x) dx \\ &\leq \|u_f\|_{L^{2^*}(\Omega)} \|f\|_{L^p(\Omega)} |\Omega|^{1/q}. \end{aligned}$$

This and the Sobolev Inequality give that

$$\|u_f\|_{L^{2^*}(\Omega)} \leq \bar{C} \|f\|_{L^p(\Omega)},$$

for a suitable $\bar{C} > 0$. Combining this with (4.55) we obtain (4.41), as desired. \square

With Theorem 4.7 at hand, we now focus on the proof of Theorem 1.5. We remark that by saying that a function $u \in H^1(\mathbb{R}^N)$ satisfies (1.6) we mean, precisely, that $u \leq \bar{C}$ a.e. in \mathbb{R}^N , $u \equiv 0$ a.e. in $\mathbb{R}^N \setminus \Omega$, and

$$\int_\Omega \langle \nabla u, \nabla \varphi \rangle dx + \frac{c_{N,s}}{2} \iint_{\mathbb{R}^N \times \mathbb{R}^N} \frac{(u(x) - u(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dx dy \leq \bar{C} \int_\Omega \varphi dx,$$

for every *non-negative* function $\varphi \in C_0^\infty(\Omega, \mathbb{R})$.

The proof of Theorem 1.5 relies on the construction of an appropriate barrier, which will be built by recursive corrections of monomial functions. In fact, the arguments provided have wider applicability and can be exploited in more general contexts as well, but for concreteness we will follow on the specific operator, boundary conditions and geometry dealt with in this article.

The first step towards the proof of Theorem 1.5 consists in an elementary computation on functions which have a convex portion in their graphs.

Lemma 4.9. *Let $s \in (0, 1)$, $d > \ell > 0$ and $v : \mathbb{R} \rightarrow \mathbb{R}$. Assume that*

$$v \in C^{1,1}((-\infty, d)) \cap L^\infty(\mathbb{R})$$

and that v is convex in $(-\infty, d)$. Then, for every $x \in (0, \ell)$,

$$(-\Delta)^s v(x) \leq \frac{2c_{1,s} \|v\|_{L^\infty(\mathbb{R})}}{s(d - \ell)^{2s}}.$$

Proof. In the principal value sense, we have that, for every $x \in (0, \ell)$,

$$\int_{2x-d}^d \frac{v(y) - v(x)}{|x - y|^{1+2s}} dy = \int_{x-d}^{d-x} \frac{v(x+z) - v(x)}{|z|^{1+2s}} dz = \int_{x-d}^{d-x} \frac{v(x+z) - v(x) - v'(x)z}{|z|^{1+2s}} dz \geq 0,$$

thanks to the convexity assumption. As a result,

$$\frac{(-\Delta)^s v(x)}{c_{1,s}} \leq \int_{\mathbb{R} \setminus (2x-d, d)} \frac{v(x) - v(y)}{|x - y|^{1+2s}} dy = \int_{\mathbb{R} \setminus (x-d, d-x)} \frac{v(x) - v(x+z)}{|z|^{1+2s}} dz.$$

We stress that, in the latter integral, we have that $|z| \geq d - x \geq d - \ell$. Hence,

$$\frac{(-\Delta)^s v(x)}{c_{1,s}} \leq 2\|v\|_{L^\infty(\mathbb{R})} \int_{\{|z| \geq d-\ell\}} \frac{dz}{|z|^{1+2s}} = \frac{2\|v\|_{L^\infty(\mathbb{R})}}{s(d-\ell)^{2s}},$$

and the proof is complete. \square

The next auxiliary result for the proof of Theorem 1.5 focuses on a calculation for a modified monomial function.

Lemma 4.10. *Let $s \in (0, 1)$, $L > 0$ and $\alpha \geq 2s$. Let also*

$$(4.56) \quad w_\alpha(x) := \begin{cases} x_+^\alpha & \text{if } x < 2L, \\ (2L)^\alpha & \text{if } x \geq 2L. \end{cases}$$

Then, there exists $C > 0$, only depending on L , s and α , such that, for all $x \in (0, L)$,

$$(4.57) \quad |(-\Delta)^s w_\alpha(x)| \leq \begin{cases} C & \text{if } \alpha > 2s, \\ C(1 + |\log x|) & \text{if } \alpha = 2s. \end{cases}$$

Proof. Up to scaling, it is not restrictive to suppose that $L := 1$. Given any point $x \in (0, 1)$, we use the substitution $z := y/x$ to see that

$$\frac{(-\Delta)^s w_\alpha(x)}{c_{1,s}} = \int_{\mathbb{R}} \frac{x^\alpha - \min\{y_+^\alpha, 2^\alpha\}}{|x - y|^{1+2s}} dy = x^{\alpha-2s} \int_{\mathbb{R}} \frac{1 - \min\{z_+^\alpha, (2/x)^\alpha\}}{|1 - z|^{1+2s}} dz,$$

where the principal value notation has been omitted for the sake of shortness. Then, we observe that

$$A_1 := \left| \int_{-\infty}^0 \frac{1 - \min\{z_+^\alpha, (2/x)^\alpha\}}{|1 - z|^{1+2s}} dz \right| = \int_{-\infty}^0 \frac{dz}{(1 - z)^{1+2s}} = \int_1^{+\infty} \frac{dt}{t^{1+2s}} = \frac{1}{2s}.$$

Similarly,

$$\begin{aligned} A_2 &:= \left| \int_2^{+\infty} \frac{1 - \min\{z_+^\alpha, (2/x)^\alpha\}}{|1 - z|^{1+2s}} dz \right| \leq \int_2^{+\infty} \frac{1 + \min\{z_+^\alpha, (2/x)^\alpha\}}{(z - 1)^{1+2s}} dz \\ &= \int_1^{+\infty} \frac{dt}{t^{1+2s}} + \int_2^{2/x} \frac{z^\alpha}{(z - 1)^{1+2s}} dz + \left(\frac{2}{x}\right)^\alpha \int_{2/x}^{+\infty} \frac{dz}{(z - 1)^{1+2s}} dz \\ &\leq C_1(1 + x^{2s-\alpha}\ell(x)), \end{aligned}$$

for some $C_1 > 0$, where

$$\ell(x) := \begin{cases} 1 + |\log x| & \text{if } \alpha = 2s, \\ 1 & \text{otherwise.} \end{cases}$$

In addition, using the principal value notation,

$$\begin{aligned} A_3 &:= \left| \int_0^2 \frac{1 - \min\{z_+^\alpha, (2/x)^\alpha\}}{|1 - z|^{1+2s}} dz \right| = \left| \int_0^2 \frac{1 - z^\alpha}{|1 - z|^{1+2s}} dz \right| \\ &= \left| \int_{-1}^1 \frac{1 - (1+t)^\alpha}{|t|^{1+2s}} dt \right| = \left| \int_{-1}^1 \frac{(1+t)^\alpha - 1 - \alpha t}{|t|^{1+2s}} dt \right| \leq C_2, \end{aligned}$$

for some $C_2 > 0$. All in all, we find that

$$\frac{|(-\Delta)^s w_\alpha(x)|}{c_{1,s}} \leq x^{\alpha-2s}(A_1 + A_2 + A_3) \leq A_1 + A_3 + x^{\alpha-2s}A_2 \leq C_3(1 + \ell(x)),$$

for some $C_3 > 0$, yielding the desired result. \square

From Lemma 4.10, we obtain the following barrier.

Lemma 4.11. *Let $s \in (0, 1)$. There exist a number $d > 0$ and a function $\beta \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, d), \mathbb{R})$ satisfying the following properties:*

- there exists $C_0 > 0$ such that for all $x \geq d$

$$(4.58) \quad \beta(x) \geq C_0;$$

- for all $x \leq 0$,

$$(4.59) \quad \beta(x) = 0;$$

- there exists $C_1 \geq 1$, independent of d , such that

$$(4.60) \quad \frac{x}{C_1} \leq \beta(x) \leq C_1 x \quad \text{for all } x \in (0, d);$$

- there exists $C_2 > 0$, independent of d , such that

$$(4.61) \quad \mathcal{L}\beta(x) \geq -C_2 \quad \text{for all } x \in (0, d).$$

Furthermore, $\beta \in H_{\text{loc}}^1(\mathbb{R})$.

Proof. We distinguish two cases, according to the value of s .

CASE I: $s \in (1/2, 1)$. In this case, we let

$$\rho(s) := \frac{2s-1}{2(1-s)} \quad \text{and} \quad J := \begin{cases} [\rho(s)] & \text{if } \rho(s) \notin \mathbb{N}, \\ \rho(s) - 1 & \text{otherwise.} \end{cases}$$

Also, for each $j \in \mathbb{N}$ with $0 \leq j \leq J+1$, we set $\alpha_j := 1 + 2j(1-s)$. We observe that, for all $j \in \{0, \dots, J\}$, we have

$$\alpha_j \leq 1 + 2J(1-s) < 1 + (2s-1) = 2s.$$

Therefore, for all $j \in \{0, \dots, J\}$, we can define $(-\Delta)^s x_+^{\alpha_j}$ and, by homogeneity, we see that, for all $x > 0$,

$$(-\Delta)^s x_+^{\alpha_j} = \kappa_j x_+^{\alpha_j-2s},$$

for a suitable $\kappa_j \in \mathbb{R}$. As a matter of fact, since $\alpha_j \geq 1$, we have that $x_+^{\alpha_j}$ is a convex function and therefore $(-\Delta)^s x_+^{\alpha_j} < 0$ in $(0, +\infty)$. From this, we get

$$(4.62) \quad \kappa_j < 0 \quad \text{for every } j = 0, \dots, J.$$

We also point out that, for every $j \in \{0, \dots, J\}$,

$$\begin{aligned} \alpha_j - 2s &= \alpha_j - 2 + 2(1-s) = 1 + 2j(1-s) - 2 + 2(1-s) \\ &= 1 + 2(j+1)(1-s) - 2 = \alpha_{j+1} - 2. \end{aligned}$$

Now we define $\{c_0, \dots, c_{J+1}\}$ as follows. We let $c_0 := 1$, and then, recursively, for every index $j \in \{1, \dots, J+1\}$,

$$(4.63) \quad c_j := -\frac{\kappa_{j-1} c_{j-1}}{\alpha_j (\alpha_j - 1)}.$$

We stress that this definition is well posed, since, if $j \in \{1, \dots, J\}$,

$$\alpha_j \geq 1 + 2(1-s) > 1.$$

From this and (4.62), it follows that

$$(4.64) \quad c_j > 0 \quad \text{for every } j = 0, \dots, J+1.$$

Hence, we consider the function $\tilde{\beta} : \mathbb{R} \rightarrow \mathbb{R}$ defined as follows

$$(4.65) \quad \tilde{\beta}(x) := \sum_{j=0}^J c_j x_+^{\alpha_j}.$$

Since $\alpha_j < 2s$ for every $0 \leq j \leq J$, it is easy to recognize that

- $\tilde{\beta}(x) \geq 0$ for all $x \in \mathbb{R}$;
- $\tilde{\beta} \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, \infty), \mathbb{R})$ and $\tilde{\beta} \in H_{\text{loc}}^1(\mathbb{R})$.

Moreover, for every $x > 0$ we have

$$\begin{aligned} \mathcal{L}\tilde{\beta}(x) &= - \sum_{j=0}^J c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + \sum_{j=0}^J c_j \kappa_j x^{\alpha_j - 2s} \\ &= - \sum_{j=1}^J c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + \sum_{j=1}^J c_{j-1} \kappa_{j-1} x^{\alpha_{j-1} - 2} + c_J \kappa_J x^{\alpha_J - 2s} \\ &= -2 \sum_{j=1}^J c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + c_J \kappa_J x^{\alpha_J - 2s}, \end{aligned}$$

where (4.63) was used in the latter line. As a consequence, taking $d \in (0, 1)$ to be chosen conveniently small in what follows, employing the notation in (4.56) with $L := 1$, and introducing the function

$$(4.66) \quad \beta_{\sharp} := \tilde{\beta} + c_{J+1} w_{\alpha_{J+1}},$$

we obtain that β_{\sharp} satisfies the following properties:

- $\beta_{\sharp} \geq 0$ on \mathbb{R} (as $c_{J+1} > 0$ and $w_{\alpha_{J+1}} \geq 0$ on \mathbb{R});
- $\beta_{\sharp} \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, d), \mathbb{R})$ and $\beta_{\sharp} \in H_{\text{loc}}^1(\mathbb{R})$.

Furthermore, if $x \in (0, d)$, we get

$$\begin{aligned} \mathcal{L}\beta_{\sharp}(x) &= -2 \sum_{j=1}^J c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + c_J \kappa_J x^{\alpha_J - 2s} - c_{J+1} \Delta w_{\alpha_{J+1}}(x) + c_{J+1} (-\Delta)^s w_{\alpha_{J+1}}(x) \\ &= -2 \sum_{j=1}^J c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + c_J \kappa_J x^{\alpha_J - 2s} - c_{J+1} \alpha_{J+1} (\alpha_{J+1} - 1) x^{\alpha_{J+1} - 2} + c_{J+1} (-\Delta)^s w_{\alpha_{J+1}}(x) \\ &= -2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} + c_{J+1} (-\Delta)^s w_{\alpha_{J+1}}(x), \end{aligned}$$

where (4.63) was used once again. From this and (4.57), we obtain that,

$$(4.67) \quad \mathcal{L}\beta_{\sharp}(x) \geq -C_{\sharp} |\log x| - 2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} = C_{\sharp} \log x - 2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2},$$

for all $x \in (0, d)$ and for some $C_{\sharp} > 0$. Now, we let

$$\tilde{W}(x) := \frac{x_+^2}{4} (3 - 2 \log x)_+ + \frac{2}{C_{\sharp}} \sum_{j=1}^{J+1} c_j x_+^{\alpha_j} \quad \text{and} \quad S(d) := \max_{(-\infty, d]} \tilde{W}.$$

Notice that

$$\lim_{d \searrow 0} \frac{S(d)}{d} = 0.$$

As a result, by possibly shrinking $d \in (0, 1)$, we can additionally suppose that

$$(4.68) \quad S(d) \leq \frac{d}{4C_{\sharp}}.$$

Then, we take a continuous function

$$(4.69) \quad W : \mathbb{R} \rightarrow [0, 2S(d)]$$

satisfying the following properties:

- (i) $W(x) = \tilde{W}(x)$ for all $x \leq d$;
- (ii) $W(x) = 0$ for all $x \geq 2d$;
- (iii) $W \in C^\infty((0, +\infty), \mathbb{R})$.

We define

$$\beta(x) := \beta_\sharp(x) - C_\sharp W(x).$$

Notice that, by the regularity of β_\sharp and W , we have that

$$\beta \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, d), \mathbb{R}) \text{ and } \beta \in H_{\text{loc}}^1(\mathbb{R}).$$

Moreover, if $d > 0$ is sufficiently small,

$$(4.70) \quad W(x) = \frac{x^2}{4}(3 - 2\log x) + \frac{2}{C_\sharp} \sum_{j=1}^{J+1} c_j x^{\alpha_j} \quad \text{for all } x \in (0, d),$$

and therefore

$$W''(x) = -\log x + \frac{2}{C_\sharp} \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2}.$$

In addition, for small $d > 0$, we have that

$$W \in C^{1,a}(-1, 2d) \text{ for some } a \in (0, 1),$$

and thus $(-\Delta)^s W \in C^{0,1+a-2s}(-1, 2d)$. As a consequence (see, e.g., [46, Propositions 2.1.7 and 2.1.8]), we obtain that, for all $x \in (0, d)$,

$$|(-\Delta)^s W(x)| \leq \tilde{C},$$

for a suitable $\tilde{C} > 0$ only depending on N, s and a (hence, in particular, \tilde{C} is independent of d). Then, we deduce from (4.67) that, for every $x \in (0, d)$,

$$\begin{aligned} \mathcal{L}\beta(x) &\geq C_\sharp \log x - 2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} - C_\sharp \mathcal{L}W(x) \\ &= C_\sharp \log x - 2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} - C_\sharp (\log x + (-\Delta)^s W(x)) + 2 \sum_{j=1}^{J+1} c_j \alpha_j (\alpha_j - 1) x^{\alpha_j - 2} \\ &= -C_\sharp (-\Delta)^s W(x) \\ &\geq -C_\sharp \tilde{C}. \end{aligned}$$

This proves (4.61). Also, (4.59) is obvious. To proceed further we observe that, for every $x \in \mathbb{R}$, we have the estimate

$$(4.71) \quad \sum_{j=0}^J c_j x_+^{\alpha_j} + c_{J+1} w_{\alpha_{J+1}}(x) \geq c_0 x_+^{\alpha_0} = x_+ \geq \min\{x_+, d\}.$$

As a result, since $\alpha_0 = 1 < \alpha_1 \leq \alpha_j$ for all $j \in \{2, \dots, J+1\}$, we have that, for every $x \in \mathbb{R}$,

$$(4.72) \quad \begin{aligned} &\beta(x) + C_\sharp W(x) \\ &= \sum_{j=0}^J c_j x_+^{\alpha_j} + c_{J+1} w_{\alpha_{J+1}} \in \left[\min\{x_+, d\}, \bar{C} \max\{x_+, x_+^{\alpha_{J+1}}\} \right], \end{aligned}$$

for a suitable $\bar{C} > 0$. From this, we conclude that

$$(4.73) \quad \beta(x) \geq \min\{x_+, d\} - C_\sharp W(x).$$

Now, if $x \geq d$, we obtain from (4.68), (4.69) and (4.73) that

$$(4.74) \quad \beta(x) \geq d - C_{\#}W(x) \geq d - 2C_{\#}S(d) \geq \frac{d}{2}.$$

If instead $x \in (-\infty, d)$, we deduce from (4.73) that

$$(4.75) \quad \beta(x) \geq x_+ - \frac{C_{\#}x_+^2}{4}(3 - 2\log x)_+ - 2 \sum_{j=1}^{J+1} c_j x_+^{\alpha_j} \geq \frac{x_+}{2},$$

by possibly redefining $d > 0$ in a conveniently small way.

We notice that (4.74) implies (4.58), as desired. In addition, (4.75) proves the first inequality in (4.60). Besides, from (4.70) and (4.72) we obtain the second inequality in (4.60), thus concluding the proof in the case $s > 1/2$.

CASE II: $s \in (0, 1/2]$. This case is indeed simpler: it suffices to rerun the preceding argument starting from (4.66), taking $\tilde{\beta} \equiv 0$ and

$$\beta_{\#} := w_1 = \begin{cases} x_+ & \text{if } x < 2, \\ 2 & \text{if } x \geq 2. \end{cases}$$

We omit any further detail. □

Corollary 4.12. *Let $s \in (0, 1)$ be fixed. There exist a number $\ell > 0$ and a non-negative function $\gamma \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, \ell), \mathbb{R})$ such that*

- for all $x \leq 0$,

$$(4.76) \quad \gamma(x) = 0;$$

- there exists $c \in (0, 1)$ such that for all $x \in (0, \ell)$

$$(4.77) \quad cx \leq \gamma(x) \leq \frac{x}{c};$$

- for all $x \geq \ell$,

$$(4.78) \quad \gamma(x) \geq 1;$$

- for all $x \in (0, \ell)$,

$$(4.79) \quad \mathcal{L}\gamma(x) \geq 1.$$

Furthermore, $\gamma \in H_{\text{loc}}^1(\mathbb{R})$.

Proof. We let β as in Lemma 4.11 and, in the notation of Lemma 4.11, we will choose $\ell \in (0, d/2)$ to be suitably small. Given $M > 0$, we define

$$\beta_*(x) := \begin{cases} 0 & \text{if } x \leq 0, \\ C_2x^2 & \text{if } x \in (0, \ell), \\ 2C_2\ell x - C_2\ell^2 & \text{if } x \in [\ell, d], \\ C_2\ell(2d - \ell) & \text{if } x \in (d, +\infty), \end{cases} \quad \text{and} \quad \gamma(x) := M(\beta(x) - \beta_*(x)),$$

where $C_2 > 0$ is as in Lemma 4.11. We point out that, by taking into account the regularity of β and β_* , it is very easy to see that

$$\gamma \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, \ell), \mathbb{R}) \text{ and } \gamma \in H_{\text{loc}}^1(\mathbb{R}).$$

Moreover, since β_* is convex in $(-\infty, d)$, by Lemma 4.9 there exists a suitable constant $C_3 > 0$ such that, for every $x \in (0, \ell)$, we have

$$(-\Delta)^s \beta_* \leq \frac{C_3\ell(2d - \ell)}{(d - \ell)^{2s}}.$$

As a consequence, as long as d and ℓ are sufficiently small we have

$$(-\Delta)^s \beta_* \leq \frac{C_3\ell(2d - \ell)}{(d - \ell)^{2s}} \leq \frac{2C_3\ell d}{(d/2)^{2s}} \leq 2^{2s+1}C_3\ell d^{1-2s} \leq 2^{2s}C_3d^{2-2s} \leq \frac{C_2}{2},$$

Therefore, for all $x \in (0, \ell)$,

$$\frac{1}{M} \mathcal{L}\gamma(x) \geq \mathcal{L}\beta(x) + \Delta\beta_*(x) - \frac{C_2}{2} \geq -C_2 + 2C_2 - \frac{C_2}{2} = \frac{C_2}{2},$$

thanks to (4.61). By choosing $M \geq \frac{2}{C_2}$, we obtain (4.79), as desired. Moreover, (4.76) follows from (4.59). In addition, by (4.60), and taking

$$\ell \leq \frac{1}{2C_1C_2},$$

if $x \in (0, \ell)$ we obtain

$$\gamma(x) \geq M \left(\frac{x}{C_1} - C_2x^2 \right) \geq M \left(\frac{1}{C_1} - C_2\ell \right) x \geq \frac{Mx}{2C_1}.$$

Similarly, recalling (4.60), we have

$$\gamma(x) \leq MC_1x.$$

These considerations imply (4.77). Furthermore, exploiting (4.58) and (4.60), we see that, for every $x \geq \ell$, one has

$$(4.80) \quad \beta(x) \geq \min \left\{ C_0, \frac{\ell}{C_1} \right\} = \frac{\ell}{C_1},$$

as long as ℓ is sufficiently small. Moreover, if $x \in [\ell, d]$,

$$(4.81) \quad \beta_*(x) \leq 2C_2\ell x \leq 2C_2\ell d.$$

Similarly, if $x > d$,

$$\beta_*(x) \leq 2C_2\ell d.$$

This and (4.81) give that, for all $x \geq \ell$,

$$\beta_*(x) \leq 2C_2\ell d \leq \frac{\ell}{2C_1},$$

provided that d is chosen sufficiently small. From this and (4.80), we get

$$\gamma(x) \geq M \left(\frac{\ell}{C_1} - \frac{\ell}{2C_1} \right) = \frac{M\ell}{2C_1} \geq 1,$$

as long as $M \geq \frac{2C_1}{\ell}$. This gives (4.78), as desired. \square

The function γ constructed in Corollary 4.12 would provide a ‘good’ barrier for the proof of Theorem 1.5 if, in addition, $\gamma \in H^1(\mathbb{R})$. In fact, since we aim to apply the weak maximum principle in Theorem 1.2 to the function

$$\pm u - M\gamma$$

(for a suitable $M \in \mathbb{R}$), and since $u \in H^1(\mathbb{R}^N)$, it is crucial to have

$$(4.82) \quad \gamma \in H^1(\mathbb{R}).$$

On the other hand, since property (4.78) shows that (4.82) cannot hold, we need to perform a suitable truncation argument. This is described in the next lemma.

Lemma 4.13. *Let $\mathcal{O} \subseteq \mathbb{R}^N$ be a bounded open set, and let $\rho > 0$ be such that $\mathcal{O} \subseteq B(0, \rho)$. Moreover, let $f \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\mathcal{O}, \mathbb{R})$. Finally, let $R > 4\rho$ and let $\varphi = \varphi_R \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ satisfy the properties*

- (i) $\varphi \equiv 1$ on $B(0, R)$;
- (ii) $0 \leq \varphi \leq 1$ on \mathbb{R}^N ;
- (iii) $\text{supp}(\varphi) \subseteq B(0, 2R)$.

Then, setting $f^* := f\varphi$, for all $x \in \mathcal{O}$ we have

$$|\mathcal{L}f^*(x) - \mathcal{L}f(x)| \leq C \left(\frac{|f(x)|}{R^{2s}} + \kappa(R) \right),$$

where $C > 0$ is a constant only depending on N, s and ρ , and

$$\kappa(R) := \int_{\{|y| \geq R\}} \frac{|f(y)|}{1 + |y|^{N+2s}} dy.$$

Proof. We first observe that, since $R > 4\rho$, we have

$$R - \rho \geq \frac{R}{2} + \rho \geq \frac{R}{2}.$$

As a consequence, if $x \in \mathcal{O} \subseteq B(0, \rho)$ and $y \in \mathbb{R}^N \setminus B(0, R)$, we have

$$|x - y| \geq |y| - |x| \geq R - \rho \geq \frac{R}{2}.$$

Moreover, one also has

$$|x - y| \geq |y| - |x| \geq \frac{|y|}{2} + \frac{R}{2} - \rho \geq \frac{\rho + |y|}{2} \geq c_\rho \cdot \frac{1 + |y|}{2},$$

with $c_\rho := \min\{\rho, 1\}$. Since $\varphi = 1$ in $B(0, R) \supseteq \mathcal{O}$, for every $x \in \mathcal{O}$ we then get

$$\begin{aligned} & \left| \int_{\mathbb{R}^N} \frac{(f(x) - f(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dy \right| \\ &= \left| \int_{\{|y| \geq R\}} \frac{(f(x) - f(y))(1 - \varphi(y))}{|x - y|^{N+2s}} dy \right| \leq \int_{\{|y| \geq R\}} \frac{|f(x) - f(y)|}{|x - y|^{N+2s}} dy \\ (4.83) \quad & \leq |f(x)| \int_{\{|y| \geq R\}} \frac{dy}{|x - y|^{N+2s}} + \int_{\{|y| \geq R\}} \frac{|f(y)|}{|x - y|^{N+2s}} dy \\ & \leq |f(x)| \int_{\{|z| \geq R/2\}} \frac{dz}{|z|^{N+2s}} + (2/c_\rho)^{N+2s} \int_{\{|y| \geq R\}} \frac{|f(y)|}{(1 + |y|)^{N+2s}} dy \\ & \leq C \left(\frac{|f(x)|}{R^{2s}} + \kappa(R) \right), \end{aligned}$$

for some $C = C(N, s, \rho) > 0$. Similarly, for every $x \in \mathcal{O}$ we have

$$\begin{aligned} & |(-\Delta)^s \varphi(x)| = c_{N,s} \left| \int_{\{|y| \geq R\}} \frac{1 - \varphi(y)}{|x - y|^{N+2s}} dy \right| \\ (4.84) \quad & \leq c_{N,s} \int_{\{|y| \geq R\}} \frac{dy}{|x - y|^{N+2s}} \leq c_{N,s} \int_{\{|z| \geq R/2\}} \frac{dz}{|z|^{N+2s}} \leq \frac{C}{R^{2s}}, \end{aligned}$$

up to renaming $C > 0$. Since, obviously, $\Delta f^* = \Delta f$ on \mathcal{O} (remind that $\varphi \equiv 1$ on $B(0, R) \subseteq \mathcal{O}$), if $x \in \mathcal{O}$ we obtain (see, e.g., [7, formula (2.11)])

$$\begin{aligned}
& |\mathcal{L}f^*(x) - \mathcal{L}f(x)| \\
&= \left| (-\Delta)^s f^*(x) - (-\Delta)^s f(x) \right| \\
&= \left| \varphi(x)(-\Delta)^s f(x) + f(x)(-\Delta)^s \varphi(x) \right. \\
(4.85) \quad & \quad \left. - c_{N,s} \int_{\mathbb{R}^N} \frac{(f(x) - f(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dy - (-\Delta)^s f(x) \right| \\
&= \left| f(x)(-\Delta)^s \varphi(x) - c_{N,s} \int_{\mathbb{R}^N} \frac{(f(x) - f(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dy \right| \\
&\leq |f(x)| |(-\Delta)^s \varphi(x)| + c_{N,s} \left| \int_{\mathbb{R}^N} \frac{(f(x) - f(y))(\varphi(x) - \varphi(y))}{|x - y|^{N+2s}} dy \right|.
\end{aligned}$$

The desired result thus follows by inserting (4.83) and (4.84) into (4.85). \square

We are finally ready to prove Theorem 1.5.

Proof of Theorem 1.5. The gist is that the function γ belongs to the space $\mathcal{C}_s(\mathbb{R}^N)$ defined in (2.9) (because $\gamma(x_1)$ grows like $x_1^{\alpha_J}$ with $\alpha_J < 2s$ as $x_1 \rightarrow +\infty$, the highest exponent coming from (4.65)) and therefore the corresponding quantity $\kappa(R)$ introduced in Lemma 4.13 goes to zero as $R \rightarrow +\infty$. As a general philosophy, the main difficulty with all of the truncations in the fractional setting is to control the errors developed by nonlocality: in this argument these errors are accounted for by the quantity $\kappa(R)$, which becomes negligible for R large. In this sense, taking care of the fact that the function γ grows slower than x_1^{2s} once we leave the “boundary layer” $\{x_1 < \ell\}$ pays off now, since it allows us to have full control of the “contributions coming from infinity”.

The technical details of the proof go as follows. Up to a rigid motion, we can assume that $p = 0$ and that

$$(4.86) \quad \Omega \subset \{x \in \mathbb{R}^N : x_1 > 0\}.$$

Moreover, if $\ell > 0$ is as in Corollary 4.12, we define $\Omega_* := \Omega \cap \{x_1 < \ell\}$. Notice that, on account of (4.86), we have

$$(4.87) \quad \Omega_* \subset \{x \in \mathbb{R}^N : x_1 \in (0, \ell)\}.$$

We now choose $\rho > 0$ in such a way that $\Omega \subseteq B(0, \rho)$, and we let $R > 0$ be such that $R > 4\rho$. Moreover, if γ is as in Corollary 4.12, we define

$$\Theta(x) := \gamma(x_1) \cdot \varphi(x) \quad \text{for any } x \in \mathbb{R}^N,$$

where $\varphi = \varphi_R \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ satisfies (i)-(ii)-(iii) in the statement of Lemma 4.13. Taking into account that $\gamma \in \mathcal{C}_s(\mathbb{R}) \cap C^2((0, \ell), \mathbb{R})$, $\varphi \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ and $\varphi \equiv 1$ on the ball $B(0, R) \supseteq \Omega_*$, it is readily seen that

$$\Theta \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega_*, \mathbb{R}).$$

Moreover, since $\gamma \in H_{\text{loc}}^1(\mathbb{R})$ and $\text{supp}(\varphi) \subseteq B(0, 2R)$, we also have

$$\Theta \in H^1(\mathbb{R}^N).$$

Finally, by combining (4.79), (4.77) and Lemma 4.13, we obtain

$$\mathcal{L}\Theta(x) \geq \mathcal{L}\gamma(x_1) - C \left(\frac{|\gamma(x_1)|}{R^{2s}} + \kappa(R) \right) \geq 1 - C \left(\frac{\ell}{c} \cdot \frac{1}{R^{2s}} + \kappa(R) \right)$$

for every $x \in \Omega_* \subseteq \{x \in \mathbb{R}^N : x_1 \in (0, \ell)\}$. In view of this last computation, by enlarging $R > 0$ if necessary, we get

$$(4.88) \quad \mathcal{L}\Theta(x) \geq \frac{1}{2} \quad \text{for every } x \in \Omega_*.$$

We then turn to use $\Theta \in H^1(\mathbb{R}^N)$ as a barrier to prove (1.7). To this end, we consider the function $v : \mathbb{R}^N \rightarrow \mathbb{R}$ defined as

$$v(x) := u(x) - 4\bar{C}\Theta(x),$$

where \bar{C} is as in (1.6). We observe that, by (4.88), we have

$$(4.89) \quad \mathcal{L}v = \mathcal{L}u - 4\bar{C}\mathcal{L}\Theta \leq \bar{C} - \frac{4\bar{C}}{2} \leq 0$$

in Ω_* . Now we claim that, for a.e. $x \in \mathbb{R}^N \setminus \Omega_*$, we have

$$(4.90) \quad v(x) \leq 0.$$

To check this, we observe that

$$\mathbb{R}^N \setminus \Omega_* \subseteq (\mathbb{R}^N \setminus \Omega) \cup \{x \in \Omega : x_1 \geq \ell\}.$$

Hence we distinguish two cases.

- $x \in \mathbb{R}^N \setminus \Omega$. In this case, since $u \equiv 0$ a.e. in $\mathbb{R}^N \setminus \Omega$, we have

$$v(x) = -4\bar{C}\Theta(x) = -4\bar{C}(\gamma(x_1) \cdot \varphi(x));$$

from this, since $\gamma \geq 0$ on \mathbb{R} (by Corollary 4.12) and $0 \leq \varphi \leq 1$ (see (ii) in Lemma 4.13), we derive that $v \leq 0$ a.e. in $\mathbb{R}^N \setminus \Omega$.

- $x \in \{x \in \Omega : x_1 \geq \ell\}$. In this case, using (4.78) and the fact that $\varphi \equiv 1$ on $B(0, R) \supseteq \Omega$, we can write

$$v(x) = u(x) - 4\bar{C}\gamma(x_1) \leq u(x) - 4\bar{C} \leq \bar{C} - 4\bar{C} \leq 0,$$

and this concludes the proof of (4.90).

From (4.89), (4.90) and Theorem 1.2 (notice that $v \in H^1(\mathbb{R}^N)$, since the same is true of both u and Θ), we conclude that $v(x) \leq 0$ for a.e. $x \in \mathbb{R}^N$. Hence, we obtain (see (4.77))

$$u(x) \leq 4\bar{C}\Theta(x) = 4\bar{C}\gamma(x_1) \leq \frac{4\bar{C}x_1}{c} \leq \frac{4}{c} \cdot \bar{C} |x|$$

for a.e. $x \in \Omega \cap B(0, \ell) \subseteq \{x \in \mathbb{R}^N : x_1 \in (0, \ell)\}$, and this establishes (1.7). \square

From Theorem 1.5, we immediately obtain the following result.

Corollary 4.14. *Let Ω be open and strictly convex, and let $f \in L^\infty(\Omega)$. Let $u_f \in \mathbb{X}(\Omega)$ be the (unique) weak solution of problem (2.1).*

Then, there exists $\ell > 0$ such that, for every $p \in \partial\Omega$, we have that

$$(4.91) \quad |u_f(x)| \leq C (\|u_f\|_{L^\infty(\mathbb{R}^N)} + \|f\|_{L^\infty(\Omega)}) |x - p|, \quad \text{for a.e. } x \in B(p, \ell).$$

Proof. Formula (4.91) follows from Theorem 4.7 and (1.7), applied to both u_f and $-u_f$, choosing

$$\bar{C} := \|u_f\|_{L^\infty(\mathbb{R}^N)} + \|f\|_{L^\infty(\Omega)}.$$

This ends the proof. \square

We point out that the term

$$\|u_f\|_{L^\infty(\mathbb{R}^N)}$$

in (4.91) can be actually reabsorbed into $\|f\|_{L^\infty(\Omega)}$, as it follows from Theorem 4.7.

As a byproduct of Theorem 1.5, we also establish Theorem 1.6:

Proof of Theorem 1.6. First of all, since $f \in L^\infty(\Omega)$, we know from Theorem 1.1 that there exists a (unique) weak solution $u_f \in \mathbb{X}(\Omega)$ of problem (2.1). Moreover, by combining Theorems 4.7 and 1.5, we infer the existence of a suitable constant $\mathbf{c} > 0$, independent of u_f , such that

- (a) $\|u_f\|_{L^\infty(\mathbb{R}^N)} \leq \mathbf{c} \|f\|_{L^\infty(\Omega)}$;
- (b) there exists $\ell > 0$ such that, for every $p \in \partial\Omega$,

$$|u_f(x)| \leq \mathbf{c} \|f\|_{L^\infty(\Omega)} \cdot |x - p| \quad \text{for a.e. } x \in \Omega \cap B(p, \ell).$$

Now, since $f \in C^k(\Omega, \mathbb{R})$ (and $k \geq \frac{N}{2} + 3$), we derive from Corollary 4.6 that there exists a unique function $\widehat{u} \in C^k(\Omega, \mathbb{R})$ such that

$$(4.92) \quad \widehat{u} \equiv u_f \quad \text{a.e. on } \Omega,$$

where $k = k_{m,N}$ is as in (4.40). In particular, $k \geq 2$. Setting

$$\mathbf{u}_f : \mathbb{R}^N \rightarrow \mathbb{R}, \quad \mathbf{u}_f(x) := \begin{cases} \widehat{u}(x), & \text{for } x \in \Omega, \\ 0, & \text{for } x \notin \Omega, \end{cases}$$

we claim that \mathbf{u}_f is a classical solution of (2.1), further satisfying (i)-(ii)-(iii). Indeed, using (4.92) and the fact that $u_f \equiv 0$ a.e. in $\mathbb{R}^N \setminus \Omega$, we have

$$\mathbf{u}_f \equiv u_f \quad \text{a.e. in } \mathbb{R}^N.$$

As a consequence, $\mathbf{u}_f \in H^1(\mathbb{R}^N)$ (hence, (i) is fulfilled) and, since u_f satisfies (a)-(b), we immediately derive that \mathbf{u}_f satisfies (ii)-(iii) (with the same constants \mathbf{c} , $\ell > 0$). In particular, from (iii) we deduce that

$$\lim_{x \rightarrow p} \mathbf{u}_f(x) = 0 \quad \text{for all } p \in \partial\Omega,$$

thus, \mathbf{u}_f being bounded, we get $\mathbf{u}_f \in \mathcal{C}_s(\mathbb{R}^N)$. Finally, since u_f is a weak solution of (2.1) and $\mathbf{u}_f \in \mathcal{C}_s(\mathbb{R}^N) \cap C^2(\Omega, \mathbb{R})$ (actually, $\mathbf{u}_f \in C^k(\Omega, \mathbb{R})$), from (i) and Remark 2.5 we conclude that \mathbf{u}_f is a classical solution of (2.1). The uniqueness of \mathbf{u}_f then follows from Corollary 3.1, and the proof of Theorem 1.6 is thereby complete. \square

APPENDIX A. FAILURE OF THE MAXIMUM PRINCIPLE

A.1. The case of $\mathcal{L}' := \Delta + (-\Delta)^s$. The following examples show that the weak maximum principle contained in (1.3) *does not hold* for the operator

$$(A.1) \quad \mathcal{L}' := \Delta + (-\Delta)^s.$$

Example A.1. Let $s \in (0, 1/2)$ be arbitrarily fixed, and let

$$f : \mathbb{R} \rightarrow \mathbb{R}, \quad f(x) := \begin{cases} x^2 - 1, & \text{if } |x| \leq 1, \\ 0, & \text{if } |x| > 1. \end{cases}$$

Clearly, $f \in C_b(\mathbb{R})$ (as $|f| \leq 1$). Moreover, setting $\Omega_0 := (-1, 1)$, we also have that $f \in C^2(\Omega_0)$.

We claim that, for every fixed $x \in \mathbb{R}$, one has

$$(A.2) \quad y \mapsto \frac{f(x) - f(y)}{|x - y|^{1+2s}} \in L^1(\mathbb{R}).$$

In order to prove (A.2), we distinguish three cases:

(i) $x \in (-1, 1)$. In this case, reminding that $s \in (0, 1/2)$, we have

$$\begin{aligned} & \int_{\mathbb{R}} \frac{|f(x) - f(y)|}{|x - y|^{1+2s}} dy \\ &= |f(x)| \int_{-\infty}^{-1} \frac{dy}{|x - y|^{1+2s}} + |f(x)| \int_1^{\infty} \frac{dy}{|x - y|^{1+2s}} + \int_{-1}^1 \frac{|x^2 - y^2|}{|x - y|^{1+2s}} dy \\ &\leq \int_{-\infty}^{-1} \frac{dy}{|x - y|^{1+2s}} + \int_1^{\infty} \frac{dy}{|x - y|^{1+2s}} + 2 \int_{-1}^1 \frac{dy}{|x - y|^{2s}} < \infty. \end{aligned}$$

(ii) $x = \pm 1$. We perform the computations when $x = 1$, being the case $x = -1$ completely analogous. In this case, we have

$$\int_{\mathbb{R}} \frac{|f(x) - f(y)|}{|1 - y|^{1+2s}} dy = \int_{-1}^1 \frac{1 - y^2}{|1 - y|^{1+2s}} dy \leq 2 \int_{-1}^1 \frac{1}{|1 - y|^{2s}} dy < \infty.$$

(iii) $x \notin [-1, 1]$. In this case, since $|x - y| \geq |x| - 1$ if $y \in (-1, 1)$, we have

$$\int_{\mathbb{R}} \frac{|f(x) - f(y)|}{|x - y|^{1+2s}} dy = \int_{-1}^1 \frac{1 - y^2}{|x - y|^{1+2s}} dy \leq \frac{1}{(|x| - 1)^{1+2s}} \int_{-1}^1 (1 - y^2) dy < \infty.$$

Summing up, the claimed (A.2) is completely established.

Now, we observe that, for any $x \in (-1, 1)$,

$$\begin{aligned} \frac{|(-\Delta)^s f(x)|}{c_{1,s}} &\leq \int_{\mathbb{R}} \frac{|f(x) - f(y)|}{|x - y|^{1+2s}} dy \\ &\leq (1 - x^2) \int_{-\infty}^{-1} \frac{dy}{|x - y|^{1+2s}} + (1 - x^2) \int_1^{\infty} \frac{dy}{|x - y|^{1+2s}} + \int_{-1}^1 \frac{|x^2 - y^2|}{|x - y|^{1+2s}} dy \\ &\leq (1 - x^2) \int_{-\infty}^{-1} \frac{dy}{|x - y|^{1+2s}} + (1 - x^2) \int_1^{\infty} \frac{dy}{|x - y|^{1+2s}} + 2 \int_{-1}^1 \frac{dy}{|x - y|^{2s}} \\ &\leq 2^{2-2s} \cdot \frac{1 - s}{s(1 - 2s)}. \end{aligned}$$

As a consequence, if $\varepsilon \in (0, 1)$ and if $f_\varepsilon(x) := f(x/\varepsilon)$, we have

$$\mathcal{L}' f_\varepsilon(x) = \frac{2}{\varepsilon^2} + \frac{1}{\varepsilon^{2s}} ((-\Delta)^s f)(x/\varepsilon) \geq \frac{2}{\varepsilon^2} \left(1 - \varepsilon^{2-2s} \cdot \frac{2^{1-2s} c_{1,s} (1 - s)}{s(1 - 2s)} \right),$$

for every $x \in \mathbb{R}$ with $|x| < \varepsilon$. If we choose ε_0 so small that

$$1 - \varepsilon_0^{2-2s} \cdot \frac{2^{1-2s} c_{1,s} (1 - s)}{s(1 - 2s)} > 0,$$

we thus see that f_{ε_0} enjoys the following properties:

- (a) $f_{\varepsilon_0} \in C^2(\Omega_{\varepsilon_0}) \cap C_b(\mathbb{R})$, where $\Omega_{\varepsilon_0} := (-\varepsilon_0, \varepsilon_0)$;
- (b) the s -Laplacian of f_{ε_0} is pointwise defined on the whole of \mathbb{R} ;
- (c) $f_{\varepsilon_0} \equiv 0$ on $\mathbb{R} \setminus \Omega_{\varepsilon_0}$ and $\mathcal{L}' f_{\varepsilon_0} > 0$ on Ω_{ε_0} .

Since, obviously, $f_{\varepsilon_0} < 0$ on Ω_{ε_0} , we conclude that a weak maximum principle as in (1.3) does not hold for $\mathcal{L}' = \Delta + (-\Delta)^s$.

Example A.2. By dropping the assumption $u \equiv 0$ in $\mathbb{R}^N \setminus \Omega$, it is possible to show that \mathcal{L}' in (A.1) violates the weak maximum principle in (1.3) for every $s \in (0, 1)$. Indeed, let $s \in (0, 1)$ and let

$$f : \mathbb{R}^N \longrightarrow \mathbb{R}, \quad f(x) := |x|^2 - 1.$$

Moreover, let $\varphi \in C_0^\infty(\mathbb{R}^N, [0, 1])$ be a cut-off function such that

$$(A.3) \quad \varphi \equiv 1 \text{ on } \Omega := B(0, 1) \text{ and } \varphi \equiv 0 \text{ on } \mathbb{R}^N \setminus B(0, 2).$$

We then set $u := f\varphi$. Obviously, $u \in C_0^\infty(\mathbb{R}^N) \subseteq C^2(\Omega) \cap C_b(\mathbb{R}^N)$. Moreover, by taking into account the properties of φ in (A.3), we see that

$$\Delta u = \Delta f = 2N \text{ in } \Omega \quad \text{and} \quad (-\Delta)^s u \in L^\infty(\mathbb{R}^N).$$

As a consequence, if $\varepsilon \in (0, 1)$ and $u_\varepsilon := u(x/\varepsilon)$, we have

$$\mathcal{L}' u_\varepsilon(x) = \frac{2N}{\varepsilon^2} + \frac{1}{\varepsilon^{2s}} ((-\Delta)^s u)(x/\varepsilon) \geq \frac{1}{\varepsilon^2} \left(2N - \varepsilon^{2-2s} \|(-\Delta)^s u\|_{L^\infty(\mathbb{R}^N)} \right),$$

for all $x \in \Omega_\varepsilon := B(0, \varepsilon)$. We now argue as in Example A.1: if ε_0 is so small that

$$2N - \varepsilon_0^{2-2s} \|(-\Delta)^s u\|_{L^\infty(\mathbb{R}^N)} > 0,$$

then the function u_{ε_0} enjoys the following properties:

- (a) $u_{\varepsilon_0} \in C_0^\infty(\mathbb{R}^N) \subseteq C^2(\Omega_{\varepsilon_0}) \cap C_b(\mathbb{R})$;
- (b) the s -Laplacian of u_{ε_0} is pointwise defined on the whole of \mathbb{R} (and it is globally bounded);
- (c) $u_{\varepsilon_0} \geq 0$ on $\mathbb{R} \setminus \Omega_{\varepsilon_0}$ and $\mathcal{L}' u_{\varepsilon_0} > 0$ on Ω_{ε_0} .

Since, obviously, $u_{\varepsilon_0} < 0$ on Ω_{ε_0} , we conclude that a weak maximum principle as in (1.3) does not hold for $\mathcal{L}' = \Delta + (-\Delta)^s$, for any $s \in (0, 1)$.

A.2. The rôle of the ‘non-local boundary conditions’. Throughout the sequel, given any $R > 0$, we adopt the simplified notation

$$B_R := B(0, R).$$

Then, we claim that for every fixed $r > 1$ there exists $v \in C^2(B_1, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$ such that

$$(A.4) \quad \begin{cases} \mathcal{L}v = 0 & \text{in } B_1, \\ \inf_{B_r \setminus B_1} v > 0, \\ \min_{B_1} v < 0. \end{cases}$$

From this, since the continuity of v implies that $v \geq 0$ on ∂B_1 , we deduce that the weak maximum principle in (1.3) *does not hold* if one requires $u \geq 0$ only on $\partial\Omega$.

To prove the existence of such a function v , we let $\phi \in C_0^\infty(\mathbb{R}^N, \mathbb{R})$ be such that

- (a) $\phi \equiv -1$ in $B_{r+3} \setminus B_{r+2}$ and $\phi \equiv 0$ outside $B_{r+4} \setminus B_{r+1}$;
- (b) $-1 \leq \phi \leq 0$ on the whole of \mathbb{R}^N .

Since $\phi \in C_0^\infty(\mathbb{R}^N, \mathbb{R}) \subseteq \mathcal{S}(\mathbb{R}^N)$, we have that $f := -\mathcal{L}\phi$ can be computed pointwise in \mathbb{R}^N and $f \in C^\infty(\mathbb{R}^N) \cap L^\infty(\mathbb{R}^N)$.

As a consequence, from Theorem 1.6 we know that there exists a unique classical solution $\mathbf{u}_\phi \in C^2(B_1, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$ of the problem

$$(A.5) \quad \begin{cases} \mathcal{L}u = f = -\mathcal{L}\phi & \text{in } B_1, \\ u \equiv 0 & \text{on } \mathbb{R}^N \setminus B_1. \end{cases}$$

We then set $w := \mathbf{u}_\phi + \phi$ and we notice that, thanks to the regularity of \mathbf{u}_ϕ and ϕ , one has that $w \in C^2(B_1, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$. Furthermore, from (A.5) we obtain that

$$(A.6) \quad \begin{cases} \mathcal{L}w = 0 & \text{in } B_1, \\ w \equiv \phi & \text{in } \mathbb{R}^N \setminus B_1. \end{cases}$$

We now claim that

$$(A.7) \quad m := \inf_{B_1} w < 0.$$

For this, we argue by contradiction and we suppose that

$$w(x) \geq 0 \quad \text{for every } x \in B_1.$$

In particular, since $w \in C(\mathbb{R}^N, \mathbb{R})$ and $w \equiv \phi \equiv 0$ on ∂B_1 , we can find an *interior* maximum point $x_0 \in B_1$ for w such that

$$(A.8) \quad w(x_0) \geq 0.$$

Thus, $\Delta w(x_0) \leq 0$ and therefore, using the exterior condition in (A.6),

$$(A.9) \quad \begin{aligned} 0 &= \frac{\mathcal{L}w(x_0)}{c_{N,s}} = \frac{(-\Delta + (-\Delta)^s)w(x_0)}{c_{N,s}} \geq \frac{(-\Delta)^s w(x_0)}{c_{N,s}} = \int_{\mathbb{R}^N} \frac{w(x_0) - w(y)}{|x_0 - y|^{N+2s}} dy \\ &= \int_{B_1} \frac{w(x_0) - w(y)}{|x_0 - y|^{N+2s}} dy + \int_{\mathbb{R}^N \setminus B_1} \frac{w(x_0) - \phi(y)}{|x_0 - y|^{N+2s}} dy \\ &\geq \int_{\mathbb{R}^N \setminus B_1} \frac{w(x_0) - \phi(y)}{|x_0 - y|^{N+2s}} dy, \end{aligned}$$

where we used that fact that $w(x_0) \geq w(y)$ for all $y \in B_1$ in the last line. Moreover, by assumption (a), we know that $\phi \equiv -1$ on $B_{r+3} \setminus B_{r+2}$. Accordingly, using also (A.8) and assumption (b) on ϕ , we conclude from (A.9) that

$$\begin{aligned}
0 &\geq \int_{B_{r+3} \setminus B_{r+2}} \frac{w(x_0) + 1}{|x_0 - y|^{N+2s}} dy + \int_{(B_{r+2} \setminus B_1) \cup (\mathbb{R}^N \setminus B_{r+3})} \frac{w(x_0) - \phi(y)}{|x_0 - y|^{N+2s}} dy \\
&\geq \int_{B_{r+3} \setminus B_{r+2}} \frac{w(x_0) + 1}{|x_0 - y|^{N+2s}} dy > 0.
\end{aligned}$$

This contradiction proves (A.7). With (A.7) at hand, we define

$$u(x) := w(x) - m,$$

and we observe that, in view of the properties of w , one has

- (1) $u \in C^2(B_1, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$ (as the same is true of w);
- (2) $\mathcal{L}u = \mathcal{L}w = 0$ pointwise on B_1 ;
- (3) $u = -m > 0$ on $B_r \setminus B_1$ (as $w \equiv \phi \equiv 0$ on $B_r \setminus B_1$);
- (4) $\inf_{B_1} u = 0$.

Thus, by making use of (1)–(4), we easily conclude that the function

$$v(x) := 2u(x) + m,$$

belongs to $C^2(B_1, \mathbb{R}) \cap \mathcal{C}_s(\mathbb{R}^N)$ and satisfies (A.4).

REFERENCES

- [1] N. Abatangelo, M. Cozzi, *An elliptic boundary value problem with fractional nonlinearity*, preprint.
- [2] N. Abatangelo, E. Valdinoci, *Getting acquainted with the fractional Laplacian*, Contemporary research in elliptic PDEs and related topics, Springer INdAM Ser. **33**, Springer, Cham (2019), 1–105.
- [3] N. Alibaud, F. del Teso, J. Endal, E.R. Jakobsen, *The Liouville theorem and linear operators satisfying the maximum principle*, preprint.
- [4] G. Barles, E. Chasseigne, A. Ciomaga, C. Imbert, *Lipschitz regularity of solutions for mixed integro-differential equations*, J. Differential Equations **252** (2012), no. 11, 6012–6060.
- [5] G. Barles, E. Chasseigne, A. Ciomaga, C. Imbert, *Large time behavior of periodic viscosity solutions for uniformly parabolic integro-differential equations*, Calc. Var. Partial Differential Equations **50** (2014), no. 1-2, 283–304.
- [6] G. Barles, C. Imbert, *Second-order elliptic integro-differential equations: viscosity solutions' theory revisited*, Ann. Inst. H. Poincaré Anal. Non Linéaire **25** (2008), no. 3, 567–585.
- [7] B. Barrios, I. Peral, F. Soria, E. Valdinoci, *A Widder's type theorem for the heat equation with nonlocal diffusion*, Arch. Ration. Mech. Anal. **213** (2014), no. 2, 629–650.
- [8] R.F. Bass, M. Kaßmann, *Harnack inequalities for non-local operators of variable order*, Trans. Amer. Math. Soc. **357** (2005), no. 2, 837–850.
- [9] R.F. Bass, M. Kaßmann, *Hölder continuity of harmonic functions with respect to operators of variable order*, Comm. Partial Differential Equations **30** (2005), no. 7-9, 1249–1259.
- [10] S. Biagi, S. Dipierro, E. Valdinoci, E. Vecchi, *Semilinear elliptic equations involving mixed local and nonlocal operators*, Proc. Roy. Soc. Edinburgh Sect. A **151** (2021), no. 5, 1611–1641.
- [11] I.H. Biswas, E.R. Jakobsen, K.H. Karlsen, *Viscosity solutions for a system of integro-PDEs and connections to optimal switching and control of jump-diffusion processes*, Appl. Math. Optim. **62** (2010), no. 1, 47–80.
- [12] I.H. Biswas, E.R. Jakobsen, K.H. Karlsen, *Difference-quadrature schemes for nonlinear degenerate parabolic integro-PDE*, SIAM J. Numer. Anal. **48** (2010), no. 3, 1110–1135.
- [13] D. Blazevski, D. del-Castillo-Negrete, *Local and nonlocal anisotropic transport in reversed shear magnetic fields: Shearless Cantori and nondiffusive transport*, Phys. Rev. E **87** (2013), 063106.
- [14] J.-M. Bony, P. Courrége, P. Priouret, *Semi-groupes de Feller sur une variété à bord compacte et problèmes aux limites intégral-différentiels du second ordre donnant lieu au principe du maximum*, Ann. Inst. Fourier (Grenoble) **18** (1968), 369–521.
- [15] X. Cabré, S. Dipierro, E. Valdinoci, *The Bernstein technique for integro-differential equations*, preprint.
- [16] X. Cabré, J. Serra, *An extension problem for sums of fractional Laplacians and 1-D symmetry of phase transitions*, Nonlinear Anal. **137** (2016), 246–265.
- [17] L.A. Caffarelli, P.E. Souganidis, *Convergence of nonlocal threshold dynamics approximations to front propagation*, Arch. Ration. Mech. Anal. **195** (2010), no. 1, 1–23.
- [18] L. Caffarelli, E. Valdinoci, *A priori bounds for solutions of a nonlocal evolution PDE*, Analysis and numerics of partial differential equations, 141–163, Springer INdAM Ser., 4, Springer, Milan, 2013.
- [19] Z.-Q. Chen, P. Kim, R. Song, Z. Vondraček, *Sharp Green function estimates for $\Delta + \Delta^{\alpha/2}$ in $C^{1,1}$ open sets and their applications*, Illinois J. Math. **54** (2010), no. 3, 981–1024.

- [20] Z.-Q. Chen, P. Kim, R. Song, Z. Vondraček, *Boundary Harnack principle for $\Delta + \Delta^{\alpha/2}$* , Trans. Amer. Math. Soc. **364** (2012), no. 8, 4169–4205.
- [21] A. Ciomaga, *On the strong maximum principle for second-order nonlinear parabolic integro-differential equations*, Adv. Differential Equations **17** (2012), no. 7-8, 635–671.
- [22] P. Courrège, *Générateur infinitésimal d'un semi-groupe de convolution sur \mathbb{R}^n , et formule de Lévy-Khinchine*, Bull. Sci. Math. **88** (1964), 3–30.
- [23] F. Dell'Oro, V. Pata, *Second order linear evolution equations with general dissipation*, preprint.
- [24] F. del Teso, J. Endal, E.R. Jakobsen, *On distributional solutions of local and nonlocal problems of porous medium type*, C. R. Math. Acad. Sci. Paris **355** (2017), no. 11, 1154–1160.
- [25] E. Di Nezza, G. Palatucci, E. Valdinoci, *Hitchhiker's guide to the fractional Sobolev spaces*, Bull. Sci. Math. **136** (2012), no. 5, 521–573.
- [26] S. Dipierro, M. Medina, I. Peral, E. Valdinoci, *Bifurcation results for a fractional elliptic equation with critical exponent in \mathbb{R}^n* , Manuscripta Math. **153** (2017), no. 1-2, 183–230.
- [27] S. Dipierro, S. Patrizi, E. Valdinoci, *Heteroclinic connections for nonlocal equations*, Math. Models Methods Appl. Sci. **29** (2019), no. 14, 2585–2636.
- [28] S. Dipierro, E. Proietti Lippi, E. Valdinoci, *(Non)local logistic equations with Neumann conditions*, preprint.
- [29] S. Dipierro, J. Serra, E. Valdinoci, *Improvement of flatness for nonlocal phase transitions*, Amer. J. Math. **142** (2020), no. 4, 1083–1160.
- [30] S. Dipierro, E. Valdinoci, *Description of an ecological niche for a mixed local/nonlocal dispersal: an evolution equation and a new Neumann condition arising from the superposition of Brownian and Lévy processes*, Phys. A **575** (2021), Paper No. 126052, 20 pp.
- [31] S. Dipierro, E. Valdinoci, V. Vespri, *Decay estimates for evolutionary equations with fractional time-diffusion*, J. Evol. Equ. **19** (2019), no. 2, 435–462.
- [32] J.M. Epstein, D.M. Goedecke, F. Yu, R.J. Morris, D.K. Wagener, G.V. Bobashev, *Controlling pandemic flu: the value of international air travel restrictions*, PLoS ONE **2** (2007), no. 5, e401.
- [33] L.C. Evans, *Partial Differential Equations*, Graduate Studies in Mathematics **19**, AMS, Providence, RI (2010).
- [34] A. Farina, E. Valdinoci, *Regularity and rigidity theorems for a class of anisotropic nonlocal operators*, Manuscripta Math. **153** (2017), no. 1-2, 53–70.
- [35] A. Farina, E. Valdinoci, *Gradient estimates for a class of anisotropic nonlocal operators*, NoDEA Nonlinear Differential Equations Appl. **26** (2019), no. 4, Paper No. 25, 13 pp.
- [36] M.E. Farquhar, T.J. Moroney, Q. Yang, I.W. Turner, K. Burrage, *Computational modelling of cardiac ischaemia using a variable-order fractional Laplacian*, preprint.
- [37] E.R. Jakobsen, K.H. Karlsen, *Continuous dependence estimates for viscosity solutions of integro-PDEs*, J. Differential Equations **212** (2005), no. 2, 278–318.
- [38] E.R. Jakobsen, K.H. Karlsen, *A “maximum principle for semicontinuous functions” applicable to integro-partial differential equations*, NoDEA Nonlinear Differential Equations Appl. **13** (2006), 137–165.
- [39] W. Hoh, *Pseudo differential operators with negative definite symbols of variable order*, Rev. Mat. Iberoamericana **16** (2000), no. 2, 219–241.
- [40] M. Jleli, B. Samet, *Liouville-type theorems for an elliptic system involving fractional Laplacian operators with mixed order*, Electron. J. Differential Equations, (2017), Paper No. 105, 11 pp.
- [41] R. de la Llave, E. Valdinoci, *A generalization of Aubry-Mather theory to partial differential equations and pseudo-differential equations*, Ann. Inst. H. Poincaré Anal. Non Linéaire **26** (2009), no. 4, 1309–1344.
- [42] A. Mimica, *Heat kernel estimates for subordinate Brownian motions*, Proc. Lond. Math. Soc. (3) **113** (2016), no. 5, 627–648.
- [43] X. Ros-Oton, J. Serra, *The Dirichlet problem for the fractional Laplacian: regularity up to the boundary*, J. Math. Pures Appl. (9) **101** (2014), no. 3, 275–302.
- [44] O. Savin, E. Valdinoci, *Density estimates for a variational model driven by the Gagliardo norm*, J. Math. Pures Appl. (9) **101** (2014), no. 1, 1–26.
- [45] R. Servadei, E. Valdinoci, *Weak and viscosity solutions of the fractional Laplace equation*, Publ. Mat. **58** (2014), no. 1, 133–154.
- [46] L.E. Silvestre, *Regularity of the obstacle problem for a fractional power of the Laplace operator*, Thesis (Ph.D.)-The University of Texas at Austin (2005), 95 pp.
- [47] L.E. Silvestre, *Hölder estimates for solutions of integro-differential equations like the fractional Laplace*, Indiana Univ. Math. J. **55** (2006), no. 3, 1155–1174.
- [48] C.J. Weiss, B.G. van Bloemen Waanders, H. Antil, *Fractional operators applied to geophysical electromagnetics*, Geophysical J. Intern. **220** (2020), no. 2, 1242–1259.

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