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MOLECULAR PREDISSOCIATION RESONANCES NEAR AN ENERGY-LEVEL CROSSING II: VECTOR FIELD INTERACTION

S. FUJIIÉ¹, A. MARTINEZ² AND T. WATANABE³

ABSTRACT. We study the resonances of a two-by-two semiclassical system of one dimensional Schrödinger operators, near an energy where the two potentials intersect transversally, one of them being bonding, and the other one anti-bonding. Assuming that the interaction is a vector-field, we obtain optimal estimates on the location and on the widths of these resonances.

Keywords: Resonances; Born-Oppenheimer approximation; eigenvalue crossing.

Subject classifications: 35P15; 35C20; 35S99; 47A75.

1. INTRODUCTION

This paper is devoted to the study of diatomic molecular predissociation resonances in the Born-Oppenheimer approximation, at energies close to that of the crossing of the electronic levels. It is a continuation of [FMW] where a method was introduced in order to overcome the difficulty of working with a 2×2 system of semiclassical operators.

In all of the work, the parameter h stands for the square-root of the inverse of the (mean-) mass of the nuclei.

In [FMW], we obtained optimal estimates both on the real parts and on the imaginary parts (widths) of the resonances (that respectively correspond to the radiation frequency and to the inverse of the life-time of the molecule), under the condition that the interaction is of the form $h(r_0(x) + hr_1(x)D_x)$, with $r_0 \neq 0$ at the point where the two electronic levels cross. However, when performing a Fechbach reduction in the Born-Oppenheimer approximation

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(see, e.g., [KMSW, MaMe, MaSo]), it appears that the interaction that comes out is a vector-field of the form $ih^2r_1(x)D_x$ (plus smaller terms), with r_1 real on the real. In that case, the result of [FMW] just says that the widths are $\mathcal{O}(h^2)$, and does not provide any lower bound on them.

Here we plan to apply the techniques introduced in [FMW] in order to obtain the asymptotic behaviour of the widths of the resonances, in the physical case of a vector-field interaction.

As in [FMW], we consider a 2×2 matrix system, the diagonal part of which consists of one-dimensional semiclassical Schrödinger operators, and we assume that the two potentials cross transversally at the origin, with value 0, and that, at this energy level, one of the two potentials admits a well, while the other one is non-trapping (see figure 1).

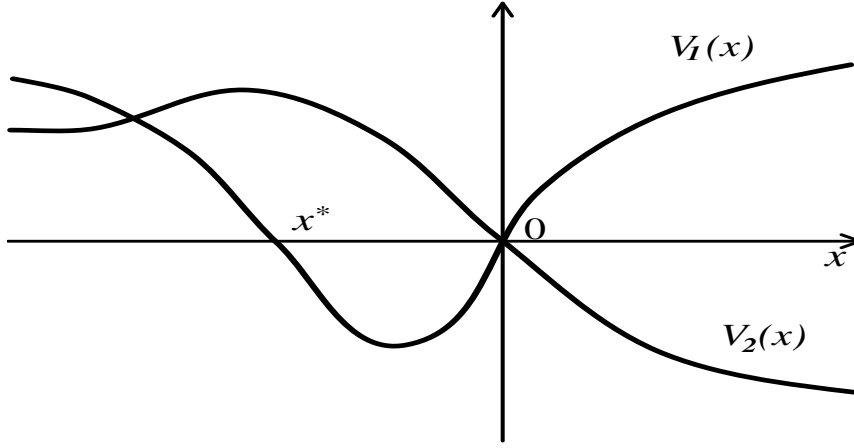


FIGURE 1. The two potentials

For such a model, we study the resonances $E = E(h)$ that have a real part $\mathcal{O}(h^{2/3})$ and an imaginary part $\mathcal{O}(h)$.

2. ASSUMPTIONS AND RESULTS

We consider a Schrödinger operator with 2×2 matrix-valued potential,

$$(2.1) \quad Pu = Eu, \quad P = \begin{pmatrix} P_1 & hW \\ hW^* & P_2 \end{pmatrix},$$

where $P_j = h^2 D_x^2 + V_j(x)$ ($j = 1, 2$) with $D_x = -i \frac{d}{dx}$, $W = W(x, hD_x)$ is a first order semiclassical differential operator, and W^* is the formal adjoint of W .

As in [FMW], we suppose the following conditions on the potentials $V_1(x), V_2(x)$ and on the interaction $W(x, hD_x)$:

(A1) $V_1(x), V_2(x)$ are real-valued analytic functions on \mathbb{R} , and extend to holomorphic functions in the complex domain,

$$\Gamma = \{x \in \mathbb{C}; |\operatorname{Im} x| < \delta_0 \langle \operatorname{Re} x \rangle\}$$

where $\delta_0 > 0$ is a constant, and $\langle t \rangle := (1 + |t|^2)^{1/2}$.

(A2) For $j = 1, 2$, V_j admit limits as $\operatorname{Re} x \rightarrow \pm\infty$ in Γ , and they satisfy,

$$\begin{aligned} \lim_{\substack{\operatorname{Re} x \rightarrow -\infty \\ x \in \Gamma}} V_1(x) &> 0; & \lim_{\substack{\operatorname{Re} x \rightarrow -\infty \\ x \in \Gamma}} V_2(x) &> 0; \\ \lim_{\substack{\operatorname{Re} x \rightarrow +\infty \\ x \in \Gamma}} V_1(x) &> 0; & \lim_{\substack{\operatorname{Re} x \rightarrow +\infty \\ x \in \Gamma}} V_2(x) &< 0. \end{aligned}$$

(A3) There exists a negative number $x^* < 0$ such that,

- $V_1 > 0$ and $V_2 > 0$ on $(-\infty, x^*)$;
- $V_1 < 0 < V_2$ on $(x^*, 0)$;
- $V_2 < 0 < V_1$ on $(0, +\infty)$,

and one has,

$$V_1'(x^*) < 0, \quad V_1'(0) =: \tau_1 > 0, \quad V_2'(0) =: -\tau_2 < 0.$$

(A4) The interaction $W(x, hD_x)$ is a differential operator of the form,

$$W(x, hD_x) = r_0(x) + ir_1(x)hD_x,$$

where $r_0(x)$ and $r_1(x)$ are bounded analytic functions on Γ , and $r_0(x)$ is real-valued on \mathbb{R} .

Notice that, in a neighborhood of $E = 0$, the scalar operator P_1 has eigenvalues, while P_2 has only essential spectrum. Hence, if the interaction W is absent, the matrix-valued operator P has embedded eigenvalues in the essential spectrum. But if W is present, it is expected that there exist, instead of embedded eigenvalues, resonances close to them in the lower half complex plane of the energy.

The resonances of P are defined, e.g., as the values $E \in \mathbb{C}$ such that the equation $Pu = Eu$ has a non identically vanishing solution such that, for some $\theta > 0$ sufficiently small, the function $x \mapsto u(xe^{i\theta})$ is in $L^2(\mathbb{R}) \oplus L^2(\mathbb{R})$ (see, e.g., [AgCo, ReSi]). We will give an equivalent definition of resonances adapted to our setting in the next section (see also [HeMa]). We denote by $\operatorname{Res}(P)$ the set of these resonances.

For $E \in \mathbb{C}$ small enough, we define the action,

$$(2.2) \quad \mathcal{A}(E) := \int_{x_1^*(E)}^{x_1(E)} \sqrt{E - V_1(t)} dt,$$

where $x_1^*(E)$ (respectively $x_1(E)$) is the unique solution of $V_1(x) = E$ close to x^* (respectively close to 0).

We fix $C_0 > 0$ arbitrarily large, and we study the resonances of P lying in the set $\mathcal{D}_h(C_0)$ given by,

$$(2.3) \quad \mathcal{D}_h(C_0) := [-C_0 h^{2/3}, C_0 h^{2/3}] - i[0, C_0 h].$$

For $h > 0$ and $k \in \mathbb{Z}$, we set,

$$(2.4) \quad \lambda_k(h) := \frac{-\mathcal{A}(0) + (k + \frac{1}{2})\pi h}{\mathcal{A}'(0)h^{2/3}}.$$

The Bohr-Sommerfeld quantization condition of eigenvalues for the scalar operator P_1 reads

$$\mathcal{A}(E) = (k + \frac{1}{2})\pi h + \mathcal{O}(h^2).$$

Then the $\lambda_k(h)h^{2/3}$'s are approximate eigenvalues of P_1 near 0. We will find resonances close to these real values.

First we recall the result obtained in [FMW]. Notice that a multiplicative factor $(\tau_1\tau_2)^{\frac{1}{3}}$ in the asymptotic formula of the imaginary part of resonances was missing in that paper. We would like to correct it on this occasion.

Theorem 2.1 ([FMW]). *Assume (A1)-(A4). For $h > 0$ small enough, one has,*

$$\text{Res}(P) \cap \mathcal{D}_h(C_0) = \{E_k(h); k \in \mathbb{Z}\} \cap \mathcal{D}_h(C_0),$$

where the $E_k(h)$'s are complex numbers that satisfy,

$$(2.5) \quad \text{Re } E_k(h) = \lambda_k(h)h^{\frac{2}{3}} - \frac{\mathcal{A}''(0)}{2\mathcal{A}'(0)}\lambda_k(h)^2h^{\frac{4}{3}} + \mathcal{O}(h^{\frac{5}{3}}),$$

$$(2.6) \quad \text{Im } E_k(h) = -\frac{2\pi^2 r_0(0)^2}{\mathcal{A}'(0)}(\tau_1\tau_2)^{\frac{1}{3}}\left(\mu_1(\lambda_k(h))^2 + \mu_2(\lambda_k(h))^2\right)h^{\frac{5}{3}} + \mathcal{O}(h^2),$$

uniformly as $h \rightarrow 0_+$. Here, the functions μ_1 and μ_2 are defined by

$$\begin{aligned} \mu_1(t) &= \int_0^\infty \text{Ai}(\tau_1^{-\frac{2}{3}}(\tau_1 y - t))\text{Ai}(-\tau_2^{-\frac{2}{3}}(\tau_2 y + t)) dy \\ \mu_2(t) &= \int_0^\infty \text{Ai}(\tau_2^{-\frac{2}{3}}(\tau_2 y - t))\text{Ai}(-\tau_1^{-\frac{2}{3}}(\tau_1 y + t)) dy, \end{aligned}$$

where Ai stands for the Airy function $\text{Ai}(x) = \frac{1}{2\pi} \int_{-\infty}^\infty e^{i(x\xi + \xi^3/3)} d\xi$.

Remark that, if $r_0(0) = 0$, this theorem gives only an estimate $\mathcal{O}(h^2)$ for the imaginary part of resonances. The following main result provides a precise asymptotic formula up to $\mathcal{O}(h^{\frac{8}{3}})$ in the case where $r_0(x)$ vanishes identically.

Theorem 2.2. *Assume moreover that $r_0(x) = 0$ and $r_1(x)$ is real on \mathbb{R} . Then the $E_k(h)$'s in the previous theorem satisfy, with $\tau_3^{-1} = \tau_1^{-1} + \tau_2^{-1}$*

$$(2.7) \quad \text{Re } E_k(h) = \lambda_k(h)h^{\frac{2}{3}} - \frac{\mathcal{A}''(0)}{2\mathcal{A}'(0)}\lambda_k(h)^2h^{\frac{4}{3}} - \frac{\mathcal{A}^{(3)}(0)}{6\mathcal{A}'(0)}\lambda_k(h)^3h^{\frac{6}{3}} + \mathcal{O}(h^{\frac{7}{3}}),$$

$$(2.8) \quad \text{Im } E_k(h) = -\frac{\pi^2 r_1(0)^2}{\mathcal{A}'(0)} \frac{\tau_3^{\frac{1}{3}}}{\tau_1 + \tau_2} \left(\text{Ai}'(-\tau_3^{-\frac{2}{3}}\lambda_k(h)) \right)^2 h^{\frac{7}{3}} + \mathcal{O}(h^{\frac{8}{3}}),$$

uniformly as $h \rightarrow 0_+$.

3. PROOF OF THEOREM 2.2

We prove Theorem 2.2. For a sufficiently small $\theta > 0$, let $I_L := (-\infty, 0]$, $I_R^\theta := F_\theta([0, +\infty))$ with $F_\theta(x) := x + i\theta f(x)$ where $f \in C^\infty(\mathbb{R}_+, \mathbb{R}_+)$, $f(x) = x$ for x large enough, $f(x) = 0$ for $x \in [0, x_\infty]$ for some $x_\infty > 0$, and f is

chosen in such a way that, for any $x \geq x_\infty$, and with some positive constant C , one has (see [FMW], Formula (3.1)),

$$(3.1) \quad \operatorname{Im} \int_{x_\infty}^{F_\theta(x)} \sqrt{E - V_2(t)} dt \geq -Ch.$$

The linear space V of solutions to the system (2.1) is of dimension four. The solutions in $L^2(I_R^\theta) \oplus L^2(I_R^\theta)$ form a two dimensional subspace $V_R = V \cap (L^2(I_R^\theta) \oplus L^2(I_R^\theta))$, and the solutions in $L^2(I_L) \oplus L^2(I_L)$ form a two dimensional subspace $V_L = V \cap (L^2(I_L) \oplus L^2(I_L))$.

Then E is a resonance if and only if the intersection $V_R \cap V_L$ is at least 1 dimensional. In other words, the quantization condition of resonances can be written in the form

$$(3.2) \quad \mathcal{W}_0(E) := \mathcal{W}(w_{1,L}, w_{2,L}, w_{1,R}, w_{2,R}) = 0,$$

where the couple $(w_{1,L}, w_{2,L})$ (resp. $(w_{1,R}, w_{2,R})$) is a basis of V_L (resp. V_R) and $\mathcal{W}(w_{1,L}, w_{2,L}, w_{1,R}, w_{2,R})$ is the Wronskian, i.e. the determinant of the 4×4 matrix

$$\begin{pmatrix} w_{1,L} & w_{2,L} & w_{1,R} & w_{2,R} \\ \partial_x w_{1,L} & \partial_x w_{2,L} & \partial_x w_{1,R} & \partial_x w_{2,R} \end{pmatrix}.$$

Such solutions $w_{1,L}, w_{2,L}, w_{1,R}, w_{2,R}$ are constructed as in [FMW] using fundamental solutions to the scalar equations $(P_j - E)u = 0$.

On I_L , and for $E \in \mathcal{D}_h(C_0)$ and $j = 1, 2$, let $u_{j,L}^\pm$ be the solutions to $(P_j - E)u = 0$ constructed in [FMW] (in particular, $u_{j,L}^-$ decays exponentially at $-\infty$, while $u_{j,L}^+$ grows exponentially, and their Wronskian $\mathcal{W}[u_{j,L}^+, u_{j,L}^-]$ is of size $h^{-\frac{2}{3}}$). We construct fundamental solutions $K_{j,L}$, $j = 1, 2$ on I_L :

$$(3.3) \quad \begin{aligned} K_{j,L}[v](x) &:= \frac{u_{j,L}^+(x)}{h^2 \mathcal{W}[u_{j,L}^+, u_{j,L}^-]} \int_{-\infty}^x u_{j,L}^-(t) v(t) dt \\ &\quad + \frac{u_{j,L}^-(x)}{h^2 \mathcal{W}[u_{j,L}^+, u_{j,L}^-]} \int_x^0 u_{j,L}^+(t) v(t) dt. \end{aligned}$$

In the same way, we construct $K_{j,R}$, $j = 1, 2$ on I_R^θ using the solutions $u_{j,R}^\pm$ ($u_{j,R}^+$ grows and $u_{j,R}^-$ decays exponentially, at ∞ along I_R^θ):

$$(3.4) \quad \begin{aligned} K_{j,R}[v](x) &:= \frac{u_{j,R}^-(x)}{h^2 \mathcal{W}[u_{j,R}^-, u_{j,R}^+]} \int_0^x u_{j,R}^+(t) v(t) dt \\ &\quad + \frac{u_{j,R}^+(x)}{h^2 \mathcal{W}[u_{j,R}^-, u_{j,R}^+]} \int_x^{+\infty} u_{j,R}^-(t) v(t) dt. \end{aligned}$$

Let $C_b^0(I_L)$ and $C_b^0(I_R^\theta)$ be the space of bounded continuous functions on I_L and I_R^θ respectively. The above operators act on these function spaces, and satisfy $(P_j - E)K_{j,L} = \operatorname{Id}$ and $(P_j - E)K_{j,R} = \operatorname{Id}$ respectively.

We have the following estimates, which are better than the elliptic case (see Proposition 3.1 and 3.2 in [FMW]).

Proposition 3.1. *As h goes to 0_+ , one has uniformly,*

$$(3.5) \quad \|hK_{2,L}W^*\|_{\mathcal{L}(C_b^0(I_L))} = \mathcal{O}(h^{\frac{2}{3}}),$$

$$(3.6) \quad \|h^2K_{1,L}WK_{2,L}W^*\|_{\mathcal{L}(C_b^0(I_L))} = \mathcal{O}(h),$$

$$(3.7) \quad \|hK_{1,R}W\|_{\mathcal{L}(C_b^0(I_R^0))} = \mathcal{O}(h^{\frac{2}{3}}),$$

$$(3.8) \quad \|h^2K_{2,R}W^*K_{1,R}W\|_{\mathcal{L}(C_b^0(I_R^0))} = \mathcal{O}(h).$$

Proof: We first prove the estimates for $K_{j,L}$. For $j = 1, 2$, we set,

$$(3.9) \quad \begin{aligned} U_j(x, t) &:= |u_{j,L}^+(x)u_{j,L}^-(t)|\mathbf{1}_{\{t < x\}} + |u_{j,L}^-(x)u_{j,L}^+(t)|\mathbf{1}_{\{t > x\}} = U_j(t, x); \\ U_j'(x, t) &:= |u_{j,L}^+(x)\partial_t u_{j,L}^-(t)|\mathbf{1}_{\{t < x\}} + |u_{j,L}^-(x)\partial_t u_{j,L}^+(t)|\mathbf{1}_{\{t > x\}}; \\ \tilde{U}_j(x, t) &:= U_j(x, t) + U_j'(x, t). \end{aligned}$$

Thanks to our choice of $K_{j,L}$, and doing an integration by parts, we see that,

$$(3.10) \quad \begin{aligned} |hK_{1,L}Wv(x)| &= \mathcal{O}(h^{\frac{2}{3}}) \left(\int_{-\infty}^0 \tilde{U}_1(x, t)|v(t)|dt + U_1(x, 0)|v(0)| \right); \\ |hK_{2,L}W^*v(x)| &= \mathcal{O}(h^{\frac{2}{3}}) \left(\int_{-\infty}^0 \tilde{U}_2(x, t)|v(t)|dt + U_2(x, 0)|v(0)| \right), \end{aligned}$$

and therefore,

$$(3.11) \quad \|hK_{2,L}W^*\| = \mathcal{O}(h^{\frac{2}{3}}) \sup_{x \in I_L} \int_{-\infty}^0 \tilde{U}_2(x, t)dt + \mathcal{O}(h^{\frac{2}{3}}) \sup_{x \in I_L} U_2(x, 0).$$

Moreover, using the asymptotics of $u_{2,L}^\pm$ and $\partial_x u_{2,L}^\pm$ on I_L , we see that $U_2(x, t) = \mathcal{O}(1)$ uniformly, and fixing some constant $C_1 > 0$ sufficiently large, we also have,

$$\tilde{U}_2(x, t) = \begin{cases} \mathcal{O}(h^{-\frac{2}{3}}) \frac{|V_2(t) - E|^{\frac{1}{4}} e^{-|\operatorname{Re} \int_t^x (V_2 - E)^{1/2}|/h}}{|V_2(x) - E|^{\frac{1}{4}}} & (x, t \leq -C_1 h^{\frac{2}{3}}), \\ \mathcal{O}(h^{-\frac{5}{6}}) |V_2(t) - E|^{\frac{1}{4}} e^{-|\operatorname{Re} \int_t^x (V_2 - E)^{1/2}|/h} & (t \leq -C_1 h^{\frac{2}{3}} \leq x \leq 0), \\ \mathcal{O}(h^{-\frac{1}{2}}) |V_2(x) - E|^{-\frac{1}{4}} e^{-|\operatorname{Re} \int_0^x (V_2 - E)^{1/2}|/h} & (x \leq -C_1 h^{\frac{2}{3}} \leq t \leq 0), \\ \mathcal{O}(h^{-\frac{2}{3}}) & (x, t \in [-C_1 h^{\frac{2}{3}}, 0]). \end{cases}$$

In particular $\tilde{U}_2(x, t) = \mathcal{O}(h^{-2/3})$ uniformly, and when $x \leq -\delta$ with $\delta > 0$ constant, there exists a positive constant α such that,

$$\int_{-\infty}^0 \tilde{U}_2(x, t)dt = \mathcal{O}(h^{-\frac{2}{3}}) \int_{-\infty}^{-\delta/2} e^{-\alpha|x-t|/h}dt + \mathcal{O}(e^{-\alpha/h}) = \mathcal{O}(h^{\frac{1}{3}}).$$

On the other hand, if δ is chosen sufficiently small and $x \in [-\delta, -C_1 h^{2/3}]$, then, there exists a (different) positive constant α such that,

$$\begin{aligned} \int_{-\infty}^0 \tilde{U}_2(x, t)dt &= \int_{-2\delta}^{-C_1 h^{2/3}} \tilde{U}_2(x, t)dt + \mathcal{O}(1) \\ &= \mathcal{O}(h^{-\frac{2}{3}}|x|^{-\frac{1}{4}}) \int_{C_1 h^{2/3}}^{2\delta} t^{\frac{1}{4}} e^{-\alpha|t^{\frac{3}{2}} - |x|^{\frac{3}{2}}|/h} dt + \mathcal{O}(1). \end{aligned}$$

Setting $t = (hs)^{2/3}$ in the integral, we obtain,

$$\int_{-\infty}^0 \tilde{U}_2(x, t) dt = \mathcal{O}(h^{-\frac{2}{3}} |x|^{-\frac{1}{4}} h^{\frac{2}{3} + \frac{1}{6}}) \int_1^\infty \frac{e^{-\alpha |s - |x|^{\frac{3}{2}}/h|}}{\sqrt{s}} ds + \mathcal{O}(1) = \mathcal{O}(1).$$

Finally, when $x \in [-C_1 h^{2/3}, 0]$, we have,

$$\begin{aligned} \int_{-\infty}^0 \tilde{U}_2(x, t) dt &= \int_{-\delta}^{-C_1 h^{2/3}} \tilde{U}_2(x, t) dt + \mathcal{O}(1) \\ &= \mathcal{O}(h^{-\frac{5}{6}}) \int_{C_1 h^{2/3}}^\delta t^{\frac{1}{4}} e^{-\alpha t^{\frac{3}{2}}/h} dt + \mathcal{O}(1) = \mathcal{O}(1). \end{aligned}$$

Thus, we have proven,

$$(3.12) \quad \sup_{x \leq 0} \int_{-\infty}^0 \tilde{U}_2(x, t) dt = \mathcal{O}(1),$$

and, by (3.11) (and the fact that $U_2 = \mathcal{O}(1)$), (3.5) follows.

Now, let us prove the estimate on $M_L := h^2 K_{1,L} W K_{2,L} W^*$. We see on the definition of $K_{1,L}$ and on (3.10) that we have,

$$\begin{aligned} |M_L v(x)| &= \mathcal{O}(h^{\frac{4}{3}}) \int_{-\infty}^0 \int_{-\infty}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) |v(s)| ds dt \\ &\quad + \mathcal{O}(h^{\frac{4}{3}}) \int_{-\infty}^0 \tilde{U}_1(x, t) U_2(t, 0) |v(0)| dt \\ &\quad + \mathcal{O}(h^{\frac{4}{3}}) U_1(x, 0) \int_{-\infty}^0 \tilde{U}_2(0, t) |v(t)| dt \\ &\quad + \mathcal{O}(h^{\frac{4}{3}}) U_1(x, 0) U_2(0, 0) |v(0)|. \end{aligned} \quad (3.13)$$

Using (3.12) and the fact that $U_j = \mathcal{O}(1)$ uniformly ($j = 1, 2$), we see that the last three terms are $\mathcal{O}(h^{4/3}) \sup_{I_L} |v|$.

In order to estimate the first term, we use the following properties of \tilde{U}_1 : For any $\delta > 0$ small enough, there exists $\alpha > 0$ constant, such that,

$$\tilde{U}_1(x, t) = \begin{cases} \mathcal{O}(h^{-\frac{2}{3}}) e^{-\alpha |t-x|/h} & (x, t \leq x^* - \delta) \\ \mathcal{O}(e^{-\alpha/h}) & (t \leq x^* - 2\delta, x \in [x^* - \delta, 0]), \\ \mathcal{O}(e^{-\alpha/h}) & (x \leq x^* - 2\delta, t \in [x^* - \delta, 0]), \\ \mathcal{O}(h^{-\frac{5}{6}} |t|^{\frac{1}{4}}) & (x \in [x^* - 4\delta, 0], t \in [-\delta, -C_1 h^{\frac{2}{3}}]), \\ \mathcal{O}(h^{-\frac{2}{3}}) & (x \in [x^* - 4\delta, 0], t \in [-C_1 h^{\frac{2}{3}}, 0] \cup [x^* - 4\delta, -\delta]), \\ \mathcal{O}(h^{-\frac{1}{2}} |x|^{-\frac{1}{4}}) & (t \in [x^* - 4\delta, 0], x \in [-\delta, -C_1 h^{\frac{2}{3}}]), \\ \mathcal{O}(h^{-\frac{2}{3}}) & (t \in [x^* - 4\delta, 0], x \in [-C_1 h^{\frac{2}{3}}, 0] \cup [x^* - 4\delta, -\delta]). \end{cases}$$

In particular $\tilde{U}_1(x, t) = \mathcal{O}(h^{-2/3})$ uniformly. Moreover, by the properties of \tilde{U}_2 , we also know that any part of the integral corresponding to $|t - s| \geq \delta$ with $\delta > 0$ constant is exponentially small.

We first consider the case $x \in (-\infty, x^* - 2\delta]$ for a small positive constant δ .

Then, we see that there exists a constant $\alpha > 0$ such that,

$$\begin{aligned} \iint_{-\infty}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) dt ds &= \mathcal{O}(h^{-\frac{4}{3}}) \int_{-\infty}^{x^* - \delta} dt \int_{-\infty}^{x^* - \delta/2} e^{-\alpha(|t-x|+|s-t|)/h} ds \\ &\quad + \mathcal{O}(e^{-\alpha/h}) \\ &= \mathcal{O}(h^{2-\frac{4}{3}}) = \mathcal{O}(h^{\frac{2}{3}}). \end{aligned}$$

Now, when $x \in [x^* - 2\delta, 0]$, and still denoting by α every new positive constant that may appear, we have,

$$\iint_{-\infty}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) dt ds = \int_{x^* - 3\delta}^0 dt \int_{x^* - 4\delta}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) ds + \mathcal{O}(e^{-\alpha/h}),$$

and,

$$\begin{aligned} \int_{x^* - 3\delta}^0 dt \int_{x^* - 4\delta}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) ds &= \mathcal{O}(h^{-\frac{2}{3}}) \int_{x^* - 3\delta}^0 dt \int_{x^* - 4\delta}^0 \tilde{U}_2(t, s) ds \\ &= \mathcal{O}(h^{-\frac{4}{3}}) \int_{x^* - 3\delta}^{-\delta} dt \int_{x^* - 4\delta}^{-\delta/2} e^{-\alpha|t-s|/h} ds + \mathcal{O}(e^{-\alpha/h}) \\ &\quad + \mathcal{O}(h^{-\frac{4}{3}}) \int_{-\delta}^{-C_1 h^{2/3}} dt \int_{-2\delta}^{-C_1 h^{2/3}} \frac{|s|^{\frac{1}{4}} e^{-\alpha|t|^{\frac{3}{2}} - |s|^{\frac{3}{2}}/h}}{|t|^{\frac{1}{2}}} ds \\ &\quad + \mathcal{O}(h^{-\frac{7}{6}}) \int_{-\delta}^{-C_1 h^{2/3}} dt \int_{-C_1 h^{2/3}}^0 \frac{e^{-\alpha|t|^{\frac{3}{2}}/h}}{|t|^{\frac{1}{4}}} ds \\ &\quad + \mathcal{O}(h^{-\frac{3}{2}}) \int_{-C_1 h^{2/3}}^0 dt \int_{-\delta}^{-C_1 h^{2/3}} |s|^{\frac{1}{4}} e^{-\alpha|s|^{\frac{3}{2}}/h} ds + \mathcal{O}(1). \end{aligned}$$

Hence,

$$\begin{aligned} (3.14) \quad \int_{x^* - 3\delta}^0 dt \int_{x^* - 4\delta}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) ds \\ = \mathcal{O}(h^{-\frac{4}{3}}) \int_{C_1 h^{2/3}}^{\delta} dt \int_{C_1 h^{2/3}}^{2\delta} \frac{s^{\frac{1}{4}} e^{-\alpha|t|^{\frac{3}{2}} - s^{\frac{3}{2}}/h}}{t^{\frac{1}{2}}} ds + \mathcal{O}(h^{-\frac{1}{2}}) \int_{C_1 h^{2/3}}^{\delta} \frac{e^{-\alpha t^{\frac{3}{2}}/h}}{t^{\frac{1}{4}}} dt \\ + \mathcal{O}(h^{-\frac{5}{6}}) \int_{C_1 h^{2/3}}^{\delta} s^{\frac{1}{4}} e^{-\alpha s^{\frac{3}{2}}/h} ds + \mathcal{O}(h^{-\frac{1}{3}}). \end{aligned}$$

For the first term, the change of variables $(t, s) \mapsto (t^{2/3}, s^{2/3})$ gives,

$$\int_{C_1 h^{2/3}}^{\delta} dt \int_{C_1 h^{2/3}}^{2\delta} \frac{s^{\frac{1}{4}} e^{-\alpha|t|^{\frac{3}{2}} - s^{\frac{3}{2}}/h}}{t^{\frac{1}{2}}} ds = \mathcal{O}(1) \iint_{C_2 h}^{\delta'} \frac{e^{-\alpha|t-s|/h}}{t^{\frac{2}{3}} s^{\frac{1}{6}}} dt ds,$$

with $C_2 := C_1^{2/3}$ and $\delta' := (2\delta)^{2/3}$. Dividing this integral in two parts, depending whether $t \leq s$ or $s \leq t$, and first integrating with respect to the

larger of the two variables, we obtain,

$$(3.15) \quad \int_{C_1 h^{2/3}}^{\delta} dt \int_{C_1 h^{2/3}}^{2\delta} \frac{s^{\frac{1}{4}} e^{-\alpha|t^{\frac{3}{2}} - s^{\frac{3}{2}}|/h}}{t^{\frac{1}{2}}} ds = \mathcal{O}(1) \int_{C_2 h}^{\delta'} dt \frac{e^{\alpha t/h}}{t^{\frac{5}{6}}} \int_t^{\delta'} e^{-\alpha s/h} ds = \mathcal{O}(h).$$

Moreover, a simple change of variable gives,

$$\int_{C_1 h^{2/3}}^{\delta} \frac{e^{-\alpha t^{\frac{3}{2}}/h}}{t^{\frac{1}{4}}} dt = \mathcal{O}(h^{\frac{1}{2}}), \quad \int_{C_1 h^{2/3}}^{\delta} s^{\frac{1}{4}} e^{-\alpha s^{\frac{3}{2}}/h} ds = \mathcal{O}(h^{\frac{5}{6}}),$$

Inserting into (3.14), we deduce that, for $x \in [x^* - 2\delta, 0]$, we have,

$$(3.16) \quad \iint_{-\infty}^0 \tilde{U}_1(x, t) \tilde{U}_2(t, s) dt ds = \mathcal{O}(h^{-\frac{1}{3}}).$$

Finally, going back to (3.13), we conclude (3.6).

The estimates (3.7), (3.8) for $K_{j,R}$ are proved similarly (x_{∞} playing the role of x^*). \square

Set $M_L := h^2 K_{1,L} W K_{2,L} W^*$ and $M_R := h^2 K_{2,R} W^* K_{1,R} W$. Thanks to Proposition 3.1, we can define the following four vector-valued functions as Neumann series for small enough h ;

$$(3.17) \quad w_{1,L} := \begin{pmatrix} \sum_{j \geq 0} M_L^j u_{1,L}^- \\ -h K_{2,L} W^* \sum_{j \geq 0} M_L^j u_{1,L}^- \end{pmatrix},$$

$$(3.18) \quad w_{2,L} := \begin{pmatrix} -\sum_{j \geq 0} M_L^j (h K_{1,L} W u_{2,L}^-) \\ u_{2,L}^- + h K_{2,L} W^* \sum_{j \geq 0} M_L^j (h K_{1,L} W u_{2,L}^-) \end{pmatrix},$$

$$(3.19) \quad w_{1,R} := \begin{pmatrix} u_{1,R}^- + h K_{1,R} W \sum_{j \geq 0} M_R^j (h K_{2,R} W^* u_{1,R}^-) \\ -\sum_{j \geq 0} M_R^j (h K_{2,R} W^* u_{1,R}^-) \end{pmatrix},$$

$$(3.20) \quad w_{2,R} := \begin{pmatrix} -h K_{1,R} W \sum_{j \geq 0} M_R^j u_{2,R}^- \\ \sum_{j \geq 0} M_R^j u_{2,R}^- \end{pmatrix}.$$

It is not difficult to see that they are solutions to the system (2.1) and that (see [FMW], Proposition 4.1),

$$w_{j,L} \in L^2(I_L) \oplus L^2(I_L) \quad ; \quad w_{j,R} \in L^2(I_R^{\theta}) \oplus L^2(I_R^{\theta}).$$

In order to get the leading term of the imaginary part of resonances, it will be necessary to compute the asymptotics of these solutions up to errors of $\mathcal{O}(h^{5/3})$. This means to compute, for example for $w_{1,L}$, two terms $u_{1,L}^- + M_L u_{1,L}^-$ for the first element, and one term $-h K_{2,L} W^* u_{1,L}^-$ for the second element just as in the elliptic interaction case.

We will compute the Wronskian $\mathcal{W}_0(E, h)$, which is independent of x , at the origin. Substituting $x = 0$ to these solutions or their derivatives, we obtain

the following asymptotic formulae just as in [FMW] (only the remainder estimates are different). For $S = L, R$, we have, uniformly as $h \rightarrow 0_+$,

$$(3.21) \quad \begin{aligned} w_{1,S}(0) &= \begin{bmatrix} u_{1,S}^-(0) + \beta_{1,S} u_{1,S}^+(0) \\ \alpha_{1,S} u_{2,S}^+(0) \end{bmatrix} + \mathcal{O}(h^{\frac{5}{3}}); \\ \tilde{\partial} w_{1,S}(0) &= \begin{bmatrix} \tilde{\partial} u_{1,S}^-(0) + \beta_{1,S} \tilde{\partial} u_{1,S}^+(0) \\ \alpha_{1,S} \tilde{\partial} u_{2,S}^+(0) \end{bmatrix} + \mathcal{O}(h^{\frac{5}{3}}), \end{aligned}$$

$$(3.22) \quad \begin{aligned} w_{2,S}(0) &= \begin{bmatrix} \alpha_{2,S} u_{1,S}^+(0) \\ u_{2,S}^-(0) + \beta_{2,S} u_{2,S}^+(0) \end{bmatrix} + \mathcal{O}(h^{\frac{5}{3}}); \\ \tilde{\partial} w_{2,S}(0) &= \begin{bmatrix} \alpha_{2,S} \tilde{\partial} u_{1,S}^+(0) \\ \tilde{\partial} u_{2,S}^-(0) + \beta_{2,S} \tilde{\partial} u_{2,S}^+(0) \end{bmatrix} + \mathcal{O}(h^{\frac{5}{3}}). \end{aligned}$$

Here, $\tilde{\partial}$ stands for $h^{2/3}\partial$ and $\alpha_{j,S}$ and $\beta_{j,S}$ are complex numbers defined by

$$(3.23) \quad \begin{aligned} \alpha_{1,L} &= \frac{-\int_{-\infty}^0 u_{2,L}^-(t)(W^* u_{1,L}^-(t))dt}{h\mathcal{W}(u_{2,L}^+, u_{2,L}^-)}; & \beta_{1,L} &= \frac{\int_{-\infty}^0 u_{1,L}^-(t)(W r_{1,L})(t)dt}{h\mathcal{W}(u_{1,L}^+, u_{1,L}^-)}; \\ \alpha_{2,L} &= \frac{-\int_{-\infty}^0 u_{1,L}^-(t)(W u_{2,L}^-(t))dt}{h\mathcal{W}(u_{1,L}^+, u_{1,L}^-)}; & \beta_{2,L} &= \frac{\int_{-\infty}^0 u_{2,L}^-(t)(W^* r_{2,L})(t)dt}{h\mathcal{W}(u_{2,L}^+, u_{2,L}^-)}; \\ \alpha_{1,R} &= \frac{-\int_0^{+\infty} u_{2,R}^-(t)(W^* u_{1,R}^-(t))dt}{h\mathcal{W}(u_{2,R}^-, u_{2,R}^+)}; & \beta_{1,R} &= \frac{\int_0^{+\infty} u_{1,R}^-(t)(W r_{1,R})(t)dt}{h\mathcal{W}(u_{1,R}^-, u_{1,R}^+)}; \\ \alpha_{2,R} &= \frac{-\int_0^{+\infty} u_{1,R}^-(t)(W u_{2,R}^-(t))dt}{h\mathcal{W}(u_{1,R}^-, u_{1,R}^+)}; & \beta_{2,R} &= \frac{\int_0^{+\infty} u_{2,R}^-(t)(W^* r_{2,R})(t)dt}{h\mathcal{W}(u_{2,R}^-, u_{2,R}^+)}, \end{aligned}$$

where we have set, for $S = L, R$,

$$r_{1,S} := hK_{2,S}W^*u_{1,S}^-, \quad r_{2,S} := hK_{1,S}Wu_{2,S}^-,$$

and where, in the case $S = R$, the integrals run over I_R^θ .

Then we can write the Wronskian $\mathcal{W}_0(E, h)$ up to $\mathcal{O}(h^{5/3})$ in terms of $\alpha_{j,S}$ and $\beta_{j,S}$ as in Section 6 of [FMW].

$$(3.24) \quad \begin{aligned} \mathcal{W}_0(E) &= -\frac{4\sqrt{2}}{\pi^2}ie^{\frac{\pi i}{4}}\left(\cos\frac{\mathcal{A}}{h}\right)(1 + \mathcal{O}(h^{\frac{4}{3}})) - \frac{4}{\pi^2}(\alpha_{1,R}\alpha_{2,L} + \alpha_{1,L}\alpha_{2,R}) \\ &+ \frac{\sqrt{2}}{\pi^2}e^{\frac{\pi i}{4}}\left(\sin\frac{\mathcal{A}}{h}\right)(4i\alpha_{1,R}\alpha_{2,R} - \alpha_{1,L}\alpha_{2,L} + 2i(\beta_{1,R} + \beta_{1,L})) + \mathcal{O}(h^{\frac{5}{3}}), \end{aligned}$$

where \mathcal{A} is the action defined in (2.2). Notice that, in (6.2) of [FMW], we used the facts $\alpha_{1,S} = \alpha_{2,S}$ and $\beta_{1,S} = \beta_{2,S}$ at the principal level.

The constants $\alpha_{j,S}$ and $\beta_{j,S}$ have the following estimates:

Proposition 3.2. *Let $E = \rho h^{2/3} \in \mathcal{D}_h(C_0)$. Then one has, as $h \rightarrow 0$,*

$$\alpha_{j,S} = \mathcal{O}(h^{\frac{2}{3}}), \quad \beta_{j,S} = \mathcal{O}(h^{\frac{4}{3}}), \quad j = 1, 2, \quad S = L, R.$$

More precisely, one has

$$\begin{aligned}
 \alpha_{1,R} &= \frac{e^{\pi i/4}}{\sqrt{2}} \pi r_1(0) h^{\frac{2}{3}} \left(\nu_{1,R}^A(\operatorname{Re} \rho) - i \nu_{1,R}^B(\operatorname{Re} \rho) \right) + \mathcal{O}(h), \\
 \alpha_{2,R} &= \frac{e^{\pi i/4}}{\sqrt{2}} \pi r_1(0) h^{\frac{2}{3}} \left(\nu_{2,R}^A(\operatorname{Re} \rho) - i \nu_{2,R}^B(\operatorname{Re} \rho) \right) + \mathcal{O}(h), \\
 \alpha_{1,L} &= 2\pi r_1(0) h^{\frac{2}{3}} \left\{ \left(\sin \frac{\mathcal{A}}{h} \right) \nu_{1,L}^A(\operatorname{Re} \rho) + \left(\cos \frac{\mathcal{A}}{h} \right) \nu_{1,L}^B(\operatorname{Re} \rho) \right\} + \mathcal{O}(h), \\
 \alpha_{2,L} &= 2\pi r_1(0) h^{\frac{2}{3}} \left\{ \left(\sin \frac{\mathcal{A}}{h} \right) \nu_{2,L}^A(\operatorname{Re} \rho) + \left(\cos \frac{\mathcal{A}}{h} \right) \nu_{2,L}^B(\operatorname{Re} \rho) \right\} + \mathcal{O}(h), \\
 \operatorname{Im} \beta_{1,R} &= \pi^2 r_1(0)^2 h^{\frac{4}{3}} \left(\nu_{1,R}^A(\operatorname{Re} \rho) \nu_{2,R}^A(\operatorname{Re} \rho) + \nu_{1,R}^B(\operatorname{Re} \rho) \nu_{2,R}^B(\operatorname{Re} \rho) \right) + \mathcal{O}(h^{\frac{5}{3}}), \\
 \operatorname{Im} \beta_{1,L} &= \mathcal{O}(h^{\frac{5}{3}}),
 \end{aligned}$$

where

$$\begin{aligned}
 \nu_{1,R}^A(t) &= \int_0^\infty \operatorname{Ai}'(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{1,R}^B(t) &= \int_0^\infty \operatorname{Ai}'(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Bi}(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{2,R}^A(t) &= \int_0^\infty \operatorname{Ai}(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}'(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{2,R}^B(t) &= \int_0^\infty \operatorname{Ai}(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Bi}'(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{1,L}^A(t) &= \int_{-\infty}^0 \operatorname{Ai}'(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{1,L}^B(t) &= \int_{-\infty}^0 \operatorname{Bi}'(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{2,L}^A(t) &= \int_{-\infty}^0 \operatorname{Ai}(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}'(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy, \\
 \nu_{2,L}^B(t) &= \int_{-\infty}^0 \operatorname{Bi}(\tau_1^{1/3}(y - \frac{t}{\tau_1})) \operatorname{Ai}'(-\tau_2^{1/3}(y + \frac{t}{\tau_2})) dy.
 \end{aligned}$$

Proof: We only prove the formula for $\alpha_{1,L}$. Other formulas can be obtained similarly.

Thanks to the exponential decay of $u_{2,L}^-$ away from 0, we obtain,

$$\alpha_{1,L} = \frac{\pi h^{\frac{2}{3}}}{2} \int_{-\delta}^0 u_{2,L}^-(t) (r_1(t) u_{1,L}^-(t))' dt + \mathcal{O}(h),$$

where $\delta > 0$ is arbitrarily small. We divide the integral into three parts, introducing a large parameter λ satisfying $\lambda h^{2/3} \rightarrow 0$;

$$\int_{-\delta}^0 u_{2,L}^-(t) (r_1(t) u_{1,L}^-(t))' dt = I_1 + I_2 + I_3,$$

where

$$I_1 = \int_{-\delta}^{-\lambda h^{2/3}} u_{2,L}^-(t) r_1(t) (u_{1,L}^-)'(t) dt, \quad I_2 = \int_{-\lambda h^{2/3}}^0 u_{2,L}^-(t) r_1(t) (u_{1,L}^-)'(t) dt,$$

$$I_3 = \int_{-\delta}^0 u_{2,L}^-(t) r_1'(t) u_{1,L}^-(t) dt.$$

On $[-\delta, -\lambda h^{2/3}]$, both $u_{2,L}^-$ and $(u_{1,L}^-)'$ are of the WKB form, one of which is exponentially small and the other of which is oscillating; for some $c > 0$,

$$|u_{2,L}^-(t)| \leq \frac{h^{\frac{1}{6}}}{|t|^{\frac{1}{4}}} e^{-c|t|^{\frac{3}{2}}/h}, \quad |(u_{1,L}^-)'(t)| \leq ch^{-\frac{5}{6}} |t|^{\frac{1}{4}},$$

and it follows that

$$|I_1| = \mathcal{O}(h^{-\frac{2}{3}}) \int_{-\delta}^{-\lambda h^{2/3}} e^{-c|t|^{\frac{3}{2}}/h} dt = \mathcal{O}(e^{-c\lambda^{\frac{3}{2}}}).$$

Taking λ larger than $((3c)^{-1} |\ln h|)^{2/3}$, we get $I_1 = \mathcal{O}(h^{1/3})$.

On the other hand, using Proposition 5.1, A.2 and A.5 of [FMW] about the asymptotic properties of the solutions $u_{2,L}^-$ and $u_{1,L}^-$ near the crossing point, we have,

$$I_2 = 4h^{-\frac{2}{3}} \int_{-\lambda h^{2/3}}^0 r_1(t) \left(\frac{\xi_1'}{\xi_2'} \right)^{-\frac{1}{2}} \text{Ai}(-h^{-\frac{2}{3}} \xi_2) \times$$

$$\left(\left(\sin \frac{\mathcal{A}}{h} \right) \text{Ai}'(h^{-\frac{2}{3}} \xi_1) + \left(\cos \frac{\mathcal{A}}{h} \right) \text{Bi}'(h^{-\frac{2}{3}} \xi_1) \right) dt + \mathcal{O}(h^{\frac{1}{3}}).$$

Making the change of variable $y := h^{-2/3}t$ and $\rho := h^{-2/3}E$ as in [FMW], we obtain,

$$I_2 = 4r_1(0) \int_{-\lambda}^0 \text{Ai}(-y - \rho) \left\{ \left(\sin \frac{\mathcal{A}}{h} \right) \text{Ai}'(y - \rho) + \left(\cos \frac{\mathcal{A}}{h} \right) \text{Bi}'(y - \rho) \right\} dy$$

$$+ \mathcal{O}(h^{\frac{2}{3}} \lambda^2) + \mathcal{O}(h^{\frac{1}{3}}).$$

Then, using the behaviour of Airy functions at $\pm\infty$, this leads to,

$$I_2 = 4r_1(0) \int_{-\infty}^0 \text{Ai}(-y - \rho) \left\{ \left(\sin \frac{\mathcal{A}}{h} \right) \text{Ai}'(y - \rho) + \left(\cos \frac{\mathcal{A}}{h} \right) \text{Bi}'(y - \rho) \right\} dy$$

$$+ \mathcal{O}(e^{-c'\lambda^{\frac{3}{2}}}) + \mathcal{O}(h^{\frac{2}{3}} \lambda^2) + \mathcal{O}(h^{\frac{1}{3}}),$$

with $c' > 0$ constant. Hence, taking $\lambda := ((3c'')^{-1} |\ln h|)^{2/3}$ with $c'' = \min\{c, c'\}$, the error is $\mathcal{O}(h^{1/3})$.

The integral I_3 contains $u_{1,L}^-$ instead of its derivative compared with I_1 and I_2 . Then one easily see that $I_3 = \mathcal{O}(h^{2/3})$.

Thus the formula for $\alpha_{1,L}$ is obtained. \square

This proposition 3.2 together with (3.24) imply that there exists a bounded complex-valued function $G(E, h)$ of $E = \rho h^{2/3} \in \mathcal{D}_h(C_0)$ and h sufficiently

small such that $E \in \mathcal{D}_h(C_0)$ is a resonance of P if and only if,

$$(3.25) \quad \cos \frac{\mathcal{A}(E)}{h} = h^{\frac{4}{3}} \left(\sin \frac{\mathcal{A}(E)}{h} \right) G(E, h).$$

More precisely, from the fact $\sin^2(\mathcal{A}/h) = 1 + \mathcal{O}(h^{8/3})$, we obtain the following asymptotic formula for the imaginary part of $G(E, h)$: As $h \rightarrow 0$,

$$\operatorname{Im} G = \pi^2 r_1(0)^2 (\nu_{1,R}^A(\operatorname{Re} \rho) + \nu_{1,L}^A(\operatorname{Re} \rho)) (\nu_{2,R}^A(\operatorname{Re} \rho) + \nu_{2,L}^A(\operatorname{Re} \rho)) + \mathcal{O}(h^{\frac{1}{3}}).$$

Finally notice that the functions

$$(3.26) \quad \begin{aligned} \nu_{1,R}^A(t) + \nu_{1,L}^A(t) &= \int_{-\infty}^{\infty} \operatorname{Ai}'(\tau_1^{\frac{1}{3}}(y - \frac{t}{\tau_1})) \operatorname{Ai}(-\tau_2^{\frac{1}{3}}(y + \frac{t}{\tau_2})) dy, \\ \nu_{2,R}^A(t) + \nu_{2,L}^A(t) &= \int_{-\infty}^{\infty} \operatorname{Ai}(\tau_1^{\frac{1}{3}}(y - \frac{t}{\tau_1})) \operatorname{Ai}'(-\tau_2^{\frac{1}{3}}(y + \frac{t}{\tau_2})) dy, \end{aligned}$$

are both proportional to the derivative of the function

$$\int_{-\infty}^{\infty} \operatorname{Ai}(\tau_1^{\frac{1}{3}}(y - \frac{t}{\tau_1})) \operatorname{Ai}(-\tau_2^{\frac{1}{3}}(y + \frac{t}{\tau_2})) dy = (\tau_1 + \tau_2)^{-\frac{1}{3}} \operatorname{Ai}(-\tau_3^{-\frac{2}{3}} t),$$

with $\tau_3^{-1} = \tau_1^{-1} + \tau_2^{-1}$, and their product is given by $\frac{\tau_3^{\frac{1}{3}}}{\tau_1 + \tau_2} \left(\operatorname{Ai}'(-\tau_3^{-\frac{2}{3}} \lambda_k) \right)^2$.

Theorem 2.2 then follows from (3.25).

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