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REGULARITY OF FLAT FREE BOUNDARIES FOR A $p(x)$ -LAPLACIAN PROBLEM WITH RIGHT HAND SIDE

FAUSTO FERRARI AND CLAUDIA LEDERMAN

ABSTRACT. We consider viscosity solutions to a one-phase free boundary problem for the $p(x)$ -Laplacian with non-zero right hand side. We apply the tools developed in [D] to prove that flat free boundaries are $C^{1,\alpha}$. Moreover, we obtain some new results for the operator under consideration that are of independent interest.

1. INTRODUCTION AND MAIN RESULTS

In this paper we study a one-phase free boundary problem governed by the $p(x)$ -Laplacian with non-zero right hand side. More precisely, we denote by

$$\Delta_{p(x)}u := \operatorname{div}(|\nabla u|^{p(x)-2}\nabla u),$$

where p is a function such that $1 < p(x) < +\infty$. Then our problem is the following:

$$(1.1) \quad \begin{cases} \Delta_{p(x)}u = f, & \text{in } \Omega^+(u) := \{x \in \Omega : u(x) > 0\}, \\ |\nabla u| = g, & \text{on } F(u) := \partial\Omega^+(u) \cap \Omega. \end{cases}$$

Here $\Omega \subset \mathbb{R}^n$ is a bounded domain, $p \in C^1(\Omega)$, $f \in C(\Omega) \cap L^\infty(\Omega)$ and $g \in C^{0,\beta}(\Omega)$, $g \geq 0$.

This problem comes out naturally from limits of a singular perturbation problem with forcing term as in [LW1], where the authors analyze solutions to (1.1), arising in the study of flame propagation with nonlocal and electromagnetic effects. On the other hand, (1.1) appears by minimizing the following functional

$$(1.2) \quad \mathcal{E}(v) = \int_{\Omega} \left(\frac{|\nabla v|^{p(x)}}{p(x)} + Q^2(x)\chi_{\{v>0\}} + f(x)v \right) dx$$

studied in [LW3], as well as in the seminal paper by Alt and Caffarelli [AC] in the case $p(x) \equiv 2$ and $f \equiv 0$. We refer also to [LW4], where (1.1) appears in the study of an optimal design problem.

We are interested in the regularity of the free boundary for viscosity solutions of (1.1). This problem has been already faced in [LW2] for weak solutions with the aid of the techniques developed in [AC].

Key words and phrases. free boundary problem, singular/degenerate operator, variable exponent spaces, regularity of the free boundary, non-zero right hand side, viscosity solutions.

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In the present work we are following the strategy introduced in the important paper by De Silva [D], that was inspired by [S], for one-phase problems and linear non-divergence operators. [D] was further extended to two-phase problems in different settings, see [DFS1, DFS2, DFS3]. The same technique was applied to the p -Laplace operator ($p(x) \equiv p$ in (1.1)) for the one phase case, with $p \geq 2$, in [LR]. See also [LT].

In the linear homogeneous case, $f \equiv 0$, (1.1) was studied for viscosity solutions in the pioneer works by Caffarelli [C1, C2]. The results in [C1, C2] have been widely generalized to different classes of homogeneous elliptic problems. See for example [CFS, FS1, FS2] for linear operators, [AF, F1, F2, Fe1, W1, W2, RT] for fully nonlinear operators and [LN1, LN2] for the p -Laplacian. See also [ART].

As already mentioned, problem (1.1) was originally studied in the linear homogeneous case in [AC], associated to (1.2). These techniques were generalized to the linear case with $f \neq 0$ in [GS, Le]. In the homogeneous case, to a quasilinear uniformly elliptic situation [ACF], to the p -Laplacian [DP], to an Orlicz setting [MW] and to the $p(x)$ -Laplacian with $p(x) \geq 2$ [FMW]. Finally, (1.1) with $1 < p(x) < \infty$ and $f \neq 0$ was dealt with in [LW2].

In this paper we show that flat free boundaries of viscosity solutions to (1.1) are $C^{1,\alpha}$. In the forthcoming work [FL] we prove that Lipschitz free boundaries of viscosity solutions to (1.1) are $C^{1,\alpha}$.

Our main result is the following (for the precise definition of viscosity solution to (1.1) we refer to Section 2)

Theorem 1.1 (Flatness implies $C^{1,\alpha}$). *Let u be a viscosity solution to (1.1) in B_1 . Assume that $0 \in F(u)$, $g(0) = 1$ and $p(0) = p_0$. There exists a universal constant $\bar{\varepsilon} > 0$ such that, if the graph of u is $\bar{\varepsilon}$ -flat in B_1 , in the direction e_n , that is*

$$(1.3) \quad (x_n - \bar{\varepsilon})^+ \leq u(x) \leq (x_n + \bar{\varepsilon})^+, \quad x \in B_1,$$

and

$$(1.4) \quad \|\nabla p\|_{L^\infty(B_1)} \leq \bar{\varepsilon}, \quad \|f\|_{L^\infty(B_1)} \leq \bar{\varepsilon}, \quad [g]_{C^{0,\beta}(B_1)} \leq \bar{\varepsilon},$$

then $F(u)$ is $C^{1,\alpha}$ in $B_{1/2}$.

In addition to the assumptions already stated above, we suppose that

$$(1.5) \quad \nabla p \in L^\infty(\Omega)$$

and that there exist positive numbers p_{\min}, p_{\max} , such that

$$(1.6) \quad 1 < p_{\min} \leq p(x) \leq p_{\max} < \infty.$$

In Theorem 1.1 the constants $\bar{\varepsilon}$ and α depend only on p_{\min} , p_{\max} and n (the dimension of the space).

The proof of Theorem 1.1 is based on an improvement of flatness, obtained via a compactness argument which linearizes the problem into a limiting one. The key tool is a geometric Harnack inequality that localizes the free boundary well, and allows the rigorous passage to the limit.

Let us point out that carrying out, for the inhomogeneous $p(x)$ -Laplace operator, the strategy devised in [D] required the development of new tools. In fact, the $p(x)$ -Laplacian is a nonlinear operator that appears naturally in divergence form from

minimization problems, i.e., in the form $\operatorname{div} A(x, \nabla u) = f(x)$, with

$$\lambda |\eta|^{p(x)-2} |\xi|^2 \leq \sum_{i,j=1^n} \frac{\partial A_i}{\partial \eta_j}(x, \eta) \xi_i \xi_j \leq \Lambda |\eta|^{p(x)-2} |\xi|^2, \quad \xi \in \mathbb{R}^n.$$

This operator is singular in the regions where $1 < p(x) < 2$ and degenerate in the ones where $p(x) > 2$.

Some results for this type of operators we needed to use to achieve our goals are available in the literature for weak solutions (in the sense of Definition 3.1 in Section 3). These results are Harnack inequality (see [Wo]) and $C^{1,\alpha}$ estimates (see [Fa] and [FZ]). However, the program followed in [D] relies on solutions of the corresponding equations in a viscosity sense (see [CIL]).

The equivalence between weak and viscosity solutions of $\Delta_{p(x)} u = f$ was proved in [JJ, JLM, MO] in the case of the p -Laplacian (i.e., for $p(x) \equiv p$) and in [JLP] in the case of the homogeneous $p(x)$ -Laplacian (i.e., for $f \equiv 0$). To our knowledge there is no such result in the literature for the inhomogeneous $p(x)$ -Laplacian.

Hence, in order to proceed with the arguments in [D], we prove in Theorem 3.2 that weak solutions of $\Delta_{p(x)} u = f$ are indeed viscosity solutions. This new result is of independent interest, since it may be applied in other contexts.

On the other hand, the approach in [D] requires the use of barriers of the type $w(x) = c_1 |x - x_0|^{-\gamma} - c_2$, together with suitable modification of them. In the present work we are able to employ the same kind of barriers. Showing that they are also appropriate to deal with the inhomogeneous $p(x)$ -Laplace operator was a nontrivial and delicate task, that we perform in Lemma 4.2. Again, the difficulty relies on the nonlinear singular/degenerate nature and x dependence of our equation and also on the presence of the logarithmic term appearing in the nondivergence form of the operator (see (3.1)).

The results in Lemma 4.2 are new even for $p(x) \equiv p$ in the range $1 < p < 2$. These barriers, which are novel in the $p(x)$ -Laplace context, are different from the ones used in the literature for this operator (see, for instance, [FMW, Wo, LW4]). Consequently, our results in Lemma 4.2 have possible applications to other situations.

We would like to stress at this stage that partial differential equations with non-standard growth have been receiving a lot of attention and that the $p(x)$ -Laplacian is a model case in this class. A list of applications of this type of operators includes the modelling of non-Newtonian fluids, for instance, electrorheological [R] or thermorheological fluids [AR]. Also non-linear elasticity [Z1], image reconstruction [AMS, CLR] and the modelling of electric conductors [Z2], to cite a few.

The fact that solutions to the inhomogeneous $p(x)$ -Laplacian are locally of class $C^{1,\alpha}$ plays a critical role in the analysis of this paper. A comprehensive account for sharp conditions for regularity of solutions of some elliptic equations with non-standard growth can be found in [AM] and [Fa].

We finally remark that our main result, Theorem 1.1, is applied in the companion paper [FL] to prove that Lipschitz free boundaries of viscosity solutions of (1.1) are $C^{1,\alpha}$.

Our work is organized as follows. In Section 2 we provide notation and basic definitions, and we also present an auxiliary result on a Neumann problem which will be used in the proof of Theorem 1.1. In Section 3 we discuss the relationship between the different notions of solutions to $\Delta_{p(x)} u = f$ we are using. In particular,

we prove Theorem 3.2 which shows that weak solutions to $\Delta_{p(x)}u = f$ are viscosity solutions of the same equation. In Section 4 we prove some auxiliary results, which include Lemma 4.2, concerning the existence of barrier functions for $\Delta_{p(x)}u = f$. Next, in Section 5 we prove a geometric Harnack inequality for problem (1.1). In Section 6 we prove an improvement of flatness lemma. Finally, in Section 7 we prove our main result, Theorem 1.1. For the sake of completeness, we also include an Appendix at the end of the paper where we introduce the Sobolev spaces with variable exponent, which are the appropriate spaces to work with weak solutions of the $p(x)$ -Laplacian.

2. BASIC DEFINITIONS, NOTATION AND PRELIMINARIES

In this section, we provide notation and basic definitions we will use throughout our work. We also present an auxiliary result on a Neumann problem that will be applied in the paper.

Notation. For any continuous function $u : \Omega \subset \mathbb{R}^n \rightarrow \mathbb{R}$ we denote

$$\Omega^+(u) := \{x \in \Omega : u(x) > 0\}, \quad F(u) := \partial\Omega^+(u) \cap \Omega.$$

We refer to the set $F(u)$ as the *free boundary* of u , while $\Omega^+(u)$ is its *positive phase* (or *side*).

Below we give the definition of viscosity solution to problem (1.1) and we deduce some consequences. In particular, we refer to the usual C -viscosity definition of sub/supersolution and solution of an elliptic PDE, see e.g. [CIL].

First we need the following standard notion.

Definition 2.1. Given $u, \varphi \in C(\Omega)$, we say that φ touches u from below (resp. above) at $x_0 \in \Omega$ if $u(x_0) = \varphi(x_0)$, and

$$u(x) \geq \varphi(x) \quad (\text{resp. } u(x) \leq \varphi(x)) \quad \text{in a neighborhood } O \text{ of } x_0.$$

If this inequality is strict in $O \setminus \{x_0\}$, we say that φ touches u strictly from below (resp. above).

Definition 2.2. Let u be a continuous nonnegative function in Ω . We say that u is a viscosity solution to (1.1) in Ω , if the following conditions are satisfied:

- (i) $\Delta_{p(x)}u = f$ in $\Omega^+(u)$ in the weak sense of Definition 3.1, see Section 3.
- (ii) For every $\varphi \in C(\Omega)$, $\varphi \in C^2(\overline{\Omega^+(\varphi)})$. If φ^+ touches u from below (resp. above) at $x_0 \in F(u)$ and $\nabla\varphi(x_0) \neq 0$, then

$$|\nabla\varphi(x_0)| \leq g(x_0) \quad (\text{resp. } \geq g(x_0)).$$

Next theorem follows as a consequence of Theorem 3.2 in Section 3.

Theorem 2.3. *Let u be a viscosity solution to (1.1) in Ω . Then the following conditions are satisfied:*

- (i) $\Delta_{p(x)}u = f$ in $\Omega^+(u)$ in the viscosity sense, that is:
 - (ia) for every $\varphi \in C^2(\Omega^+(u))$ and for every $x_0 \in \Omega^+(u)$, if φ touches u from above at x_0 and $\nabla\varphi(x_0) \neq 0$, then $\Delta_{p(x_0)}\varphi(x_0) \geq f(x_0)$, that is, u is a viscosity subsolution;
 - (ib) for every $\varphi \in C^2(\Omega^+(u))$ and for every $x_0 \in \Omega^+(u)$, if φ touches u from below at x_0 and $\nabla\varphi(x_0) \neq 0$, then $\Delta_{p(x_0)}\varphi(x_0) \leq f(x_0)$, that is, u is a viscosity supersolution.

- (ii) For every $\varphi \in C(\Omega)$, $\varphi \in C^2(\overline{\Omega^+(\varphi)})$. If φ^+ touches u from below (resp. above) at $x_0 \in F(u)$ and $\nabla\varphi(x_0) \neq 0$, then

$$|\nabla\varphi(x_0)| \leq g(x_0) \quad (\text{resp. } \geq g(x_0)).$$

Remark 2.4. If $p(x) \equiv p$ or $f \equiv 0$, then any function satisfying the conditions of Theorem 2.3 is a solution to (1.1) in the sense of Definition 2.2 (see Remark 3.3).

We introduce also the notion of comparison sub/supersolution.

Definition 2.5. We say that $v \in C(\Omega)$ is a strict (comparison) subsolution (resp. supersolution) to (1.1) in Ω if $v \in C^2(\overline{\Omega^+(v)})$, $\nabla v \neq 0$ in $\overline{\Omega^+(v)}$ and the following conditions are satisfied:

- (i) $\Delta_{p(x)}v > f$ (resp. $< f$) in $\Omega^+(v)$;
(ii) If $x_0 \in F(v)$, then

$$|\nabla v(x_0)| > g(x_0) \quad (\text{resp. } |\nabla v(x_0)| < g(x_0)).$$

Notice that by the implicit function theorem, according to our definition, the free boundary of a comparison sub/supersolution is C^2 .

As a consequence of the previous discussion we have

Lemma 2.6. Let u be a viscosity solution to (1.1) in Ω . If v is a strict (comparison) subsolution to (1.1) in Ω and $u \geq v^+$ in Ω then $u > v$ in $\Omega^+(v) \cup F(v)$. Analogously, if v is a strict (comparison) supersolution to (1.1) in Ω and $v \geq u$ in Ω then $v > u$ in $\Omega^+(u) \cup F(u)$.

Notation. From now on $B_\rho(x_0) \subset \mathbb{R}^n$ will denote the open ball of radius ρ centered at x_0 , and $B_\rho = B_\rho(0)$. A positive constant depending only on the dimension n , p_{\min} , p_{\max} will be called a universal constant. We will use c , c_i to denote small universal constants and C , C_i to denote large universal constants.

The rest of the section is devoted to the study of the linearized problem associated with our free boundary problem (1.1). That is, the classical Neumann problem for a constant coefficient linear operator. Precisely, we consider the following boundary value problem:

$$(2.1) \quad \begin{cases} \mathcal{L}_{p_0}\tilde{u} = 0 & \text{in } B_\rho \cap \{x_n > 0\}, \\ \tilde{u}_n = 0 & \text{on } B_\rho \cap \{x_n = 0\}. \end{cases}$$

Here $1 < p_{\min} \leq p_0 \leq p_{\max} < \infty$, \tilde{u}_n denotes the derivative in the e_n direction of \tilde{u} and

$$(2.2) \quad \mathcal{L}_{p_0}u := \Delta u + (p_0 - 2)\partial_{nn}u.$$

Theorem 1.1 will follow via a compactness argument combined with regularity properties of solutions to (2.1), namely Theorem 2.9.

We use the notion of viscosity solution to (2.1). We recall standard notions and a regularity result for viscosity solutions to (2.1).

Definition 2.7. Let \tilde{u} be a continuous function on $B_\rho \cap \{x_n \geq 0\}$. We say that \tilde{u} is a viscosity solution to (2.1) if given a quadratic polynomial $P(x)$ touching \tilde{u} from below (resp. above) at $\bar{x} \in B_\rho \cap \{x_n \geq 0\}$,

- (i) if $\bar{x} \in B_\rho \cap \{x_n > 0\}$ then $\mathcal{L}_{p_0}P \leq 0$ (resp. $\mathcal{L}_{p_0}P \geq 0$), i.e. $\mathcal{L}_{p_0}\tilde{u} = 0$ in the viscosity sense in $B_\rho \cap \{x_n > 0\}$;

(ii) if $\bar{x} \in B_\rho \cap \{x_n = 0\}$ then $P_n(\bar{x}) \leq 0$ (resp. $P_n(\bar{x}) \geq 0$).

Remark 2.8. Notice that in the definition above we can choose polynomials P that touch \tilde{u} strictly from above/below. Also, it suffices to verify that (ii) holds for polynomials \tilde{P} with $\mathcal{L}_{p_0}\tilde{P} > 0$ (see [D]).

We will use the following regularity result for viscosity solutions to the linearized problem (2.1). For the proof we refer to Theorem 7.4 in [MS].

Theorem 2.9. *Let \tilde{u} be a viscosity solution to (2.1) in $B_{1/2} \cap \{x_n \geq 0\}$. Then, $\tilde{u} \in C^2(B_{1/2} \cap \{x_n \geq 0\})$ and it is a classical solution to (2.1).*

Moreover, if $\|\tilde{u}\|_\infty \leq 1$, then there exists a constant $\bar{C} > 0$, depending only on n, p_{\min} and p_{\max} , such that

$$(2.3) \quad |\tilde{u}(x) - \tilde{u}(0) - \nabla \tilde{u}(0) \cdot x| \leq \bar{C}r^2 \quad \text{in } B_r \cap \{x_n \geq 0\},$$

for all $r \leq 1/4$.

3. DIFFERENT NOTIONS OF SOLUTIONS TO $p(x)$ -LAPLACIAN

In this section we discuss the relationship between the different notions of solutions to $\Delta_{p(x)}u = f$ we are using, namely weak and viscosity solutions.

We start by observing that direct calculations show that, for C^2 functions u such that $\nabla u(x) \neq 0$,

$$(3.1) \quad \begin{aligned} \Delta_{p(x)}u &= \operatorname{div}(|\nabla u|^{p(x)-2}\nabla u) \\ &= |\nabla u(x)|^{p(x)-2} \left(\Delta u + (p(x) - 2)\Delta_\infty^N u + \langle \nabla p(x), \nabla u(x) \rangle \log |\nabla u(x)| \right), \end{aligned}$$

where

$$\Delta_\infty^N u := \left\langle D^2 u(x) \frac{\nabla u(x)}{|\nabla u(x)|}, \frac{\nabla u(x)}{|\nabla u(x)|} \right\rangle$$

denotes the normalized ∞ -Laplace operator.

First we need (see the Appendix for the definition of Sobolev spaces with variable exponent)

Definition 3.1. Assume that $1 < p_{\min} \leq p(x) \leq p_{\max} < \infty$ with $p(x)$ Lipschitz continuous in Ω and $\|\nabla p\|_{L^\infty} \leq L$, for some $L > 0$ and $f \in L^\infty(\Omega)$.

We say that u is a weak solution to $\Delta_{p(x)}u = f$ in Ω if $u \in W^{1,p(\cdot)}(\Omega)$ and, for every $\varphi \in C_0^\infty(\Omega)$, there holds that

$$-\int_\Omega |\nabla u(x)|^{p(x)-2} \nabla u \cdot \nabla \varphi \, dx = \int_\Omega \varphi f(x) \, dx.$$

We next prove

Theorem 3.2. *Let p and f be as in Definition 3.1. Assume moreover that $f \in C(\Omega)$ and $p \in C^1(\Omega)$.*

Let $u \in W^{1,p(\cdot)}(\Omega) \cap C(\Omega)$ be a weak solution to $\Delta_{p(x)}u = f$ in Ω . Then u is a viscosity solution to $\Delta_{p(x)}u = f$ in Ω .

Proof. Let us show that u is a viscosity supersolution to $\Delta_{p(x)}u = f$ in Ω .

Step I. We will first prove the result under the extra assumption that $f \in W^{1,\infty}(\Omega)$ and $p \in C^{1,\beta}(\Omega)$, for some $0 < \beta < 1$.

In fact, let $v \in C^2(\Omega)$ such that v touches u from below at $x_0 \in \Omega$, with $\nabla v(x_0) \neq 0$. We will show that

$$(3.2) \quad \Delta_{p(x_0)}v(x_0) \leq f(x_0).$$

Let us fix $r > 0$ such that $\overline{B_r(x_0)} \subset \Omega$. From Theorem 1.1 in [Fa] we know that $u \in C^{1,\alpha}$ in $\overline{B_r(x_0)}$, for some $0 < \alpha < 1$. We can assume that $\alpha \leq \beta$.

Since v touches u from below at x_0 , we know that $\nabla u(x_0) = \nabla v(x_0) \neq 0$. Then, we can choose r small enough so that

$$c_1 \leq |\nabla u(x)| \leq C_1 \quad \text{in } B_r(x_0), \quad (c_1, C_1 \text{ positive constants}).$$

Now, arguing as in Theorem 3.2 in [CL] we deduce that $u \in W_{\text{loc}}^{2,2}(B_r(x_0))$ and it is a solution to the linear uniformly elliptic equation

$$\Delta_{p(x)} u = \sum_{i,j=1}^n a_{ij}(x) u_{x_i x_j} + \sum_{i=1}^n b_i(x) u_{x_i} = f \quad \text{in } B_r(x_0)$$

where

$$a_{ij}(x) = |\nabla u|^{p(x)-2} \left(\delta_{ij} + (p(x) - 2) \frac{u_{x_i} u_{x_j}}{|\nabla u|^2} \right),$$

and

$$b_i(x) = |\nabla u|^{p(x)-2} \left(p_{x_i}(x) \log |\nabla u| \right),$$

with

$$\beta_1 |\xi|^2 \leq \sum_{i,j=1}^n a_{ij}(x) \xi_i \xi_j \leq \beta_2 |\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \quad \forall x \in B_r(x_0),$$

for β_1, β_2 positive constants. It follows (see, for instance, Theorem 9.19 in [GT]) that $u \in C^{2,\alpha}$ in $B_r(x_0)$.

Since v touches u from below at x_0 , we have $\nabla u(x_0) = \nabla v(x_0)$ and $D^2 u(x_0) \geq D^2 v(x_0)$ and then,

$$\begin{aligned} f(x_0) &= \Delta_{p(x_0)} u(x_0) \\ &= \sum_{i,j=1}^n |\nabla u(x_0)|^{p(x_0)-2} \left(\delta_{ij} + (p(x_0) - 2) \frac{u_{x_i}(x_0) u_{x_j}(x_0)}{|\nabla u(x_0)|^2} \right) u_{x_i x_j}(x_0) \\ &\quad + \sum_{i=1}^n |\nabla u(x_0)|^{p(x_0)-2} \left(p_{x_i}(x_0) \log |\nabla u(x_0)| \right) u_{x_i}(x_0) \\ &\geq \Delta_{p(x_0)} v(x_0). \end{aligned}$$

That is, (3.2) holds.

Step II. We now assume that f and p are as in the statement and we will show that u is a viscosity supersolution to $\Delta_{p(x)} u = f$ in Ω .

Again, let $v \in C^2(\Omega)$ such that v touches u from below at $x_0 \in \Omega$, with $\nabla v(x_0) \neq 0$. We will show that

$$(3.3) \quad \Delta_{p(x_0)} v(x_0) \leq f(x_0).$$

Assume that $\Delta_{p(x_0)} v(x_0) > f(x_0)$. Then, there exist $r > 0$ and $\sigma > 0$ small such that

$$(3.4) \quad \begin{aligned} |\nabla v(x)| &> \sigma \quad \text{in } B_r(x_0), \\ \Delta_{p(x)} v(x) &> f(x) + \sigma \quad \text{in } B_r(x_0). \end{aligned}$$

We now take $p_k \in C^{1,\beta}(\overline{B_r(x_0)})$, for some $0 < \beta < 1$, with $\frac{1}{2}(1 + p_{\min}) \leq p_k(x) \leq p_{\max}$, $p_k \leq p$ in $B_r(x_0)$ and $\|\nabla p_k\|_{L^\infty} \leq 2L$, and $f_k \in W^{1,\infty}(B_r(x_0))$, $\|f_k\|_{L^\infty} \leq 2\|f\|_{L^\infty}$, such that

$$(3.5) \quad \begin{aligned} f_k &\rightarrow f \quad \text{uniformly on } \overline{B_r(x_0)}, \\ p_k &\rightarrow p \quad \text{and } \nabla p_k \rightarrow \nabla p \quad \text{uniformly on } \overline{B_r(x_0)}. \end{aligned}$$

Let $u_k \in W^{1,p_k(\cdot)}(B_r(x_0))$ be the (weak) solutions to

$$\begin{aligned} \Delta_{p_k(x)} u_k &= f_k \text{ in } B_r(x_0), \\ u_k &= u \text{ on } \partial B_r(x_0). \end{aligned}$$

Using Theorem 4.1 in [FZ] and Theorem 1.2 in [Fa], we get that $u_k \in C^{1,\alpha}$ in $\overline{B_r(x_0)}$, for some $0 < \alpha < 1$, $\|u_k\|_{C^{1,\alpha}(\overline{B_r(x_0)})} \leq C$ and

$$(3.6) \quad u_k \rightarrow u \quad \text{uniformly on } \overline{B_r(x_0)}.$$

Moreover, from the results in *Step I* we know that, for every k , u_k is a viscosity supersolution to $\Delta_{p_k(x)} u_k = f_k$ in $B_r(x_0)$.

We fix $\varepsilon > 0$ and define

$$\tilde{v}(x) = v(x) - \varepsilon|x - x_0|^2.$$

Since there holds (3.4), we can choose ε small enough so that

$$(3.7) \quad \begin{aligned} |\nabla \tilde{v}(x)| &> \frac{\sigma}{2} \quad \text{in } B_r(x_0), \\ \Delta_{p(x)} \tilde{v}(x) &> f(x) + \frac{\sigma}{2} \quad \text{in } B_r(x_0). \end{aligned}$$

Now, from (3.5) and (3.7), we get

$$(3.8) \quad \Delta_{p_k(x)} \tilde{v}(x) > f_k(x) + \frac{\sigma}{4} \quad \text{in } B_r(x_0), \quad \text{if } k \geq k_0.$$

We now take $0 < \delta < \frac{\varepsilon}{4}r^2$. Recalling (3.6), we can choose $k \geq k_0$ such that

$$|u_k - u| < \delta \quad \text{in } \overline{B_r(x_0)},$$

so that we have

$$\begin{aligned} u_k + \delta &> \tilde{v} \quad \text{in } \overline{B_r(x_0)}, \\ u_k(x_0) - \delta &< \tilde{v}(x_0). \end{aligned}$$

We now take

$$\bar{t} = \inf \left\{ t \in \mathbb{R} / u_k + t \geq \tilde{v} \quad \text{in } \overline{B_r(x_0)} \right\}.$$

Then, $|\bar{t}| \leq \delta$ and

$$(3.9) \quad \begin{aligned} u_k &\geq \tilde{v} - \bar{t} \quad \text{in } \overline{B_r(x_0)}, \\ u_k(\bar{x}) &= \tilde{v}(\bar{x}) - \bar{t}, \quad \text{for some } \bar{x} \in \overline{B_r(x_0)}. \end{aligned}$$

Suppose $\bar{x} \in \partial B_r(x_0)$. Then,

$$u_k(\bar{x}) = \tilde{v}(\bar{x}) - \bar{t} = v(\bar{x}) - \varepsilon r^2 - \bar{t} \leq u(\bar{x}) - \varepsilon r^2 + \delta \leq u_k(\bar{x}) + 2\delta - \varepsilon r^2,$$

a contradiction since we have chosen $\delta < \frac{\varepsilon}{4}r^2$.

Then $\bar{x} \in B_r(x_0)$ and (3.9) says that $\tilde{v} - \bar{t}$ touches u_k from below at \bar{x} . Since $\nabla \tilde{v}(\bar{x}) \neq 0$, we get

$$\Delta_{p_k(\bar{x})} \tilde{v}(\bar{x}) \leq f_k(\bar{x}).$$

This contradicts (3.8) and we conclude that (3.3) holds. So u is a viscosity supersolution to $\Delta_{p(x)}u = f$ in Ω .

The proof that u is a viscosity subsolution to $\Delta_{p(x)}u = f$ in Ω follows similarly. \square

Remark 3.3. As already mentioned in the Introduction, the equivalence between weak and viscosity solutions to the $p(x)$ -Laplacian with right hand side $f \equiv 0$ was proved in [JLP]. On the other hand, this equivalence, in case $p(x) \equiv p$ and $f \not\equiv 0$ was dealt with in [JJ] and [MO]. See also [JLM] for the case $p(x) \equiv p$ and $f \equiv 0$.

We also obtain the following result that will be used in the proof of Lemma 5.1

Proposition 3.4. *Let p and f be as in Definition 3.1. Let $B_{2r}(x_0) \subset\subset \Omega$.*

Let $u \in W^{1,p(\cdot)}(\Omega) \cap L^\infty(\Omega)$ be a weak solution to $\Delta_{p(x)}u = f$ in Ω such that

$$c_1 \leq |\nabla u(x)| \leq C_1 \quad \text{in } B_{2r}(x_0), \quad c_1, C_1 \text{ positive constants.}$$

Then, $u \in W^{2,n}(B_r(x_0))$ and it is a strong solution to the linear uniformly elliptic equation

$$\sum_{i,j=1}^n a_{ij}(x)u_{x_i x_j} + \sum_{i=1}^n b_i(x)u_{x_i} = f \quad \text{in } B_r(x_0)$$

where

$$a_{ij}(x) = |\nabla u|^{p(x)-2} \left(\delta_{ij} + (p(x) - 2) \frac{u_{x_i} u_{x_j}}{|\nabla u|^2} \right),$$

and

$$b_i(x) = |\nabla u|^{p(x)-2} \left(p_{x_i}(x) \log |\nabla u| \right),$$

with

$$\beta_1 |\xi|^2 \leq \sum_{i,j=1}^n a_{ij}(x) \xi_i \xi_j \leq \beta_2 |\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \quad \forall x \in B_r(x_0),$$

for β_1, β_2 positive constants, depending only on $c_1, C_1, p_{\min}, p_{\max}$.

Proof. We take $f_k \in W^{1,\infty}(B_{2r}(x_0))$, $\|f_k\|_{L^\infty} \leq 2\|f\|_{L^\infty}$, such that

$$f_k \rightarrow f \quad \text{in } L^1(B_{2r}(x_0)).$$

Let $u_k \in W^{1,p(\cdot)}(B_{2r}(x_0))$ be the (weak) solutions to

$$\Delta_{p(x)} u_k = f_k \quad \text{in } B_{2r}(x_0),$$

$$u_k = u \quad \text{on } \partial B_{2r}(x_0).$$

Using Theorem 4.1 in [FZ] and Theorem 1.2 in [Fa], we get that $u_k \in C^{1,\alpha}$ in $\overline{B_{2r}(x_0)}$, for some $0 < \alpha < 1$, $\|u_k\|_{C^{1,\alpha}(\overline{B_{2r}(x_0)})} \leq C$ and

$$u_k \rightarrow u, \quad \nabla u_k \rightarrow \nabla u \quad \text{uniformly on } \overline{B_{2r}(x_0)}.$$

Then, for k large,

$$\frac{c_1}{2} \leq |\nabla u_k(x)| \leq 2C_1 \quad \text{in } B_{2r}(x_0).$$

Now, arguing as in Theorem 3.2 in [CL], we deduce that, for k large, $u_k \in W_{\text{loc}}^{2,2}(B_{2r}(x_0))$ and it is a solution to the linear uniformly elliptic equation

$$\sum_{i,j=1}^n a_{ij}^k(x)(u_k)_{x_i x_j} + \sum_{i=1}^n b_i^k(x)(u_k)_{x_i} = f_k \quad \text{in } B_{2r}(x_0)$$

where

$$a_{ij}^k(x) = |\nabla u_k|^{p(x)-2} \left(\delta_{ij} + (p(x) - 2) \frac{(u_k)_{x_i} (u_k)_{x_j}}{|\nabla u_k|^2} \right),$$

and

$$b_i^k(x) = |\nabla u_k|^{p(x)-2} \left(p_{x_i}(x) \log |\nabla u_k| \right),$$

with

$$\beta_1 |\xi|^2 \leq \sum_{i,j=1}^n a_{ij}^k(x) \xi_i \xi_j \leq \beta_2 |\xi|^2, \quad \forall \xi \in \mathbb{R}^N, \quad \forall x \in B_{2r}(x_0),$$

for β_1, β_2 positive constants, depending only on $c_1, C_1, p_{\min}, p_{\max}$. Moreover, $a_{ij}^k \in C^\alpha(\overline{B_{2r}(x_0)})$ and $\|b_i^k\|_{L^\infty(B_{2r}(x_0))} \leq \tilde{C}$.

It follows (see, for instance, Lemma 9.16 and Theorem 9.11 in [GT]) that

$$u_k \in W_{\text{loc}}^{2,n}(B_{2r}(x_0)) \cap L^\infty(B_{2r}(x_0)) \quad \text{and} \quad \|u_k\|_{W^{2,n}(B_r(x_0))} \leq \tilde{C},$$

for some positive constant \tilde{C} . Then, passing to the limit $k \rightarrow \infty$, we get the desired result. \square

4. AUXILIARY RESULTS

In this section we prove some results that will be of use in our main theorem. Namely, a Harnack inequality for an auxiliary problem of $p(x)$ -Laplacian type and an existence result of barrier functions for the $p(x)$ -Laplacian operator.

In the next result we assume for simplicity that $\|f\|_{L^\infty(\Omega)} \leq 1$, but a similar result holds for any $f \in L^\infty(\Omega)$. We have

Lemma 4.1. *Assume that $1 < p_{\min} \leq p(x) \leq p_{\max} < \infty$ with $p(x)$ Lipschitz continuous in Ω and $\|\nabla p\|_{L^\infty} \leq L$, for some $L > 0$. Let $x_0 \in \Omega$ and $0 < R \leq 1$ such that $\overline{B_{4R}(x_0)} \subset \Omega$. Let $v \in W^{1,p(\cdot)}(\Omega) \cap L^\infty(\Omega)$ be a nonnegative solution to*

$$(4.1) \quad \operatorname{div}(|\nabla v + e|^{p(x)-2}(\nabla v + e)) = f \quad \text{in } \Omega,$$

where $f \in L^\infty(\Omega)$ with $\|f\|_{L^\infty(\Omega)} \leq 1$ and $e \in \mathbb{R}^n$ with $|e| = 1$. Then, there exists C such that

$$(4.2) \quad \sup_{B_R(x_0)} v \leq C \left[\inf_{B_R(x_0)} v + R \left(\|f\|_{L^\infty(B_{4R}(x_0))}^{\frac{1}{p_{\max}-1}} + C \right) \right].$$

The constant C depends only on $n, p_{\min}, p_{\max}, \|v\|_{L^\infty(B_{4R}(x_0))}$ and L .

Proof. We define $A : \Omega \times \mathbb{R}^n \rightarrow \mathbb{R}^n$

$$A(x, \xi) = |\xi + e|^{p(x)-2}(\xi + e).$$

Then equation (4.1) takes the form

$$\operatorname{div} A(x, \nabla v) = f(x) \quad \text{in } \Omega.$$

We first observe that, for every $\xi \in \mathbb{R}^n$,

$$|A(x, \xi)| = |\xi + e|^{p(x)-1} \leq C_1 |\xi|^{p(x)-1} + C_1,$$

where C_1 depends only on p_{\max} . On the other hand, for every $\xi \in \mathbb{R}^n$,

$$(4.3) \quad \begin{aligned} \langle A(x, \xi), \xi \rangle &= |\xi + e|^{p(x)-2} \langle \xi + e, \xi \rangle \\ &= |\xi + e|^{p(x)} - |\xi + e|^{p(x)-2} \langle \xi + e, e \rangle \\ &\geq |\xi + e|^{p(x)} - |\xi + e|^{p(x)-1}. \end{aligned}$$

Now, if $|\xi + e| \leq 2$, we get from (4.3)

$$(4.4) \quad \begin{aligned} \langle A(x, \xi), \xi \rangle &\geq |\xi + e|^{p(x)} - 2^{p(x)-1} \\ &\geq C_2 |\xi|^{p(x)} - C_3, \end{aligned}$$

where C_2 and C_3 depend only on p_{\max} . If $|\xi + e| > 2$, we obtain from (4.3)

$$(4.5) \quad \begin{aligned} \langle A(x, \xi), \xi \rangle &\geq |\xi + e|^{p(x)} - |\xi + e|^{p(x)-1} \\ &= |\xi + e|^{p(x)} (1 - |\xi + e|^{-1}) \\ &\geq \frac{1}{2} |\xi + e|^{p(x)} \geq C_4 |\xi|^{p(x)} - \frac{1}{2}, \end{aligned}$$

where C_4 depends only on p_{\max} . Then, from (4.4) and (4.5) we deduce

$$\langle A(x, \xi), \xi \rangle \geq C_5 |\xi|^{p(x)} - C_6,$$

where C_5 and C_6 depend only on p_{\max} . Now the result follows from Theorem 1.1 in [Wo]. \square

We now continue with a technical result concerning the existence of barrier functions for the $p(x)$ -Laplacian operator.

Lemma 4.2. *Let $x_0 \in B_1$ and $0 < \bar{r}_1 < \bar{r}_2 \leq 1$. Assume that $1 < p_{\min} \leq p(x) \leq p_{\max} < \infty$ and $\|\nabla p\|_{L^\infty} \leq \varepsilon^{1+\theta}$, for some $0 < \theta \leq 1$. Let c_0, c_1, c_2 be positive constants and let $c_3 \in \mathbb{R}$.*

There exist positive constants $\gamma \geq 1$, \bar{c} , ε_0 and ε_1 such that the functions

$$w(x) = c_1 |x - x_0|^{-\gamma} - c_2,$$

$$v(x) = q(x) + \frac{c_0}{2} \varepsilon (w(x) - 1), \quad q(x) = x_n + c_3$$

satisfy, for $\bar{r}_1 \leq |x - x_0| \leq \bar{r}_2$,

$$(4.6) \quad \Delta_{p(x)} w \geq \bar{c}, \quad \text{for } 0 < \varepsilon \leq \varepsilon_0,$$

$$(4.7) \quad \frac{1}{2} \leq |\nabla v| \leq 2, \quad \Delta_{p(x)} v > \varepsilon^2, \quad \text{for } 0 < \varepsilon \leq \varepsilon_1.$$

Here $\gamma = \gamma(n, p_{\min}, p_{\max})$, $\bar{c} = \bar{c}(p_{\min}, p_{\max}, c_1)$, $\varepsilon_0 = \varepsilon_0(n, p_{\min}, p_{\max}, \bar{r}_1, c_1)$, $\varepsilon_1 = \varepsilon_1(n, p_{\min}, p_{\max}, \bar{r}_1, c_0, c_1, \theta)$.

Proof. Without loss of generality we can assume that $x_0 = 0$. We will divide the proof into five steps.

Step 1. For simplicity, we assume first that $c_1 = 1$. Let us fix $p \in \mathbb{R}$, $1 < p_{\min} \leq p \leq p_{\max} < \infty$ and $\gamma > 0$. Let us consider $x \in \mathbb{R}^n \setminus \{0\}$.

Then, $w(x) = |x|^{-\gamma} - c_2$ and $\nabla w = -\gamma |x|^{-\gamma-2} x$, so that

$$\frac{\nabla w}{|\nabla w|} = -\frac{x}{|x|}.$$

Moreover

$$(4.8) \quad \begin{aligned} D^2 w &= \gamma(\gamma + 2) |x|^{-\gamma-2} \frac{x}{|x|} \otimes \frac{x}{|x|} - \gamma |x|^{-\gamma-2} I \\ &= \gamma |x|^{-\gamma-2} \left((\gamma + 2) \frac{x}{|x|} \otimes \frac{x}{|x|} - I \right). \end{aligned}$$

As a consequence

$$(4.9) \quad \text{Tr}(D^2w) = \gamma|x|^{-\gamma-2}((\gamma+2)-n).$$

Thus

$$(4.10) \quad \begin{aligned} \Delta_p w &= |\nabla w|^{p-2} \left(\Delta w + (p-2) \langle D^2 w \frac{\nabla w}{|\nabla w|}, \frac{\nabla w}{|\nabla w|} \rangle \right) \\ &= \gamma^{p-1} |x|^{-(\gamma+1)(p-2)} |x|^{-\gamma-2} \left((\gamma+2)-n + (p-2) \langle [(\gamma+2) \frac{x}{|x|} \otimes \frac{x}{|x|} - I] \frac{x}{|x|}, \frac{x}{|x|} \rangle \right) \\ &= \gamma^{p-1} |x|^{-\gamma(p-1)-p} (\gamma+2-n + (p-2)(\gamma+1)) \\ &= \gamma^{p-1} |x|^{-\gamma(p-1)-p} (\gamma(p-1) + p - n) \\ &\geq \gamma^{p-1} |x|^{-\gamma(p-1)-p} (\gamma(p_{\min}-1) + p_{\min} - n) \geq \gamma^{p-1} |x|^{-\gamma(p-1)-p}, \end{aligned}$$

if $\gamma > 0$ is such that

$$(4.11) \quad \gamma(p_{\min}-1) + p_{\min} - n \geq 1.$$

On the other hand,

$$D^2v = \frac{c_0}{2} \varepsilon D^2w.$$

Then, for x such that $\nabla v(x) \neq 0$,

$$(4.12) \quad \begin{aligned} \Delta_p v &= |\nabla v|^{p-2} \left(\Delta v + (p-2) \langle D^2 v \frac{\nabla v}{|\nabla v|}, \frac{\nabla v}{|\nabla v|} \rangle \right) \\ &= \frac{c_0}{2} \varepsilon |\nabla v|^{p-2} \left(\Delta w + (p-2) \langle D^2 w \frac{\nabla v}{|\nabla v|}, \frac{\nabla v}{|\nabla v|} \rangle \right) \\ &= \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \left\{ (\gamma+2)-n + (p-2) \langle [(\gamma+2) \frac{x}{|x|} \otimes \frac{x}{|x|} - I] \frac{\nabla v}{|\nabla v|}, \frac{\nabla v}{|\nabla v|} \rangle \right\} \\ &= \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \left\{ (\gamma+2)-n + (p-2) \left[(\gamma+2) \langle \frac{x}{|x|}, \frac{\nabla v}{|\nabla v|} \rangle^2 - 1 \right] \right\} \\ &= \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \left\{ (\gamma+2) \left[1 + (p-2) \langle \frac{x}{|x|}, \frac{\nabla v}{|\nabla v|} \rangle^2 \right] - n - p + 2 \right\}. \end{aligned}$$

We also observe that

$$(4.13) \quad 0 \leq \langle \frac{x}{|x|}, \frac{\nabla v}{|\nabla v|} \rangle^2 \leq 1.$$

Hence, in case $p_{\min} \leq p \leq 2$, it follows from (4.12)

$$(4.14) \quad \begin{aligned} \Delta_p v &\geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \{ (\gamma+2)(1+p-2) - n - p + 2 \} \\ &= \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \{ (\gamma+2)(p-1) - n - p + 2 \} \\ &\geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} \{ (\gamma+2)(p_{\min}-1) - n \} \geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2}, \end{aligned}$$

if $\gamma > 0$ is such that

$$(4.15) \quad (\gamma+2)(p_{\min}-1) - n \geq 1.$$

Moreover, in case $2 < p \leq p_{\max}$, it follows from (4.12)

$$(4.16) \quad \begin{aligned} \Delta_p v &\geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} (\gamma + 2 - n - p + 2) \\ &\geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2} (\gamma + 4 - n - p_{\max}) \geq \frac{c_0}{2} \gamma \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2}, \end{aligned}$$

if $\gamma > 0$ is such that

$$(4.17) \quad \gamma + 4 - n - p_{\max} \geq 1.$$

We now fix

$$(4.18) \quad \gamma = \gamma(n, p_{\min}, p_{\max}) = \max \left\{ 1, \frac{1 + n - p_{\min}}{p_{\min} - 1}, \frac{1 + n}{p_{\min} - 1} - 2, n + p_{\max} - 3 \right\}.$$

Then, $\gamma = \gamma(n, p_{\min}, p_{\max}) \geq 1$ and γ satisfies (4.11), (4.15) and (4.17). Hence we obtain from (4.10), (4.14) and (4.16) that for every $p \in [p_{\min}, p_{\max}]$ and $x \in \mathbb{R}^n \setminus \{0\}$

$$(4.19) \quad \Delta_p w \geq |x|^{-\gamma(p-1)-p},$$

$$(4.20) \quad \Delta_p v \geq \frac{c_0}{2} \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2}, \quad \text{if } \nabla v(x) \neq 0.$$

Step 2. We now assume that $c_1 > 0$ is arbitrary. We fix $\gamma = \gamma(n, p_{\min}, p_{\max}) \geq 1$ as above, given by (4.18). It is not hard to see that similar computations as those in Step 1, but with $c_1 > 0$ arbitrary, imply that for every $p \in [p_{\min}, p_{\max}]$ and $x \in \mathbb{R}^n \setminus \{0\}$

$$(4.21) \quad \Delta_p w \geq c_1^{p-1} |x|^{-\gamma(p-1)-p},$$

$$(4.22) \quad \Delta_p v \geq \frac{c_0}{2} c_1 \varepsilon |\nabla v|^{p-2} |x|^{-\gamma-2}, \quad \text{if } \nabla v(x) \neq 0.$$

Step 3. We now observe that there holds

$$\nabla v = e_n + \frac{c_0}{2} \varepsilon \nabla w.$$

Then, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$\begin{aligned} \left| |\nabla v| - 1 \right| &= \left| |\nabla v| - |e_n| \right| \leq \left| \nabla v - e_n \right| = \left| \frac{c_0}{2} \varepsilon \nabla w \right| \\ &= \frac{c_0}{2} c_1 \varepsilon \gamma |x|^{-\gamma-1} \leq \frac{c_0}{2} c_1 \varepsilon \bar{r}_1^{-\gamma-1} \leq \frac{1}{2}, \end{aligned}$$

if we let $\varepsilon \leq \bar{\varepsilon}_1 = \bar{\varepsilon}_1(n, p_{\min}, p_{\max}, \bar{r}_1, c_0, c_1)$ and therefore,

$$(4.23) \quad \frac{1}{2} \leq |\nabla v| \leq 2, \quad \text{for } \varepsilon \leq \bar{\varepsilon}_1.$$

So the first assertion in (4.7) follows.

Step 4. We now consider $p(x)$ a Lipschitz continuous function such that $1 < p_{\min} \leq p(x) \leq p_{\max} < \infty$.

We first observe that, for any $R \geq 1$,

$$\begin{aligned} t^{p(x)-1} |\log t| &\leq t^{p_{\min}-1} |\log t| \leq C_1(p_{\min}), \quad \text{if } 0 < t < 1, \\ t^{p(x)-1} |\log t| &\leq t^{p_{\max}-1} |\log t| \leq R^{p_{\max}-1} \log R \quad \text{if } 1 \leq t \leq R, \end{aligned}$$

so that

$$(4.24) \quad t^{p(x)-1} |\log t| \leq C_2(p_{\min}, p_{\max}, R), \quad \text{if } 0 < t \leq R.$$

It then follows from (4.23) and (4.24) that, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.25) \quad |\nabla v|^{p(x)-1} |\log |\nabla v|| \leq C_3(p_{\min}, p_{\max}), \quad \text{if } \varepsilon \leq \bar{\varepsilon}_1.$$

We also have, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$|\nabla w| = c_1 \gamma |x|^{-\gamma-1} \leq c_1 \gamma \bar{r}_1^{-\gamma-1},$$

so using again (4.24), we get, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.26) \quad |\nabla w|^{p(x)-1} |\log |\nabla w|| \leq C_4(n, p_{\min}, p_{\max}, \bar{r}_1, c_1).$$

Step 5. We now assume that $p(x)$ satisfies moreover that $\|\nabla p\|_{L^\infty} \leq \varepsilon^{1+\theta}$, for some $0 < \theta < 1$. Then, from (4.25) we obtain, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.27) \quad \left| |\nabla v|^{p(x)-2} \langle \nabla p(x), \nabla v \rangle \log |\nabla v| \right| \leq |\nabla v|^{p(x)-1} |\log |\nabla v|| \|\nabla p\|_{L^\infty} \leq \varepsilon^{1+\theta} C_3,$$

if $\varepsilon \leq \bar{\varepsilon}_1$. Hence, from (4.22), (4.27) and (4.23), for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.28) \quad \begin{aligned} \Delta_{p(x)} v &= |\nabla v|^{p(x)-2} (\Delta v + (p(x) - 2) \langle D^2 v, \frac{\nabla v}{|\nabla v|}, \frac{\nabla v}{|\nabla v|} \rangle + \langle \nabla p(x), \nabla v \rangle \log |\nabla v|) \\ &\geq \frac{c_0}{2} c_1 \varepsilon |\nabla v|^{p(x)-2} |x|^{-\gamma-2} - \varepsilon^{1+\theta} C_3 \\ &\geq \frac{c_0}{2} c_1 \varepsilon C_5 |x|^{-\gamma-2} - \varepsilon^{1+\theta} C_3 \geq \frac{c_0}{2} c_1 \varepsilon C_5 - \varepsilon^{1+\theta} C_3 = \varepsilon \left(\frac{c_0}{2} c_1 C_5 - \varepsilon^\theta C_3 \right), \end{aligned}$$

if $\varepsilon \leq \bar{\varepsilon}_1$, where we have used that $\bar{r}_2 \leq 1$ and $C_5 = C_5(p_{\min}, p_{\max})$, $C_5 = \min\{(\frac{1}{2})^{p_{\max}-2}, 2^{p_{\min}-2}\}$. We conclude that, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$\Delta_{p(x)} v \geq \varepsilon \left(\frac{c_0}{2} c_1 C_5 - \varepsilon^\theta C_3 \right) \geq \varepsilon \frac{c_0}{4} c_1 C_5 > \varepsilon^2,$$

if moreover $\varepsilon \leq \tilde{\varepsilon}_1 = \tilde{\varepsilon}_1(p_{\min}, p_{\max}, c_0, c_1, \theta)$. That is, the second assertion in (4.7) follows.

Finally, from (4.26) we obtain, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.29) \quad \left| |\nabla w|^{p(x)-2} \langle \nabla p(x), \nabla w \rangle \log |\nabla w| \right| \leq |\nabla w|^{p(x)-1} |\log |\nabla w|| \|\nabla p\|_{L^\infty} \leq \varepsilon^{1+\theta} C_4.$$

Hence, from (4.21) and (4.29), for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$(4.30) \quad \begin{aligned} \Delta_{p(x)} w &= |\nabla w|^{p(x)-2} (\Delta w + (p(x) - 2) \langle D^2 w, \frac{\nabla w}{|\nabla w|}, \frac{\nabla w}{|\nabla w|} \rangle + \langle \nabla p(x), \nabla w \rangle \log |\nabla w|) \\ &\geq c_1^{p(x)-1} |x|^{-\gamma(p(x)-1)-p(x)} - \varepsilon^{1+\theta} C_4 \geq 2\bar{c} - \varepsilon C_4, \end{aligned}$$

if $\varepsilon \leq 1$. Here we have used that $\bar{r}_2 \leq 1$ and we have denoted $\bar{c} = \bar{c}(p_{\min}, p_{\max}, c_1) = \frac{1}{2} \min\{c_1^{p_{\min}-1}, c_1^{p_{\max}-1}\}$. We conclude that, for $\bar{r}_1 \leq |x| \leq \bar{r}_2$,

$$\Delta_{p(x)} w \geq \bar{c},$$

if $\varepsilon \leq \varepsilon_0 = \varepsilon_0(n, p_{\min}, p_{\max}, \bar{r}_1, c_1)$. This proves (4.6) and finishes the proof. \square

5. GEOMETRIC REGULARITY RESULTS

In this section we prove a Harnack type inequality for a solution u to problem (1.1), following the approach in [D]. We will argue assuming that

$$(5.1) \quad \|f\|_{L^\infty(\Omega)} \leq \varepsilon^2, \quad \|g - 1\|_{L^\infty(\Omega)} \leq \varepsilon^2, \quad \|\nabla p\|_{L^\infty(\Omega)} \leq \varepsilon^{1+\theta}, \quad \|p - p_0\|_{L^\infty(\Omega)} \leq \varepsilon,$$

holds, for $0 < \varepsilon < 1$, for some constant $0 < \theta \leq 1$.

The proof of Harnack inequality is based on the following lemma.

Lemma 5.1. *Let u be a solution to (1.1)–(5.1) in B_1 . There exists a universal constant $\bar{\varepsilon}$ such that if $0 < \varepsilon \leq \bar{\varepsilon}$ and u satisfies*

$$(5.2) \quad q^+(x) \leq u(x) \leq (q(x) + \varepsilon)^+, \quad x \in B_1, \quad q(x) = x_n + \sigma, \quad |\sigma| < \frac{1}{20},$$

and in $x_0 = \frac{1}{10}e_n$,

$$u(x_0) \geq (q(x_0) + \frac{\varepsilon}{2})^+,$$

then

$$(5.3) \quad u \geq (q + c\varepsilon)^+ \quad \text{in } \overline{B}_{\frac{1}{2}},$$

for some universal $0 < c < 1$. Analogously, if

$$(5.4) \quad u(x_0) \leq (q(x_0) + \frac{\varepsilon}{2})^+,$$

then

$$(5.5) \quad u \leq (q + (1 - c)\varepsilon)^+ \quad \text{in } \overline{B}_{\frac{1}{2}}.$$

Proof. The proof follows the original one in [D] adapted by the dichotomy discussed in [LR]. We will prove the first statement.

From (5.2) we have that $u \geq q$ in B_1 .

We also notice that $B_{1/20}(x_0) \subset B_1^+(u)$. Then,

$$(5.6) \quad \Delta_{p(x)} u = f \quad \text{in } B_{1/20}(x_0).$$

Thus, by Theorem 1.1 in [Fa], $u \in C^{1,\alpha}$ in $\overline{B}_{1/40}(x_0)$, where $\alpha = \alpha(p_{\min}, p_{\max}, n) \in (0, 1)$ and $\|u\|_{C^{1,\alpha}(\overline{B}_{1/40}(x_0))} \leq C$, with $C = C(p_{\min}, p_{\max}, n) \geq 1$. Here we have used (5.1) and also that (5.2) implies that $\|u\|_{L^\infty(B_1)} \leq 3$.

We will consider two cases:

Case (i). Suppose $|\nabla u(x_0)| < \frac{1}{4}$. We choose $r_1 > 0$, $r_1 = r_1(p_{\min}, p_{\max}, n) \leq 1/40$ such that $|\nabla u(x)| \leq \frac{1}{2}$ in $B_{r_1}(x_0)$. In addition, there exists a constant $0 < r_2 = r_2(r_1) = r_2(p_{\min}, p_{\max}, n) < r_1$ such that $(x - r_2 e_n) \in B_{r_1}(x_0)$, for every $x \in B_{r_1/2}(x_0)$. We observe that $\tilde{v} = u - q$ satisfies

$$(5.7) \quad \operatorname{div}(|\nabla \tilde{v} + e_n|^{p(x)-2}(\nabla \tilde{v} + e_n)) = f \quad \text{in } B_{\frac{1}{20}}(x_0).$$

We now apply Lemma 4.1 to the function $\tilde{v} = u - q$ in $B_{4r_3}(x_0)$, where $r_3 = \min\{\frac{r_1}{4}, \frac{r_2}{8}\}$. In particular we obtain from (4.2) that

$$u(x) - q(x) \geq C^{-1}(u(x_0) - q(x_0)) - r_3 \geq \frac{\varepsilon}{2C} - r_3,$$

for $x \in B_{r_3}(x_0)$. Here $C = C(n, p_{\min}, p_{\max})$ is a universal constant because $\|f\|_{L^\infty(B_1)} \leq \varepsilon^2$, see (5.1), and $\|\tilde{v}\|_{L^\infty(B_1)} \leq 2$.

On the other hand, for all $x \in B_{r_3}(x_0)$ we obtain

$$\begin{aligned} \frac{\varepsilon}{2C} - r_3 &\leq u(x) - q(x) = u((x - r_2 e_n) + r_2 e_n) - q((x - r_2 e_n) + r_2 e_n) \\ &= u((x - r_2 e_n) + r_2 e_n) - q(x - r_2 e_n) - r_2 \leq u(x - r_2 e_n) - q(x - r_2 e_n) + \frac{r_2}{2} - r_2. \end{aligned}$$

As a consequence, denoting $c_0 = C^{-1}$ and $\bar{x}_0 := x_0 - r_2 e_n$, we get for all $x \in B_{r_3}(\bar{x}_0)$

$$(5.8) \quad \frac{c_0}{2} \varepsilon = \frac{\varepsilon}{2C} \leq \frac{\varepsilon}{2C} - r_3 + \frac{r_2}{2} = \frac{\varepsilon}{2C} - r_3 - \frac{r_2}{2} + r_2 \leq u(x) - q(x).$$

Let us define the function $w : \bar{D} \rightarrow \mathbb{R}$, $D := B_{\frac{4}{5}}(\bar{x}_0) \setminus \bar{B}_{r_3}(\bar{x}_0)$ as

$$w(x) = c \left(|x - \bar{x}_0|^{-\gamma} - \left(\frac{4}{5}\right)^{-\gamma} \right),$$

for $\gamma = \gamma(n, p_{\min}, p_{\max}) \geq 1$ given in Lemma 4.2 (see (4.18)). We choose $c = c(n, p_{\min}, p_{\max}) > 0$ in such a way that

$$w = \begin{cases} 0, & \text{on } \partial B_{\frac{4}{5}}(\bar{x}_0) \\ 1, & \text{on } \partial B_{r_3}(\bar{x}_0). \end{cases}$$

As usual, we define for every $x \in \bar{B}_{\frac{4}{5}}(\bar{x}_0)$

$$v(x) = q(x) + c_0 \frac{\varepsilon}{2} (w(x) - 1)$$

and for $t \geq 0$ we set

$$v_t(x) = v(x) + t, \quad x \in \bar{B}_{\frac{4}{5}}(\bar{x}_0).$$

We extend w to 1 in $B_{r_3}(\bar{x}_0)$, so that it results

$$v_0(x) = v(x) \leq q(x) \leq u(x), \quad x \in \bar{B}_{\frac{4}{5}}(\bar{x}_0).$$

Let

$$\bar{t} = \sup\{t \geq 0 : v_t \leq u \text{ in } \bar{B}_{\frac{4}{5}}(\bar{x}_0)\}.$$

Claim: $\bar{t} \geq \frac{c_0 \varepsilon}{2}$.

Assuming that the previous Claim holds, we obtain from the definition of v that, in $B_{\frac{4}{5}}(\bar{x}_0)$, the inequality

$$u(x) \geq v(x) + \bar{t} \geq q(x) + \frac{c_0 \varepsilon}{2} w(x)$$

is satisfied.

On the other hand, $B_{\frac{1}{2}} \subset B_{\frac{3}{5}}(\bar{x}_0)$ and since

$$w(x) \geq \begin{cases} c \left(\left(\frac{3}{5}\right)^{-\gamma} - \left(\frac{4}{5}\right)^{-\gamma} \right), & B_{\frac{3}{5}}(\bar{x}_0) \setminus B_{r_3}(\bar{x}_0), \\ 1, & B_{r_3}(\bar{x}_0), \end{cases}$$

we conclude that, in $B_{\frac{1}{2}}$,

$$u(x) - q(x) \geq c_1 \varepsilon,$$

with $0 < c_1 = c_1(n, p_{\min}, p_{\max}) < 1$ universal, as desired.

We now have to prove the Claim. We argue by contradiction assuming that $\bar{t} < \frac{c_0 \varepsilon}{2}$. Let $y_0 \in \bar{B}_{\frac{4}{5}}(\bar{x}_0)$ be the contact point between $v_{\bar{t}}$ and u , where

$$v_{\bar{t}}(y_0) = u(y_0).$$

We will prove that $y_0 \in \overline{B_{r_3}}(\bar{x}_0)$. In fact, recalling that w vanishes on $\partial B_{\frac{4}{5}}(\bar{x}_0)$ and from the definition of $v_{\bar{t}}$, we obtain

$$v_{\bar{t}} = q - \frac{c_0}{2}\varepsilon + \bar{t} < u \quad \text{on } \partial B_{\frac{4}{5}}(\bar{x}_0),$$

because $u \geq q$ and $\bar{t} < \frac{c_0\varepsilon}{2}$.

We can apply Lemma 4.2 to v . Hence, there exists $\varepsilon_1 = \varepsilon_1(n, p_{\min}, p_{\max}, \theta)$ a universal constant such that

$$\frac{1}{2} \leq |\nabla v_{\bar{t}}| = |\nabla v| \leq 2,$$

$$\Delta_{p(x)} v_{\bar{t}} = \Delta_{p(x)} v > \varepsilon^2 \geq f,$$

for every $0 < \varepsilon \leq \varepsilon_1$ and for every $x \in D = B_{\frac{4}{5}}(\bar{x}_0) \setminus \overline{B_{r_3}}(\bar{x}_0)$.

On the other hand, from the definition of $v_{\bar{t}}$, we have

$$(5.9) \quad |\nabla v_{\bar{t}}| \geq |(v_{\bar{t}})_n| = |1 + \frac{c_0}{2}\varepsilon w_n|,$$

where $(v_{\bar{t}})_n$ and w_n denote the partial derivatives with respect to x_n of $v_{\bar{t}}$ and w .

Let us show that $w_n > \hat{c}$ in $\{v_{\bar{t}} \leq 0\} \cap D$, for $\hat{c} > 0$ universal.

In fact, whenever $0 < \varepsilon \leq \varepsilon_2$, for ε_2 universal, we have

$$\{v_{\bar{t}} \leq 0\} \cap D \subset \{q \leq \frac{c_0\varepsilon}{2}\} = \{x_n \leq \frac{c_0\varepsilon}{2} - \sigma\} \subset \{x_n \leq \frac{5}{80}\}.$$

On the other hand,

$$\nabla w = -\gamma c |x - \bar{x}_0|^{-\gamma-2} (x - \bar{x}_0) = -\gamma c |x - \bar{x}_0|^{-\gamma-1} \frac{x - \bar{x}_0}{|x - \bar{x}_0|}.$$

Moreover, denoting $\nu_x = \frac{x - \bar{x}_0}{|x - \bar{x}_0|}$, we observe that, in $\{v_{\bar{t}} \leq 0\} \cap D$, we have $-\langle \nu_x, e_n \rangle > 0$ since

$$x_n - (\bar{x}_0)_n = x_n - \frac{1}{10} + r_2 \leq -\frac{1}{80} \quad \text{in } \{x_n \leq \frac{5}{80}\}.$$

In particular, there holds in $\{v_{\bar{t}} \leq 0\} \cap D$

$$w_n = \langle \nabla w, e_n \rangle = -\gamma \langle \nu_x, e_n \rangle c |x - \bar{x}_0|^{-\gamma-1} \geq c \gamma \frac{1}{80} \frac{5}{4} \left(\frac{4}{5}\right)^{-1-\gamma} = \hat{c} > 0.$$

Thus, from (5.9) we deduce that

$$|\nabla v_{\bar{t}}| \geq 1 + \frac{c_0}{2}\varepsilon w_n \geq 1 + \frac{c_0}{2}\hat{c}\varepsilon$$

in $\{v_{\bar{t}} \leq 0\} \cap D$, which implies, for ε sufficiently small,

$$|\nabla v_{\bar{t}}| > 1 + \varepsilon^2 \geq g,$$

on $F(v_{\bar{t}}) \cap D$. Then $v_{\bar{t}}$ is a strict subsolution to (1.1) in D touching u at y_0 . Hence $y_0 \in \overline{B_{r_3}}(\bar{x}_0)$ and this generates a contradiction with (5.8), because

$$u(y_0) = v_{\bar{t}}(y_0) = v(y_0) + \bar{t} = q(y_0) + \bar{t} < q(y_0) + c_0\varepsilon.$$

Case (ii). Now suppose $|\nabla u(x_0)| \geq \frac{1}{4}$. By exploiting the $C^{1,\alpha}$ regularity of u in $\overline{B_{\frac{1}{40}}}(x_0)$, we know that u is Lipschitz continuous in $\overline{B_{\frac{1}{40}}}(x_0)$, as well as there exist a constant $0 < r_0 = r_0(n, p_{\min}, p_{\max})$, with $8r_0 \leq \frac{1}{40}$, and $C = C(n, p_{\min}, p_{\max}) > 1$ such that

$$\frac{1}{8} \leq |\nabla u| \leq C \quad \text{in } B_{8r_0}(x_0).$$

In addition, since (5.6) holds, it follows by Proposition 3.4, that $u \in W^{2,n}(B_{4r_0}(x_0))$ and it is a solution to the linear uniformly elliptic equation

$$\mathcal{L}h = f \quad \text{in } B_{4r_0}(x_0),$$

where

$$\begin{aligned} \mathcal{L}h &= \text{Tr}(A(x)D^2h(x)) + \langle b, \nabla h(x) \rangle, \\ A(x) &:= |\nabla u|^{p(x)-2} \left(I + (p(x) - 2) \frac{\nabla u(x)}{|\nabla u(x)|} \otimes \frac{\nabla u(x)}{|\nabla u(x)|} \right), \end{aligned}$$

and

$$b(x) := |\nabla u|^{p(x)-2} \log |\nabla u(x)| \nabla p(x).$$

Hence $A \in C^{0,\alpha}(\overline{B}_{4r_0}(x_0))$, $b \in C(\overline{B}_{4r_0}(x_0))$ and \mathcal{L} has universal ellipticity constants (depending only on n, p_{\min}, p_{\max}). Moreover, $\|b\|_{L^\infty(B_{4r_0}(x_0))} \leq C\varepsilon^{1+\theta}$, C universal, because $\|\nabla p\|_{L^\infty(B_1)} \leq \varepsilon^{1+\theta}$ (see (5.1)).

In this way, we conclude that $u - q$ satisfies

$$\text{Tr}(A(x)D^2h(x)) + \langle b, \nabla h(x) \rangle = f - \langle b, e_n \rangle \quad \text{in } B_{4r_0}(x_0).$$

Then, applying Harnack's inequality (see, for instance, [GT], Chap. 9) and recalling again (5.1), we obtain

$$\begin{aligned} (5.10) \quad u(x) - q(x) &\geq C_1(u(x_0) - q(x_0)) - C_2(\|f\|_{L^\infty(B_{4r_0}(x_0))} + \|b\|_{L^\infty(B_{4r_0}(x_0))}) \\ &\geq C_1 \frac{\varepsilon}{2} - C_2(\varepsilon^2 + C\varepsilon^{1+\theta}) \geq \frac{c_0}{2}\varepsilon, \end{aligned}$$

for every $x \in B_{r_0}(x_0)$, for $0 < \varepsilon \leq \varepsilon_3$. Here ε_3 , C_1 , C_2 and c_0 are positive universal constants. At this point, we can repeat the same argument of Case (i) around the point x_0 , considering the annulus $B_{\frac{4}{5}}(x_0) \setminus \overline{B}_{r_0}(x_0)$. This completes the proof. \square

The next result is the main tool in Theorem 1.1.

Theorem 5.2 (Harnack inequality). *There exists a universal constant $\bar{\varepsilon}$, such that if u solves (1.1)–(5.1), and for some point $x_0 \in \Omega^+(u) \cup F(u)$,*

$$(5.11) \quad (x_n + a_0)^+ \leq u(x) \leq (x_n + b_0)^+ \quad \text{in } B_r(x_0) \subset \Omega,$$

with

$$b_0 - a_0 \leq \varepsilon r, \quad \varepsilon \leq \bar{\varepsilon},$$

then

$$(x_n + a_1)^+ \leq u(x) \leq (x_n + b_1)^+ \quad \text{in } B_{r/40}(x_0),$$

with

$$a_0 \leq a_1 \leq b_1 \leq b_0, \quad b_1 - a_1 \leq (1 - c)\varepsilon r,$$

and $0 < c < 1$ universal.

Proof. Assume without loss of generality that $x_0 = 0, r = 1$.

We call $q(x) = x_n + a_0$. Assumption (5.11) gives that

$$(5.12) \quad q^+(x) \leq u(x) \leq (q(x) + \varepsilon)^+ \quad \text{in } B_1,$$

since $b_0 \leq a_0 + \varepsilon$. We distinguish three cases.

Case 1. $|a_0| < 1/20$. We now distinguish two cases: $u(\hat{x}_0) \geq (q(\hat{x}_0) + \frac{\varepsilon}{2})^+$ or $u(\hat{x}_0) \leq (q(\hat{x}_0) + \frac{\varepsilon}{2})^+$, where $\hat{x}_0 = \frac{1}{10}e_n$.

Assume that

$$u(\hat{x}_0) \geq (q(\hat{x}_0) + \frac{\varepsilon}{2})^+, \quad \hat{x}_0 = \frac{1}{10}e_n,$$

(the other case is treated similarly). Then, by Lemma 5.1, if $\varepsilon \leq \bar{\varepsilon}$,

$$(q(x) + c\varepsilon)^+ \leq u(x) \quad \text{in } \overline{B_{\frac{1}{2}}},$$

for $0 < c < 1$ universal, which gives the desired improvement.

Case 2. $a_0 \leq -1/20$. In this case it follows from (5.12) that, for $\varepsilon < 1/40$, 0 belongs to the zero phase of $(q(x) + \varepsilon)^+$, which implies that 0 belongs to the zero phase of u . A contradiction.

Case 3. $a_0 \geq 1/20$. In this case it follows from (5.11) that

$$B_{1/20} \subset B_1^+(u).$$

Then, denoting $\hat{u} = u - a_0$, we have

$$(5.13) \quad \Delta_{p(x)} u = \Delta_{p(x)} \hat{u} = f \quad \text{in } B_{1/20}.$$

Observing that $\|\hat{u}\|_{L^\infty(B_1)} \leq 2$ and recalling (5.1), we obtain from the application of Theorem 1.1 in [Fa] to \hat{u} , that $u \in C^{1,\alpha}$ in $\overline{B_{1/40}}$, where $\alpha = \alpha(p_{\min}, p_{\max}, n) \in (0, 1)$ and $\|\nabla u\|_{C^\alpha(\overline{B_{1/40}})} \leq C$, with $C = C(p_{\min}, p_{\max}, n) \geq 1$.

We now distinguish two cases: $u(0) - q(0) \geq \frac{\varepsilon}{2}$ or $u(0) - q(0) \leq \frac{\varepsilon}{2}$.

Assume that

$$u(0) - q(0) \geq \frac{\varepsilon}{2},$$

(the other case is treated similarly). We will proceed as in the proof of Lemma 5.1.

If $|\nabla u(0)| < \frac{1}{4}$, we argue as in Case (i) of Lemma 5.1, taking $\bar{x}_0 = -r_2 e_n$. Here $r_2 > 0$ is universal, chosen as in that lemma, and such that we also have

$$B_{1/40} \subset\subset B_{r_4}(\bar{x}_0) \subset\subset B_{1/20},$$

for an appropriate chosen universal $r_4 > 0$. We now take r_3 universal as in Lemma 5.1, let

$$D := B_{r_4}(\bar{x}_0) \setminus \overline{B_{r_3}(\bar{x}_0)},$$

and define w in D as in that lemma. Then, arguing as in that proof, we obtain

$$(5.14) \quad u(x) - q(x) \geq c_1 \varepsilon \quad \text{in } B_{1/40},$$

with $0 < c_1 < 1$, if $\varepsilon \leq \bar{\varepsilon}$, $\bar{\varepsilon}$ and c_1 universal.

If $|\nabla u(0)| \geq \frac{1}{4}$, we proceed as in Case (ii) of Lemma 5.1 and we consider the barrier w in

$$D := B_{1/20} \setminus \overline{B_{r_0}},$$

with $r_0 > 0$ universal and small. We obtain again (5.14), thus completing the proof. \square

From Theorem 5.2, with the same arguments employed in [D], we obtain the following estimate that will be crucial in the improvement of flatness procedure.

Corollary 5.3. *Let u be as in Theorem 5.2 satisfying (5.11) for $r = 1$. Then in $B_1(x_0)$, $\tilde{u}_\varepsilon(x) = \frac{u(x) - x_n}{\varepsilon}$ has a Hölder modulus of continuity at x_0 , outside the ball of radius $\varepsilon/\bar{\varepsilon}$, i.e., for all $x \in (\Omega^+(u) \cup F(u)) \cap B_1(x_0)$, with $|x - x_0| \geq \varepsilon/\bar{\varepsilon}$,*

$$|\tilde{u}_\varepsilon(x) - \tilde{u}_\varepsilon(x_0)| \leq C|x - x_0|^\gamma.$$

Here $\bar{\varepsilon}$ is as in Theorem 5.2, and C and $0 < \gamma < 1$ are universal.

6. IMPROVEMENT OF FLATNESS

In this section we present the main improvement of flatness lemma. Theorem 1.1 will then be obtained by applying this lemma in an iterative way.

Lemma 6.1 (Improvement of flatness). *Let u satisfy (1.1) in B_1 and*

$$(6.1) \quad \|f\|_{L^\infty(B_1)} \leq \varepsilon^2, \quad \|g-1\|_{L^\infty(B_1)} \leq \varepsilon^2, \quad \|\nabla p\|_{L^\infty(B_1)} \leq \varepsilon^{1+\theta}, \quad \|p-p_0\|_{L^\infty(B_1)} \leq \varepsilon,$$

for $0 < \varepsilon < 1$, for some constant $0 < \theta \leq 1$. Suppose that

$$(6.2) \quad (x_n - \varepsilon)^+ \leq u(x) \leq (x_n + \varepsilon)^+ \quad \text{in } B_1, \quad 0 \in F(u).$$

If $0 < r \leq r_0$ for r_0 universal, and $0 < \varepsilon \leq \varepsilon_0$ for some ε_0 depending on r , then

$$(6.3) \quad (x \cdot \nu - r\varepsilon/2)^+ \leq u(x) \leq (x \cdot \nu + r\varepsilon/2)^+ \quad \text{in } B_r,$$

with $|\nu| = 1$ and $|\nu - e_n| \leq \tilde{C}\varepsilon$ for a universal constant \tilde{C} .

Proof. We divide the proof of this lemma into 3 steps. We will use the following notation:

$$\Omega_\rho(u) := (B_1^+(u) \cup F(u)) \cap B_\rho.$$

Step 1: Compactness. Fix $r \leq r_0$ with r_0 universal (the precise r_0 will be given in Step 3). Assume by contradiction that we can find a sequence $\varepsilon_k \rightarrow 0$ and a sequence u_k of solutions to (1.1) in B_1 with right hand side f_k , exponent p_k and free boundary condition g_k satisfying (6.1) with $\varepsilon = \varepsilon_k$, such that u_k satisfies (6.2), i.e.,

$$(6.4) \quad (x_n - \varepsilon_k)^+ \leq u_k(x) \leq (x_n + \varepsilon_k)^+ \quad \text{for } x \in B_1, \quad 0 \in F(u_k),$$

but u_k does not satisfy the conclusion (6.3) of the lemma.

Set

$$\tilde{u}_k(x) = \frac{u_k(x) - x_n}{\varepsilon_k}, \quad x \in \Omega_1(u_k).$$

Then, (6.4) gives

$$(6.5) \quad -1 \leq \tilde{u}_k(x) \leq 1 \quad \text{for } x \in \Omega_1(u_k).$$

From Corollary 5.3, it follows that the function \tilde{u}_k satisfies

$$(6.6) \quad |\tilde{u}_k(x) - \tilde{u}_k(y)| \leq C|x - y|^\gamma,$$

for C and $0 < \gamma < 1$ universal and

$$|x - y| \geq \varepsilon_k/\bar{\varepsilon}, \quad x, y \in \Omega_{1/2}(u_k).$$

From (6.4) it clearly follows that $F(u_k)$ converges to $B_1 \cap \{x_n = 0\}$ in the Hausdorff distance. This fact and (6.6) together with Ascoli-Arzelà give that, as $\varepsilon_k \rightarrow 0$, the graphs of the \tilde{u}_k over $\Omega_{1/2}(u_k)$ converge (up to a subsequence) in the Hausdorff distance to the graph of a Hölder continuous function \tilde{u} over $B_{1/2} \cap \{x_n \geq 0\}$.

Step 2: Limiting Solution. We now show that \tilde{u} solves the following linearized problem

$$(6.7) \quad \begin{cases} \mathcal{L}_{p_0} \tilde{u} = 0 & \text{in } B_{1/2} \cap \{x_n > 0\}, \\ \tilde{u}_n = 0 & \text{on } B_{1/2} \cap \{x_n = 0\}, \end{cases}$$

in the sense of Definition 2.7. Here \mathcal{L}_{p_0} is as in (2.2).

Let $P(x)$ be a quadratic polynomial touching \tilde{u} at $\bar{x} \in B_{1/2} \cap \{x_n \geq 0\}$ strictly from below. We need to show that

- (i) if $\bar{x} \in B_{1/2} \cap \{x_n > 0\}$ then $\mathcal{L}_{p_0} P \leq 0$;
- (ii) if $\bar{x} \in B_{1/2} \cap \{x_n = 0\}$ then $P_n(\bar{x}) \leq 0$.

Since $\tilde{u}_k \rightarrow \tilde{u}$ in the sense specified above, there exist points $x_k \in \Omega_{1/2}(u_k)$, $x_k \rightarrow \bar{x}$ and constants $c_k \rightarrow 0$ such that

$$(6.8) \quad \tilde{u}_k(x_k) = P(x_k) + c_k$$

and

$$(6.9) \quad \tilde{u}_k \geq P + c_k \quad \text{in a neighborhood of } x_k.$$

From the definition of \tilde{u}_k , (6.8) and (6.9) read

$$u_k(x_k) = Q_k(x_k)$$

and

$$u_k(x) \geq Q_k(x) \quad \text{in a neighborhood of } x_k,$$

where

$$Q_k(x) = \varepsilon_k(P(x) + c_k) + x_n.$$

For notational simplicity we will drop the sub-index k from Q_k .

We first notice that

$$(6.10) \quad \nabla Q = \varepsilon_k \nabla P + e_n,$$

thus,

$$(6.11) \quad \nabla Q(x_k) \neq 0, \quad \text{for } k \text{ large.}$$

We now distinguish two cases.

(i) If $\bar{x} \in B_{1/2} \cap \{x_n > 0\}$ then $x_k \in B_{1/2}^+(u_k)$ (for k large). Since Q touches u_k from below at x_k , and $\nabla Q(x_k) \neq 0$, we get

$$\begin{aligned} \varepsilon_k^2 &\geq f_k(x_k) \\ &\geq \Delta_{p_k(x_k)} Q(x_k) \\ &= |\nabla Q(x_k)|^{p_k(x_k)-2} \Delta Q + |\nabla Q(x_k)|^{p_k(x_k)-4} (p_k(x_k) - 2) \sum_{i,j=1}^n Q_{x_i}(x_k) Q_{x_j}(x_k) Q_{x_i x_j} \\ &\quad + |\nabla Q(x_k)|^{p_k(x_k)-2} \langle \nabla p_k(x_k), \nabla Q(x_k) \rangle \log |\nabla Q(x_k)| \\ &= \varepsilon_k |\nabla Q(x_k)|^{p_k(x_k)-2} \Delta P + \varepsilon_k |\nabla Q(x_k)|^{p_k(x_k)-4} (p_k(x_k) - 2) \sum_{i,j=1}^n Q_{x_i}(x_k) Q_{x_j}(x_k) P_{x_i x_j} \\ &\quad + |\nabla Q(x_k)|^{p_k(x_k)-2} \langle \nabla p_k(x_k), \nabla Q(x_k) \rangle \log |\nabla Q(x_k)|. \end{aligned}$$

Using that $|\nabla p_k(x_k)| \leq \varepsilon_k$, we obtain

$$\begin{aligned} \varepsilon_k &\geq |\nabla Q(x_k)|^{p_k(x_k)-2} \Delta P + |\nabla Q(x_k)|^{p_k(x_k)-4} (p_k(x_k) - 2) \sum_{i,j=1}^n Q_{x_i}(x_k) Q_{x_j}(x_k) P_{x_i x_j} \\ &\quad - |\nabla Q(x_k)|^{p_k(x_k)-1} |\log |\nabla Q(x_k)||. \end{aligned}$$

Now, passing to the limit $k \rightarrow \infty$ and recalling that

$$\nabla Q(x_k) \rightarrow e_n, \quad p_k(x_k) \rightarrow p_0, \quad \varepsilon_k \rightarrow 0,$$

we conclude that $\mathcal{L}_{p_0} P \leq 0$ as desired.

(ii) If $\bar{x} \in B_{1/2} \cap \{x_n = 0\}$, as observed in Remark 2.8, we can assume that $\mathcal{L}_{p_0}P > 0$. We claim that for k large enough, $x_k \in F(u_k)$. Otherwise $x_{k_j} \in B_{1/2}^+(u_{k_j})$ for a subsequence $k_j \rightarrow \infty$ and as in case (i), passing to the limit, we get

$$\mathcal{L}_{p_0}P \leq 0,$$

a contradiction. Thus, $x_k \in F(u_k)$ for k large.

Since Q^+ touches u_k from below at $x_k \in F(u_k)$ and (6.11) holds

$$|\nabla Q(x_k)| \leq g_k(x_k) \leq 1 + \varepsilon_k^2,$$

which, by (6.10), gives

$$|\nabla Q(x_k)|^2 = \varepsilon_k^2 |\nabla P(x_k)|^2 + 1 + 2\varepsilon_k P_n(x_k) \leq 1 + 3\varepsilon_k^2.$$

Thus, after division by ε_k ,

$$\varepsilon_k |\nabla P(x_k)|^2 - 3\varepsilon_k + 2P_n(x_k) \leq 0.$$

Passing to the limit as $k \rightarrow \infty$, we obtain $P_n(\bar{x}) \leq 0$ as desired.

Step 3: Improvement of flatness. From the previous step, \tilde{u} solves (6.7) and from (6.5),

$$-1 \leq \tilde{u}(x) \leq 1 \quad \text{in } B_{1/2} \cap \{x_n \geq 0\}.$$

From Theorem 2.9 and the bound above we find that, for the given r ,

$$|\tilde{u}(x) - \tilde{u}(0) - \nabla \tilde{u}(0) \cdot x| \leq C_0 r^2 \quad \text{in } B_r \cap \{x_n \geq 0\},$$

if $r_0 \leq 1/4$, for a universal constant C_0 . In particular, since $\tilde{u}(0) = 0$ and also $\tilde{u}_n(0) = 0$, we obtain

$$x' \cdot \tilde{\nu} - C_0 r^2 \leq \tilde{u}(x) \leq x' \cdot \tilde{\nu} + C_0 r^2 \quad \text{in } B_r \cap \{x_n \geq 0\},$$

where $x' = (x_1, \dots, x_{n-1})$, $\tilde{\nu} = \nabla_{x'} \tilde{u}(0)$ and $|\tilde{\nu}| \leq C_0$. Therefore, for k large enough we get

$$x' \cdot \tilde{\nu} - C_1 r^2 \leq \tilde{u}_k(x) \leq x' \cdot \tilde{\nu} + C_1 r^2 \quad \text{in } \Omega_r(u_k),$$

for a universal constant C_1 . From the definition of \tilde{u}_k the inequality above reads

$$(6.12) \quad \varepsilon_k x' \cdot \tilde{\nu} + x_n - \varepsilon_k C_1 r^2 \leq u_k(x) \leq \varepsilon_k x' \cdot \tilde{\nu} + x_n + \varepsilon_k C_1 r^2 \quad \text{in } \Omega_r(u_k).$$

We next set

$$\nu_k = \frac{1}{\sqrt{1 + \varepsilon_k^2 |\tilde{\nu}|^2}} (e_n + \varepsilon_k(\tilde{\nu}, 0)).$$

Then,

$$|\nu_k| = 1, \quad |\nu_k - e_n| \leq \tilde{C} \varepsilon_k,$$

and

$$\nu_k = e_n + \varepsilon_k(\tilde{\nu}, 0) + \varepsilon_k^2 \tau, \quad |\tau| \leq \tilde{C},$$

with \tilde{C} universal. We now deduce from (6.12)

$$x \cdot \nu_k - \varepsilon_k^2 \tilde{C} r - \varepsilon_k C_1 r^2 \leq u_k(x) \leq x \cdot \nu_k + \varepsilon_k^2 \tilde{C} r + \varepsilon_k C_1 r^2 \quad \text{in } \Omega_r(u_k).$$

If we fix r_0 satisfying $C_1 r_0 \leq 1/4$ and we take k large enough so that $\varepsilon_k \tilde{C} \leq 1/4$, we get

$$x \cdot \nu_k - \varepsilon_k r/2 \leq u_k(x) \leq x \cdot \nu_k + \varepsilon_k r/2 \quad \text{in } \Omega_r(u_k).$$

Recalling (6.4), we obtain for large k

$$(x \cdot \nu_k - \varepsilon_k r/2)^+ \leq u_k(x) \leq (x \cdot \nu_k + \varepsilon_k r/2)^+ \quad \text{in } B_r,$$

thus u_k satisfies the conclusion (6.3) of the lemma, a contradiction. \square

7. REGULARITY OF THE FREE BOUNDARY

In this section we finally prove our main result, namely, Theorem 1.1.

Proof of Theorem 1.1. Let u be a viscosity solution to (1.1) in B_1 with $0 \in F(u)$, $g(0) = 1$ and $p(0) = p_0$. Consider the sequence

$$u_k(x) = \frac{1}{\rho_k} u(\rho_k x), \quad x \in B_1,$$

with $\rho_k = \bar{r}^k$, $k = 0, 1, \dots$, for a fixed \bar{r} such that

$$\bar{r}^\beta \leq 1/4, \quad \bar{r} \leq r_0,$$

with r_0 the universal constant in Lemma 6.1, taking $\theta = 1$ in (6.1).

Each u_k is a solution to (1.1) with right hand side $f_k(x) = \rho_k f(\rho_k x)$, exponent $p_k(x) = p(\rho_k x)$, and free boundary condition $g_k(x) = g(\rho_k x)$. For the chosen \bar{r} , by taking $\bar{\varepsilon} = \varepsilon_0(\bar{r})^2$, the assumption (6.1) holds for $\varepsilon = \varepsilon_k = 2^{-k} \varepsilon_0(\bar{r})$. Indeed, in B_1 , in view of (1.4),

$$\begin{aligned} |f_k(x)| &\leq \|f\|_\infty \rho_k \leq \bar{\varepsilon} \bar{r}^k \leq \varepsilon_k^2, \\ |g_k(x) - 1| &= |g(\rho_k x) - g(0)| \leq [g]_{0,\beta} \rho_k^\beta \leq \bar{\varepsilon} \bar{r}^{k\beta} \leq \varepsilon_k^2, \\ |\nabla p_k(x)| &\leq \|\nabla p\|_\infty \rho_k \leq \bar{\varepsilon} \bar{r}^k \leq \varepsilon_k^2, \\ |p_k(x) - p_0| &= |p(\rho_k x) - p(0)| \leq \|\nabla p\|_\infty \rho_k \leq \bar{\varepsilon} \bar{r}^k \leq \varepsilon_k^2. \end{aligned}$$

The hypothesis (1.3) guarantees that for $k = 0$ also the flatness assumption (6.2) in Lemma 6.1 is satisfied by u_0 . Then it easily follows, by applying inductively Lemma 6.1, that each u_k is ε_k -flat in B_1 in the sense of (6.2), in the direction ν_k , with $|\nu_k| = 1$, $|\nu_k - \nu_{k+1}| \leq \tilde{C} \varepsilon_k$ ($\nu_0 = e_n$). Now, a standard iteration argument gives the desired statement. \square

APPENDIX A. LEBESGUE AND SOBOLEV SPACES WITH VARIABLE EXPONENT

Let $p : \Omega \rightarrow [1, \infty)$ be a measurable bounded function, called a variable exponent on Ω , and denote $p_{\max} = \text{esssup } p(x)$ and $p_{\min} = \text{essinf } p(x)$. The variable exponent Lebesgue space $L^{p(\cdot)}(\Omega)$ is defined as the set of all measurable functions $u : \Omega \rightarrow \mathbb{R}$ for which the modular $\varrho_{p(\cdot)}(u) = \int_\Omega |u(x)|^{p(x)} dx$ is finite. The Luxemburg norm on this space is defined by

$$\|u\|_{L^{p(\cdot)}(\Omega)} = \|u\|_{p(\cdot)} = \inf\{\lambda > 0 : \varrho_{p(\cdot)}(u/\lambda) \leq 1\}.$$

This norm makes $L^{p(\cdot)}(\Omega)$ a Banach space.

There holds the following relation between $\varrho_{p(\cdot)}(u)$ and $\|u\|_{L^{p(\cdot)}(\Omega)}$:

$$\begin{aligned} \min \left\{ \left(\int_\Omega |u|^{p(x)} dx \right)^{1/p_{\min}}, \left(\int_\Omega |u|^{p(x)} dx \right)^{1/p_{\max}} \right\} &\leq \|u\|_{L^{p(\cdot)}(\Omega)} \\ &\leq \max \left\{ \left(\int_\Omega |u|^{p(x)} dx \right)^{1/p_{\min}}, \left(\int_\Omega |u|^{p(x)} dx \right)^{1/p_{\max}} \right\}. \end{aligned}$$

Moreover, the dual of $L^{p(\cdot)}(\Omega)$ is $L^{p'(\cdot)}(\Omega)$ with $\frac{1}{p(x)} + \frac{1}{p'(x)} = 1$.

$W^{1,p(\cdot)}(\Omega)$ denotes the space of measurable functions u such that u and the distributional derivative ∇u are in $L^{p(\cdot)}(\Omega)$. The norm

$$\|u\|_{1,p(\cdot)} := \|u\|_{p(\cdot)} + \|\nabla u\|_{p(\cdot)}$$

makes $W^{1,p(\cdot)}(\Omega)$ a Banach space.

The space $W_0^{1,p(\cdot)}(\Omega)$ is defined as the closure of the $C_0^\infty(\Omega)$ in $W^{1,p(\cdot)}(\Omega)$.

For further details on these spaces, see [DHHR], [KR], [RR] and their references.

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